

Dark matter from dark gauge groups

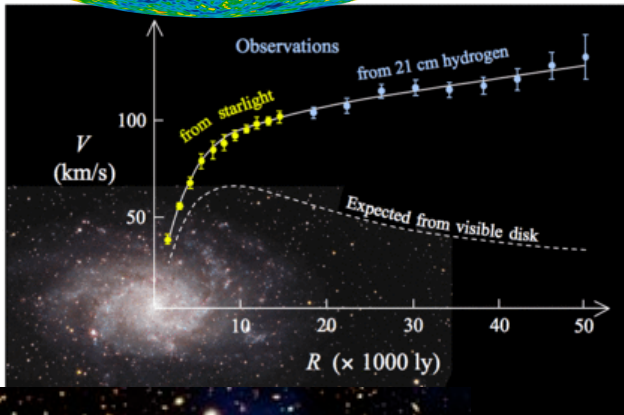
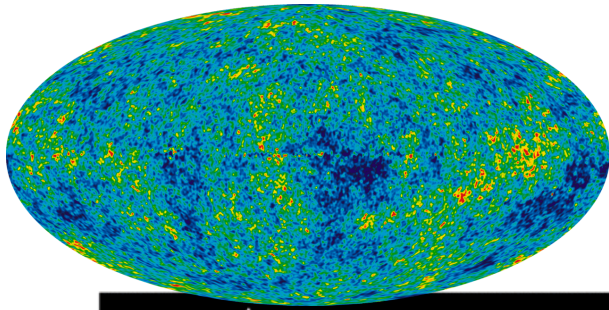
Christian Gross



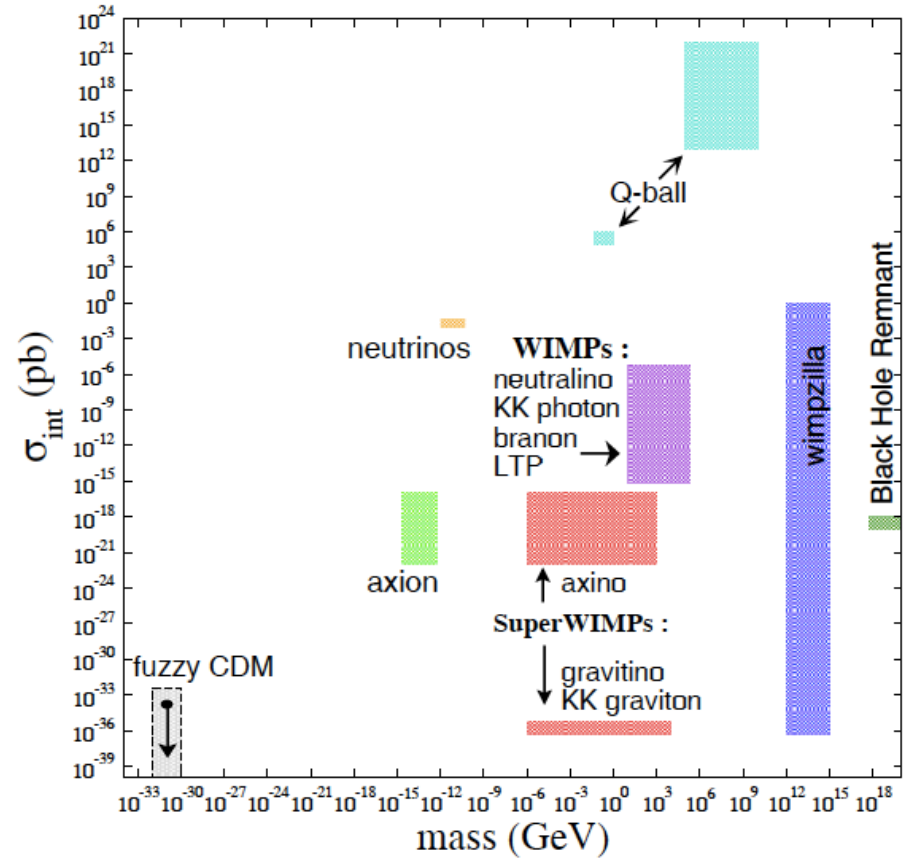
Planck 2017, Warsaw

*based on work in collaboration with G. Arcadi, O. Lebedev,
Y. Mambrini, S. Pokorski, T. Toma*

We know DM exists:



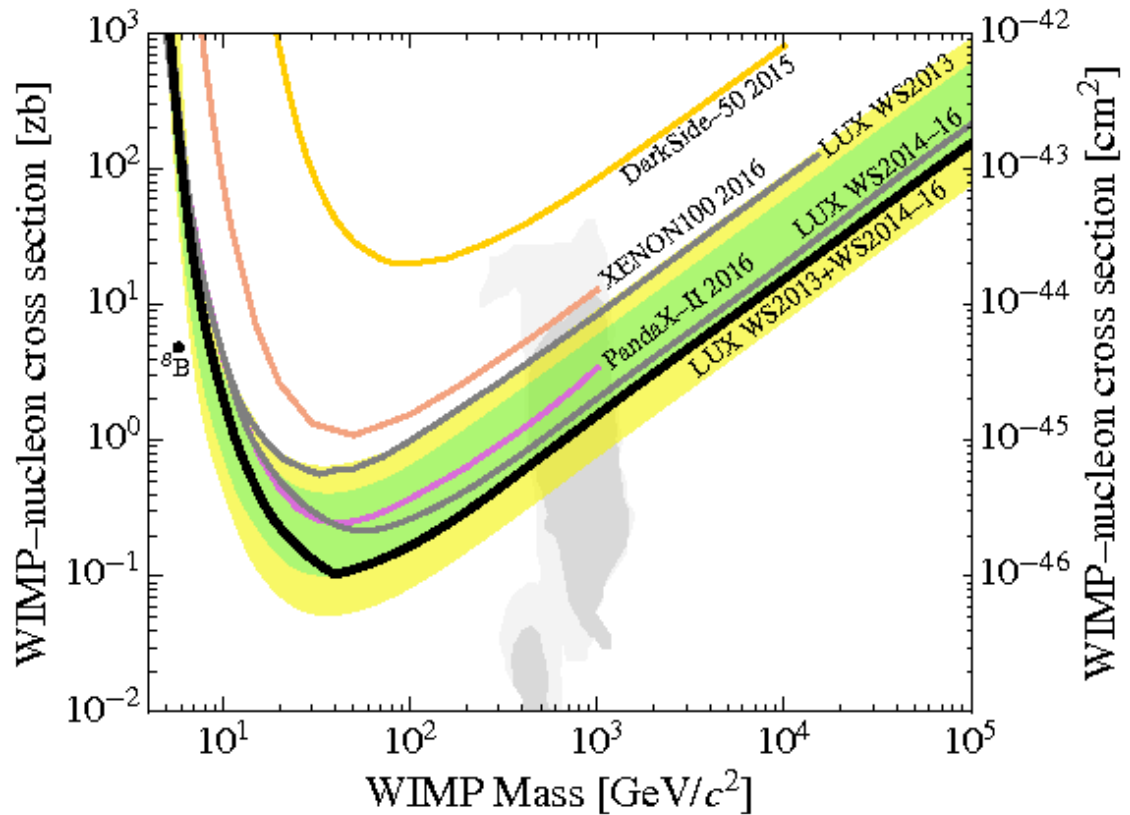
....but we have no clue
what kind of particle(s)
DM consists of:



[figure taken from E.-K. Park,
contribution to DMSAG
report, July 18, 2007]

why WIMPs are popular

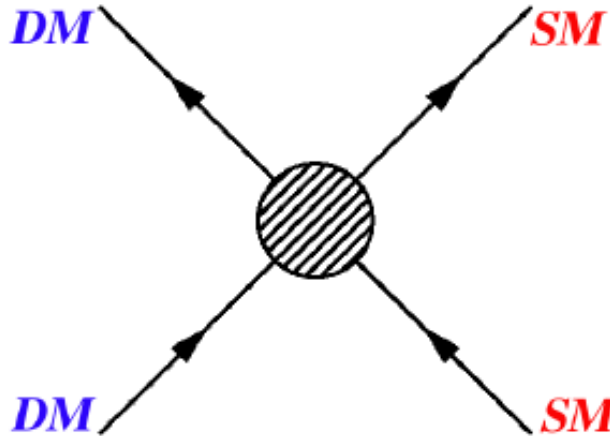
- The WIMP miracle
→ link to TeV scale BSM physics
- Huge efforts to search for WIMPs in direct detection experiments



[LUX, 1608.07648]

why simple WIMP DM models are under pressure

thermal freeze-out (early Univ.)
indirect detection (now)



production at colliders

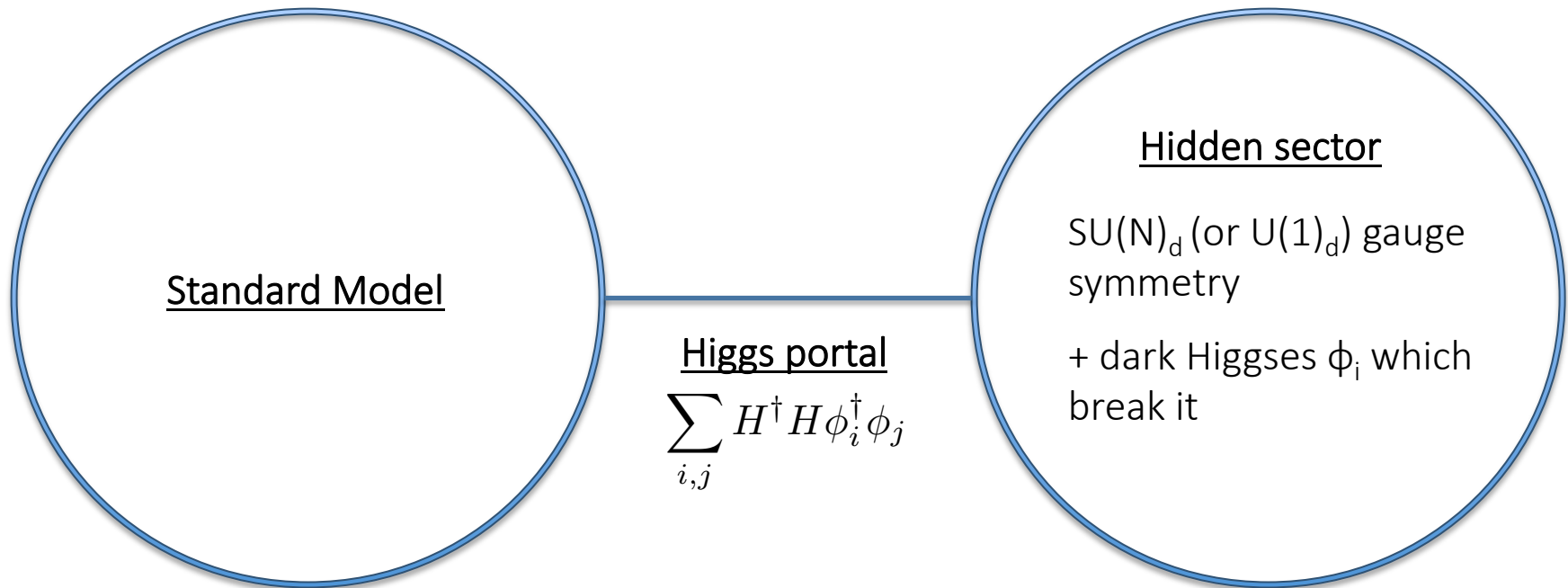


direct detection



- Reason for tension in simplest models:
 - DD limits require small coupling
 - small $\langle \sigma v \rangle$
 - WIMPs overabundant due to $\Omega \propto 1/\langle \sigma v \rangle$
- Need to break the 'crossing relation', e.g. by:
 - resonant DM annihilation
 - additional annihilation channels into the dark sector
 - cancellation among different direct detection diagrams
 - ...

WIMP DM from hidden gauge groups



plan for the rest of the talk:

- ① why the massive gauge fields are stable
- ② three ways to naturally reconcile direct detection limits and relic abundance

stability of $U(1)_d$ vector dark matter

$$\mathcal{L}_{\text{dark}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + |D_{\mu}\phi|^2 - V(\phi)$$

complex scalar

A_{μ} is stable due to charge-conjugation symmetry

$$Z_2: \begin{cases} \phi \rightarrow \phi^* \\ A_{\mu} \rightarrow -A_{\mu} \end{cases}$$

[Lebedev, Lee, Mambrini, 2011]

stability of $SU(2)_d$ vector dark matter

$$\mathcal{L}_{\text{dark}} = -\frac{1}{4} \sum_{a=1}^3 F_{\mu\nu}^a F^{a\mu\nu} + |D_\mu \phi|^2 - V(\phi)$$

↖ $SU(2)_d$ doublet

$A_\mu^a \rightarrow -A_\mu^a$ is not a symmetry,
due to triple gauge boson vertex

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- since the generators are $\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$,

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which can be shown to be a remnant of $SU(N)_d$ gauge symmetry

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[CG, Lebedev, Mambrini, 2015]

- this $Z_2 \times Z_2$ ensures stability of the three A_μ^a ;
it enlarges to a custodial $SO(3) \rightarrow$ the A_μ^a have same masses

[Hambye, 2008]

stability of $SU(N \geq 3)_d$ vector dark matter

[CG, Lebedev, Mambrini, 2015]

- $SU(N \geq 3)_d$ broken completely by $N-1$ dark Higgs N -plets
- $N(N-1)/2$ physical CP-even scalars
 $N(N-3)/2+1$ " CP-odd "
- assuming CP invariance, obtain $Z_2 \times Z_2$ symmetry,
as in the $SU(2)_d$ case
($Z_2 \times Z_2$ is part of a global $U(1) \times Z_2$)
- DM: two degenerate vectors + $\left\{ \begin{array}{l} \text{one slightly lighter vector} \\ \text{or} \\ \text{one CP-odd scalar} \end{array} \right.$

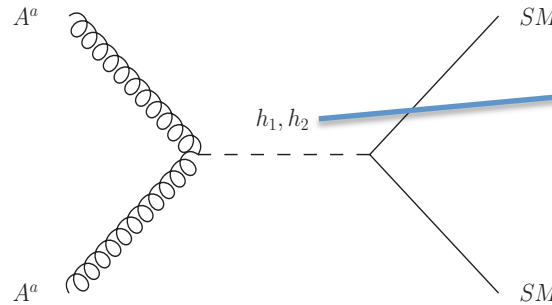
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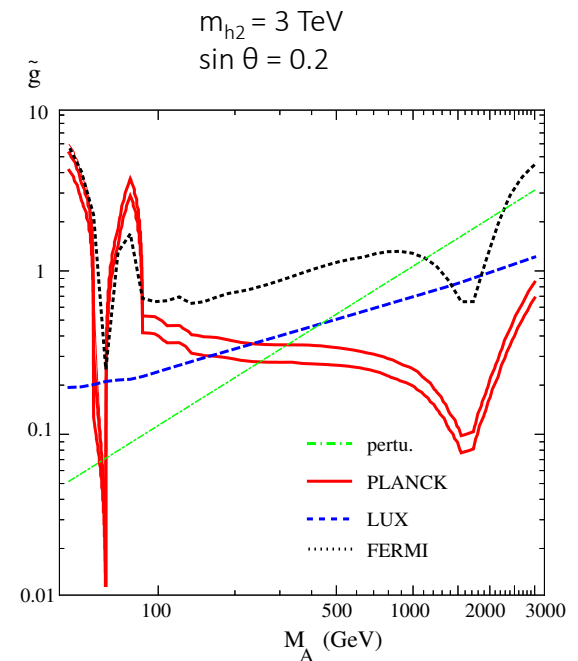
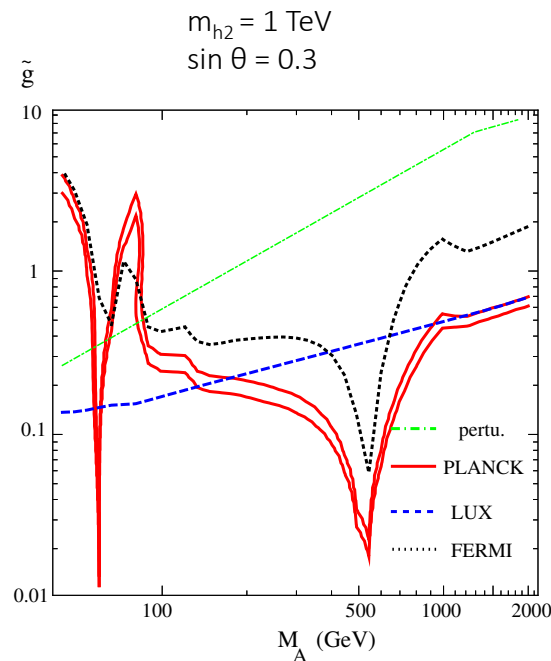
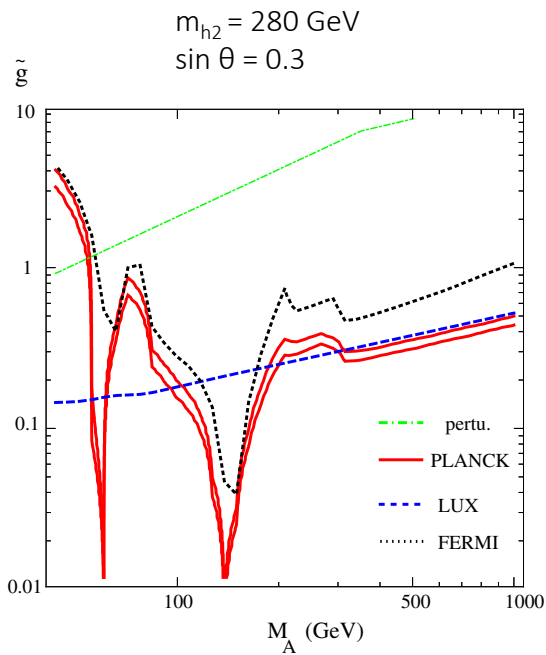
a) DM annihilation via (broad) resonances

[CG, Lebedev, Mambrini, 2015]

DM annihilation
via s-channel Higgses:

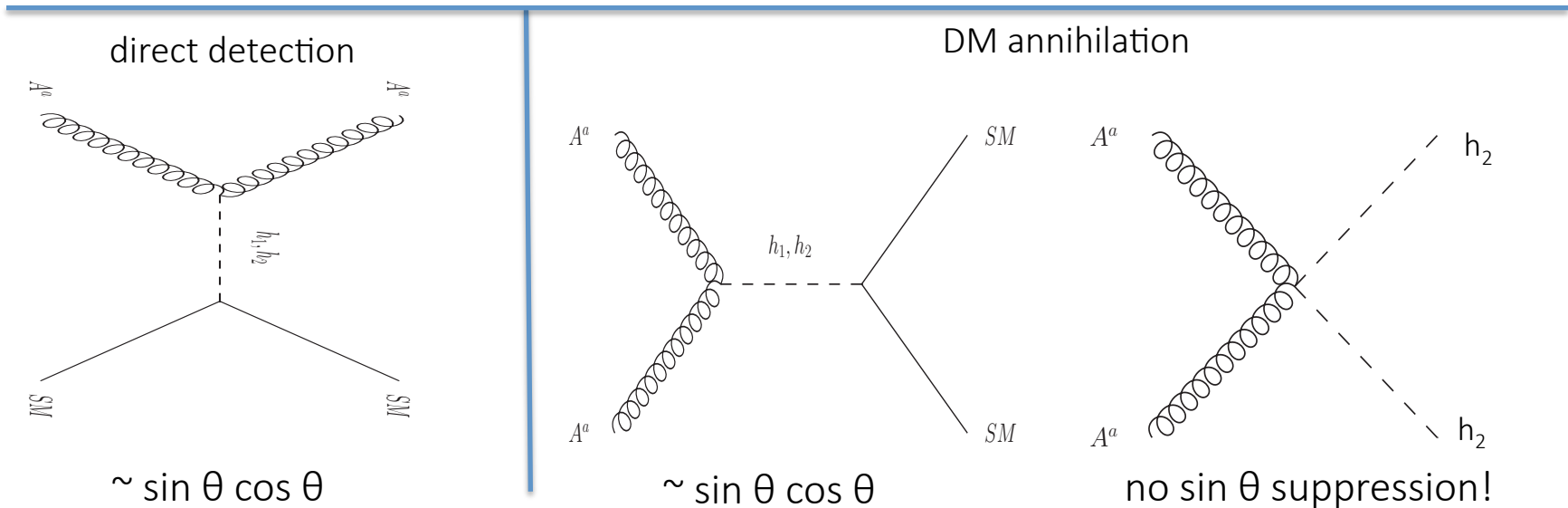


h_1, h_2 : mass eigenstates
of dark Higgs and
 $SU(2)_L$ -Higgs,
with mixing $\sin \theta$



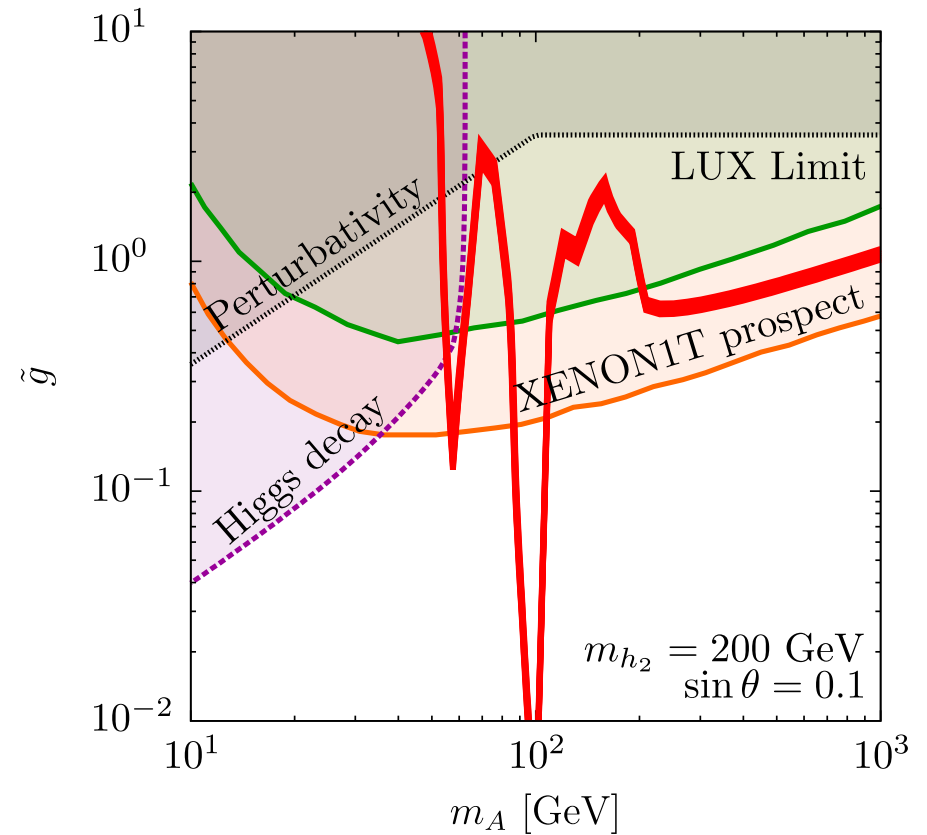
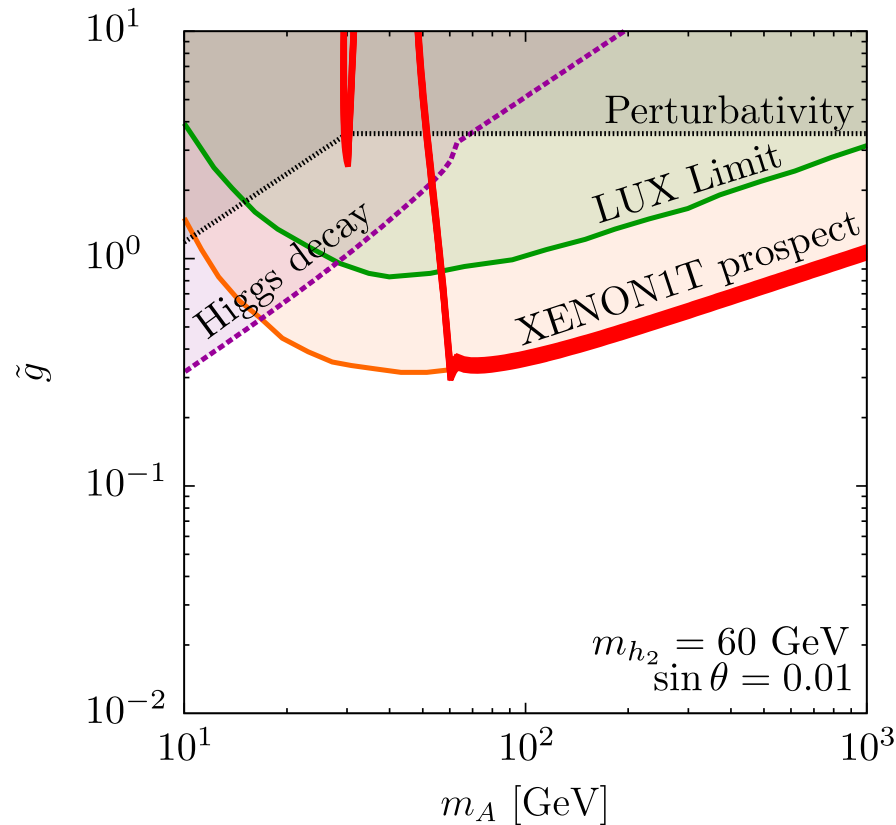
b) DM annihilation mostly into hidden sector

- The basic mechanism is known as secluded DM: [Pospelov, Ritz, Voloshin, 2007]
DM annihilation into dark sector states may provide an efficient extra annihilation channel. This breaks the correlation between DD and annihilation cross section.
- The scalar sector of spontaneously broken hidden gauge groups automatically contains such dark sector states into which DM may annihilate: the hidden sector scalars.



b) DM annihilation mostly into hidden sector

[Arcadi,CG,Lebedev,Pokorski,Toma; 2016]

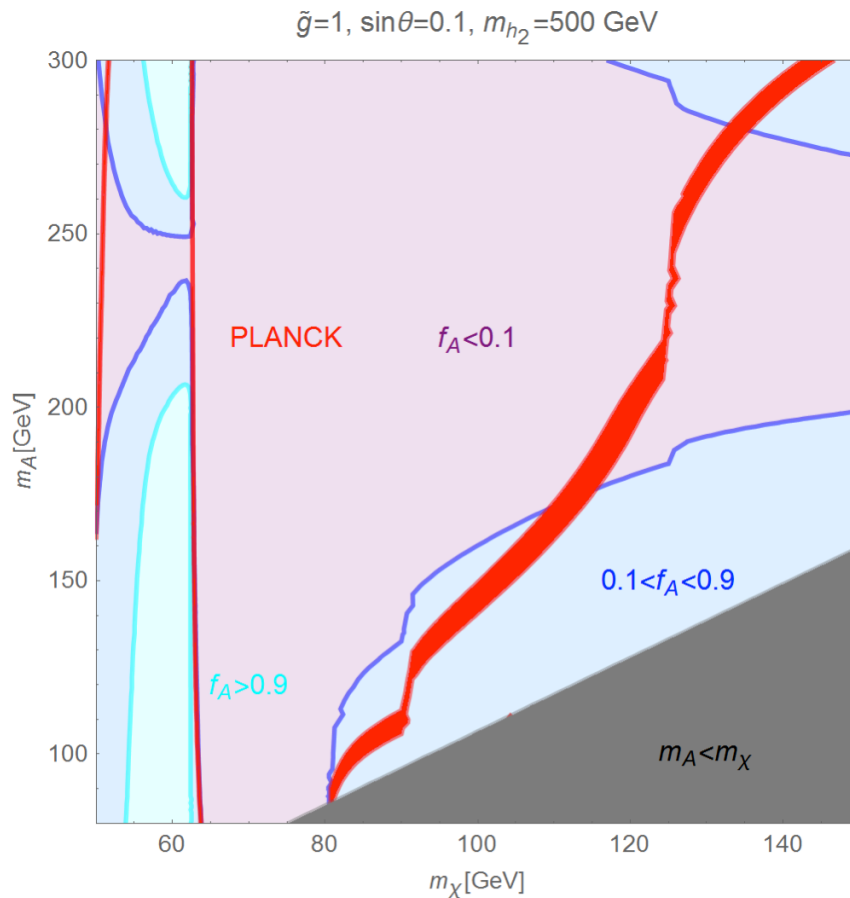


- The coupling required to obtain the correct relic abundance drops dramatically as soon as the dark annihilation channel opens up.
- The smaller $\sin \theta$ is, the stronger this effect is.

c) cancellation among direct detection diagrams

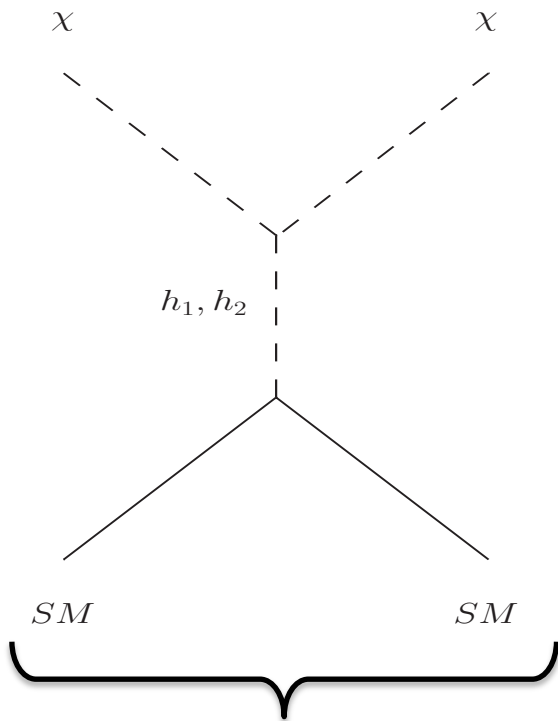
[Arcadi,CG,Lebedev,Mambrini,Pokorski,Toma; 2016]

- for $SU(N \geq 3)_d$ DM can consist of vectors and a CP-odd scalar χ
- Both components may be dominant, depending on parameters:

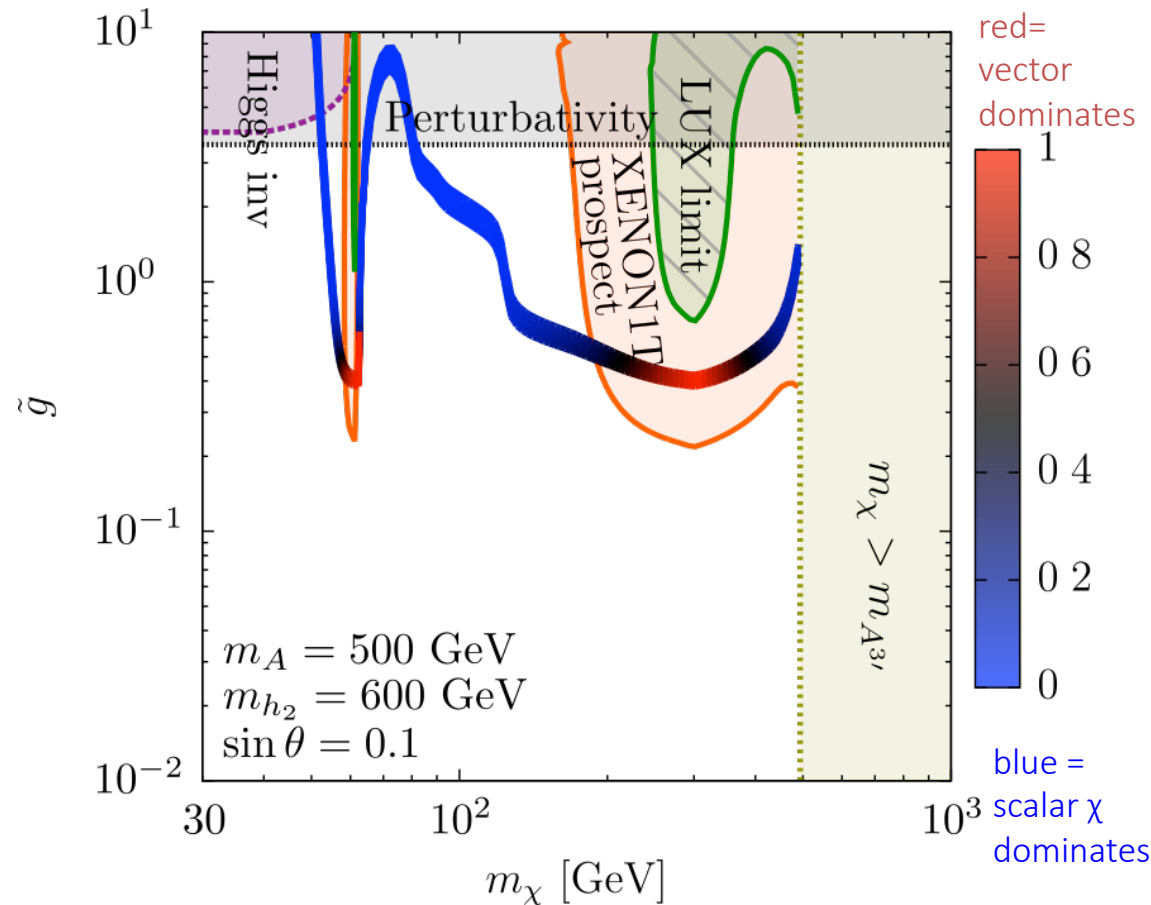


- shown: relative contribution of vector DM to total DM:
 $f_A = \Omega_A / \Omega_{\text{total}}$
- since, very roughly,
 $\Omega_{\text{total}} \propto 1/\langle\sigma v\rangle_A + 1/\langle\sigma v\rangle_\chi$
component with smaller $\langle\sigma v\rangle$ dominates
- this explains e.g. behaviour when one of the components annihilates resonantly

c) cancellation among direct detection diagrams



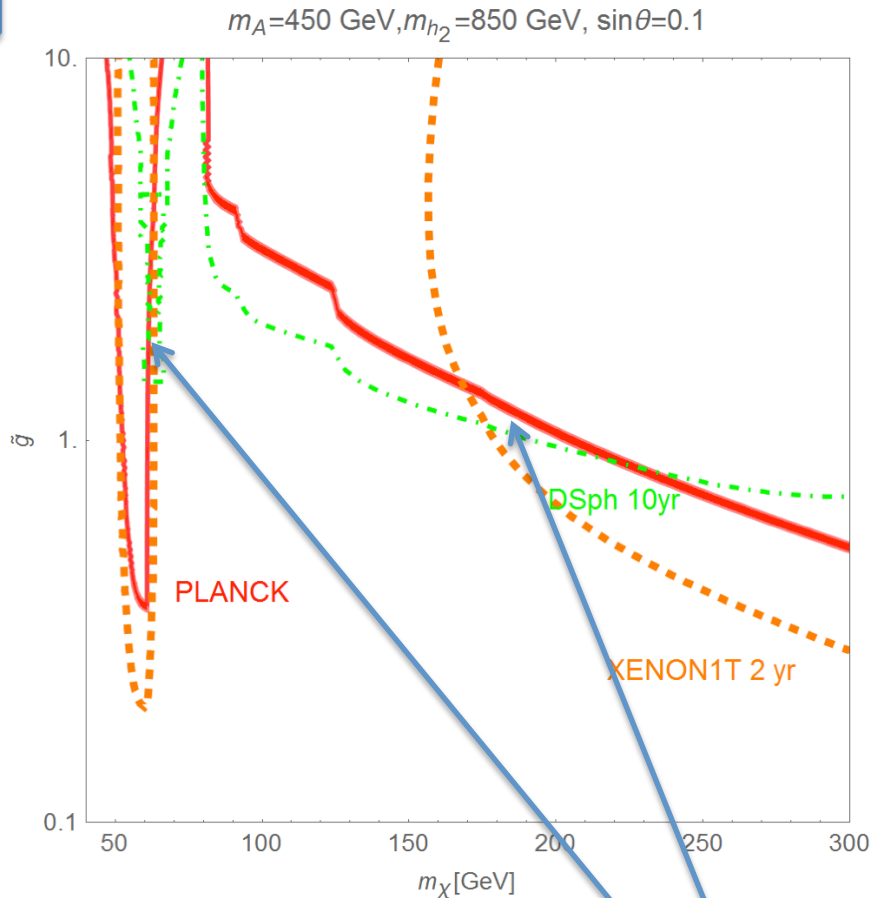
couplings lead to exact tree-level cancellation in scattering of χ -DM on nuclei, independently of the mass of h_2



→ WIMPs can be completely invisible in direct detection

side remark:

in the mixed scalar-vector DM case, one might be able to see one DM component in indirect detection and the other in direct detection



basic point:

- χ is invisible in DD, but might be observable in ID (if its DM fraction is large enough)
- on the other hand: vector component could be visible in DD, but is hardly visible in ID because main annihilation (often) into dark sector ($A A \rightarrow \chi \chi$)

two regions where testing this scenario might be possible with future data

Summary

- Gauge fields of a spontaneously broken dark $SU(N)_d$ are viable DM candidates
- Stability of DM is due to a $Z_2 \times Z_2$ symmetry that automatically arises for minimal CP-conserving Higgs sectors
- Several ways to reconcile relic abundance and direct detection limits:
 - annihilation via (broad) resonances
 - annihilation dominantly into dark sector
 - cancellation among direct detection diagram