
A Higgs Portal Update

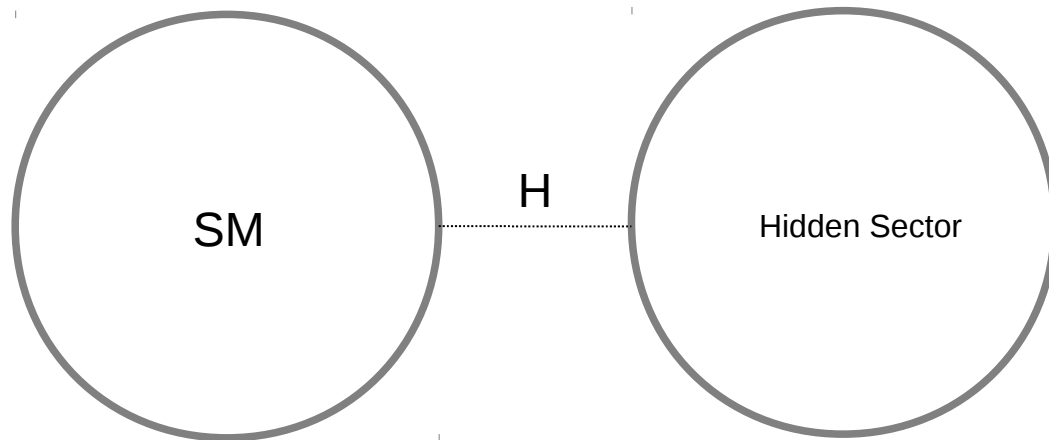
Oleg Lebedev



University of Helsinki

-
- the Higgs and the hidden sector
 - the Higgs and dark matter
 - the Higgs and inflation
-

The Higgs and the hidden sector



Special role of the Higgs :

Silveira, Zee '85
Veltman, Yndurain '89
...

$|H|^2$ = the only gauge and Lorentz-inv. dim-2 operator

$$L = a |H|^2 S^2 + b |H|^2 S$$

(S = “hidden” scalar)

b=0 (S has hidden charge):

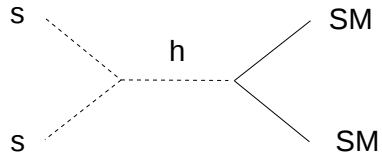
$$L = a |H|^2 S^2$$

“S” is stable and couples weakly to SM

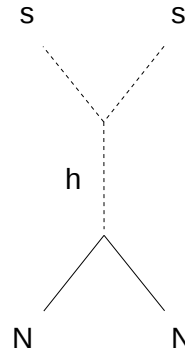


DARK MATTER (?)

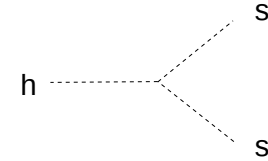
Dark matter:



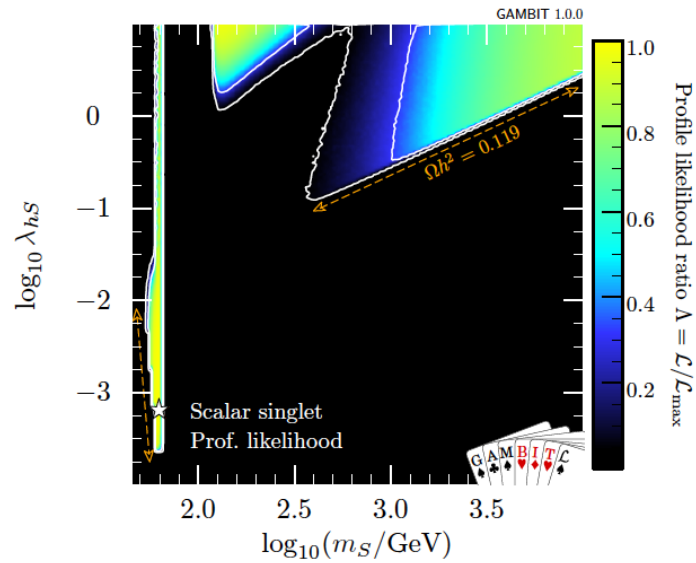
DM annihilation



DM direct detection



Higgs decay



GAMBIT collaboration
1705.07931

white contour = 2σ bound

(Pseudo) Goldstone dark matter

Add a complex scalar S , require softly broken $U(1)$ symmetry:

$$V = V_0 + V_{\text{soft}} ,$$

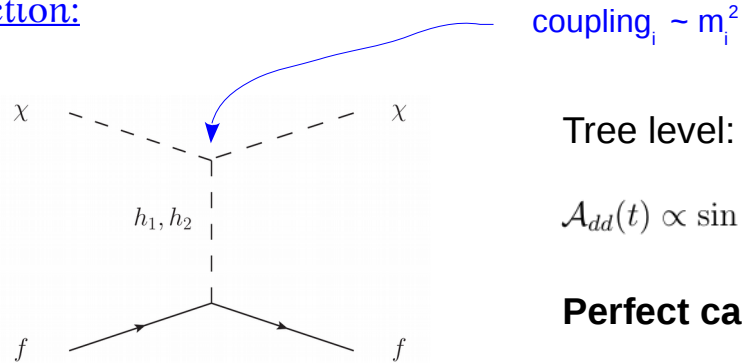
$$V_0 = -\frac{\mu_H^2}{2} |H|^2 - \frac{\mu_S^2}{2} |S|^2 + \frac{\lambda_H}{2} |H|^4 + \lambda_{HS} |H|^2 |S|^2 + \frac{\lambda_S}{2} |S|^4 ,$$

$$V_{\text{soft}} = -\frac{\mu_S'^2}{4} S^2 + \text{h.c.}$$

All parameters are real $\Rightarrow \langle S \rangle = \text{real} , \quad S \rightarrow S^* \text{ symmetry}$

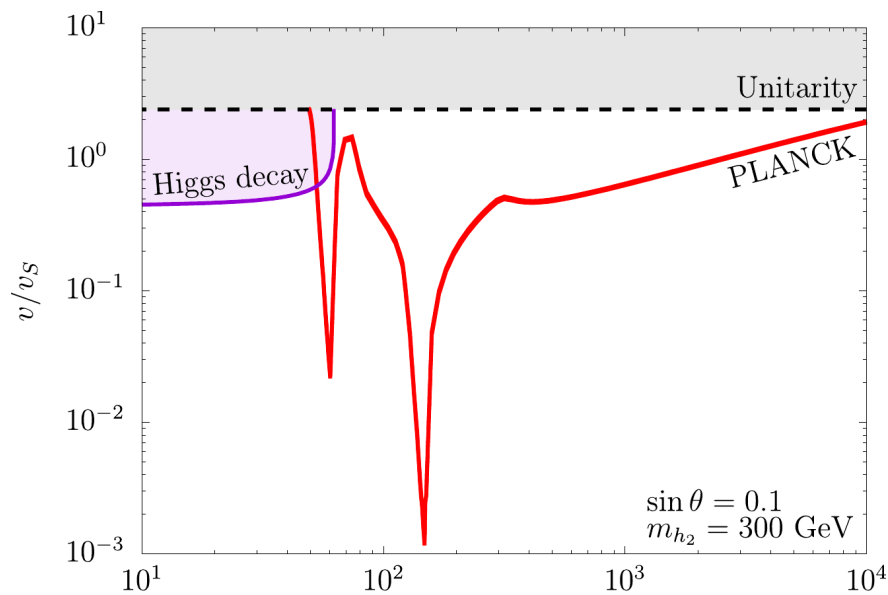
$\text{Im } S = \text{Dark Matter}$

\sim Goldstone boson

Direct detection:

Tree level:

$$\mathcal{A}_{dd}(t) \propto \sin \theta \cos \theta \left(\frac{m_2^2}{t - m_2^2} - \frac{m_1^2}{t - m_1^2} \right) \rightarrow 0$$

Perfect cancellation for any parameter choice !

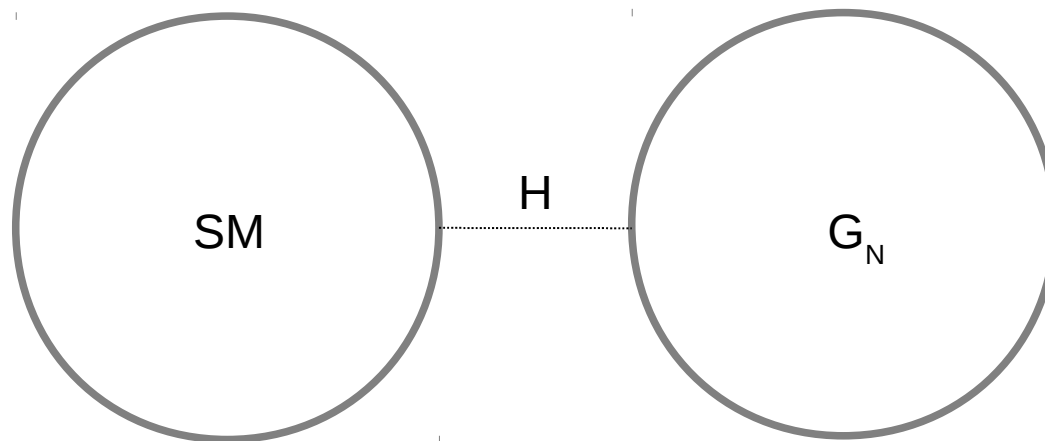
Direct detection = loop-suppressed

Annihilation = unsuppressed

**Excellent WIMP**

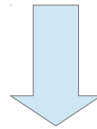
(from 60 GeV to 10 TeV)

The Higgs and "secluded" dark matter



$$V_{\text{portal}} \sim \bar{H} H \bar{S} S$$

Minimal $SU(N)$ breaking



automatic vector dark matter

E.g. $SU(2)$ gauge field DM

Minimal SU(N) Higgsing:

Need N-1 fundamentals, gauge them to

$$\underbrace{\begin{pmatrix} 0 \\ 0 \\ \vdots \\ 0 \\ a \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ \vdots \\ b_1 \\ b_2 \end{pmatrix} \dots \begin{pmatrix} 0 \\ z_1 \\ \vdots \\ \vdots \\ z_{N-1} \end{pmatrix}}_{N-1}$$

Unbroken symmetry: $U_1 = \text{const} \times \text{diag}(e^{iq}, 1, \dots, 1)$

Lightest gauge fields transforming under U_1 are **stable**

For SU(3) and larger groups,

$$m_{\text{DM}} > m_{\text{unstable}}$$

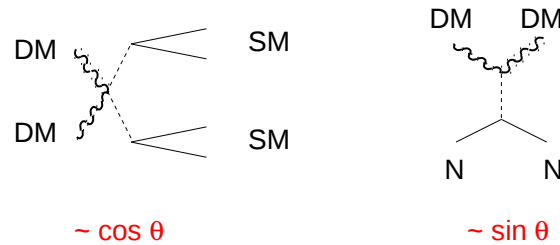
if CP is broken



"secluded" DM

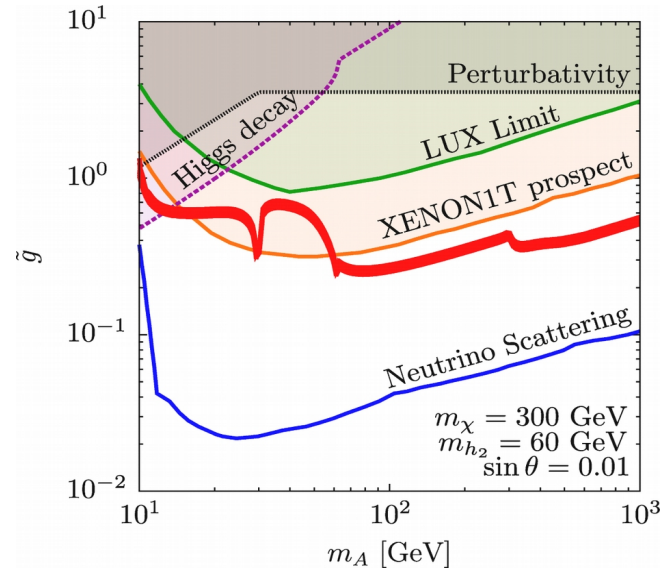
à la Pospelov et al. '07

(also Hambye '08)



Arcadi, Gross, OL, Pokorski, Toma '16

SU(3):



correct relic density

- DM annihilation efficient
- Direct detection suppressed



Higgs portal DM = viable WIMP

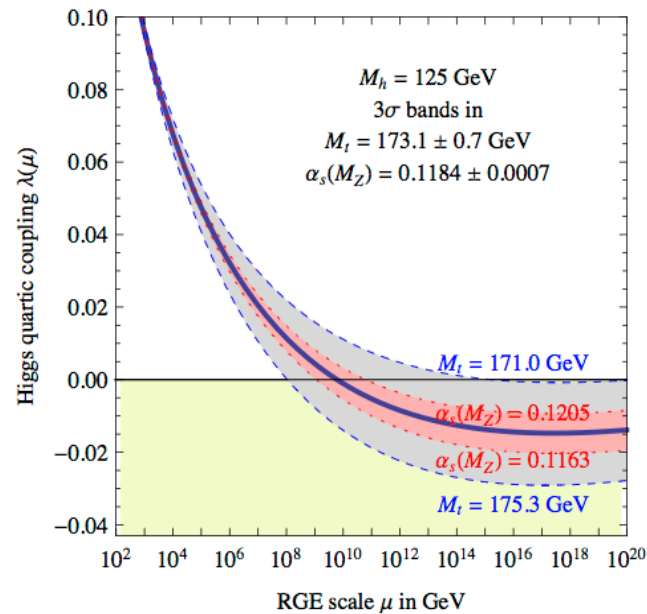
(especially if there's more than one state in the hidden sector)

The Higgs and inflation

Buttazzo et al.'13

SM stability bound:

$$m_h > (129.6 \pm 1.5) \text{ GeV}$$

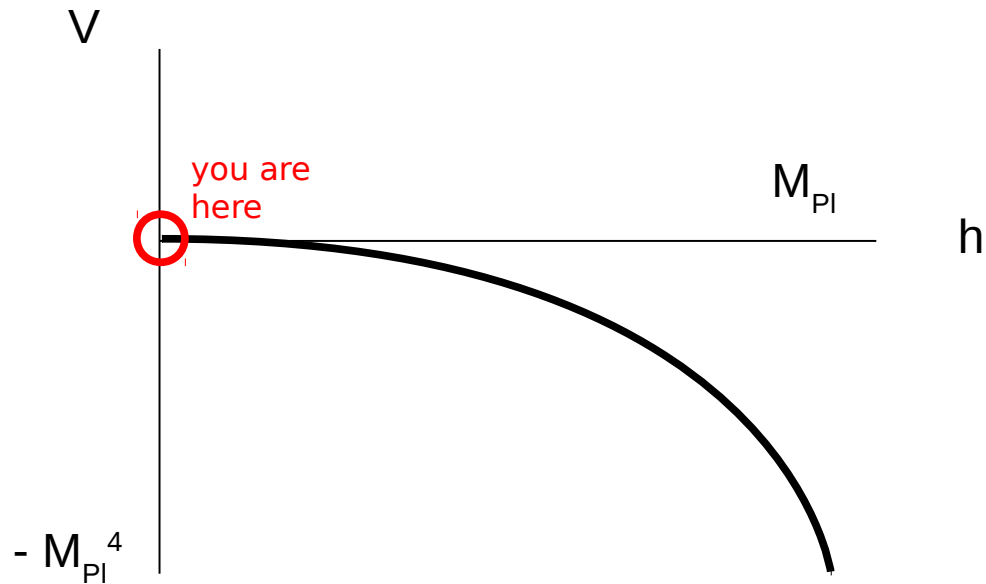


(not settled : Alekhin et al. '12
Bezrukov et al. '12)

$$h \gg \Lambda \sim 10^{10} \text{ GeV}$$



$$V \sim \frac{1}{4} \lambda(h) h^4, \quad \lambda(h) < 0$$



$$\Lambda = 10^{-8} M_{\text{Pl}}$$

,

$$\text{barrier} = 10^{-32} M_{\text{Pl}}^4$$

Problems :

- how did the Universe end up at $h \sim 0$?
- why did it stay there during inflation ?

Solutions :

- modify the Higgs potential during inflation
 - just modify the Higgs potential
-

Vacuum stabilization via Higgs-inflaton mixing

Renormalizable Higgs–inflaton/gravity interactions:

$$\begin{aligned} -\mathcal{L}_{h\phi} &= \lambda_{h\phi} H^\dagger H \phi^2 + 2\sigma H^\dagger H \phi, \\ -\mathcal{L}_{hR} &= \xi_h H^\dagger H \hat{R}. \end{aligned}$$

always leads to Higgs-inflaton mixing

Two mass eigenstates $h_{1,2}$ with mixing angle θ :

$$2\lambda_h v^2 = m_1^2 \cos^2 \theta + m_2^2 \sin^2 \theta$$



λ_h increases for $m_2 > m_1$!

The set-up :

$$S = \int d^4x \sqrt{-\hat{g}} \left[\frac{1}{2} \Omega^2 \hat{R} - \frac{1}{2} \hat{g}^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} \hat{g}^{\mu\nu} \partial_\mu h \partial_\nu h - V(\phi, h) \right]$$

$$\Omega^2 = 1 + \xi_\phi \phi^2 + \xi_h h^2 ,$$

$$V(\phi, h) = \frac{\lambda_h}{4} h^4 - \frac{\mu_h^2}{2} h^2 + \frac{\lambda_{h\phi}}{2} h^2 \phi^2 + \sigma h^2 \phi + \frac{\lambda_\phi}{4} \phi^4 + \frac{b_3}{3} \phi^3 - \frac{\mu_\phi^2}{2} \phi^2 + b_1 \phi ,$$

Inflation is driven by ϕ with $\xi_\phi \gg 1$ (à la Bezrukov-Shaposhnikov) :

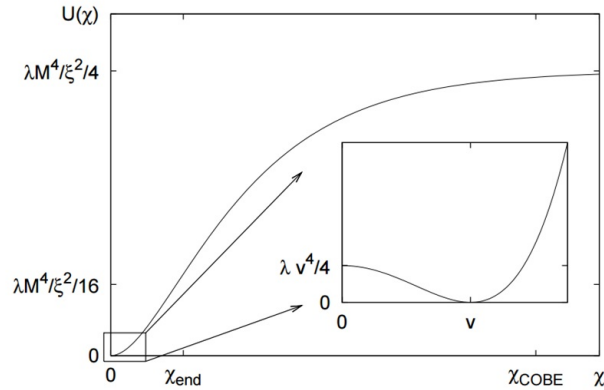


Fig. 1. Effective potential in the Einstein frame.

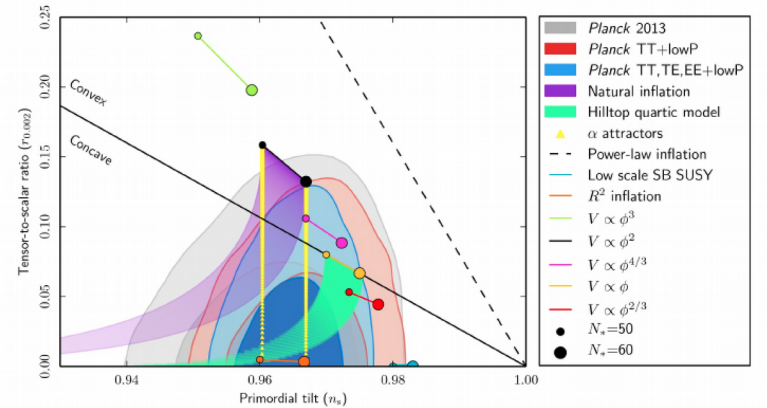


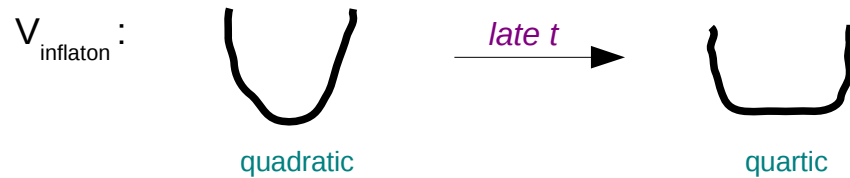
Fig. 12. Marginalized joint 68% and 95% CL regions for n_s and $r_{0.002}$ from *Planck* in combination with other data sets, compared to the theoretical predictions of selected inflationary models.

High energy constraints :

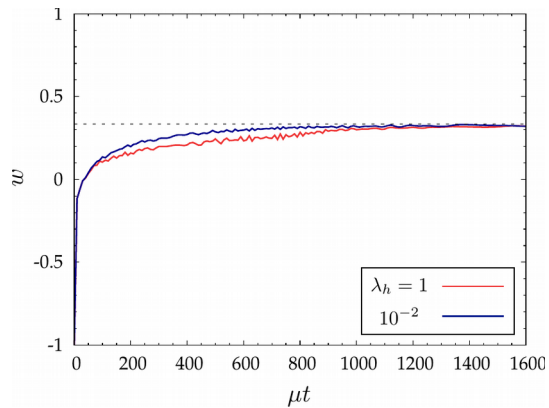
$$\lambda_\phi < 10^{-5} \quad \Rightarrow \quad \text{no unitarity issues}$$
$$\lambda_{h\phi} < 10^{-2} \quad \Rightarrow \quad \text{no large rad. corrections}$$

(also assume a single scale \Rightarrow TeV dimensionful parameters)

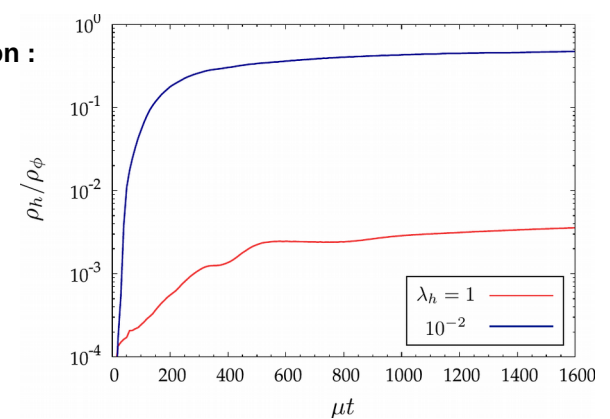
Preheating :



**Equation
of state:**



**Higgs
production :**



Reheating :

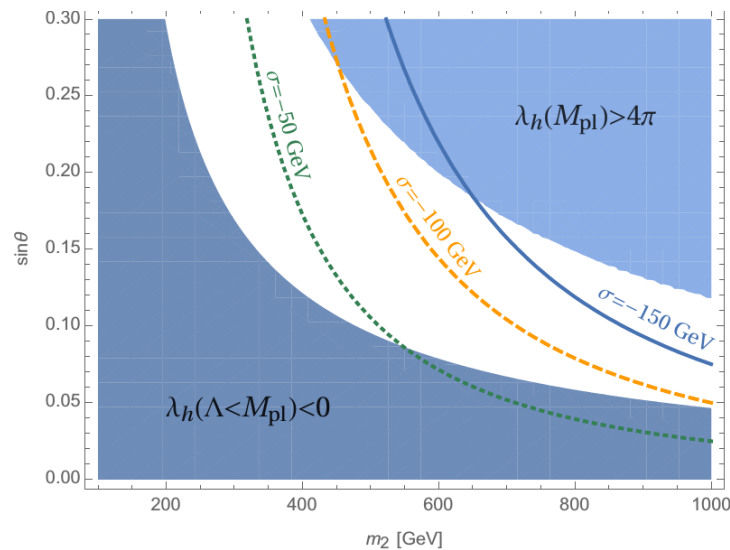
$$\phi + \phi \rightarrow h + h$$



$$T_{\text{reh}} > 10^{12} \text{ GeV}$$

ϕ freezes-out and decays $\phi \rightarrow h + h$

Vacuum stability :



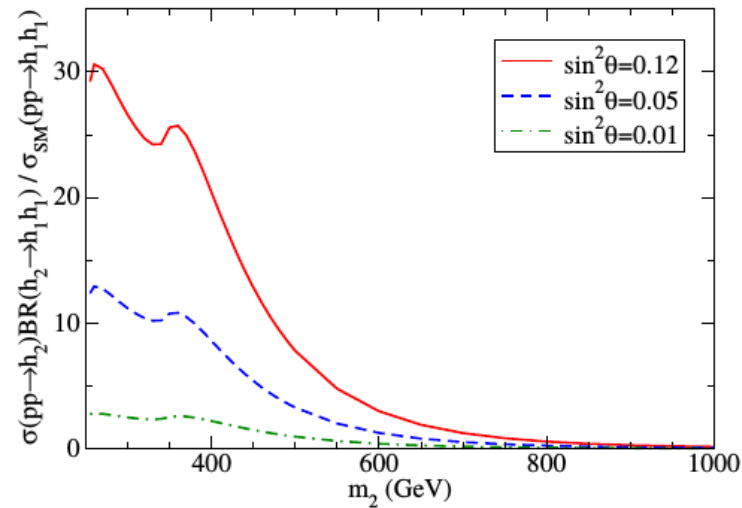
$$\sigma' v = \frac{\sin 2\theta}{4} (m_1^2 - m_2^2)$$

the trilinear interaction is
crucial for the mixing

Inflaton search at LHC :

- Universal Higgs coupling reduction
- Heavy Higgs–like resonance
- Resonant decay $h_2 \rightarrow h_1 h_1$

Lewis, Sullivan '17



Conclusion

- Higgs portal WIMP DM is viable:
 - *Goldstone DM*
 - *Secluded DM*
- Higgs—inflaton mixing
 - *vacuum stability*
 - *inflaton at LHC*