

Global effective field theory for top physics at lepton colliders

Gauthier Durieux
(DESY Hamburg)

preliminary results
based on current work with Cen Zhang (Brookhaven),
Martín Perelló, Marcel Vos (Valencia, ATLAS/ILC)

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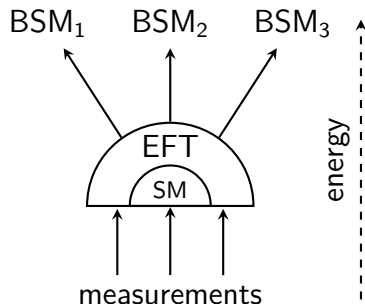


Introduction

The new physics effective field theory (aka SM EFT)

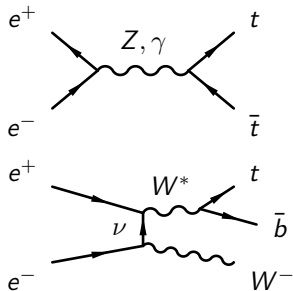
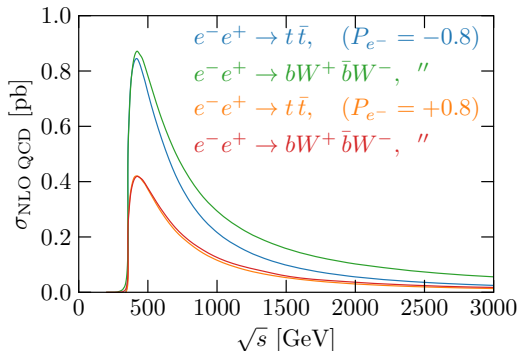
provides a systematic parametrization of the theory space
in direct vicinity of the SM

- ▶ in a low-energy limit
- ▶ through a proper QFT
- ▶ fully general when global



Aiming at a global EFT analysis of $e^+e^- \rightarrow t\bar{t} \rightarrow bW^+ \bar{b}W^-$

- Including four-fermion operators, notably
- Examining the impact of – NLO QCD corrections
 - off-shell top effect
- Studying various observables
 - + the impact of $\sqrt{s}$ and beam polarization



Up-sector EFT

[Grzadkowski et al '10]

Two-quark operators:

Scalar: $O_{u\varphi} \equiv \bar{q}u \tilde{\varphi} \varphi^\dagger \varphi,$

Vector: $O_{\varphi q}^1 \equiv \bar{q}\gamma^\mu q \varphi^\dagger \overleftrightarrow{D}_\mu \varphi \equiv O_{\varphi q}^+ + O_{\varphi q}^V - O_{\varphi q}^A,$

$O_{\varphi q}^3 \equiv \bar{q}\gamma^\mu \tau^I q \varphi^\dagger \overleftrightarrow{D}_\mu^I \varphi \equiv O_{\varphi q}^+ - O_{\varphi q}^V + O_{\varphi q}^A$ (CC also)

$O_{\varphi u} \equiv \bar{u}\gamma^\mu u \varphi^\dagger \overleftrightarrow{D}_\mu \varphi \equiv O_{\varphi q}^V + O_{\varphi q}^A$

$O_{\varphi ud} \equiv \bar{u}\gamma^\mu d \tilde{\varphi}^\dagger \overleftrightarrow{D}_\mu \varphi,$ (CC only, m_b int.)

Tensor: $O_{uB} \equiv \bar{q}\sigma^{\mu\nu} u \tilde{\varphi} B_{\mu\nu}, \equiv O_{uA} - \tan\theta_W O_{uZ}$

$O_{uW} \equiv \bar{q}\sigma^{\mu\nu} \tau^I u \tilde{\varphi} W_{\mu\nu}^I, \equiv O_{uA} + \cotan\theta_W O_{uZ}$ (CC also)

$O_{dW} \equiv \bar{q}\sigma^{\mu\nu} \tau^I d \tilde{\varphi} W_{\mu\nu}^I,$ (CC only, m_b int.)

$O_{uG} \equiv \bar{q}\sigma^{\mu\nu} T^A u \tilde{\varphi} G_{\mu\nu}^A.$ (NLO only)

Two-quark–two-lepton operators:

Scalar: $O_{lequ}^S \equiv \bar{l}e \varepsilon \bar{q}u,$ (CC also, m_e int.)

$O_{ledq} \equiv \bar{l}e \bar{d}q,$ (CC only, m_e int.)

Vector: $O_{lq}^1 \equiv \bar{l}\gamma_\mu l \bar{q}\gamma^\mu q \equiv O_{lq}^+ + O_{lq}^V - O_{lq}^A,$

$O_{lq}^3 \equiv \bar{l}\gamma_\mu \tau^I l \bar{q}\gamma^\mu \tau^I q \equiv O_{lq}^+ - O_{lq}^V + O_{lq}^A,$ (CC also)

$O_{lu} \equiv \bar{l}\gamma_\mu l \bar{u}\gamma^\mu u \equiv O_{lq}^V + O_{lq}^A,$

$O_{eq} \equiv \bar{e}\gamma^\mu e \bar{q}\gamma_\mu q \equiv O_{eq}^V - O_{eq}^A,$

$O_{eu} \equiv \bar{e}\gamma_\mu e \bar{u}\gamma^\mu u \equiv O_{eq}^V + O_{eq}^A,$

Tensor: $O_{lequ}^T \equiv \bar{l}\sigma_{\mu\nu} e \varepsilon \bar{q}\sigma^{\mu\nu} u.$ (CC also, m_e int.)

Anomalous vertices

$$t\bar{t}\gamma : \quad \gamma_\mu \overbrace{(F_{1V}^\gamma + \gamma_5 F_{1A}^\gamma)}^{\sim \phi} + \frac{\sigma_{\mu\nu} i q^\nu}{2m_t} \overbrace{(F_{2V}^\gamma + i\gamma_5 F_{2A}^\gamma)}^{\sim \text{Re,Im}\{C_{uA}\}}$$

$$t\bar{t}Z : \quad \gamma_\mu \overbrace{(F_{1V}^Z + \gamma_5 F_{1A}^Z)}^{\sim C_{\varphi q}^V, C_{\varphi q}^A} + \frac{\sigma_{\mu\nu} i q^\nu}{2m_t} \overbrace{(F_{2V}^Z + i\gamma_5 F_{2A}^Z)}^{\sim \text{Re,Im}\{C_{uZ}\}}$$

$$t\bar{b}W : \quad \gamma_\mu \overbrace{(F_{1V}^W + \gamma_5 F_{1A}^W)}^{\sim C_{\varphi ud}, C_{\varphi q}^+ - \frac{1}{2}(C_{\varphi q}^V - C_{\varphi q}^A)} + \frac{\sigma_{\mu\nu} i q^\nu}{2m_t} \overbrace{(F_{2V}^W + i\gamma_5 F_{2A}^W)}^{\sim s_W^2 C_{uA} + s_W c_W C_{uZ} \pm C_{dW}^*}$$

Insufficiencies:

- Miss four-fermion operators
- Conflict with gauge invariance
Do not allow for radiative corrections to be computed
- Complex couplings where the tree-level EFT prescribes real ones
- Hide correlations induced by gauge invariance
Preclude the combination of measurements in various sectors

Two ideas

for global EFT analyses

Statistically optimal observables

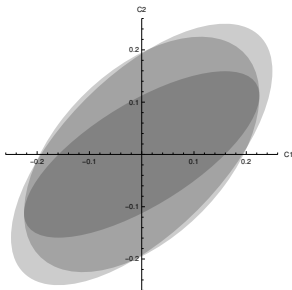
[Atwood, Soni '92]

[Diehl, Nachtmann '94]

minimize the one-sigma ellipsoid in EFT parameter space.

(joint efficient set of estimators, saturating the Rao-Cramér-Fréchet bound: $V^{-1} = I$)

For small C_i , with a phase-space distribution $\sigma(\Phi) = \sigma_0(\Phi) + \sum_i C_i \sigma_i(\Phi)$, the statistically optimal set of observables is: $O_i(\Phi) = \sigma_i(\Phi)/\sigma_0(\Phi)$.



e.g. $\sigma(\phi) = 1 + \cos(\phi) + C_1 \sin(\phi) + C_2 \sin(2\phi)$

1. asymmetries: $O_i \sim \text{sign}\{\sin(i\phi)\}$

2. moments: $O_i \sim \sin(i\phi)$

3. statistically optimal: $O_i \sim \frac{\sin(i\phi)}{1 + \cos \phi}$

$\Rightarrow$ area ratios 1.9 : 1.7 : 1

Previous applications in $e^+e^- \rightarrow t\bar{t}$:

[Grzadkowski, Hioki '00] [Janot '15] [Khiem et al '15]

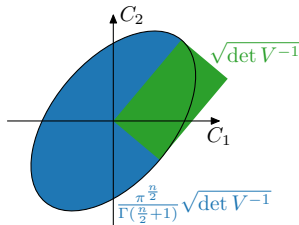
Global determinant parameter (GDP)

[GD, Grojean, Gu, Wang, '17]

In a n -dimensional Gaussian fit,
with covariance matrix V ,

$$\text{GDP} \equiv \sqrt[n]{\det V^{-1}}$$

provides a geometric average
of the constraints strengths.



Interestingly, GDP ratios are operator-basis independent!

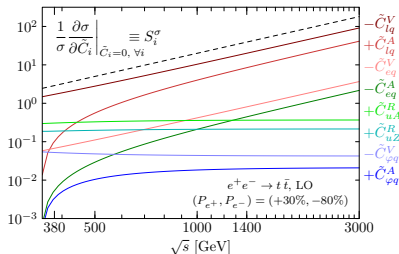
- as the volume scales linearly with coefficient normalization
- as the volume is invariant under rotations

$\Rightarrow$ conveniently assess constraint strengthening.

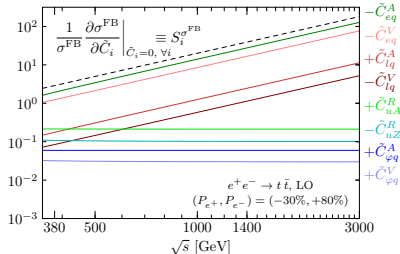
Operator sensitivities

σ and A^{FB} sensitivities

Total cross section (left pol.):



FB-integrated cross section (right pol.):



Few features:

- quadratic energy growth for four-fermion operators
- no growth for two-fermion operators (dipoles included)
- p -wave $\beta = \sqrt{1 - 4m_t^2/s}$ suppression of axial vectors at threshold
- enhanced sensitivity of axial vector operators in σ^{FB}
- sensitivity sign flip for $C_{\varphi q}^V$ and C_{uZ}^R when polarization is reversed
- etc.

Helicity amplitude decomposition in $bW^+\bar{b}W^-$

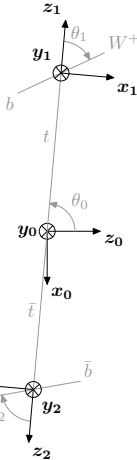
[Jacob,Wick '59]

Production amplitudes: $++ : A_1 \sim \frac{2m_t}{\sqrt{s}} V + \sqrt{s} (D - \beta\tilde{D})$
 $-- : A_2 \sim \frac{2m_t}{\sqrt{s}} V + \sqrt{s} (D + \beta\tilde{D})$
 $+- : A_3 \sim (V + \beta A) + 2m_t D$
 $-+ : A_4 \sim (V - \beta A) + 2m_t D$

[Schmidt '95]

In terms of $\Omega = \{\theta_0, \theta_1, \phi_1, \theta_2, \phi_2\}$ helicity angles:

$\frac{d\sigma}{d\Omega} \propto$	$+3/4$	$(A_3 ^2 + A_4 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$(1 + \cos^2 \theta_0)$				
	$+3/4$	$(A_3 ^2 - A_4 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$(1 + \cos^2 \theta_0)$				
	$+3/4$	$(A_3 ^2 - A_4 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$(1 + \cos^2 \theta_0)$	$\cos \theta_1$	$\cos \theta_2$		
	$+3/4$	$(A_3 ^2 + A_4 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$(1 + \cos^2 \theta_0)$	$\cos \theta_1$	$\cos \theta_2$		
	$-3/2$	$(A_3 ^2 - A_4 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$\cos \theta_0$		$\cos \theta_2$		
	$-3/2$	$(A_3 ^2 + A_4 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\cos \theta_0$		$\cos \theta_2$		
	$-3/2$	$(A_3 ^2 + A_4 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$\cos \theta_0$	$\cos \theta_1$			
	$-3/2$	$(A_3 ^2 - A_4 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\cos \theta_0$	$\cos \theta_1$	$\cos \theta_2$		
	$+3/2$	$(A_1 ^2 + A_2 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$\sin^2 \theta_0$				
	$-3/2$	$(A_1 ^2 - A_2 ^2)$	$ a_2 ^2 + a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin^2 \theta_0$		$\cos \theta_2$		
	$+3/2$	$(A_1 ^2 - A_2 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 + b_3 ^2$	$\sin^2 \theta_0$	$\cos \theta_1$			
	$-3/2$	$(A_1 ^2 + A_2 ^2)$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin^2 \theta_0$	$\cos \theta_1$	$\cos \theta_2$		
	$+3/2$	$\sqrt{2} \operatorname{Re}\{A_1^* A_4\}$	$ a_2 ^2 - a_4 ^2$	$ b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\cos \phi_1$	
	$+3/2$	$\sqrt{2} \operatorname{Re}\{A_1^* A_4\}$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\cos \phi_1$	
	$+3/2$	$\sqrt{2} \operatorname{Re}\{A_2^* A_3\}$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\cos \phi_1$	
	$+3/2$	$\sqrt{2} \operatorname{Re}\{A_2^* A_3\}$	$ a_2 ^2 - a_4 ^2$	$ b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\cos \phi_1$	
	$-3/2$	$\sqrt{2} \operatorname{Re}\{A_2^* A_4\}$	$ a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\cos \phi_2$	
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	$-3/2$	$\sqrt{2} \operatorname{Re}\{A_1^* A_3\}$	$ a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$(1 - \cos \theta_1)$	$\sin \theta_2$	$\cos \phi_2$	
	-3	$\operatorname{Re}\{A_1^* A_2\}$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin^2 \theta_0$	$\sin \theta_1$	$\sin \theta_2$	$\cos(\phi_1 + \phi_2)$	
	$-3/2$	$\operatorname{Re}\{A_2^* A_4\}$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin^2 \theta_0$	$\sin \theta_1$	$\sin \theta_2$	$\cos(\phi_1 - \phi_2)$	
	$+3/2$	$\sqrt{2} \operatorname{Im}\{A_1^* A_4\}$	$ a_2 ^2 - a_4 ^2$	$ b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\sin \phi_1$	
	$+3/2$	$\sqrt{2} \operatorname{Im}\{A_1^* A_4\}$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\sin \phi_1$	
	$-3/2$	$\sqrt{2} \operatorname{Im}\{A_2^* A_3\}$	$ a_2 ^2 - a_4 ^2$	$ b_1 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\sin \phi_1$	
	$-3/2$	$\sqrt{2} \operatorname{Im}\{A_2^* A_3\}$	$ a_2 ^2 - a_4 ^2$	$ b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\sin \phi_1$	
	$+3/2$	$\sqrt{2} \operatorname{Im}\{A_2^* A_4\}$	$ a_4 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$	
	$+3/2$	$\sqrt{2} \operatorname{Im}\{A_2^* A_4\}$	$ a_2 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$(1 - \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$	
	$-3/2$	$\sqrt{2} \operatorname{Im}\{A_1^* A_3\}$	$ a_2 ^2$	$ b_1 ^2 - b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$	



Helicity amplitude decomposition in $bW^+\bar{b}W^-$

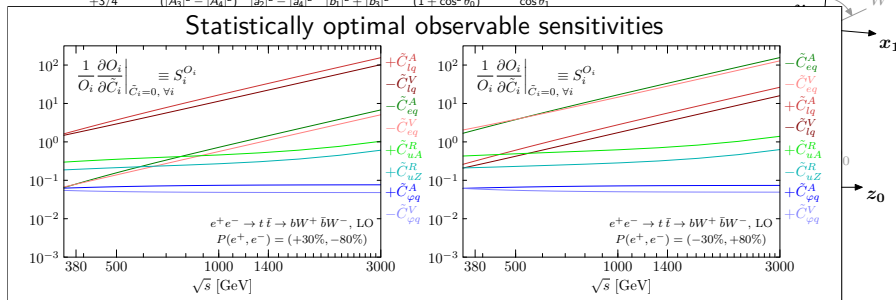
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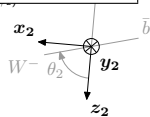
[Schmidt '95]

In terms of $\Omega = \{\theta_0, \theta_1, \phi_1, \theta_2, \phi_2\}$ helicity angles:

$$\begin{array}{cccccc} +3/4 & (|A_3|^2 + |A_4|^2) & |a_2|^2 + |a_4|^2 & |b_1|^2 + |b_3|^2 & (1 + \cos^2 \theta_0) & \\ +3/4 & (|A_3|^2 - |A_4|^2) & |a_2|^2 + |a_4|^2 & |b_1|^2 - |b_3|^2 & (1 + \cos^2 \theta_0) & \\ +3/4 & (|A_3|^2 - |A_4|^2) & |a_2|^2 - |a_4|^2 & |b_1|^2 + |b_3|^2 & (1 + \cos^2 \theta_0) & \cos \theta_1 \cos \theta_2 \end{array}$$



$$\begin{array}{cccccc} +3/2 & \sqrt{2} & \text{Im}\{A_1^* A_4\} & |a_2|^2 - |a_4|^2 & |b_3|^2 & \sin \theta_0 (1 + \cos \theta_0) & \sin \theta_1 & (1 + \cos \theta_2) & \sin \phi_1 \\ +3/2 & \sqrt{2} & \text{Im}\{A_1^* A_4\} & |a_2|^2 - |a_4|^2 & |b_1|^2 & \sin \theta_0 (1 + \cos \theta_0) & \sin \theta_1 & (1 - \cos \theta_2) & \sin \phi_1 \\ -3/2 & \sqrt{2} & \text{Im}\{A_2^* A_3\} & |a_2|^2 - |a_4|^2 & |b_1|^2 & \sin \theta_0 (1 - \cos \theta_0) & \sin \theta_1 & (1 + \cos \theta_2) & \sin \phi_1 \\ -3/2 & \sqrt{2} & \text{Im}\{A_2^* A_3\} & |a_2|^2 - |a_4|^2 & |b_3|^2 & \sin \theta_0 (1 - \cos \theta_0) & \sin \theta_1 & (1 - \cos \theta_2) & \sin \phi_1 \\ +3/2 & \sqrt{2} & \text{Im}\{A_3^* A_4\} & |a_4|^2 & |b_1|^2 - |b_3|^2 & \sin \theta_0 (1 + \cos \theta_0) & (1 + \cos \theta_1) & \sin \theta_2 & \sin \phi_2 \\ +3/2 & \sqrt{2} & \text{Im}\{A_3^* A_4\} & |a_2|^2 & |b_1|^2 - |b_3|^2 & \sin \theta_0 (1 + \cos \theta_0) & (1 - \cos \theta_1) & \sin \theta_2 & \sin \phi_2 \\ -3/2 & \sqrt{2} & \text{Im}\{A_3^* A_4\} & |a_2|^2 & |b_1|^2 - |b_3|^2 & \sin \theta_0 (1 - \cos \theta_0) & (1 + \cos \theta_1) & \sin \theta_2 & \sin \phi_2 \\ -3/2 & \sqrt{2} & \text{Im}\{A_3^* A_4\} & |a_2|^2 & |b_1|^2 - |b_3|^2 & \sin \theta_0 (1 - \cos \theta_0) & (1 - \cos \theta_1) & \sin \theta_2 & \sin \phi_2 \end{array}$$



Benchmark analysis

resonant $e^+e^- \rightarrow t\bar{t} \rightarrow bW^+ \bar{b}W^-$

$m_b/m_t \rightarrow 0$

analytically at LO

with perfect detector

and statistical uncertainties only

500 fb^{-1} at $\sqrt{s} = 500 \text{ GeV}$

1 ab^{-1} at $\sqrt{s} = 1 \text{ TeV}$

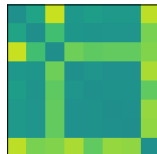
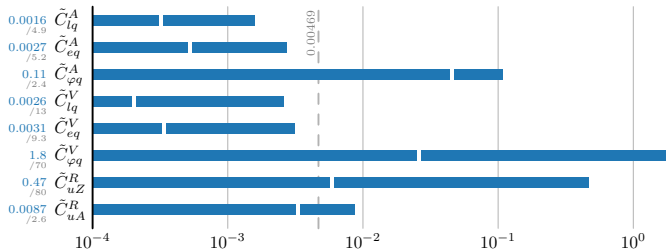
70% with $P(e^+, e^-) = (+0.3, -0.8)$

30% with $P(e^+, e^-) = (-0.3, +0.8)$

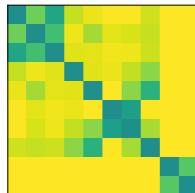
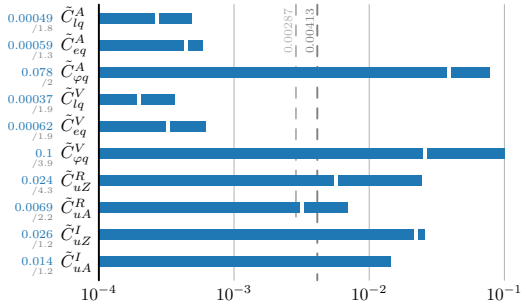
Global constraints

in TeV^{-2}

$\sigma + A^{\text{FB}}$:



Statistically optimal observables:

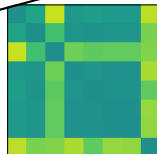
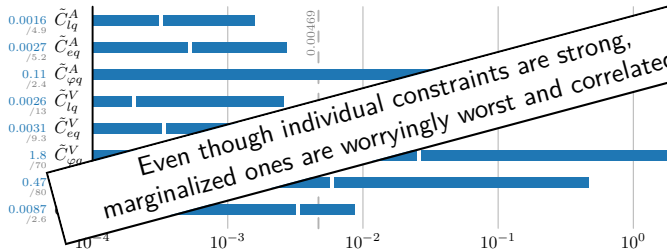


factor of 1.6 improvement
of the 8-coefficient's GDP
with linear coefficient
dependence only

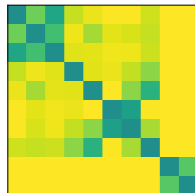
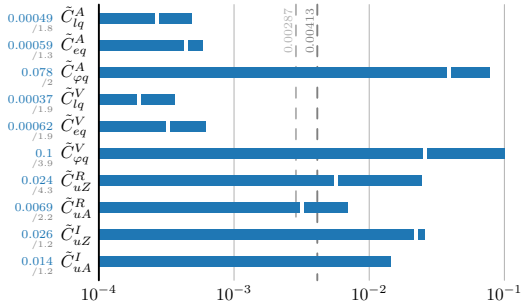
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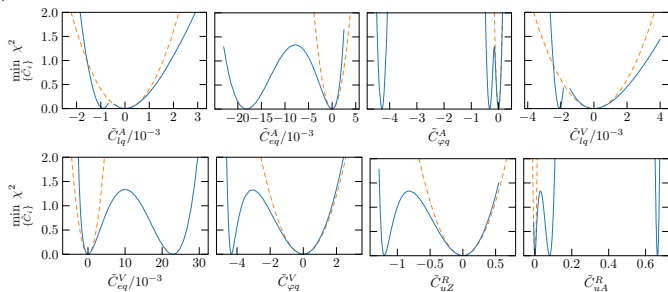
Statistically optimal observables:



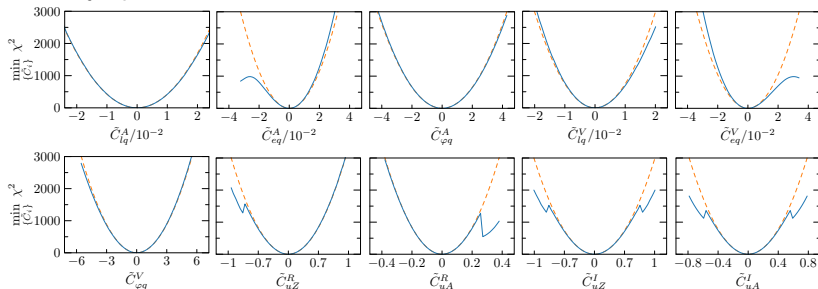
factor of 1.6 improvement
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Adding quadratic coefficient dependences

$\sigma + A^{\text{FB}}$:

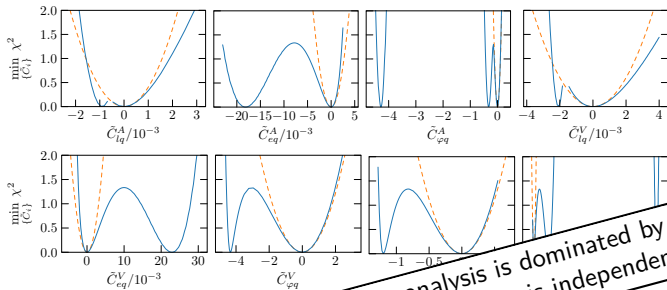


Statistically optimal observables:



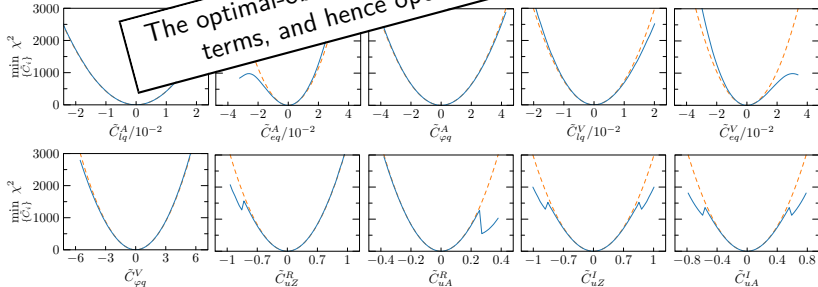
Adding quadratic coefficient dependences

$\sigma + A^{\text{FB}}$:



Statistically optimal observable

The optimal-observable analysis is dominated by linear terms, and hence operator basis independent.



Run parameters optimization,
GDP-based

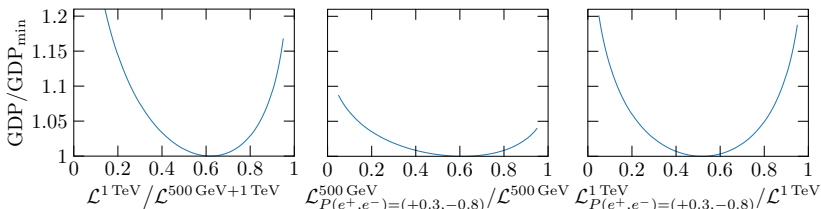
Examples of run parameters optimization

Given 1.5 ab^{-1} to share between two energies and polarizations, the optimal repartition is:

$\sqrt{s} = 500 \text{ GeV}$	570 fb^{-1}	61% with $P(e^+, e^-) = (+0.3, -0.8)$
1 TeV	930 fb^{-1}	52%

→ GDP is 1.02 times better than the benchmark one

for the optimal observable analysis with all 11 coefficients.



Same performances require 5.6 ab^{-1} with only $\sqrt{s} = 380 + 500 \text{ GeV}$:

$\sqrt{s} = 380 \text{ GeV}$	1.7 ab^{-1}	59%	with $P(e^+, e^-) = (+0.3, -0.8)$
500 GeV	3.9 ab^{-1}	53%	

Summary

The EFT parametrizes systematically the theory space in direct vicinity of the standard model.

A global analysis of future-lepton-collider constraints on the top EFT is ongoing:

- including notably four-fermion operators,
- examining NLO QCD and off-shell top effects.

Statistically optimal observables greatly help

- constraining all directions in the effective-theory space,
- producing basis-independent limits.

Global determinant parameter ratios assess the strengthening of global constraints, basis independently.

Backup

NLO in QCD for $e^+e^- \rightarrow bW^+\bar{b}W^-$

For various beam polarizations and center-of-mass energies:

pol	$\sqrt{s}$ [GeV]	σ_{SM} [fb]	$ A_{eq} $	$ A_{eq} $	$ A_{eq} $	$ V_{eq} $	$ V_{eq} $	σ_I [fb]	$ R_{eq} $	$ R_{eq} $	$ I_{eq} $	$ I_{eq} $	$ R_{eq} $
00	300	$2.92^{+0.23}_{-0.15}$	$0.353^{+0.23}_{-0.15}$	$-0.0856^{+0.04}_{-0.04}$	$0.14^{+0.23}_{-0.15}$	$-0.621^{+0.23}_{-0.15}$	$-0.303^{+0.23}_{-0.15}$	$-0.136^{+0.23}_{-0.15}$	$0.349^{+0.23}_{-0.15}$	$0.32^{+0.23}_{-0.15}$	$-0.000225^{+0.0002}_{-0.0002}$	$-0.000125^{+0.0001}_{-0.0001}$	$0.00021^{+0.0002}_{-0.0002}$
00	380	825^{+2}_{-1}	77.1^{+3}_{-1}	-55.6^{+2}_{-1}	53.1^{+2}_{-1}	-993^{+2}_{-1}	-635^{+2}_{-1}	-62.9^{+2}_{-1}	118^{+2}_{-1}	323^{+2}_{-1}	$0.107^{+0.001}_{-0.001}$	$-0.434^{+0.001}_{-0.001}$	$0.25^{+0.001}_{-0.001}$
00	500	$669^{+0.6}_{-0.5}$	258^{+1}_{-1}	-233^{+1}_{-1}	49.2^{+1}_{-1}	-1230^{+1}_{-1}	-750^{+1}_{-1}	-102^{+1}_{-1}	102^{+1}_{-1}	263^{+1}_{-1}	$-2.08^{+0.01}_{-0.01}$	$1.78^{+0.01}_{-0.01}$	$0.715^{+0.01}_{-0.01}$
00	1000	$221^{+0.9}_{-0.8}$	936^{+2}_{-1}	$-475^{+0.2}_{-0.2}$	$15.0^{+0.2}_{-0.2}$	$-1070^{+0.2}_{-0.2}$	$-940^{+0.2}_{-0.2}$	$-15.5^{+0.2}_{-0.2}$	$36.1^{+0.2}_{-0.2}$	$87.3^{+0.2}_{-0.2}$	$0.392^{+0.001}_{-0.001}$	$-8.74^{+0.001}_{-0.001}$	$0.907^{+0.001}_{-0.001}$
00	1400	$132^{+0.6}_{-0.5}$	$391^{+0.9}_{-0.8}$	$-412^{+0.3}_{-0.3}$	$8.89^{+0.3}_{-0.3}$	$-1460^{+0.3}_{-0.3}$	$-816^{+0.3}_{-0.3}$	$-8.89^{+0.3}_{-0.3}$	$21.6^{+0.3}_{-0.3}$	$59.8^{+0.3}_{-0.3}$	$1.8^{+0.01}_{-0.01}$	$0.257^{+0.001}_{-0.001}$	$0.414^{+0.001}_{-0.001}$
00	3000	$40.2^{+0.1}_{-0.1}$	$1080^{+0.2}_{-0.2}$	$-70.4^{+0.2}_{-0.2}$	$1.28^{+0.01}_{-0.01}$	$-1270^{+0.01}_{-0.01}$	$-689^{+0.01}_{-0.01}$	$-0.717^{+0.01}_{-0.01}$	$10^{+0.01}_{-0.01}$	$14.2^{+0.01}_{-0.01}$	$1.85^{+0.001}_{-0.001}$	$2.15^{+0.001}_{-0.001}$	$-0.261^{+0.001}_{-0.001}$
+-	300	$2.73^{+0.3}_{-0.1}$	$0.351^{+0.2}_{-0.1}$	—	$0.126^{+0.2}_{-0.1}$	$-0.62^{+0.2}_{-0.1}$	—	$-0.14^{+0.2}_{-0.1}$	$0.376^{+0.2}_{-0.1}$	$0.241^{+0.2}_{-0.1}$	$6.25e-06^{+0.0001}_{-0.0001}$	$2.7e-06^{+0.0001}_{-0.0001}$	$0.000197^{+0.0001}_{-0.0001}$
+-	380	579^{+2}_{-1}	73^{+1}_{-1}	—	36.1^{+2}_{-1}	-968^{+2}_{-1}	—	-165^{+2}_{-1}	165^{+2}_{-1}	198^{+2}_{-1}	$-0.324^{+0.001}_{-0.001}$	$0.185^{+0.001}_{-0.001}$	—
+-	500	$469^{+0.4}_{-0.3}$	$287^{+0.5}_{-0.4}$	—	$11.7^{+0.4}_{-0.3}$	$-1270^{+0.4}_{-0.3}$	—	$-0.2^{+0.4}_{-0.3}$	$130^{+0.4}_{-0.3}$	$164^{+0.4}_{-0.3}$	$-0.442^{+0.001}_{-0.001}$	$0.554^{+0.001}_{-0.001}$	—
+-	1000	$160^{+0.5}_{-0.4}$	$470^{+0.8}_{-0.7}$	—	$11.8^{+0.6}_{-0.5}$	$-1450^{+0.6}_{-0.5}$	—	$-15.5^{+0.6}_{-0.5}$	$44.9^{+0.6}_{-0.5}$	$52.6^{+0.6}_{-0.5}$	$-0.663^{+0.001}_{-0.001}$	$5.09^{+0.001}_{-0.001}$	$0.587^{+0.001}_{-0.001}$
+-	1400	$84.9^{+0.9}_{-0.8}$	$507^{+1.0}_{-0.9}$	—	$7.57^{+0.9}_{-0.8}$	$-1230^{+0.9}_{-0.8}$	—	$-7.76^{+0.9}_{-0.8}$	$22.2^{+0.9}_{-0.8}$	$29.5^{+0.9}_{-0.8}$	$-1.22^{+0.01}_{-0.01}$	$-2.38^{+0.01}_{-0.01}$	$0.281^{+0.01}_{-0.01}$
+-	3000	$23.3^{+0.3}_{-0.3}$	$356^{+0.4}_{-0.4}$	—	$0.574^{+0.01}_{-0.01}$	$-1270^{+0.01}_{-0.01}$	—	$-1.08^{+0.01}_{-0.01}$	$6.28^{+0.01}_{-0.01}$	$1.93^{+0.01}_{-0.01}$	$8.36^{+0.001}_{-0.001}$	$0.197^{+0.001}_{-0.001}$	—
++	300	$0.218^{+0.4}_{-0.3}$	—	$-0.085e-2^{+0.001}_{-0.001}$	$0.0147^{+0.001}_{-0.001}$	$-0.302^{+0.001}_{-0.001}$	$-0.00343^{+0.001}_{-0.001}$	$-0.0259^{+0.001}_{-0.001}$	$0.0799^{+0.001}_{-0.001}$	$3.38e-06^{+0.0001}_{-0.0001}$	$-7.78e-06^{+0.0001}_{-0.0001}$	$1.84e-05^{+0.0001}_{-0.0001}$	—
++	380	$249^{+0.4}_{-0.4}$	—	$-51.6^{+0.4}_{-0.4}$	$16.2^{+0.4}_{-0.4}$	$-649^{+0.4}_{-0.4}$	—	$-41.8^{+0.4}_{-0.4}$	$124^{+0.4}_{-0.4}$	$0.0946^{+0.001}_{-0.001}$	$-0.0633^{+0.001}_{-0.001}$	$0.000205^{+0.0001}_{-0.0001}$	—
++	500	$203^{+0.5}_{-0.4}$	—	$-213^{+0.5}_{-0.4}$	$15.8^{+0.5}_{-0.4}$	$-810^{+0.5}_{-0.4}$	—	$-76.7^{+0.5}_{-0.4}$	$34.2^{+0.5}_{-0.4}$	$99.7^{+0.5}_{-0.4}$	$0.316^{+0.001}_{-0.001}$	$0.187^{+0.001}_{-0.001}$	$0.255^{+0.001}_{-0.001}$
++	1000	$63.4^{+0.9}_{-0.8}$	—	$-327^{+0.9}_{-0.8}$	$8.88^{+0.9}_{-0.8}$	$-810^{+0.9}_{-0.8}$	—	$-0.24^{+0.9}_{-0.8}$	$34.1^{+0.9}_{-0.8}$	$10.2^{+0.9}_{-0.8}$	$-0.832^{+0.01}_{-0.01}$	$0.255^{+0.001}_{-0.001}$	$0.39^{+0.001}_{-0.001}$
++	1400	$33.8^{+0.9}_{-0.7}$	—	$-493^{+0.9}_{-0.7}$	$2.86^{+0.9}_{-0.7}$	$-850^{+0.9}_{-0.7}$	—	$-0.208^{+0.9}_{-0.7}$	$-4.75^{+0.9}_{-0.7}$	$10.2^{+0.9}_{-0.7}$	$0.448^{+0.001}_{-0.001}$	$0.475^{+0.001}_{-0.001}$	$0.258^{+0.001}_{-0.001}$
++	3000	—	—	$-424^{+0.01}_{-0.01}$	$0.226^{+0.01}_{-0.01}$	—	$0.146^{+0.01}_{-0.01}$	$-1.5^{+0.01}_{-0.01}$	$-497^{+0.01}_{-0.01}$	$2.52^{+0.01}_{-0.01}$	$2110^{+0.001}_{-0.001}$	$-0.0453^{+0.001}_{-0.001}$	—

(MG5_aMC@NLO, complex mass scheme, $m_b/m_t \rightarrow 0$,
EFT dependence of the total width not included)