

Exact $SU(5)$ Yukawa matrix unification in the General Flavour Violating MSSM

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based on MI, K. Kowalska, JHEP 1504 (2015) 120, arXiv: 1412.8651v2

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1. $SU(5)$ Yukawa matrix unification
2. Minimal Supersymmetric Standard Model
3. chirally-enhanced SUSY threshold corrections
4. off-diagonal soft terms help $\rightarrow$ General Flavour Violating MSSM
5. Phenomenology of Yukawa unification in the GFV MSSM:
 - ▶ 2nd + 3rd generation
 - ▶ 1st + 2nd + 3rd generation

Unification - SU(5) model: matter & Higgs sector

Georgi, Glashow, 1974

$$\underbrace{(\bar{3}, 1, \frac{1}{3})}_{d_R^*} \oplus \underbrace{(1, 2, -\frac{1}{2})}_l = \underbrace{\bar{5}}_{\psi_{\bar{5}}}$$
$$\underbrace{(3, 2, \frac{1}{6})}_q \oplus \underbrace{(\bar{3}, 1, -\frac{2}{3})}_{u_R^*} \oplus \underbrace{(1, 1, 1)}_{e_R^*} = \underbrace{10}_{\psi_{10}},$$

$$W \ni \psi_{10} \mathbf{Y}^{de} \psi_{\bar{5}} H_{\bar{5}} + \psi_{10} \mathbf{Y}^u \psi_{10} H_5$$

Unification - SU(5) model: matter & Higgs sector

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$$W \ni \psi_{10} \mathbf{Y}^{de} \psi_{\bar{5}} H_{\bar{5}} + \psi_{10} \mathbf{Y}^u \psi_{10} H_5$$

$$Y_{ii}^{d, MSSM} = Y_{ii}^{e, MSSM}$$

Gauge coupling unification

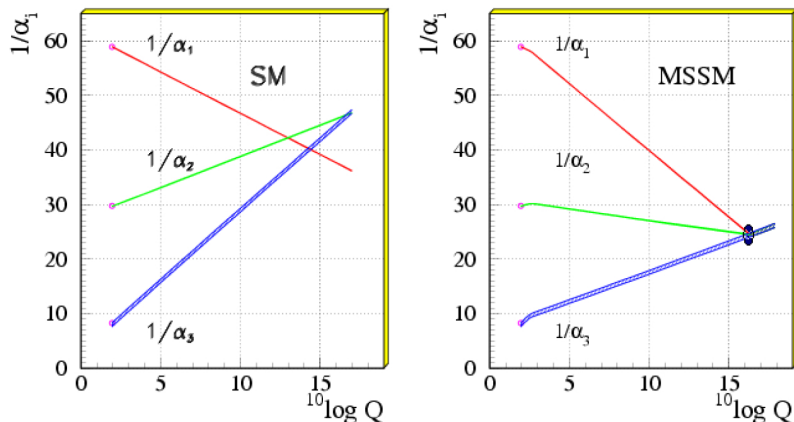
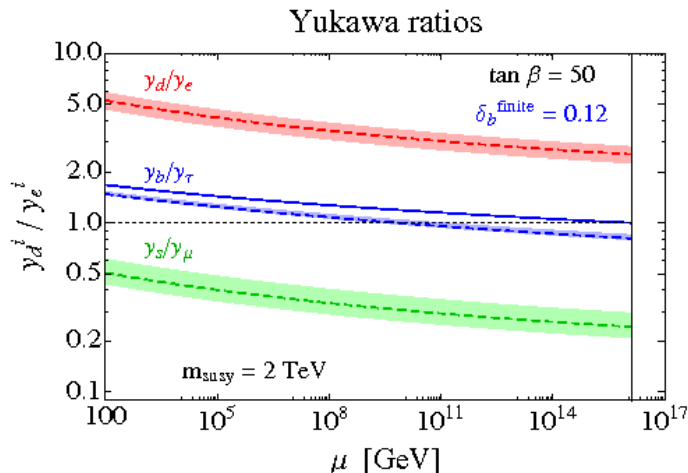


Figure : Gauge coupling unification in non-SUSY GUTs on the left vs. SUSY GUTs on the right using the LEP data (1991)
arXiv: hep-ph/0012288

Yukawa couplings at the GUT scale



Elor, Hall,
Pinner,
Ruderman,
JHEP 1210
(2012) 111,
arXiv:1206.5301

2nd generation: $Y_\mu(M_{\text{GUT}}) \approx 3 Y_s(M_{\text{GUT}})$
1st generation: $Y_e(M_{\text{GUT}}) \approx 1/3 Y_d(M_{\text{GUT}})$

Change the boundary condition at the high scale

- ▶ additional Higgs fields, e.g.

H. Georgi and C. Jarlskog, Phys. Lett. B86 (1979) 297

$$H_5, \quad H_{\bar{5}}, \quad \boxed{H_{45}} \rightarrow Y_\mu = 3Y_s, \quad Y_e = 1/3Y_d$$

- ▶ correction $O(1)$ from higher-dim. operators

D. Emmanuel-Costa and S. Wiesenfeldt, Nucl. Phys. B 661 (2003) 62

S. Antusch and M. Spinrath, Phys. Rev. D 79 (2009) 095004

S. Antusch, S.F. King and M. Spinrath, Phys. Rev. D 89 (2014) 055027

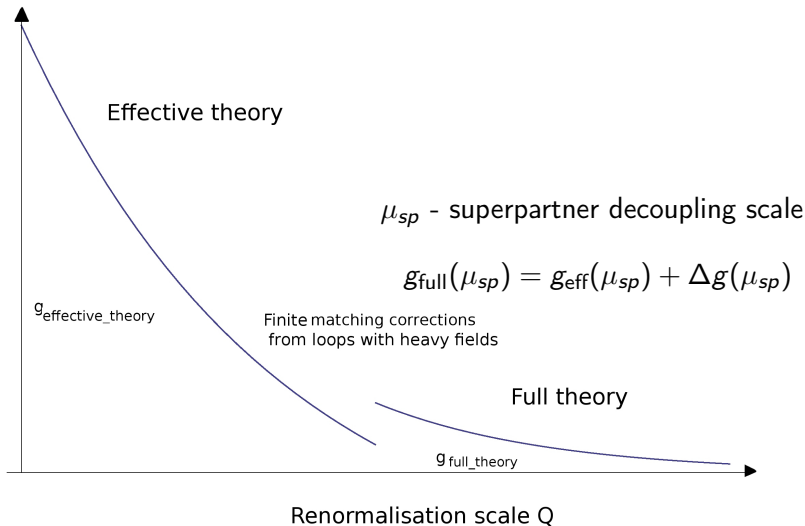
$$W = W_Y + W_{HO} \rightarrow \begin{aligned} Y_d^{ij} &= Y_{de}^{ij} + \frac{1}{M_{\text{PL}}} F^{ij} \\ Y_e^{ij} &= Y_{de}^{ij} + \frac{1}{M_{\text{PL}}} G^{ij} \end{aligned}$$

Manipulate the boundary condition between SM and MSSM - play with threshold corrections

- ▶ Diaz-Cruz, Murayama, Pierce, Phys.Rev.D 65:075011, 2002
(particular ansatz using A -terms for unification)
- ▶ Ts. Enkhbat, arXiv:0909.5597
(general diagonal A -terms)
- ▶ MI, Eur.Phys.J. C75 (2015) 51
(update - new exp results, broader $\tan \beta$ range, weaker impact on flavour observables)

Threshold corrections

Renormalised constant g



A. Crivellin, L. Hofer, J. Rosiek, JHEP 1107 (2011) 017

$$v_f Y_{ii}^{f MSSM} = v_f Y_{ii}^{f SM} - \Sigma_{ii}^f(Y_j^{f'}, \dots).$$

A. Crivellin, L. Hofer, J. Rosiek, *JHEP* 1107 (2011) 017

$$v_f Y_{ii}^{f\,MSSM} = v_f Y_{ii}^{f\,SM} - \Sigma_{ii}^f(Y_j^{f'}, \dots).$$

$$m_i^{d(\ell)\,SM} - v_d Y_{ii}^{d(\ell)\,MSSM} = \Sigma_{ii\mathcal{X}}^{d(\ell)\,LR} + \epsilon_i^{d(\ell)} v_u Y_{ii}^{d(\ell)(0)} + O\left(\frac{v^2}{M_{SUSY}}\right),$$

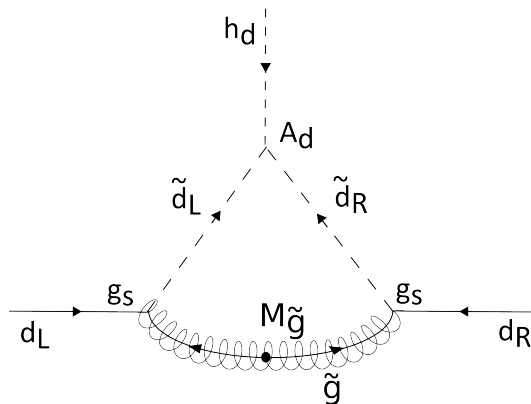
A. Crivellin, L. Hofer, J. Rosiek, *JHEP* 1107 (2011) 017

$$v_f Y_{ii}^{f MSSM} = v_f Y_{ii}^{f SM} - \Sigma_{ii}^f(Y_j^{f'}, \dots).$$

$$m_i^{d(\ell) SM} - v_d Y_{ii}^{d(\ell) MSSM} = \Sigma_{ii}^{d(\ell) LR} + \epsilon_i^{d(\ell)} v_u Y_{ii}^{d(\ell)(0)} + O\left(\frac{v^2}{M_{SUSY}}\right),$$

$$Y_{ii}^{d(\ell) MSSM} = \frac{m_i^{d(\ell) SM} - \Sigma_{ii}^{d(\ell) LR}}{v_d(1 + \tan \beta \cdot \epsilon_i^{d(\ell)})}.$$

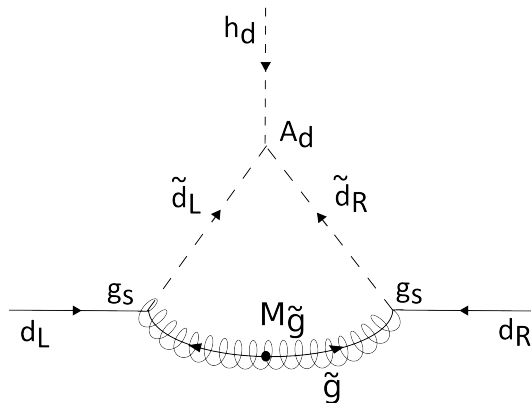
Threshold corrections - example diagrams



- ▶ Diaz-Cruz, Murayama, Pierce, Phys.Rev.D 65:075011, 2002
- ▶ Ts. Enkhbat, arXiv:0909.5597
- ▶ Ml, Eur.Phys.J. C75 (2015) 51

$$(\Sigma_{ii}^d)^{\tilde{g}} \sim \alpha_S m_{\tilde{g}} (v_d A_{ii}^d - v_d Y_{ii}^d \mu \tan \beta)$$

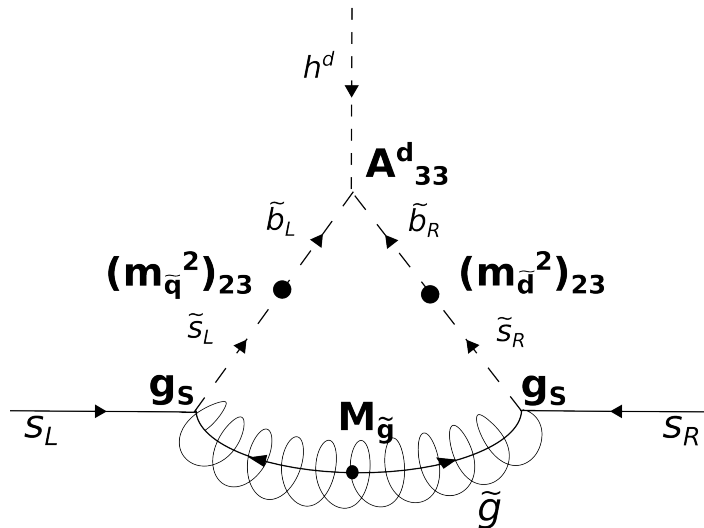
Threshold corrections - example diagrams



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- ▶ Ml, Eur.Phys.J. C75 (2015) 51

$A_s \sim m_{\tilde{s}}$ required for strange-muon unification
 $\Rightarrow$ MSSM vacuum metastable

Threshold corrections - example diagrams



$$(\Sigma_{22}^d)^{\tilde{g}} \sim \alpha_s M_{\tilde{g}} v_d (A_{33}^d - Y_b \mu \tan \beta) (m_{\tilde{q}}^2)_{23} (m_{\tilde{d}}^2)_{23}$$

$$(m_{\tilde{l}}^2)_{ij} = (m_{\tilde{d}}^2)_{ij} \equiv (m_{\text{dl}}^2)_{ij}$$

$$(m_{\tilde{q}}^2)_{ij} = (m_{\tilde{u}}^2)_{ij} = (m_{\tilde{e}}^2)_{ij} \equiv (m_{\text{ue}}^2)_{ij}$$

$$A_{ij}^d = A_{ij}^e \equiv A_{ij}^{\text{de}}$$

$$A_{ij}^u$$

$$M_1 = M_2 = M_3 \equiv M_{1/2},$$

$$\tan \beta = \frac{v_u}{v_d}$$

$$m_{\text{H}_u}^2, \quad m_{\text{H}_d}^2$$

BayesFITSv.3.2 A. Fowlie, M. Kazana, K. Kowalska, S. Munir, L. Roszkowski, E. M. Sessolo, S. Trojanowski, Y. L. S. Tsai [arXiv:1206.0264], K. Kowalska [arXiv:1406.0710]

MultiNest v2.7

F. Feroz, M. P. Hobson and M. Bridges, [arXiv:0809.3437]

SUSY_Flavor v2.10

A. Crivellin, J. Rosiek,
P. H. Chankowski, A. Dedes,
S. Jaeger and P. Tanedo
[arXiv:1203.5023]

SPheno v3.3.3

W. Porod and F. Staub [arXiv:1104.1573]

HIGGSBounds v4.0.0 HIGGSSIGNALS v1.0.0

P. Bechtle et al. [arXiv:0811.4169]
[arXiv:1102.1898], [arXiv:1311.0055],
[arXiv:1305.1933]

DarkSUSY v5.0.6

P. Gondolo, J. Edsjo,
P. Ullio, L. Bergstrom,
M. Schelke, E. A. Baltz,
[astro-ph/0406204]

Ranges of input parameters

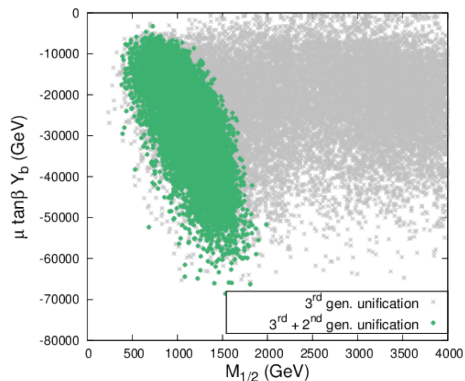
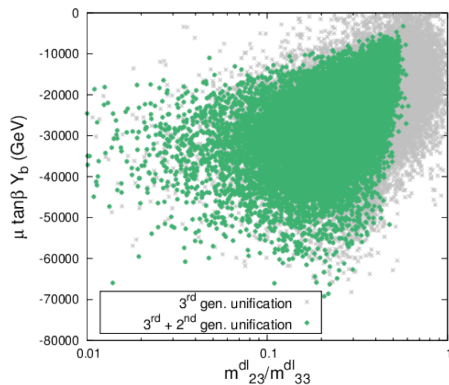
Parameter	Scanning Range
$M_{1/2}$	[100, 4000] GeV
m_{H_u}	[100, 8000] GeV
m_{H_d}	[100, 8000] GeV
$\tan \beta$	[3, 35]
$\text{sgn } \mu$	-1
A_{33}^{de}	[-5000, 5000] GeV
A_{33}^u	[-9000, 9000] GeV
A_{11}^{de}/A_{33}^{de}	[-0.00028, 0.00028]
A_{22}^{de}/A_{33}^{de}	[-0.065, 0.065]
A_{22}^u/A_{33}^u	[-0.005, 0.005]
$m_{ii}^{dl}, i = 1, 2, 3$	[100, 7000] GeV
m_{23}^{dl}/m_{33}^{dl}	[0, 1]
$m_{ii}^{ue}, i = 1, 2, 3$	[100, 7000] GeV

$$m_{ij}^{dl} \equiv \sqrt{(m_{dl}^2)_{ij}}, \quad m_{ij}^{ue} \equiv \sqrt{(m_{ue}^2)_{ij}}.$$



3rd + 2nd family Yukawa unification

relevant GFV parameter: m_{23}^{dl}



Experimental constraints

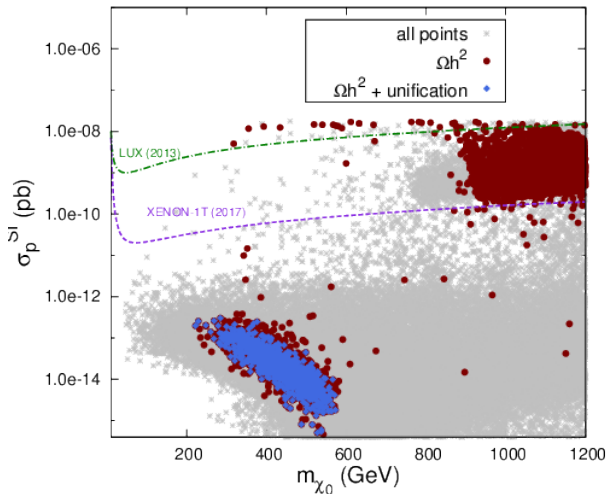
Measurement	Mean or range	Error [exp., th.]
$\Omega_\chi h^2$	0.1199	[0.0027, 10%]
m_h (by CMS)	125.7 GeV	[0.4, 3.0] GeV
$\sin^2 \theta_{\text{eff}}$	0.23155	[0.00012, 0.00015]
M_W	80.385 GeV	[0.015, 0.015] GeV
$\text{BR}(\bar{B} \rightarrow X_s \gamma) \times 10^4$	3.43	[0.22, 0.23]
$\text{BR}(B_s \rightarrow \mu^+ \mu^-) \times 10^9$	2.8	[0.7, 0.23]
$\text{BR}(B_d \rightarrow \mu^+ \mu^-) \times 10^{10}$	3.9	[1.6, 0.2]
$\Delta M_{B_s} \times 10^{11}$	1.1691 GeV	[0.0014, 0.1580] GeV
$\Delta M_{B_d} \times 10^{13}$	3.357 GeV	[0.033, 0.340] GeV
$\Delta M_{B_d} / \Delta M_{B_s} \times 10^2$	2.87	[0.02, 0.14]
$\sin(2\beta)_{\text{exp}}$	0.682	[0.019, 0.003]
$\text{BR}(B_u \rightarrow \tau \nu) \times 10^4$	1.14	[0.27, 0.07]
$\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \times 10^{10}$	1.73	[1.15, 0.04]
$ d_n \times 10^{26}$	< 2.9 e cm	[0, 30%]
$\epsilon_K \times 10^3$	2.228	[0.011, 0.17]



Experimental constraints - Lepton Flavour Violation

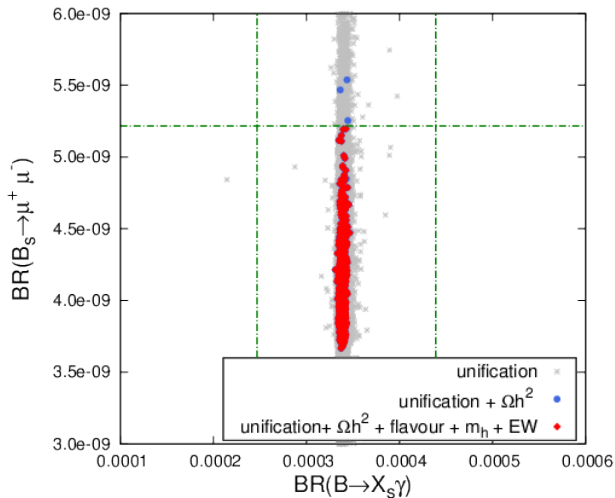
$$\begin{array}{l|l} \text{BR}(\mu^+ \rightarrow e^+ \gamma) \times 10^{13} & < 5.7 \\ \text{BR}(\tau^\pm \rightarrow e^\pm \gamma) \times 10^8 & < 3.3 \\ \text{BR}(\tau^\pm \rightarrow \mu^\pm \gamma) \times 10^8 & < 4.4 \\ \text{BR}(\mu^+ \rightarrow e^+ e^+ e^-) \times 10^{12} & < 1.0 \\ \text{BR}(\tau^\pm \rightarrow e^\pm e^+ e^-) \times 10^8 & < 2.7 \\ \text{BR}(\tau^\pm \rightarrow \mu^\pm \mu^+ \mu^-) \times 10^8 & < 2.1 \end{array}$$

3rd + 2nd family unification: Dark matter



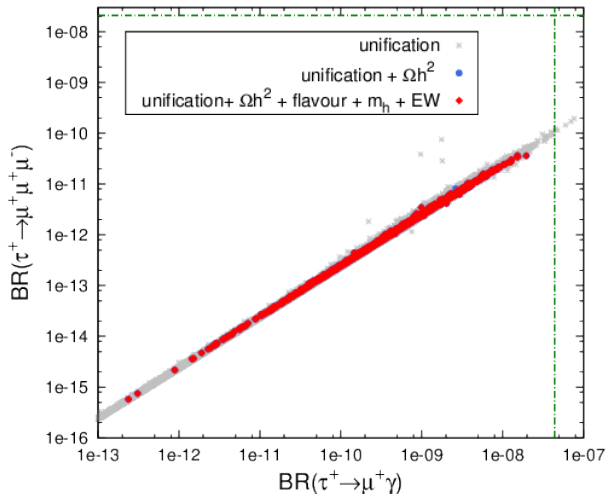
only bino DM

3rd + 2nd family unification: Flavour observables



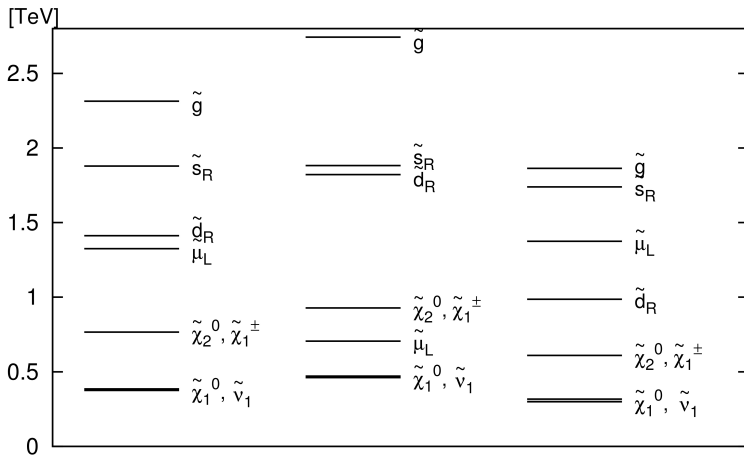
dashed lines - 3σ experimental limits

3rd + 2nd family unification: Flavour observables

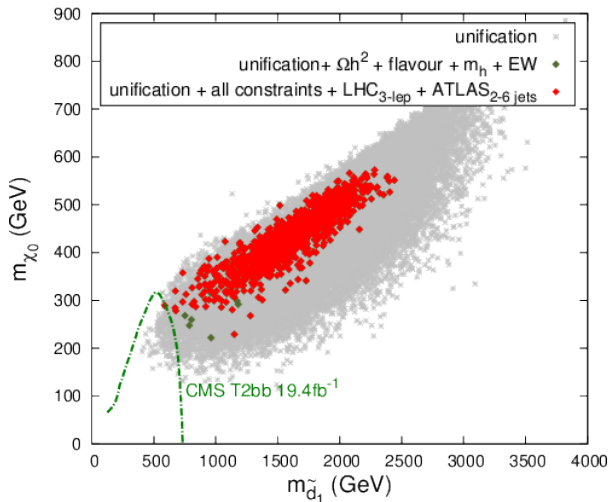


dashed lines - 3σ experimental limits

3rd + 2nd family unification: typical spectra

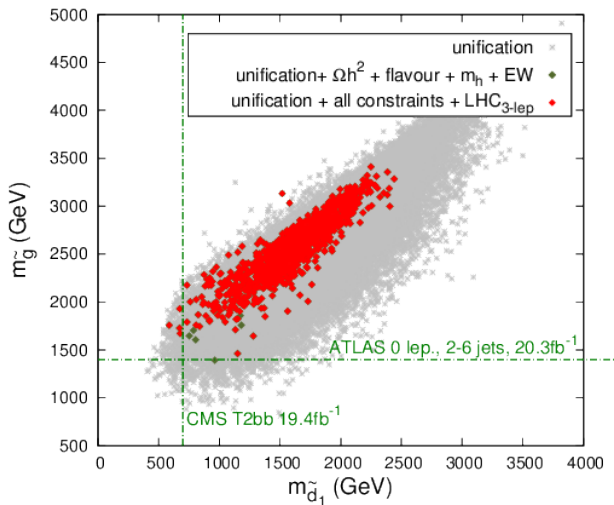


3rd + 2nd family unification: LHC SUSY searches



dashed line - 95% CL CMS sbottom production search

3rd + 2nd family unification: LHC SUSY searches

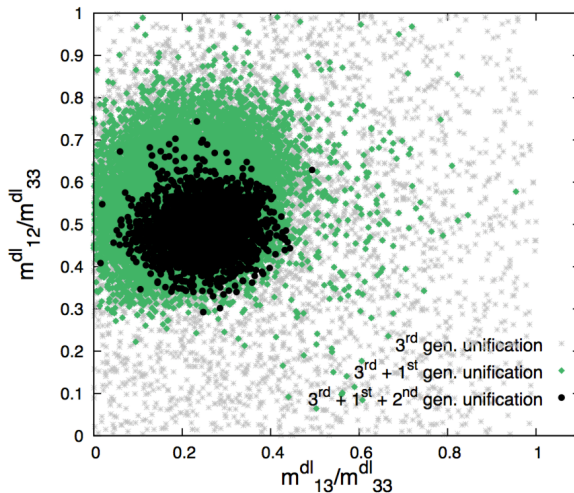


3rd + 2nd + 1st family Yukawa unification: GFV_{123}

Parameter	Scanning Range
$M_{1/2}$	[100, 4000] GeV
m_{H_u}	[100, 8000] GeV
m_{H_d}	[100, 8000] GeV
$\tan \beta$	[3, 45]
$\text{sgn } \mu$	-1
A_{33}^{de}	[0, 5000] GeV
A_{33}^u	[-9000, 9000] GeV
A_{11}^{de}/A_{33}^{de}	[-0.00028, 0.00028]
A_{22}^{de}/A_{33}^{de}	[-0.065, 0.065]
A_{22}^u/A_{33}^u	[-0.005, 0.005]
$A_{ij}^{de}/A_{33}^{de}, i \neq j$	[-0.5, 0.5]
$m_{ii}^{dl}, i = 1, 2, 3$	[100, 7000] GeV
m_{23}^{dl}/m_{33}^{dl}	[0, 1]
m_{13}^{dl}/m_{33}^{dl}	[0, 1]
m_{12}^{dl}/m_{33}^{dl}	[0, 1]
$m_{ii}^{ue}, i = 1, 2, 3$	[100, 7000] GeV

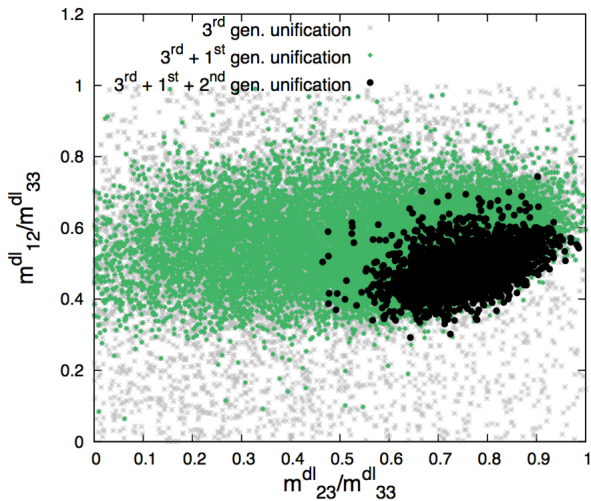
3rd + 2nd + 1st family Yukawa unification

relevant GFV parameters: m_{23}^{dl} , m_{13}^{dl} , m_{12}^{dl} , A_{12}^{de}



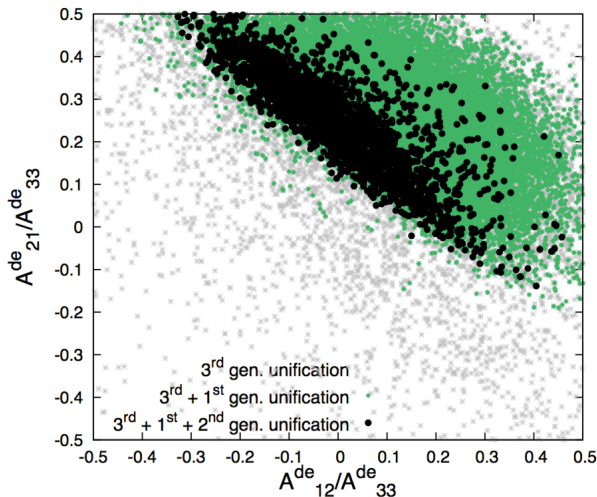
3rd + 2nd + 1st family Yukawa unification

relevant GFV parameters: m_{23}^{dl} , m_{13}^{dl} , m_{12}^{dl} , A_{12}^{de}



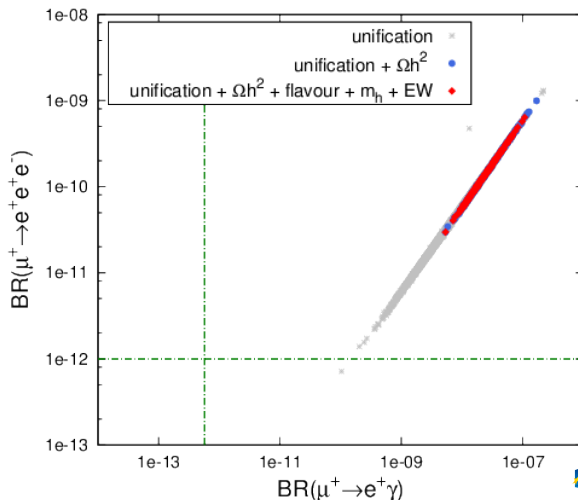
3rd + 2nd + 1st family Yukawa unification

relevant GFV parameters: m_{23}^{dl} , m_{13}^{dl} , m_{12}^{dl} , A_{12}^{de}



3rd + 2nd + 1st family unification: LFV

- ▶ consistent with quark flavour observables
- ▶ **strongly disfavoured** by the Lepton Flavour Violating observables



- ▶ Are there other regions consistent with Yukawa unification?
- ▶ Could the exclusion of GFV_{123} Yukawa unification be avoided?
e.g. much higher SUSY masses,
an $SU(5)$ GUT scenario with $m_{\tilde{l}} \neq m_{\tilde{d}}$
- ▶ Could two-loop threshold corrections be any relevant?
- ▶ $Y_d = Y_e$ in a GFV_{23} -like scenario without vacuum metastability?

Non-trivial flavour structure of the MSSM

can facilitate the SU(5) Yukawa matrix unification

- ▶ Unification of the 2nd and 3rd generation phenomenologically allowed (relevant parameter: $(m_{dl}^2)_{23}$)
- ▶ Full unification of all three generations is strongly disfavoured by the limits on LFV (problems with: $(m_{dl}^2)_{12}$, $A_{12/21}^{de}$)

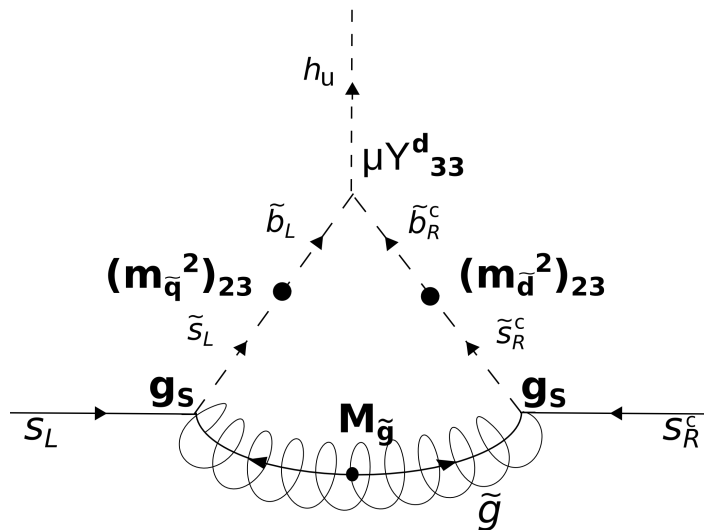
Supplementary slides



INNOVATIVE ECONOMY
NATIONAL COMPETITIVENESS REPORT



Threshold corrections - example diagrams



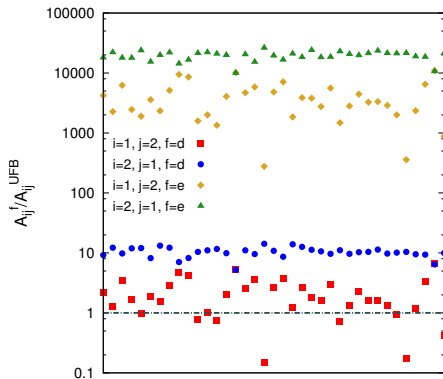
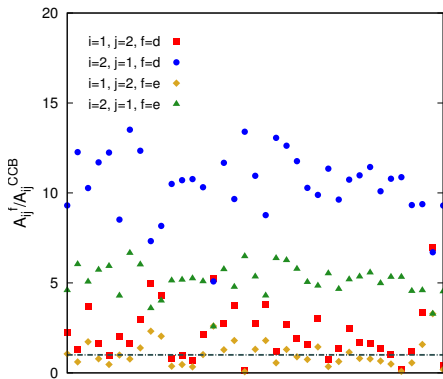
In the down-squark sector, Tree-level formulae for the CCB and UFB bounds in the down-squark sector:

$$(v_d/\sqrt{2})A_{ij}^d \leq m_k^d[(m_{\tilde{q}}^2)_{ii} + (m_{\tilde{d}}^2)_{jj} + m_{H_d}^2 + \mu^2]^{1/2}, \quad k = \text{Max}(i, j)$$

$$(v_d/\sqrt{2})A_{ij}^d \leq m_k^d[(m_{\tilde{q}}^2)_{ii} + (m_{\tilde{d}}^2)_{jj} + (m_{\tilde{l}}^2)_{ii} + (m_{\tilde{e}}^2)_{jj}]^{1/2}$$

J. A. Casas and S. Dimopoulos, [hep-ph/9606237]

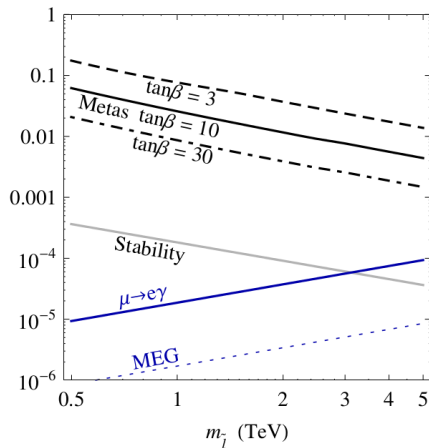
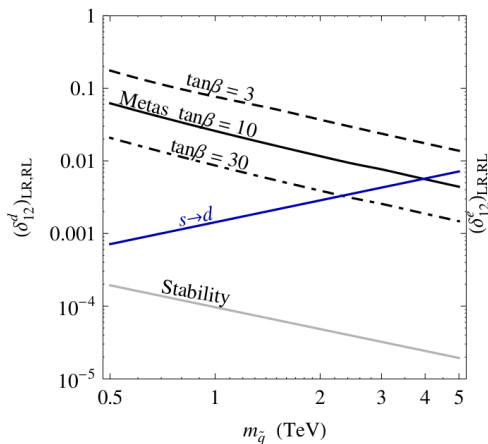
EW vacuum stability



EW vacuum CCB (a) and UFB (b) upper bounds (dashed) on the elements $A_{12/21}^{d,e}$

EW vacuum stability

J. h. Park, [arXiv:1011.4939]:



metastability bounds are 2-3 orders of magnitude weaker.

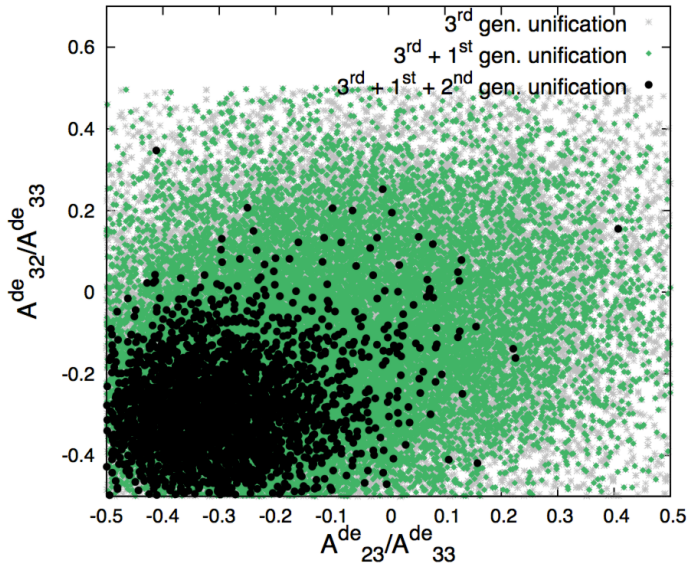
Constants values

we scanned over $(m_t^{\text{pole}}, m_b^{\overline{MS}}(m_b), \alpha_{\text{em}}^{-1}(M_Z)$ and $\alpha_s^{\overline{MS}}(M_Z))$ ($\bar{\rho}$, $\bar{\eta}$, A , λ)

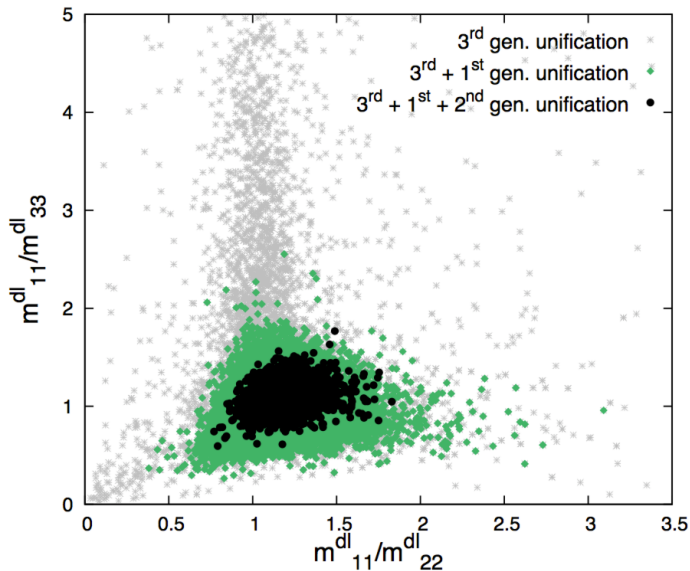
m_t^{pole}		$m_b^{\overline{MS}}(m_b)$		$\alpha_s^{\overline{MS}}(M_Z)$		$\alpha_{\text{em}}^{-1}(M_Z)$	
$173.34 \pm 0.76 \text{ GeV}$		$4.18 \pm 0.03 \text{ GeV}$		0.1184 ± 0.0007		127.944 ± 0.015	
$m_u^{\overline{MS}}$	$m_d^{\overline{MS}}$	$m_s^{\overline{MS}}$	$m_c^{\overline{MS}}(m_c)$	m_e^{pole}	m_μ^{pole}	m_τ^{pole}	M_Z^{pole}
2.3 MeV	4.8 MeV	95 MeV	1.275 GeV	511 keV	106 MeV	1.777 GeV	91.19 GeV
$\bar{\rho}$		$\bar{\eta}$		A		λ	
0.159 ± 0.045		0.363 ± 0.049		0.802 ± 0.020		0.22535 ± 0.0006	

Table : Standard Model parameters (PDG 2014) used in our numerical calculations. The light (u , d , s) quark masses are $\overline{MS}$ -renormalized at 2 GeV.

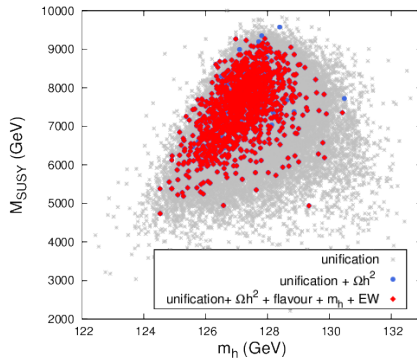
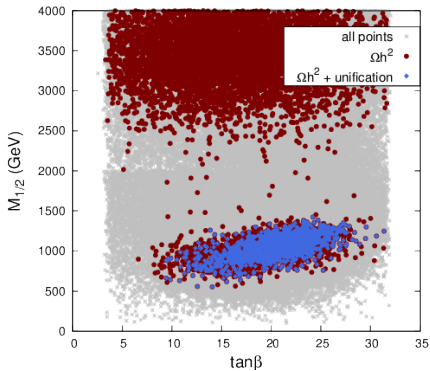
Yukawa unification



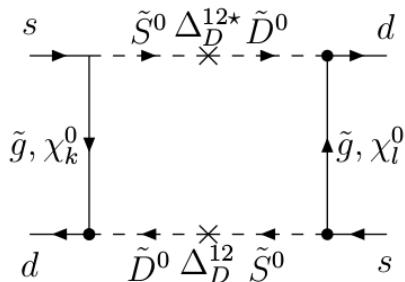
Yukawa unification



Dark matter & Higgs mass



Kaon and B mixing

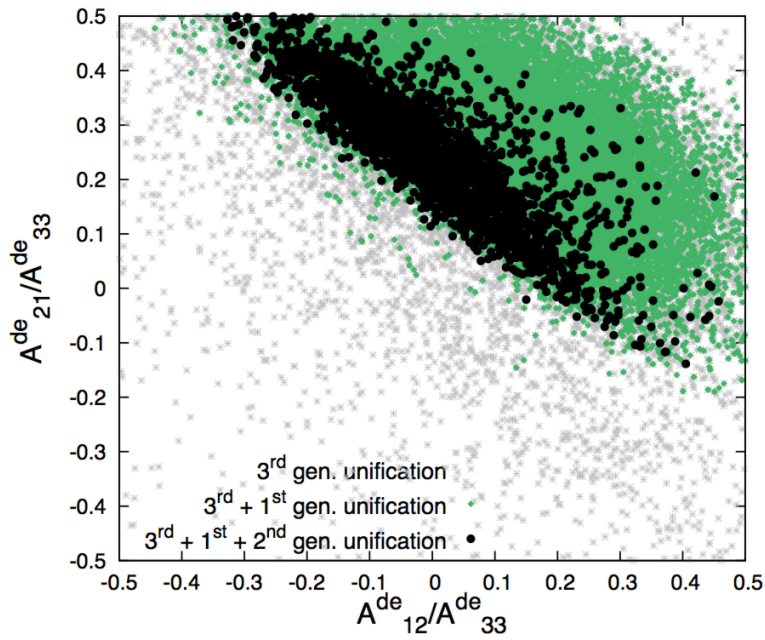


$$\Delta M_{B_{d(s)}} = 2 \left| \langle \bar{B}_{d(s)}^0 | H_{\text{eff}}^{\Delta B=2} | B_{d(s)}^0 \rangle \right|$$

$$\varepsilon_K = \frac{\exp(i\pi/4)}{\sqrt{2}\Delta M_K} \Im \langle \bar{K}^0 | H_{\text{eff}}^{\Delta S=2} | K^0 \rangle$$

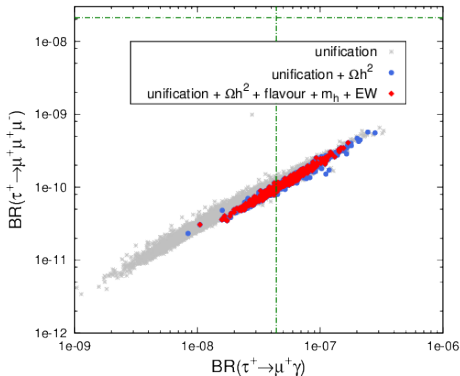
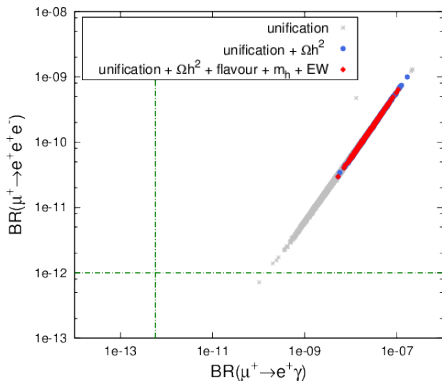
$\Delta_D^{12} = m_{12}^d$ in super-CKM basis

Misiak, Pokorski, Rosiek,
hep-ph/9703442



3rd + 2nd + 1st family unification: LFV

- ▶ consistent with quark flavour observables
- ▶ **strongly disfavoured** by the Lepton Flavour Violating observables



Parameter	Scanning Range	Parameter	Scanning Range
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m_{H_d}	[100, 8000] GeV	m_{13}^{dl}/m_{33}^{dl}	[0, 1]
$\tan \beta$	[3, 45]	m_{12}^{dl}/m_{33}^{dl}	[0, 1]
$\text{sgn } \mu$	-1	$m_{ii}^{ue}, i = 1, 2, 3$	[100, 7000] GeV
A_{33}^{de}	[0, 5000] GeV	Table : Ranges of the input SUSY parameters used in our initial scan. Several omitted soft SUSY-breaking parameters at the GUT scale (namely A_{11}^u as well as A_{ij}^u and m_{ij}^{ue} for $i \neq j$) have been set to zero.	
A_{33}^u	[-9000, 9000] GeV		
A_{11}^{de}/A_{33}^{de}	[-0.00028, 0.00028]		
A_{22}^{de}/A_{33}^{de}	[-0.065, 0.065]		
A_{22}^u/A_{33}^u	[-0.005, 0.005]		
$A_{ij}^{de}/A_{33}^{de}, i \neq j$	[-0.5, 0.5]		

Minimal Supersymmetric Standard Model

Superfields	Fermions	Scalars
$Q = \begin{pmatrix} U_L \\ D_L \end{pmatrix}$	$q = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$	$\tilde{q} = \begin{pmatrix} \tilde{u}_L \\ \tilde{d}_L \end{pmatrix}$
U_R	u_R	$\tilde{u}_R$
D_R	d_R	$\tilde{d}_R$
$L = \begin{pmatrix} N \\ E_L \end{pmatrix}$	$l = \begin{pmatrix} \nu \\ e_L \end{pmatrix}$	$\tilde{l} = \begin{pmatrix} \tilde{\nu} \\ \tilde{e}_L \end{pmatrix}$
E_R	e_R	$\tilde{e}_R$
H_d	$\tilde{h}_d = \begin{pmatrix} \tilde{h}_d^0 \\ \tilde{h}_d^- \end{pmatrix}$	$h_d = \begin{pmatrix} h_d^0 \\ h_d^- \end{pmatrix}$
H_u	$\tilde{h}_u = \begin{pmatrix} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{pmatrix}$	$h_u = \begin{pmatrix} h_u^+ \\ h_u^0 \end{pmatrix}$



Yukawa unification - anatomy of the problem

Yukawa interactions in the superpotential of the minimal $SU(5)$ SUSY GUT:

$$\mathcal{W} \ni \psi_{10} \mathbf{Y}^{\text{de}} \psi_5 \overline{H}_5 + \psi_{10} \mathbf{Y}^{\text{u}} \psi_{10} H_5 \quad (0.1)$$

Here $\overline{H}_5$ and H_5 are two Higgs superfields that couple to model's matter fields. The masses of known fermions are thus given by only two independent 3×3 matrices $\mathbf{Y}_{\text{de}}$ and $\mathbf{Y}_{\text{u}}$

the superpotential of MSSM:

$$\mathcal{W}_{MSSM} = QY^u U_R H_u + QY^d D_R H_d + LY^e E_R H_d + \mu H_d H_u.$$

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$$\mathcal{W}_{MSSM} = QY^u U_R H_u + QY^d D_R H_d + LY^e E_R H_d + \mu H_d H_u.$$

the soft supersymmetry-breaking terms:

$$\begin{aligned} \mathcal{L}_{soft}^{MSSM} = & -\frac{1}{2}[m_{\tilde{g}}(\tilde{G}^a)^T C \tilde{G}^a + m_{\tilde{W}}(\tilde{W}^I)^T C \tilde{W}^I + m_{\tilde{B}}\tilde{B}^T C \tilde{B} + h.c.] - m_{h_d}^2 h_d^\dagger \\ & - \tilde{q}^\dagger(\mathbf{m}_{\tilde{q}}^2)\tilde{q} - (\tilde{u}_R)^\dagger(\mathbf{m}_{\tilde{u}}^2)(\tilde{u}_R) - (\tilde{d}_R)^\dagger(\mathbf{m}_{\tilde{d}}^2)(\tilde{d}_R) - \tilde{l}^\dagger(\mathbf{m}_{\tilde{l}}^2)\tilde{l} - (\tilde{e}_R)^\dagger(\mathbf{m}_{\tilde{e}}^2) \\ & + \tilde{q}\mathbf{A}^u \tilde{u}_R h_u + \tilde{q}\mathbf{A}^d \tilde{d}_R h_d + \tilde{l}\mathbf{A}^e \tilde{e}_R h_d + B\mu h_d h_u + \text{h.c.} \end{aligned}$$

Problem's anatomy in SU(5)

In SM and MSSM the fermion masses are independent parameters and are given by 3 Yukawa matrices:

$$Y^u \rightarrow m_u, m_c, m_t$$

$$Y^d \rightarrow m_d, m_s, m_b$$

$$Y^e \rightarrow m_e, m_\mu, m_\tau$$

In the minimal SU(5) Grand Unified Theory the symmetry requires:

$$Y_d = Y_e, Y_s = Y_\mu, Y_b = Y_\tau$$

flavour mixing (CKM matrix can be included in) $\mathbf{Y}_u$