
Freeze-in at stronger coupling

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- *particle production during and after inflation*
- *Planck-suppressed operators*
- *freeze-in dark matter*

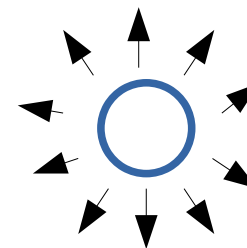
Main references : [2306.13061](#) [hep-ph] Cosme, Costa, OL

[2210.02293](#) [hep-ph] OL

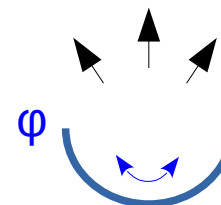
Non-thermal relics / DM have **memory** !

Production mechanisms (*all add up*):

- during inflation



- via inflaton oscillations



- inflaton decay



- thermal emission (freeze-in)



Decoupled scalar production during inflation

Scalar “s” with

$$V(s) = \frac{1}{2}m_s^2 s^2 + \frac{1}{4}\lambda_s s^4$$

$$\lambda_s \ll 1, \quad m_s \ll H$$

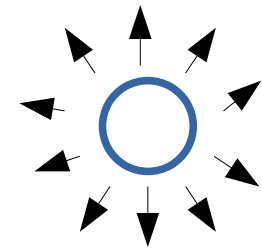
Starobinsky-Yokoyama equilibrium distribution of de Sitter fluctuations:

$$P(s) \propto \exp \left[-8\pi^2 V(s) / (3H^4) \right]$$

$$\langle s^2 \rangle \simeq 0.1 \times \frac{H_{\text{end}}^2}{\sqrt{\lambda_s}}$$

Mean field:

$$\bar{s} \equiv \sqrt{\langle s^2 \rangle}$$



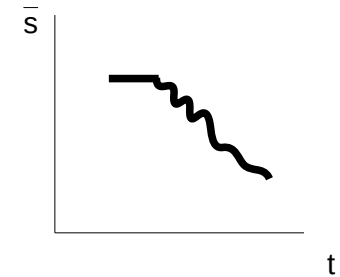
scalar
fluctuation
generation

Effective mass:

$$m_{\text{eff}}^2 = m_s^2 + 3\lambda_s \bar{s}^2$$

Evolution:

$$\bar{s}_{\text{end}} \xrightarrow{a^0} \bar{s}_{\text{osc}} \xrightarrow{a^{-1}} \bar{s}_{\text{dust}}$$



frozen \rightarrow oscillates in s^4 potential \rightarrow oscillates in s^2 potential

$$H > m_{\text{eff}}$$

$$H \sim m_{\text{eff}}$$

$$m_s \sim m_{\text{eff}}$$

Relic number density (*non-rel.*) = energy density / particle mass :

$$n \simeq m_s^3 / \lambda_s$$

Constraints

Require

$$Y \leq 4.4 \times 10^{-10} \frac{\text{GeV}}{m_s}$$

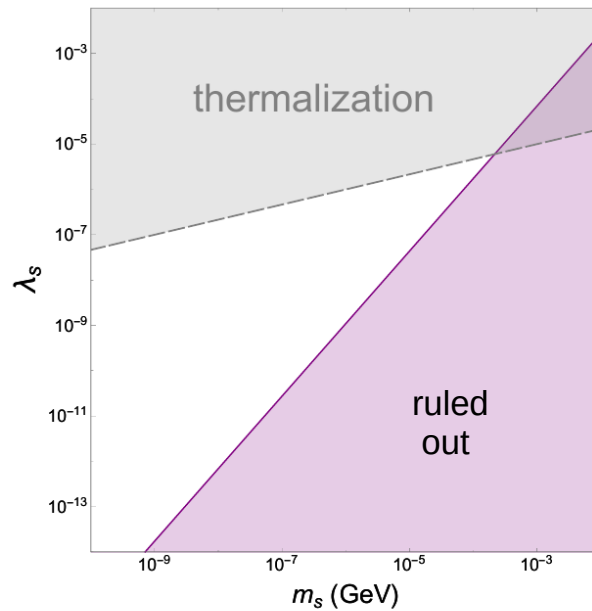
instant reheating
or φ^4



$$m_s \lambda_s^{-5/8} \lesssim 10^{-7} \left(\frac{M_{\text{Pl}}}{H_{\text{end}}} \right)^{3/2} \text{ GeV}$$

Hubble rate at the end of inflation

Very strong constraint :



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Starobinsky-Yokoyama
fluctuations

$$H_{\text{end}} \sim 10^{14} \text{ GeV}$$

$$\Delta_{\text{NR}} = 1$$

In general, the abundance depends on duration of the *non-relativistic* expansion period (ϕ^2 pot.):

$$H_{\text{end}} \xrightarrow{a^{-3/2}} H_{\text{reh}} \quad \Delta_{\text{NR}} \equiv \left(\frac{H_{\text{end}}}{H_{\text{reh}}} \right)^{1/2} > 1$$

Dilutes the energy in the condensate \rightarrow weaker constraint

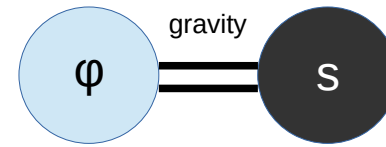
$$m_s \lambda_s^{-3/4} \lesssim 10^{-7} \Delta_{\text{NR}} \left(\frac{M_{\text{Pl}}}{H_{\text{end}}} \right)^{3/2} \text{ GeV}$$

$$H_{\text{end}} \sim 10^{14} \text{ GeV} \quad \rightarrow \quad m_s \ll \Delta_{\text{NR}} \text{ GeV}$$

Only particles far below the GeV scale are allowed for $\Delta_{\text{NR}} = 1$

Quantum gravity effects

Induce gauge invariant operators
(with unknown coefficients)



Dim-6 gravity-induced couplings:

Also induced by **classical** gravity!

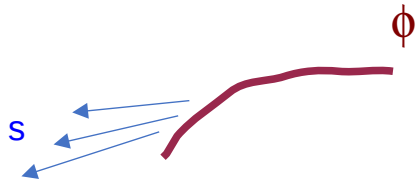
$$\Delta\mathcal{L}_6 = \frac{C_1}{M_{\text{Pl}}^2} (\partial_\mu\phi)^2 s^2 + \frac{C_2}{M_{\text{Pl}}^2} (\phi\partial_\mu\phi)(s\partial^\mu s) + \frac{C_3}{M_{\text{Pl}}^2} (\partial_\mu s)^2 \phi^2 - \frac{C_4}{M_{\text{Pl}}^2} \phi^4 s^2 - \frac{C_5}{M_{\text{Pl}}^2} \phi^2 s^4$$

Main operators for on-shell fields contributing to s-pair production:

$$\mathcal{O}_3 = \frac{1}{M_{\text{Pl}}^2} (\partial_\mu s)^2 \phi^2 \quad , \quad \mathcal{O}_4 = \frac{1}{M_{\text{Pl}}^2} \phi^4 s^2$$

(supplemented with dim-4 $\mathcal{O}_{\text{renorm}} = \frac{m_\phi^2}{M_{\text{Pl}}^2} \phi^2 s^2$ and 4-DM op $\frac{C_5}{M_{\text{Pl}}^2} \phi^2 s^4$)

Particle production:



\mathcal{O}_4 dominates

$$\Gamma = \frac{C_4^2}{4\pi M_{\text{Pl}}^4} \sum_{n=1}^{\infty} |\hat{\zeta}_n|^2 \quad \hat{\zeta}_n = \sum_{m=-\infty}^{\infty} \zeta_{n-m} \zeta_m$$

$$\dot{n} + 3Hn = 2\Gamma$$

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$$\Delta_{\text{NR}} \equiv \left(\frac{H_{\text{end}}}{H_{\text{reh}}} \right)^{1/2}$$

$$|C_4| < 10^{-3} \Delta_{\text{NR}}^{1/2} \frac{H_{\text{end}}^{5/4} M_{\text{Pl}}^{11/4}}{\phi_0^4} \sqrt{\frac{\text{GeV}}{m_s}}$$

$$\phi_0 \sim M_{\text{Pl}} \text{ and } H_{\text{end}} \sim 10^{14} \text{ GeV}$$



$$|C_4| < \text{few} \times 10^{-9} \Delta_{\text{NR}}^{1/2} \sqrt{\frac{\text{GeV}}{m_s}}$$

$$|C_3| \lesssim 10^{-1} \Delta_{\text{NR}}^{1/2} \sqrt{\frac{\text{GeV}}{m_s}}$$

Higher dim operators:

$$\mathcal{O}^{(p)} = \frac{\phi^p s^2}{M_{\text{Pl}}^{p-2}}$$

$$|C^{(p)}| < 10^{-3} \Delta_{\text{NR}}^{1/2} \frac{H_{\text{end}}^{5/4} M_{\text{Pl}}^{p-5/4}}{\phi_0^p} \sqrt{\frac{\text{GeV}}{m_s}}$$



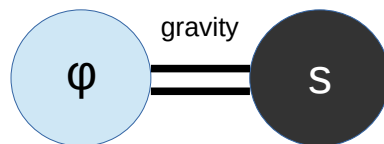
Planck-suppressed operators are very efficient in particle production!

$$\frac{\phi^4 s^2}{M_{\text{Pl}}^2}, \quad \frac{\phi^6 s^2}{M_{\text{Pl}}^4}, \quad \frac{\phi^8 s^2}{M_{\text{Pl}}^6}, \quad \dots$$

Main observation :

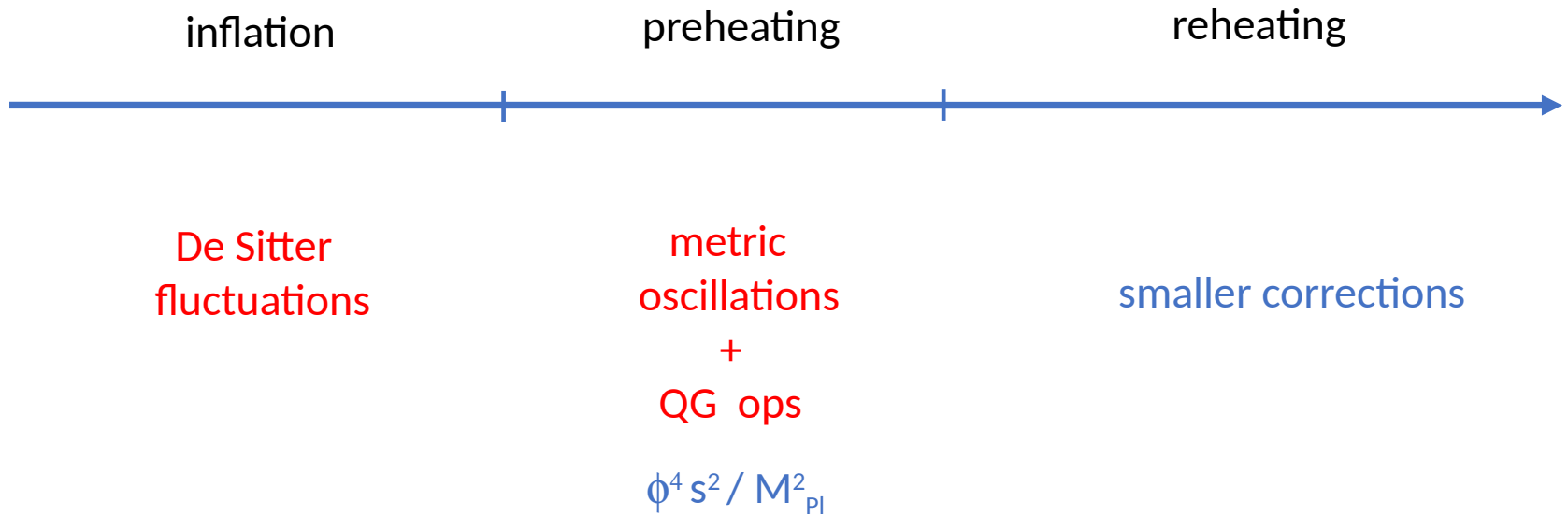
*Planck—suppressed (“gravity--induced”) operators
with small Wilson coefficients
can account for all of the dark matter !*

Non-thermal DM model building is highly **UV sensitive** :



- abundance is additive (“memory”)
- need to control quantum gravity
- **predictivity ?**

Irreducible gravity background for Freeze-in :



The problem is not to produce DM, but to get rid of it !

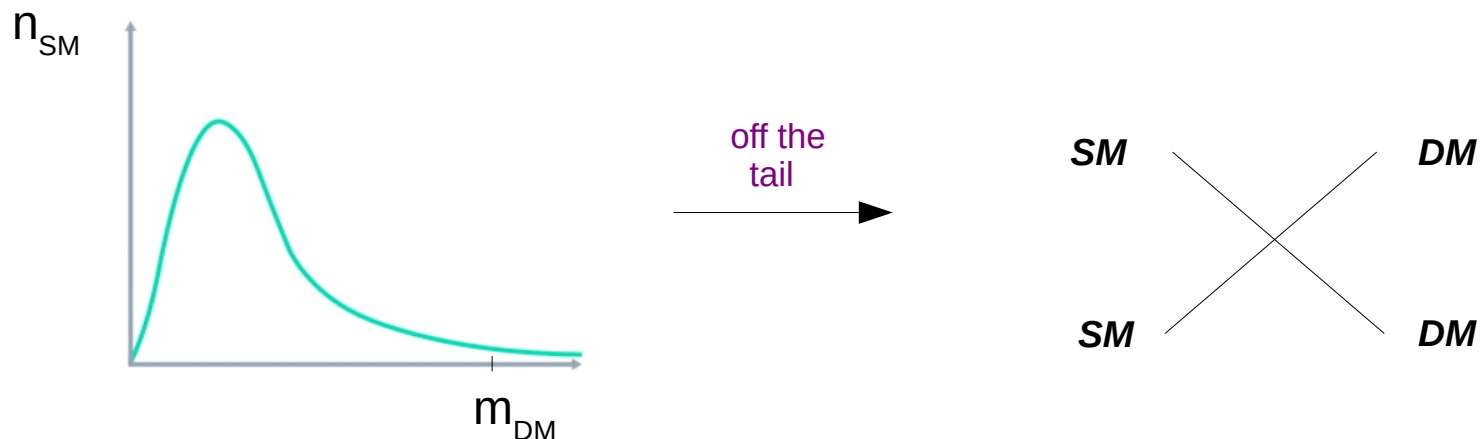
Freeze-in at stronger coupling

Need to get rid of the gravitational background → matter dominated expansion
 ϕ^2 local inflaton potential



relics are diluted, low reheating temperature T_R

What if $T_R < m_{DM}$?



Boltzmann-suppressed DM production requires a stronger coupling → observable !

Simplest model = Higgs portal DM

$$V(s) = \frac{1}{2} \lambda_{hs} s^2 H^\dagger H + \frac{1}{2} m_s^2 s^2$$

Boltzmann equation:

$$\dot{n} + 3Hn = \Gamma(h_i h_i \rightarrow ss) - \Gamma(ss \rightarrow h_i h_i)$$

No annihilation:

$$\Gamma(h_i h_i \rightarrow ss) \simeq \frac{\lambda_{hs}^2 T^3 m_s}{27 \pi^4} e^{-2m_s/T} \quad \rightarrow \quad \lambda_{hs} \simeq 3 \times 10^{-11} e^{m_s/T_R} \sqrt{\frac{T_R}{m_s}}$$

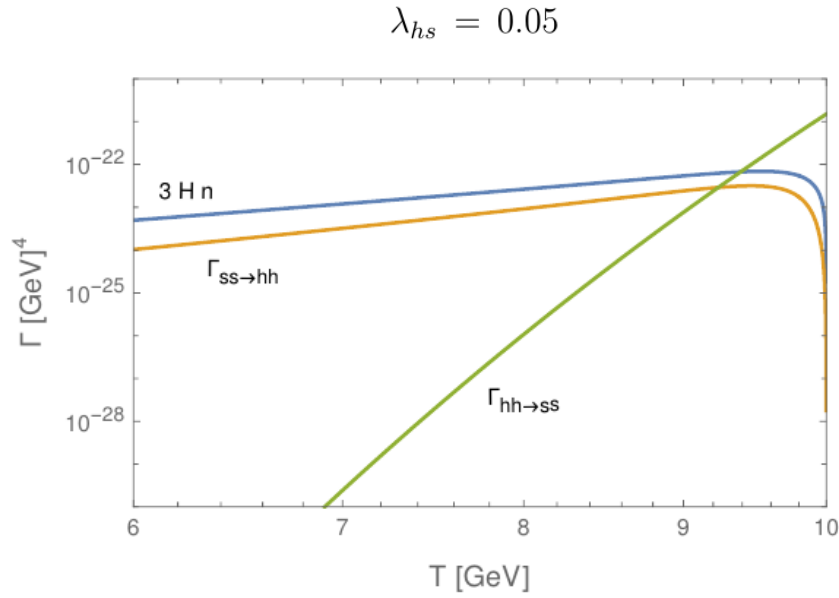
With annihilation:

$$\Gamma(ss \rightarrow h_i h_i) = \sigma(ss \rightarrow h_i h_i) v_r n^2, \quad \sigma(ss \rightarrow h_i h_i) v_r = 4 \times \frac{\lambda_{hs}^2}{64 \pi m_s^2}$$

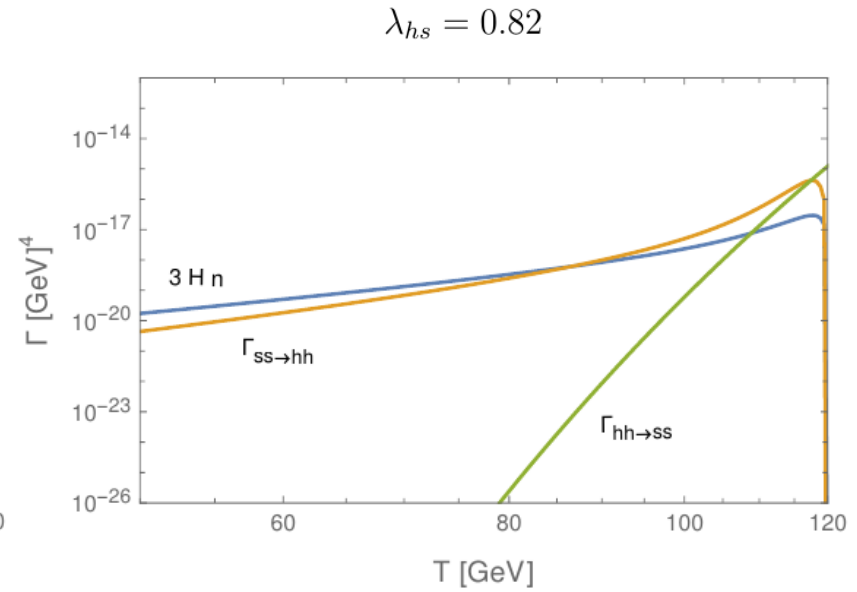
No thermalization:

$$\Gamma(h_i h_i \rightarrow ss) \neq \Gamma(ss \rightarrow h_i h_i)$$

Reaction rates:

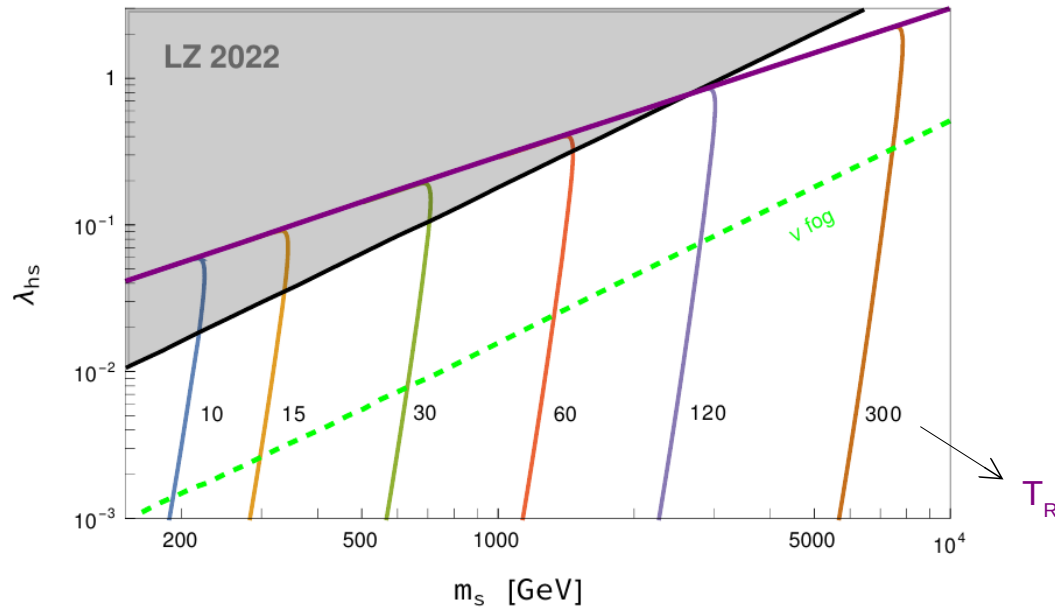


$$T_R = 10 \text{ GeV}, m_s = 223.5 \text{ GeV}$$



$$T_R = 120 \text{ GeV}, m_s = 2.81 \text{ TeV}$$

Correct relic density (solid lines):



Already probed
by direct detection!

As coupling \uparrow
FI \rightarrow FO

CONCLUSION

- *gravitational DM production = strong background*
- *low T_R dilutes gravitational relics*
- *freeze-in for $m_{DM} > T_R$ requires a significant coupling*
- *freeze-in already probed by direct detection (LZ2022)*
- *further probes by XENONnT, DARWIN*