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Revisiting the $b \rightarrow c\ell\nu$ anomalies with charged Higgs boson

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In collaboration with

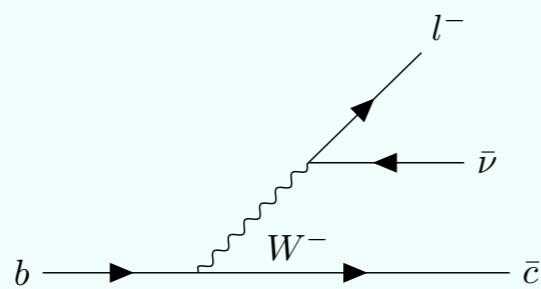
Nilakshi Das and Rupak Dutta

PLANCK 2023, 22-26 May, Warsaw

Anomalies in B decays

In Standard Model (SM), three flavours of charged leptons has the same gauge coupling strength, known as lepton flavour universality (LFU).

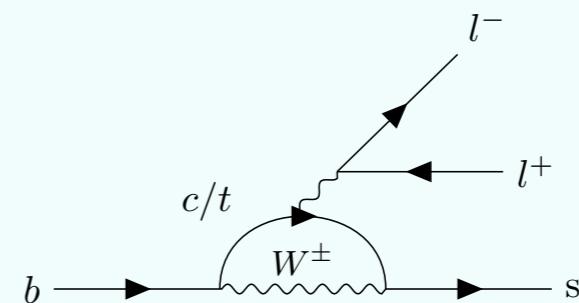
Violation of LFU is first observed in semi-leptonic B decays :



Charged current (main focus)

e.g. $B \rightarrow D^* l \nu, B_c \rightarrow J/\psi l \nu \dots$

Observables: $R_D, R_{D^{(*)}}, R_{J/\psi}, P_\tau^{D^*} \dots$



Neutral current

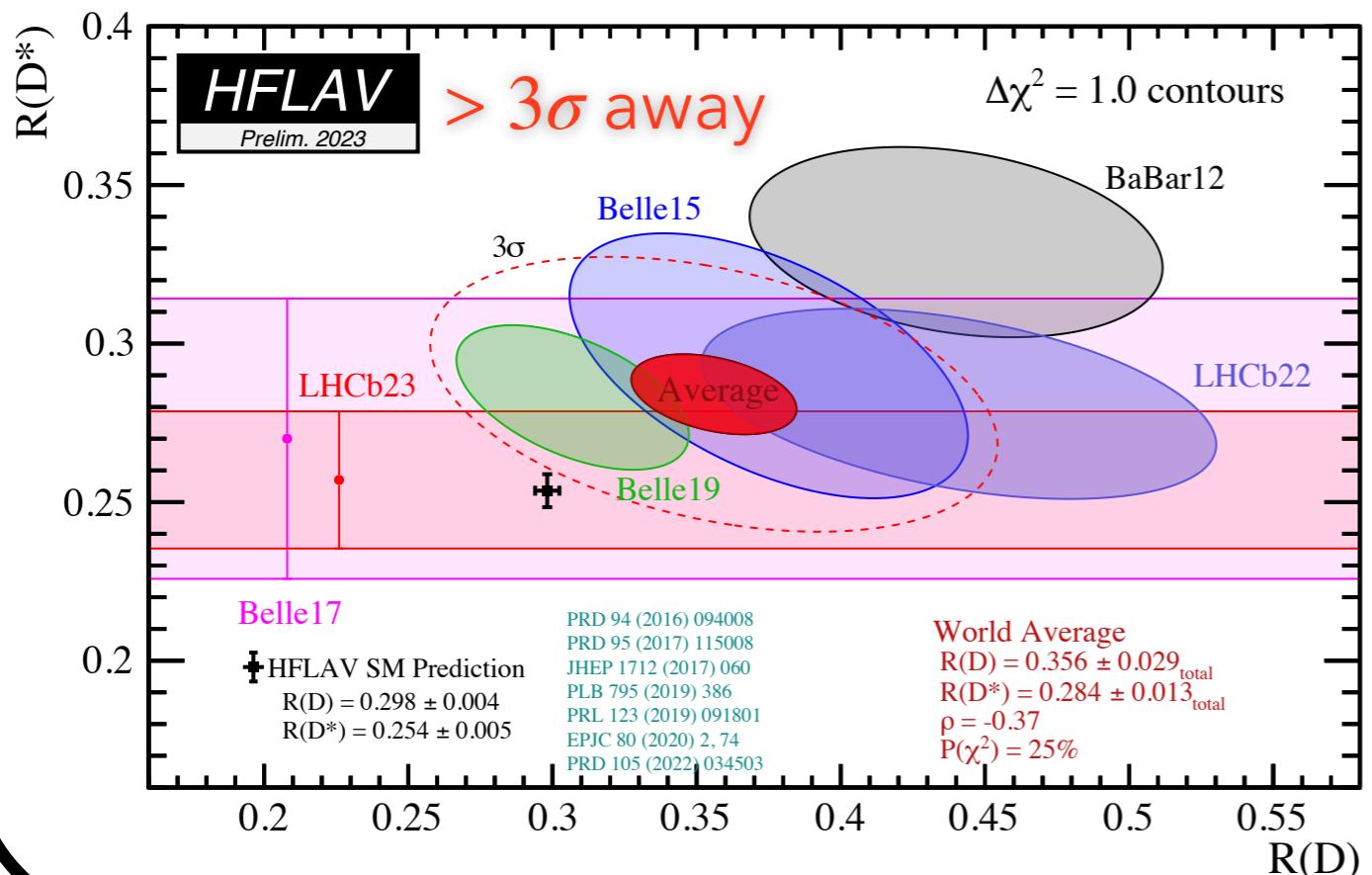
e.g. $B \rightarrow K^* l^+ l^-, B_s \rightarrow \phi \mu^+ \mu^- \dots$

Observables: $R_K, R_{K^{(*)}}, P'_5 \dots$

...Talk by Monica Pepe Altarelli
LHCb-PAPER-2022-046

Present status of $b \rightarrow c\tau\nu$ anomalies

$$R_D = \frac{\mathcal{B}(B \rightarrow D\tau\nu_\tau)}{\mathcal{B}(B \rightarrow Dl\nu_l)} \quad R_{D^*} = \frac{\mathcal{B}(B \rightarrow D^*\tau\nu_\tau)}{\mathcal{B}(B \rightarrow D^*l\nu_l)}$$



$$P_\tau^{D^*} = \frac{\Gamma(B \rightarrow D^* \tau^{\lambda=+1/2} \nu_\tau) - \Gamma(B \rightarrow D^* \tau^{\lambda=-1/2} \nu_\tau)}{\Gamma(B \rightarrow D^* \tau \nu_\tau)}$$

$$P_\tau^{D^* \text{ SM}} = -0.497 \pm 0.013 \quad 2210.10751$$

$$P_\tau^{D^* \text{ ex}} = -0.38 \pm 0.51^{+0.21}_{-0.16}$$

Iguro et.al.
Report No.: P3H-22-103

Phys. Rev. Lett. 118,
211801 (2017)

$$R_{J/\psi} = \frac{\mathcal{B}(B_c \rightarrow J/\psi \tau\nu_\tau)}{\mathcal{B}(B_c \rightarrow J/\psi \mu\nu_\mu)}$$

Phys. Rev. Lett. 125,
222003 (2020)

$$R_{J/\psi}^{SM} = 0.258 \pm 0.004$$

Phys. Rev. Lett. 120,
121801 (2018)

$$F_L^{D^*} = \frac{\Gamma(B \rightarrow D_L^* \tau\nu_\tau)}{\Gamma(B \rightarrow D^* \tau\nu_\tau)}$$

Iguro et.al.
Report No.: P3H-22-103

$$F_L^{D^* \text{ SM}} = 0.46 \pm 0.04$$

$$F_L^{D^* \text{ ex}} = 0.60 \pm 0.08 \pm 0.035$$

Report: BELLE-CONF-1805
1903.03102

$$R_{\Lambda_c} = \frac{\mathcal{B}(\Lambda_b \rightarrow \Lambda_c \tau\nu_\tau)}{\mathcal{B}(\Lambda_b \rightarrow \Lambda_c \mu\nu_\mu)}$$

Phys. Rev. D 99, 055008
(2019)

$$R_{\Lambda_c}^{SM} = 0.324 \pm 0.004$$

$$R_{\Lambda_c}^{ex} = 0.242 \pm 0.026 \pm 0.040 \pm 0.059$$

Phys. Rev. Lett. 128,
191803 (2022)

New physics explanations

Model independent approach:

$$\mathcal{H}_{eff} = \frac{4G_F}{\sqrt{2}} V_{cb} (1 + C_{LL}^V) \mathcal{O}_{C_{LL}^V} + C_{RL}^V \mathcal{O}_{C_{RL}^V} + C_{LL}^S \mathcal{O}_{C_{LL}^S} + C_{RL}^S \mathcal{O}_{C_{RL}^S} + C_{LL}^T \mathcal{O}_{C_{LL}^T} + \\ C_{LR}^V \widetilde{\mathcal{O}}_{C_{LR}^V} + C_{RR}^V \widetilde{\mathcal{O}}_{C_{RR}^V} + C_{LR}^S \widetilde{\mathcal{O}}_{C_{LR}^S} + C_{RR}^S \widetilde{\mathcal{O}}_{C_{RR}^S} + C_{RR}^T \widetilde{\mathcal{O}}_{C_{RR}^T} + h.c.$$

$$\mathcal{O}_{C_{LL}^S} = (\bar{c}_R b_L)(\bar{\tau}_R \nu_{\tau L})$$

$$\mathcal{O}_{C_{RL}^S} = (\bar{\tau}_R \nu_{\tau L})(\bar{c}_L b_R)$$

$$\widetilde{\mathcal{O}}_{C_{LL}^S} = (\bar{\tau}_L \nu_{\tau R})(\bar{c}_R b_L)$$

$$\widetilde{\mathcal{O}}_{C_{RL}^S} = (\bar{\tau}_L \nu_{\tau R})(\bar{c}_L b_R)$$

Cirigliano et.al.

Phys. Rev. D 85, 054512

Model dependent approach:

W' boson

Charged Higgs boson

Leptoquark

Model independent approach - χ^2 fit

χ^2 fit using iminuit

Observables : R_D , $R_{D^{(*)}}$, $R_{J/\psi}$, $P_\tau^{D^*}$, R_{Λ_c} , $F_L^{D^*}$

1D scenario (Real) : $C_{LL}^S = 0.15$ $C_{RL}^S = 0.16$ $C_{RR}^S/C_{LR}^S = -0.51$

1D scenario (Complex) :

$C_{LL}^S = -0.67 - i 0.84$ $C_{RL}^S = 0.16 - i 6.4 \times 10^{-6}$ $C_{RR}^S/C_{LR}^S = -0.49 - i 0.14$

$(C_{LL}^S, C_{RL}^S) = (-0.37 + i 0, 0.48 + i 0)$

$(C_{LL}^S, C_{RR}^S/C_{LR}^S) = (-0.67 - i 0.5, -0.63 - i 0.26)$

2D scenario :

$(C_{RL}^S, C_{RR}^S/C_{LR}^S) = (0.16 - i 1.6 \times 10^{-4}, \sim 0 + i 0)$

$(C_{RR}^S, C_{LR}^S) = (-0.88 - i 0.62, 0.50 + i 0.45)$

Contour Plots

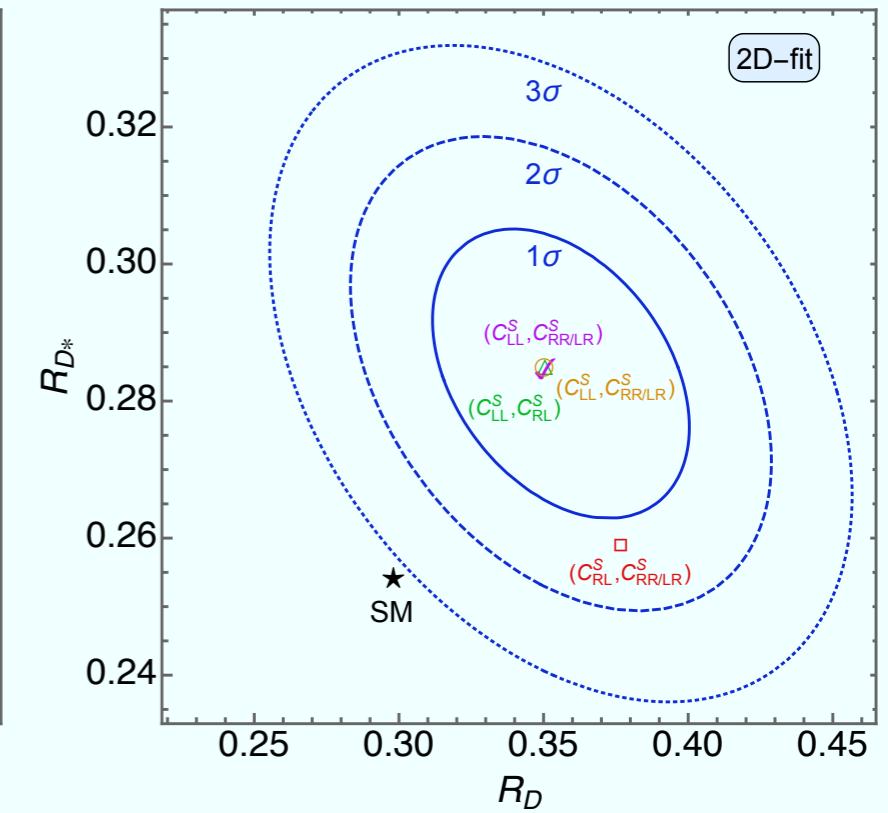
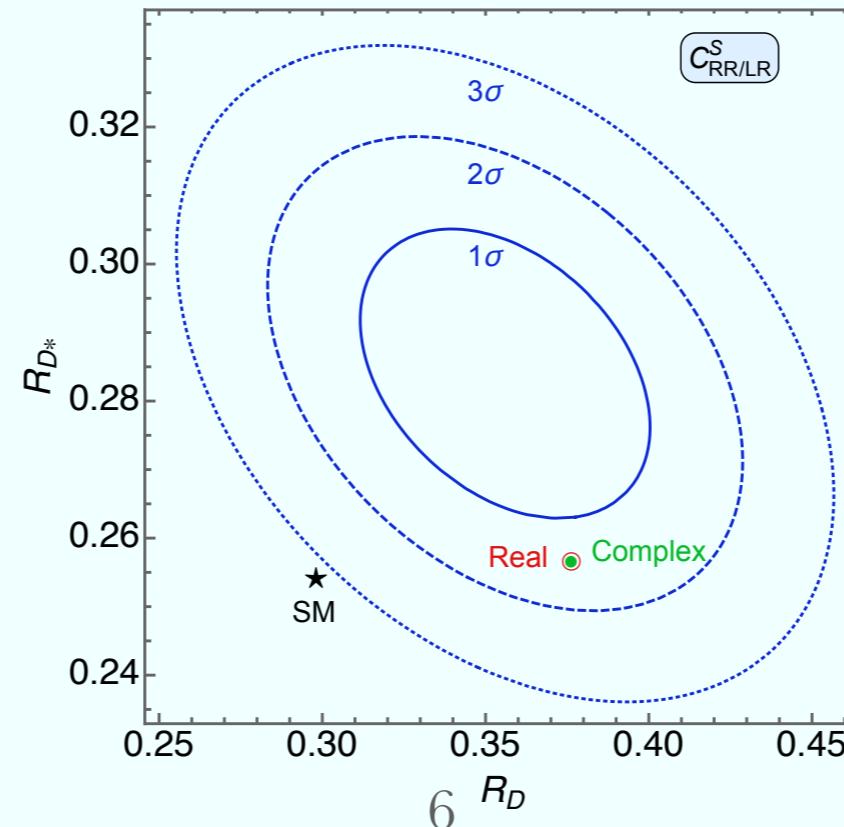
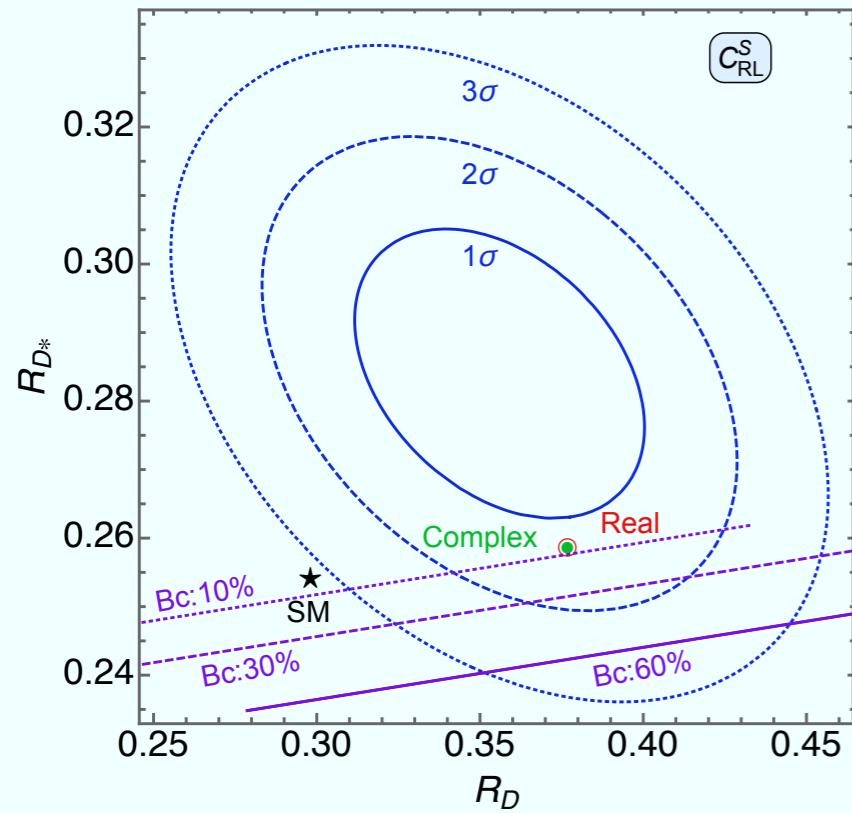
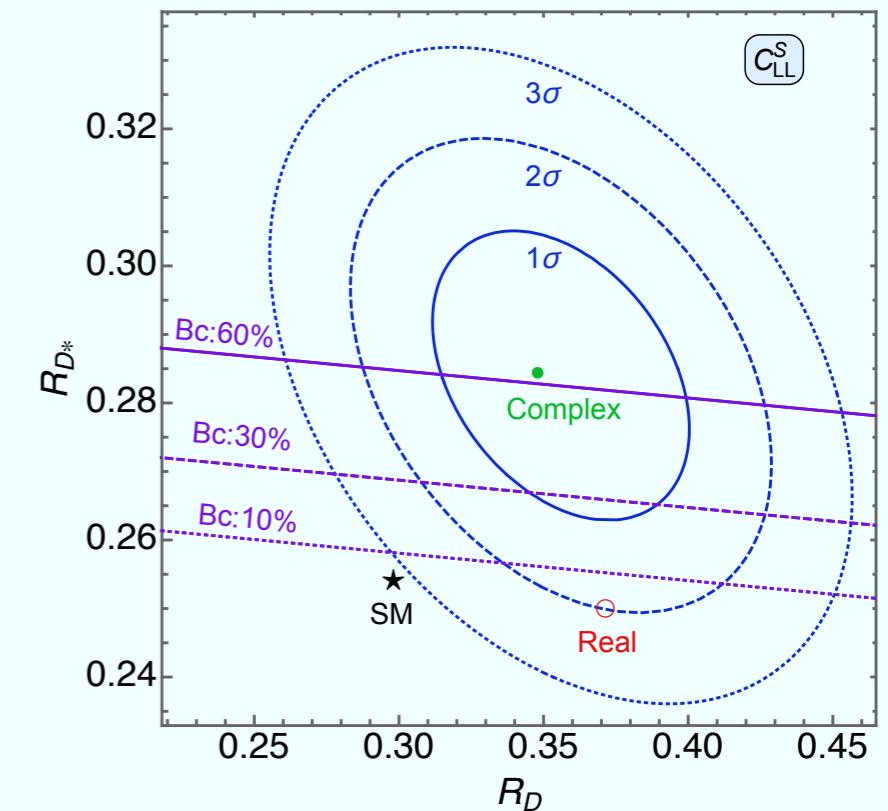
$R_D - R_{D^*}$ plane

Blue: $1\sigma, 2\sigma, 3\sigma$ contours

★ : SM prediction, ○ : Real fit, ● : Complex fit

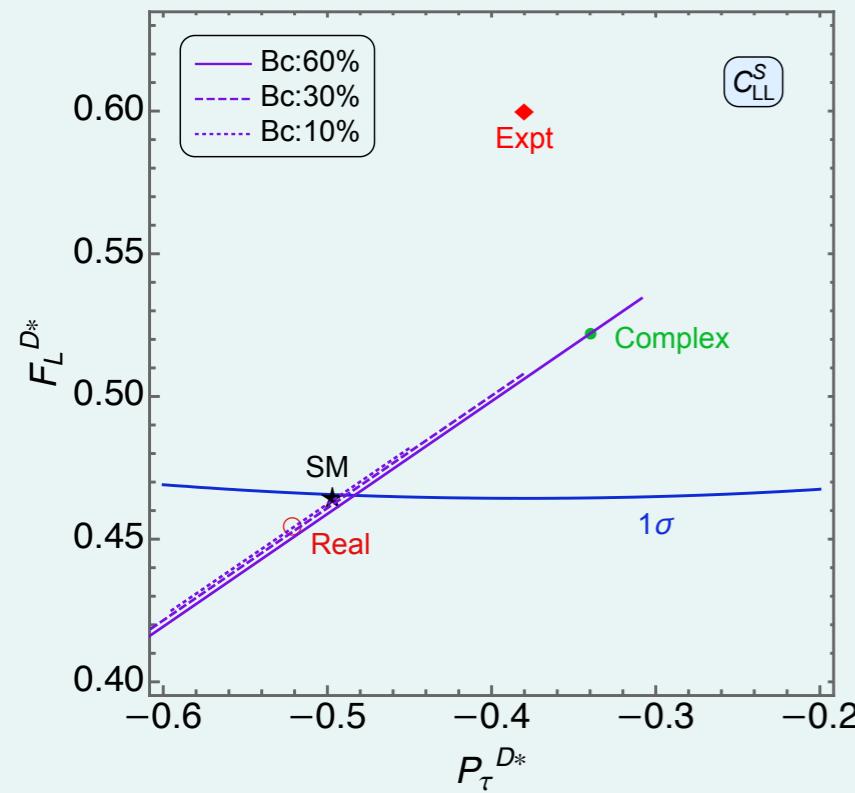
Purple: $\mathcal{B}(B_c \rightarrow \tau \bar{\nu}_\tau) = B_c = 60\%, 30\%, 10\%$

Blanke et.al. Phys. Rev. D 99, 075006 (2019)
Aebischer et.al. JHEP 07, 130 (2021)

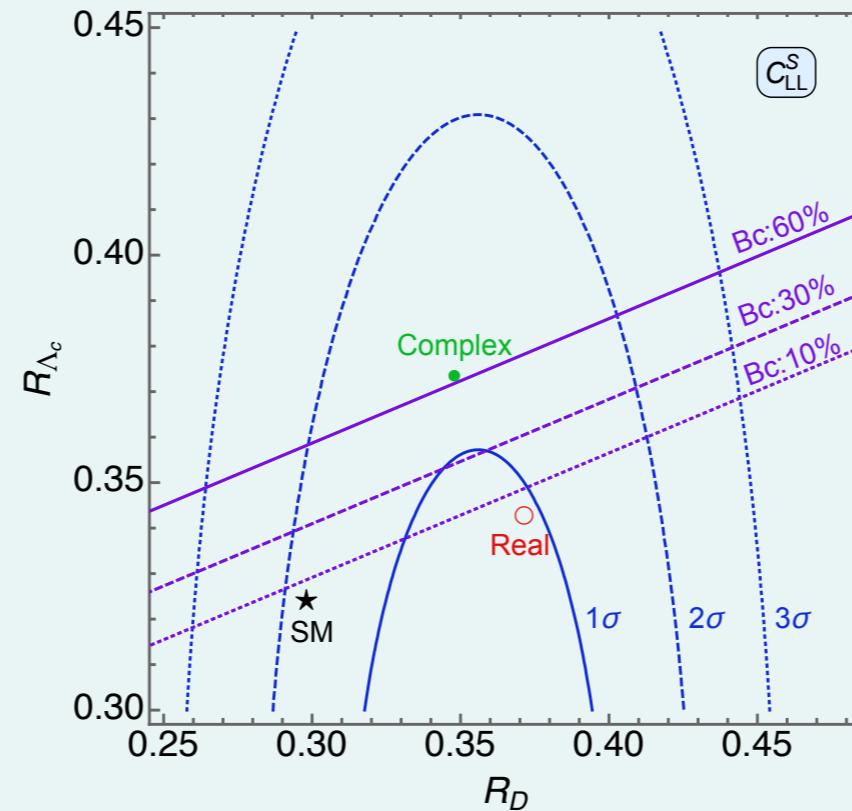


Contour Plots

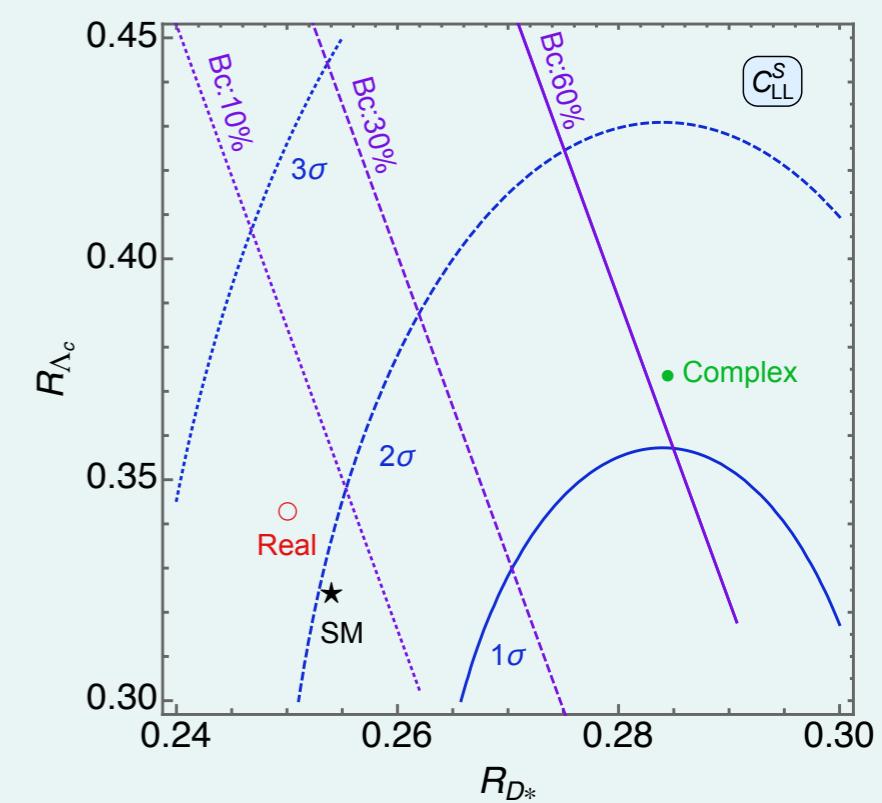
$F_L^{D^*} - P_\tau^{D^*}$ plane



$R_D - R_{\Lambda_c}$ plane

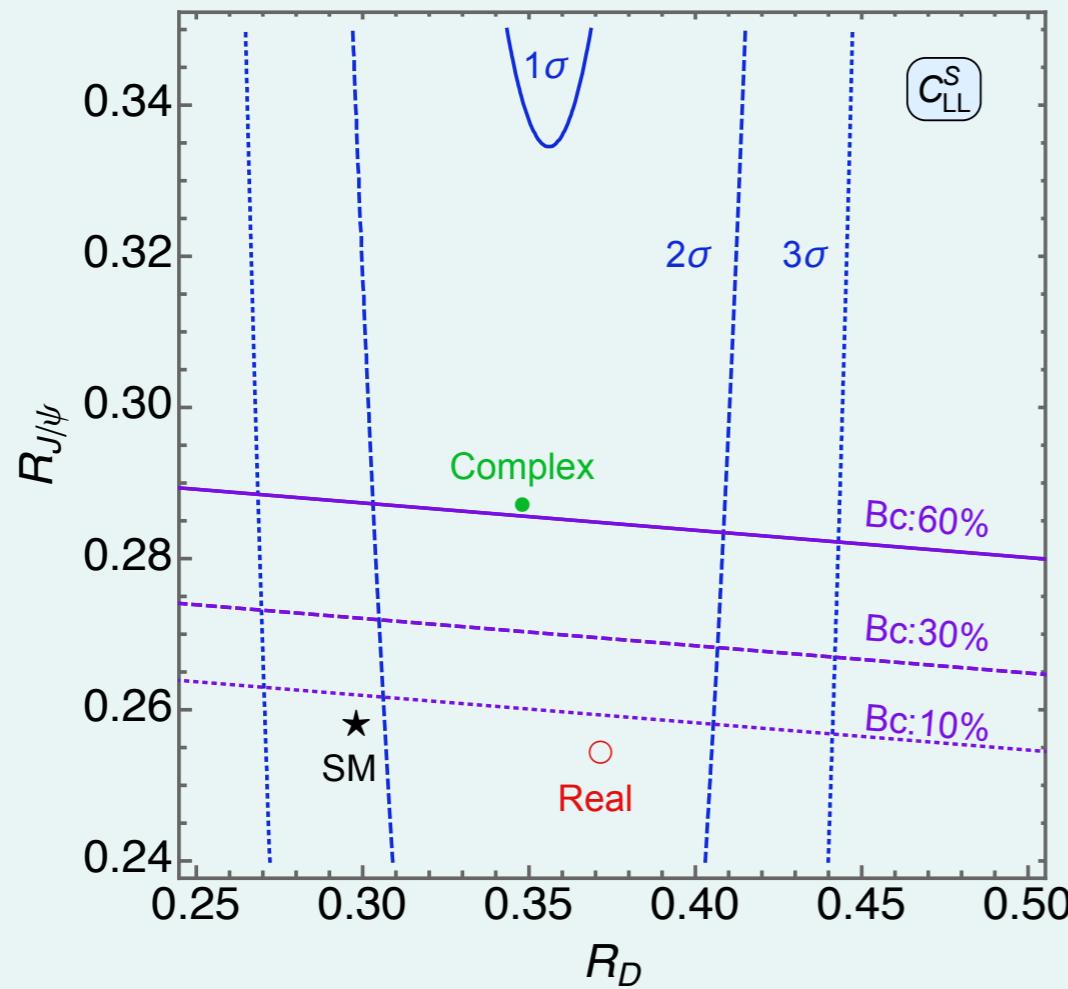


$R_{D^*} - R_{\Lambda_c}$ plane

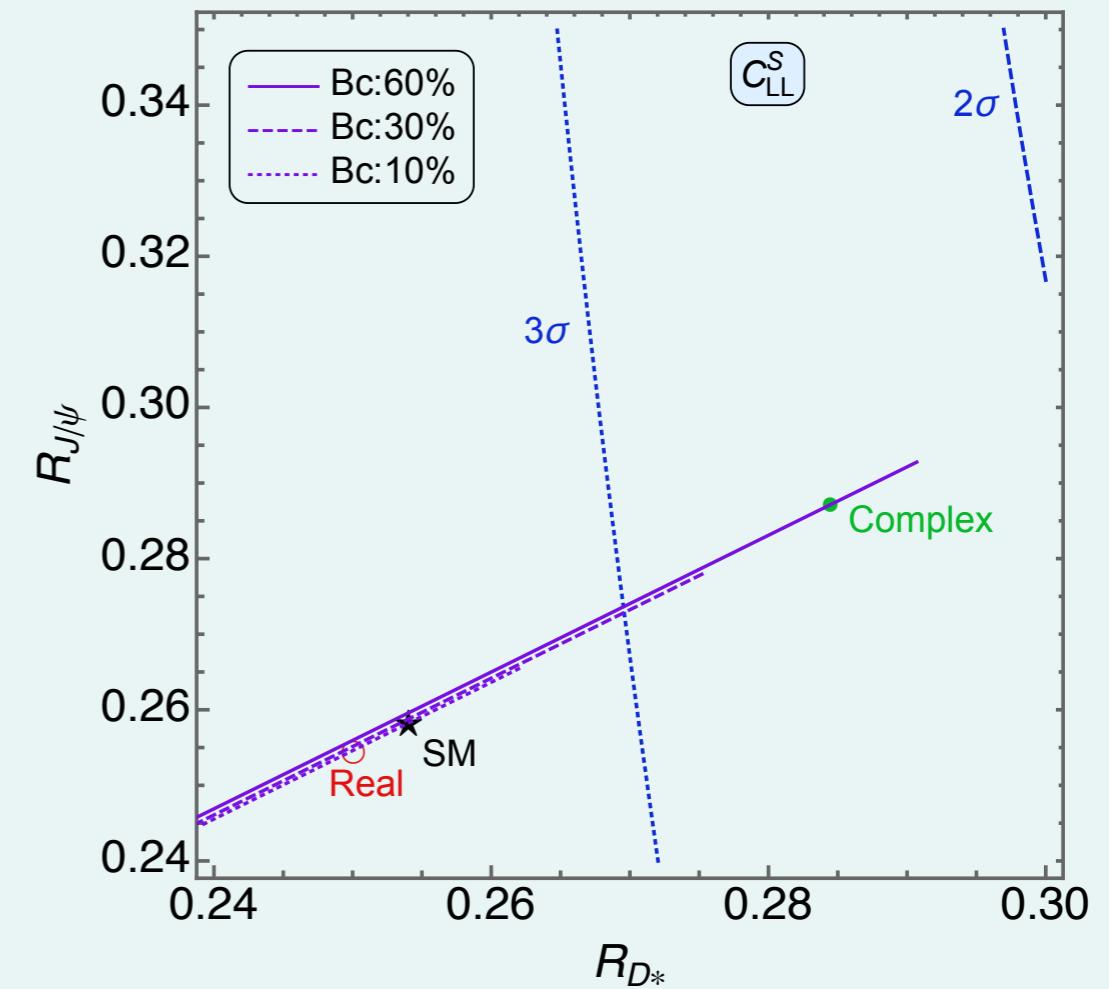


Contour Plots

$R_D - R_{J/\psi}$ plane



$R_{D^*} - R_{J/\psi}$ plane



Model dependent approach

Charged Higgs boson in generic two Higgs doublet model (2HDM) :

$$\mathcal{L}_{\text{int}} = y_{bc}^u H^- (\bar{b} P_R c) - y_{bc}^d H^- (\bar{b} P_L c) - y_{\tau\nu} H^- (\bar{\tau} P_L \nu_\tau) + h.c.$$

$$y_{bc}^u = y_{bc}$$

$$y_{bc}^d \text{ constrained by } B_s - \bar{B}_s \text{ mixing}$$

$$\epsilon_{RL}^S$$

Phys.Rev.D 105 (2022) 9, 095011

Effective Hamiltonian for the $b \rightarrow c\tau\nu$ transition :

$$\mathcal{H}_{\text{eff}} = 2\sqrt{2}G_F V_{cb} [(\bar{c}_R \gamma_\mu b_L)(\bar{\tau}_R \gamma^\mu \nu_\tau) + C_{LL}^S (\bar{c}_R b_L)(\bar{\tau}_R \nu_\tau)].$$

At the m_{H^\pm} scale,

$$C_{LL}^S = \frac{y_{bc} y_{\tau\nu}}{2\sqrt{2}G_F V_{cb} m_{H^\pm}^2} \xleftarrow{\hspace{-1cm}} \text{Collider search}$$

Collider Analysis : b-veto category

Signal process : $bc \rightarrow H^\pm \rightarrow \tau\nu$

Dominant background process : $pp \rightarrow W^\pm \rightarrow \tau\nu$

Other background processes : Drell-Yan (DY) + jets,
Misid. τ_h ($j \rightarrow \tau_h$ fake), $t\bar{t}$, VV ($V = W/Z$) + jets, single t

Trigger cuts:

$$N_{\tau_h} = 1, N_\ell = 0$$

$$\text{b-veto : } N_{\text{b-jets}} = 0$$

$$N_j \leq 2$$

Selection cuts:

$$\Delta\phi(p_{T,\tau_h}, \not{p}_T) \geq 2.4$$

$$0.7 \leq p_{T,\tau_h}/\not{p}_T \leq 1.3$$

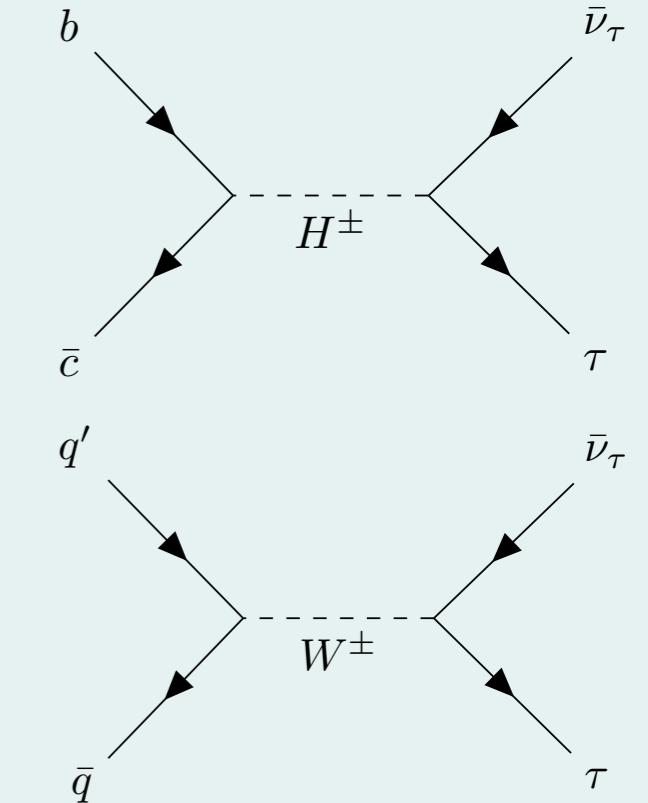
Optimised selection cuts:

$$p_{T,\tau_h} \geq [50, 50, 70, 80, 90, 110] \text{ GeV}$$

$$m_T \geq [100, 110, 150, 170, 200, 220] \text{ GeV}$$

$$E_T \geq [50, 50, 60, 80, 90, 100] \text{ GeV}$$

$$\text{for } m_{H^\pm} = [180, 200, 250, 300, 350, 400] \text{ GeV}$$

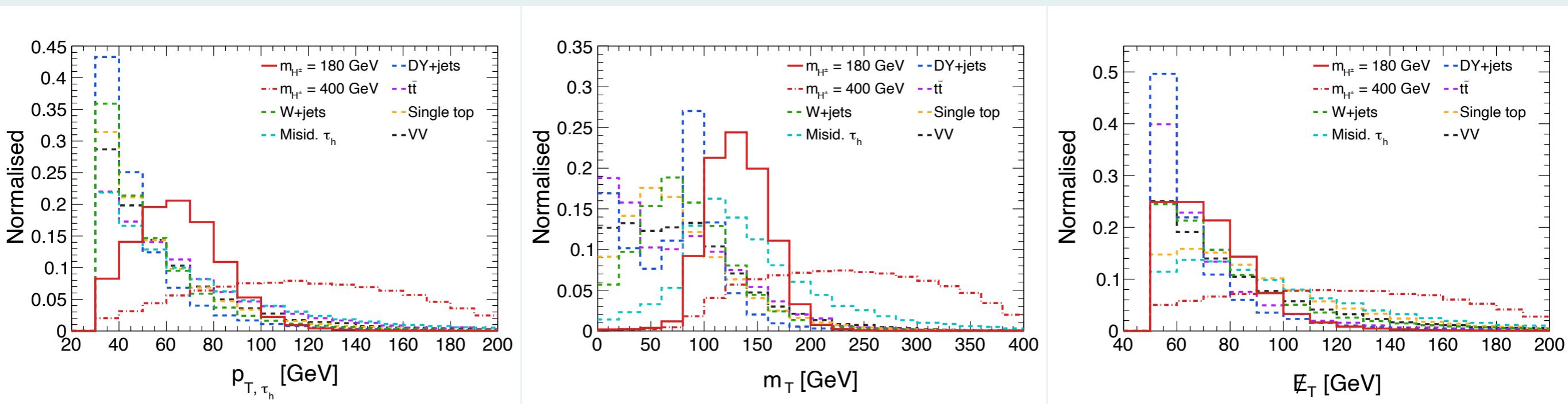
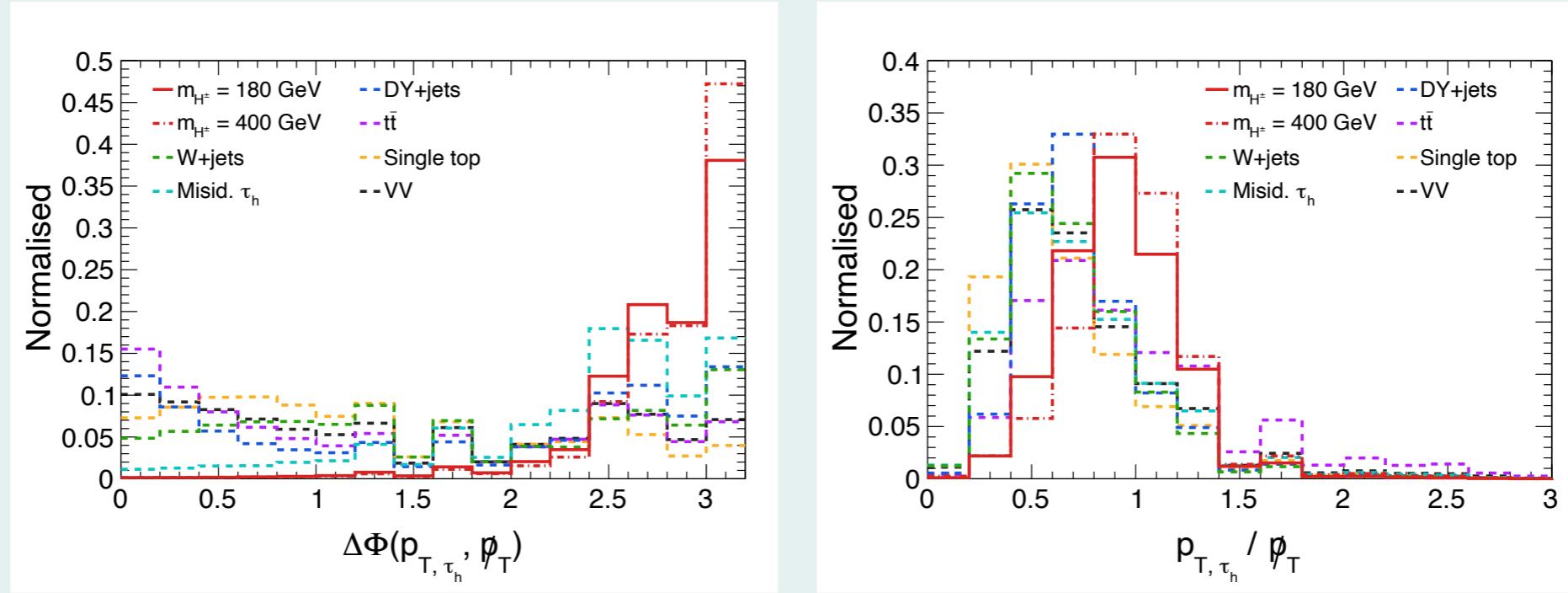


Signal Significance:

$$\sigma_s = \frac{S}{\sqrt{B}}$$

Collider Analysis : b-veto category

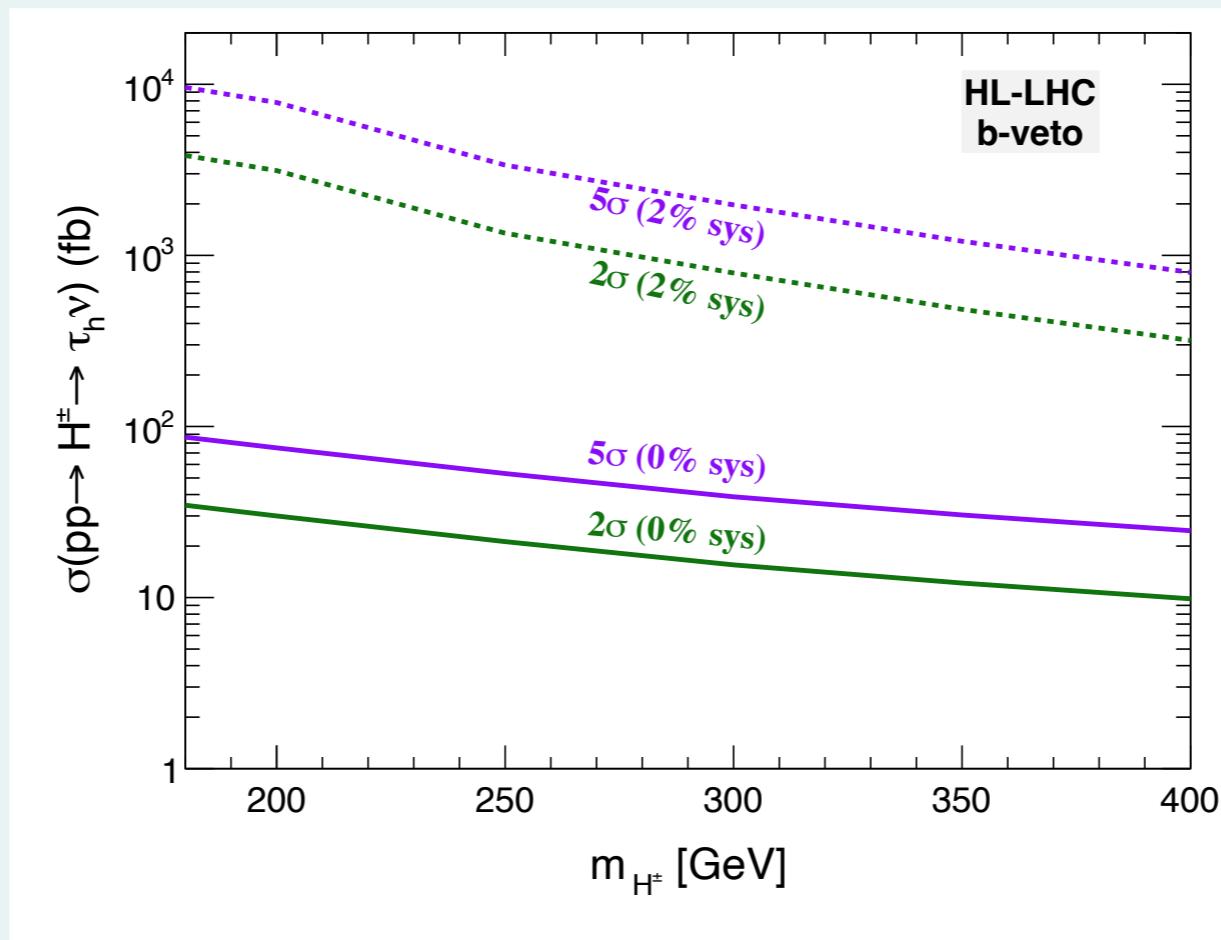
Kinematic variables



Collider Analysis : b-veto category

Upper limit on the signal production cross-section :

$$\sigma(pp \rightarrow H^\pm \rightarrow \tau_h \nu)_{\text{UL}} = \frac{N \cdot \sqrt{B}}{\epsilon \cdot \mathcal{L}}$$

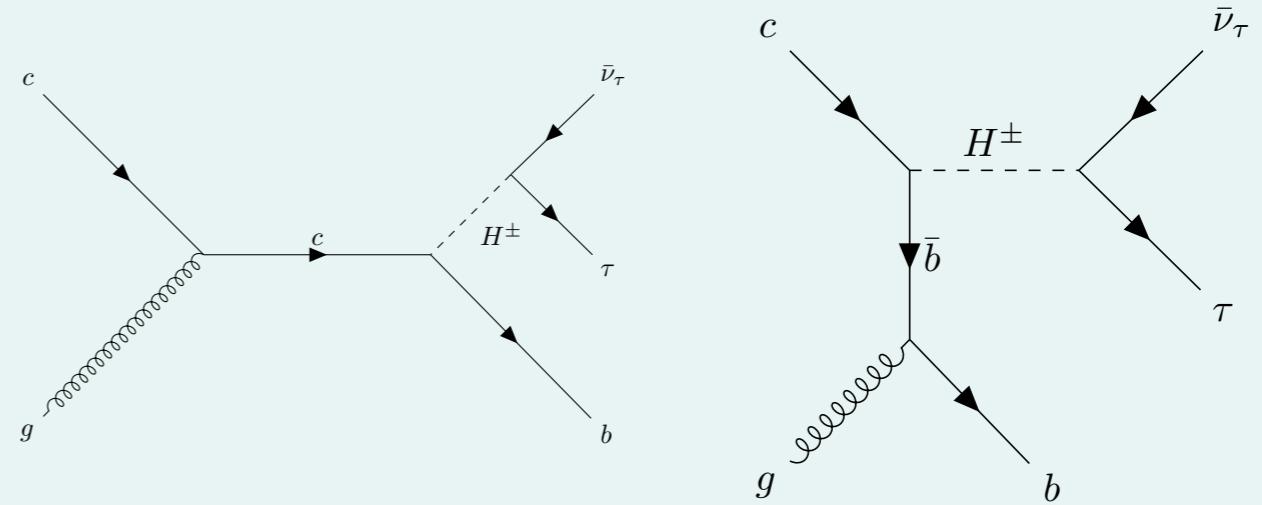


Collider Analysis : b-tag category

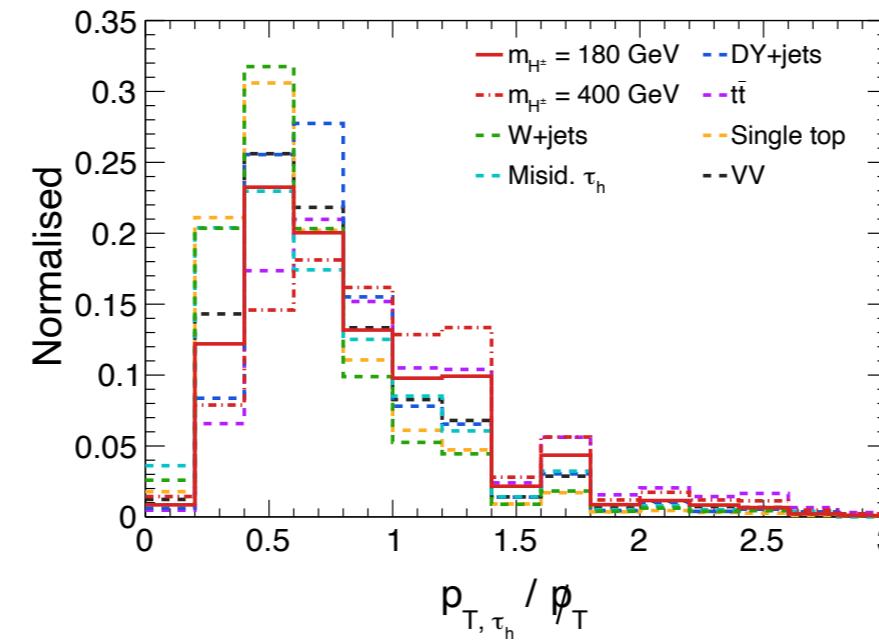
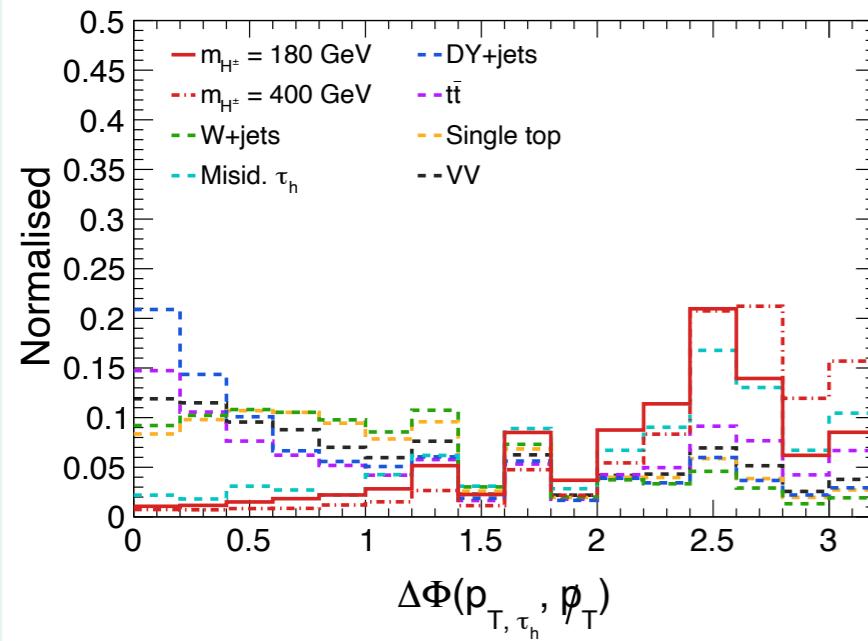
Signal process : $gc \rightarrow bH^\pm \rightarrow b\tau\nu$

b-tag : $N_{b\text{-jets}} = 1$

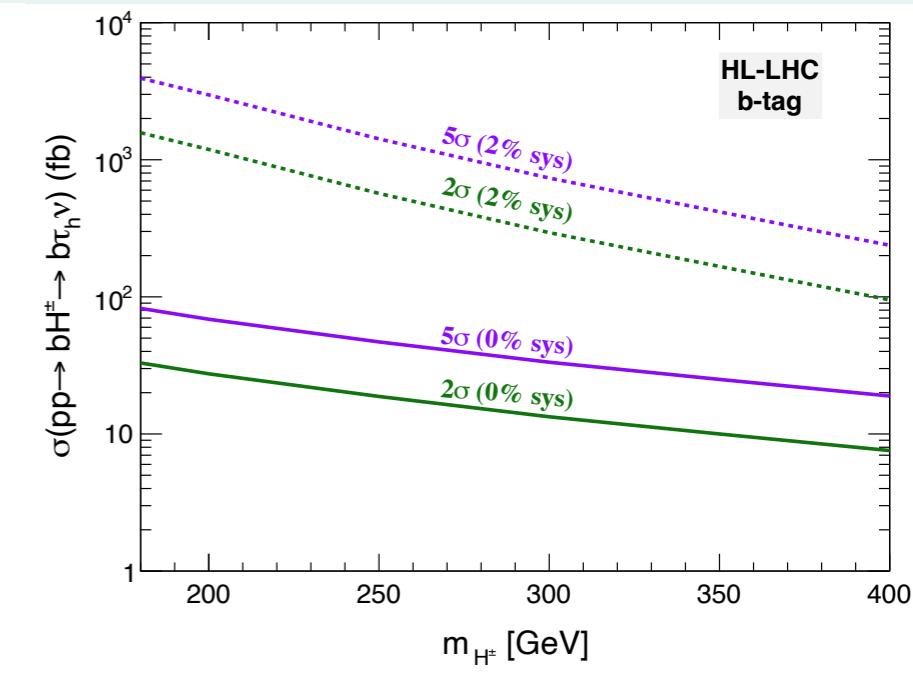
Optimise over $(p_{T,\tau_h}, \Delta\phi(p_{T,\tau_h}, p_T), m_T, E_T)$



Kinematic variables



Upper limit

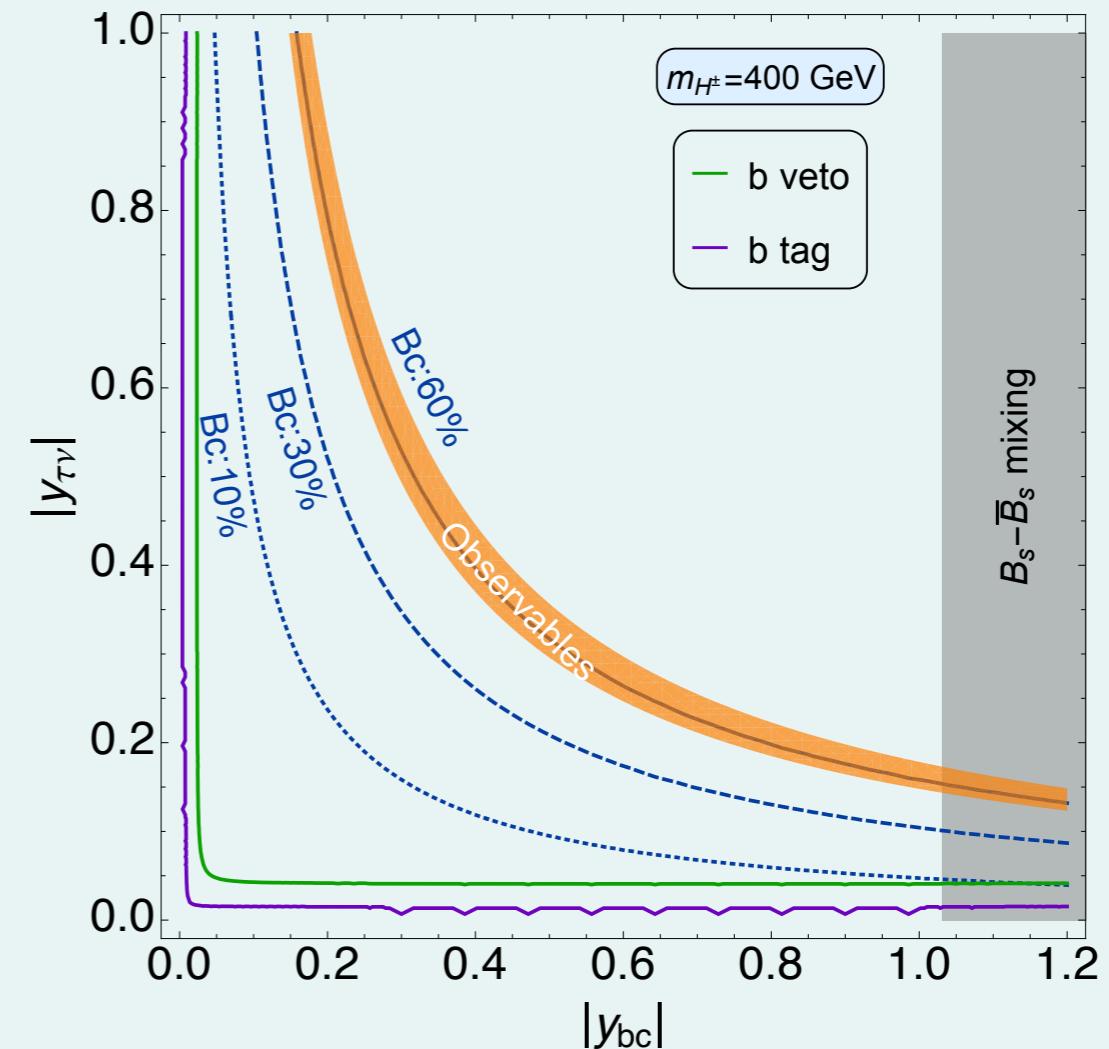
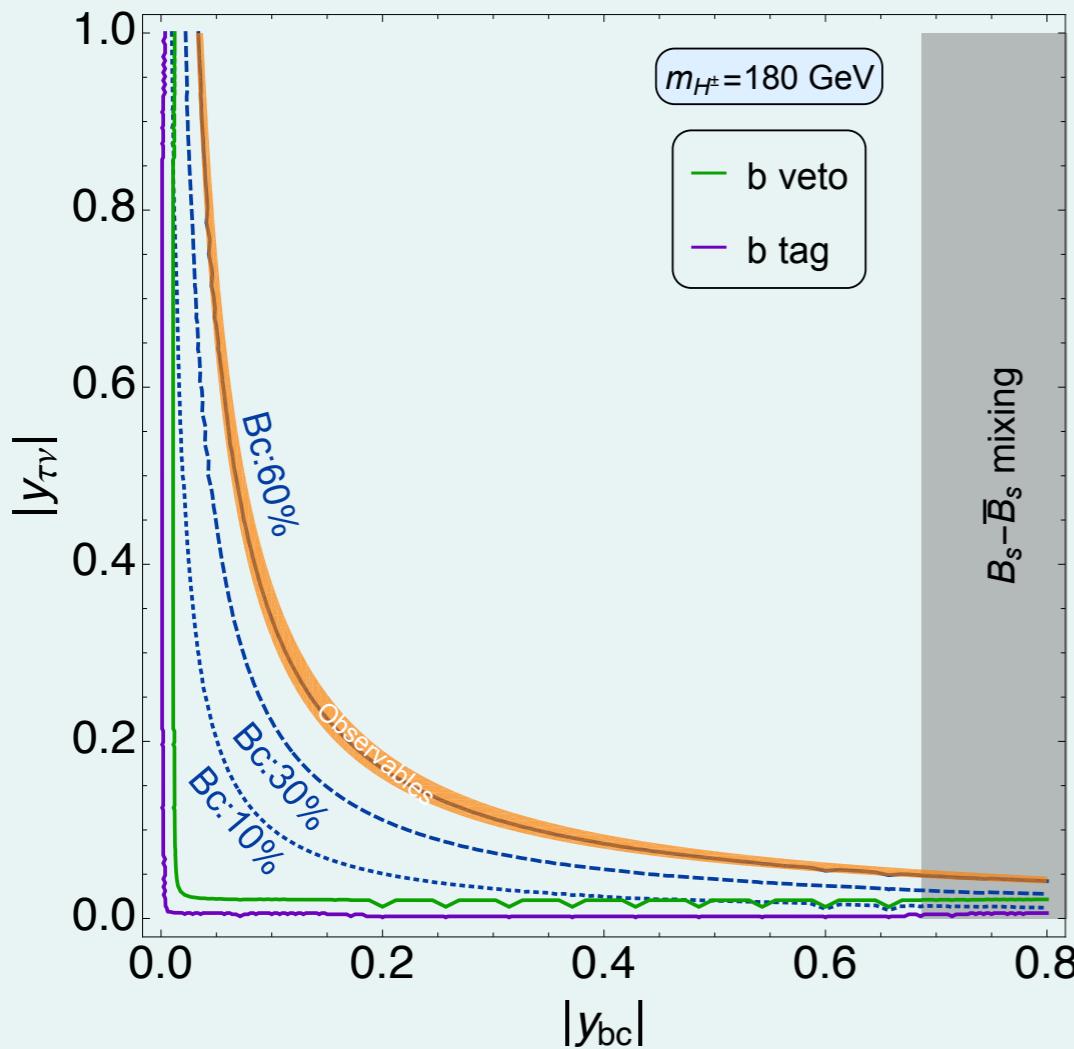


Constraining the Yukawa plane

Orange: 1σ favoured region by all six observable measurements

Green: Constraints from $bc \rightarrow H^\pm \rightarrow \tau\nu$

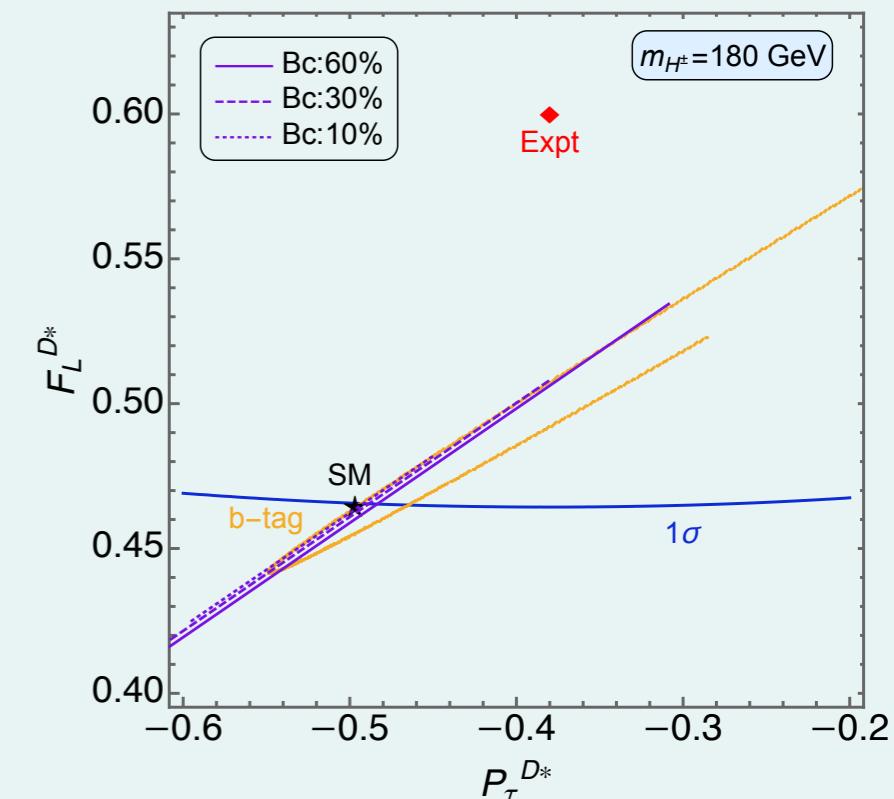
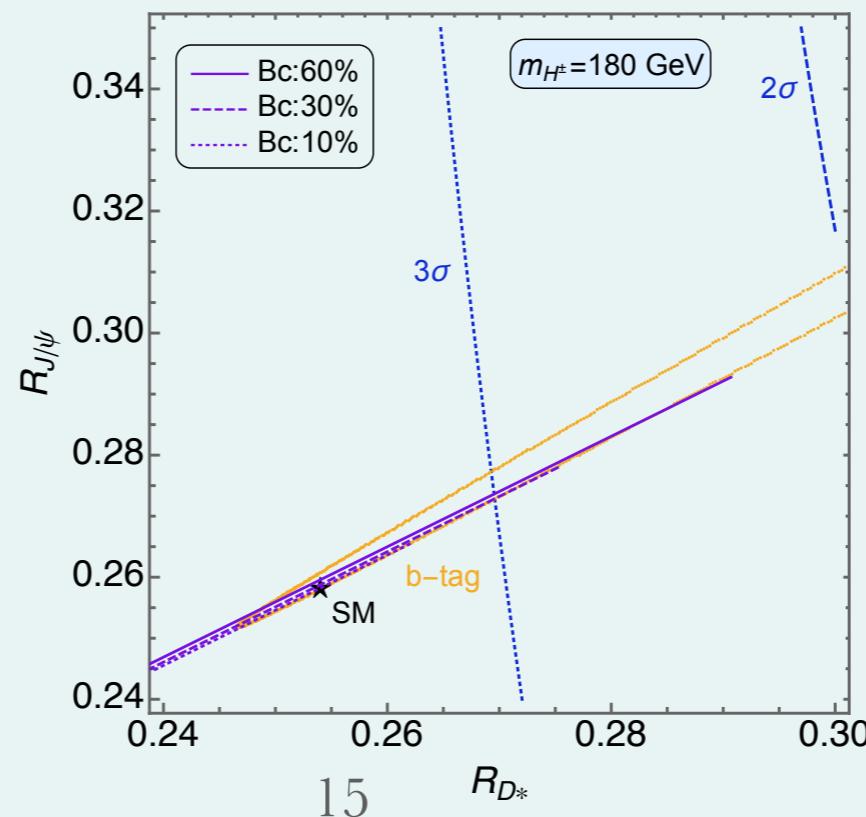
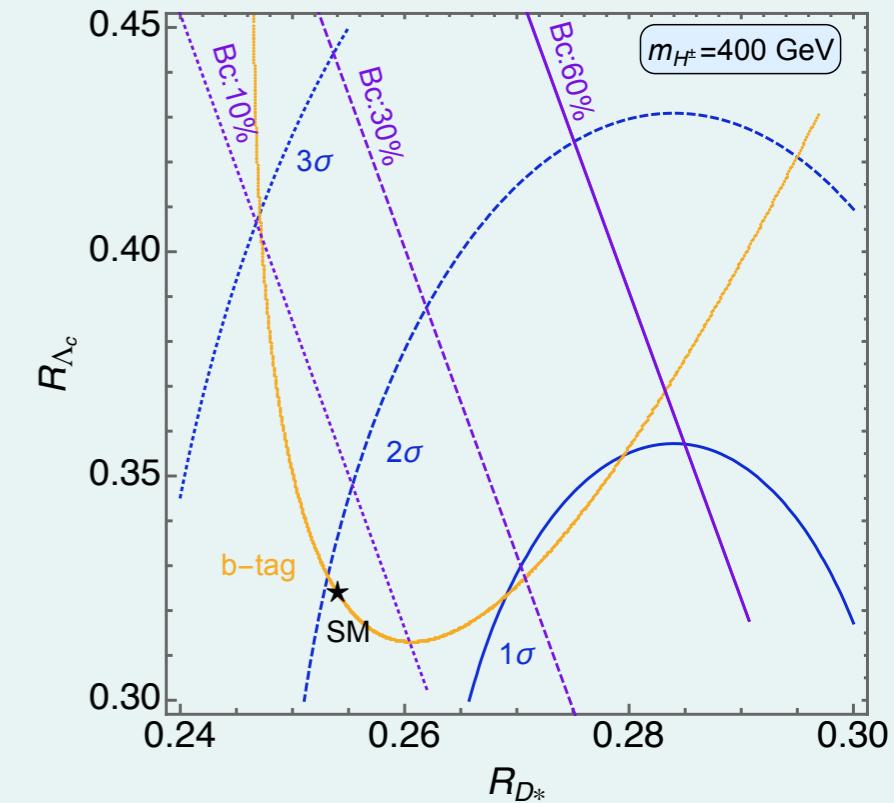
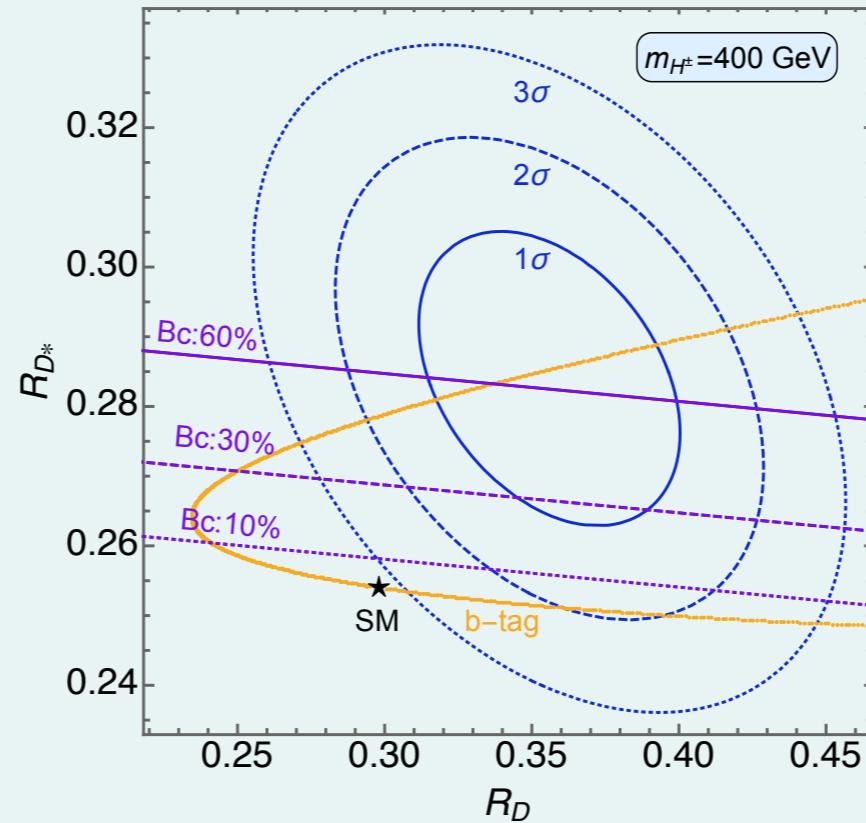
Purple: Constraints from $gc \rightarrow bH^\pm \rightarrow b\tau\nu$



Collider constraint on flavour observables

Preliminary

Orange: reach of HL-LHC
in the b-tag category



Right-handed Neutrino
scalar coupling!

Summary

- ▶ We look at the $b \rightarrow c\tau\nu$ transition anomalies in both the model independent and model dependent way.
- ▶ We fit the scalar couplings to the six $b \rightarrow c\tau\nu$ observables, assuming either real or complex, in 1D and 2D scenario.
- ▶ We show how these scalar coupling fits are in various 2D planes of flavour observables.
- ▶ A collider search in the $\tau\nu$ and $b\tau\nu$ final state is performed, where the b-tag channel gives stronger constraints on the charged Higgs production cross-section.
- ▶ Finally, we use these collider limits to estimate the reach of HL-LHC for interpreting these six anomalies.