

Charged Composite Scalar Dark Matter

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PLANCK 2017

University of Warsaw

May 24, 2017

based on 1705.xxxxx with R. Balkin, M. Ruhdorfer and A. Weiler

Motivation

- In viable composite Higgs models, Higgs doublet arises as set of (approximate) Goldstone bosons

$$SO(5) \xrightarrow[H]{f} SO(4)$$

- Simple, attractive option for DM: extra Goldstone scalar as WIMP
- Mass and interactions dictated by global symmetry + explicit breaking

Frigerio, Pomarol, Riva, Urbano 2012

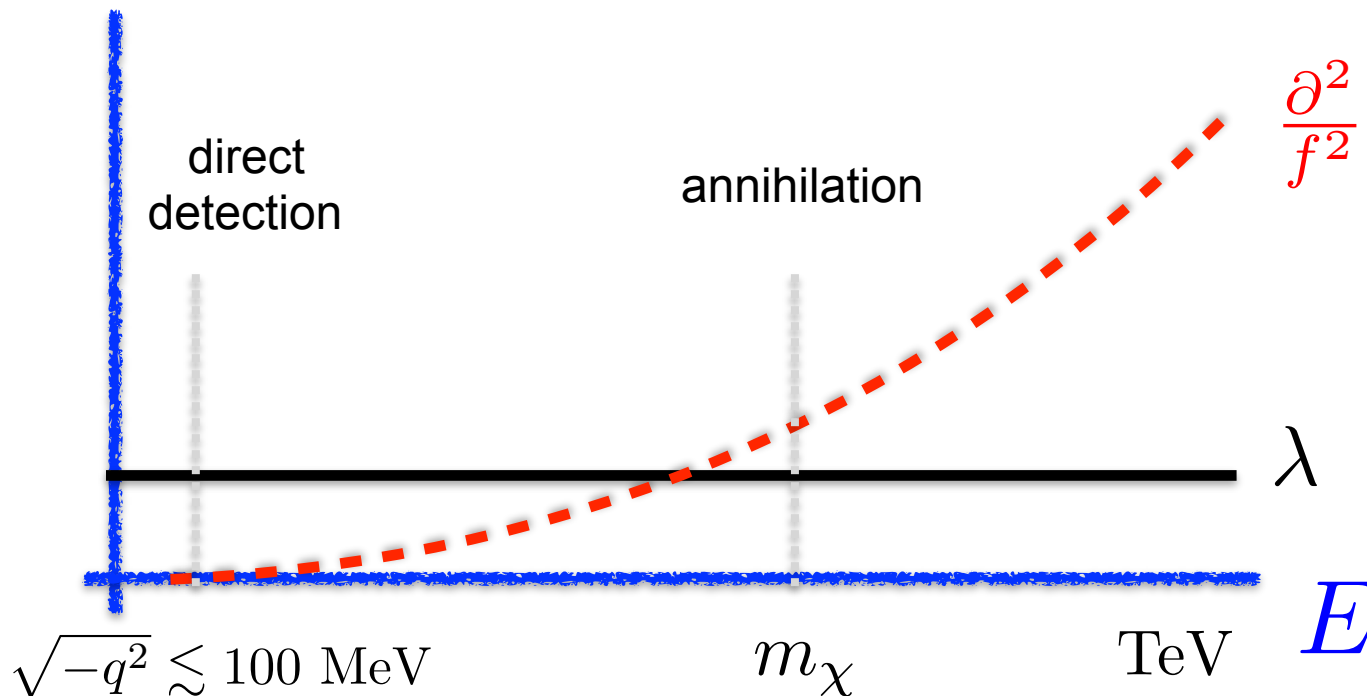
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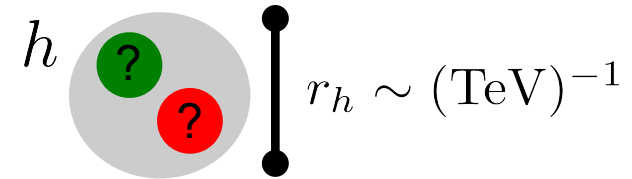
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$$g_{\text{DM-SM}}^2(E)$$



A pseudo-Goldstone composite Higgs

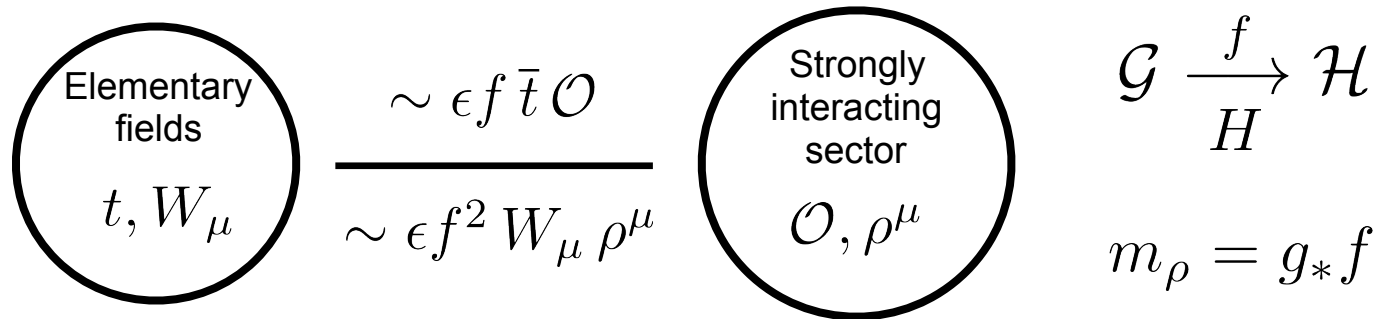
- The Higgs is a bound state of new degrees of freedom
The description of the theory changes above $\sim \text{TeV}$,
Higgs mass naturally 'screened'



$$h \quad \left[r_h \sim (\text{TeV})^{-1} \right]$$

- Take analogy with QCD further: **Higgs** as (approximate) **Goldstone** boson, like pions

Agashe, Contino,
Pomarol 2004



$$\begin{array}{c} \text{Elementary} \\ \text{fields} \\ t, W_\mu \end{array} \quad \begin{array}{c} \sim \epsilon f \bar{t} \mathcal{O} \\ \hline \sim \epsilon f^2 W_\mu \rho^\mu \end{array} \quad \begin{array}{c} \text{Strongly} \\ \text{interacting} \\ \text{sector} \\ \mathcal{O}, \rho^\mu \end{array}$$

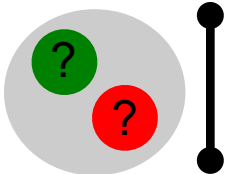
$$\mathcal{G} \xrightarrow[H]{f} \mathcal{H}$$

$$m_\rho = g_* f$$

Kaplan, 1991

A pseudo-Goldstone composite Higgs

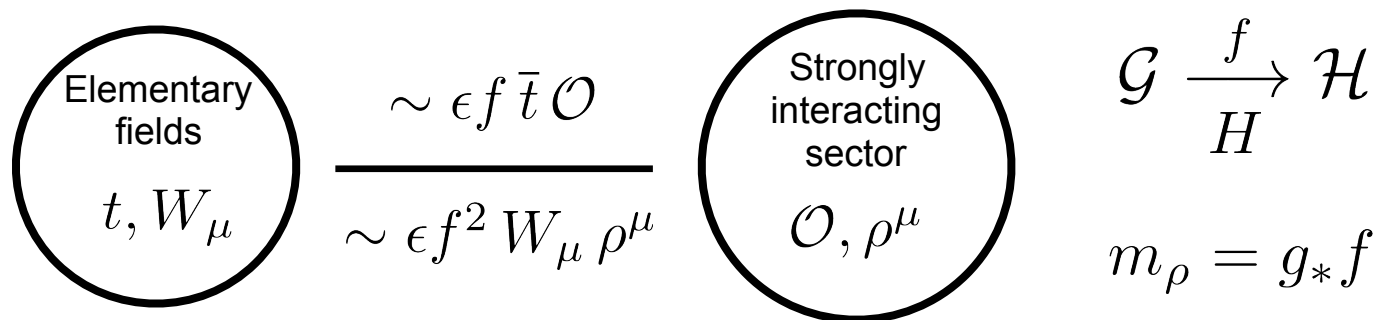
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- At tree level, the Higgs doublet is exact Goldstone, e.g. $SO(5) \xrightarrow[H]{f} SO(4)$
- Breaking of global sym by ϵ generates radiative potential,
dominated by top + vectorlike fermions, the **top partners**
- Coupling $\sim \epsilon f \bar{t} \mathcal{O}$ implies top partners are charged under QCD

Composite Scalar Dark Matter

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- Extend minimal coset $SO(5)/SO(4)$. **Symmetry stabilizing the DM?**

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Gripaios et al. 2009,
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$$SO(6)/SO(5) \rightarrow (H, \eta) \sim 4 + 1_0$$

DM can be stabilized by parity, $\eta \xrightarrow{P_\eta} -\eta$ $P_\eta = \text{diag}(1, 1, 1, 1, -1, 1)$

But $P_\eta \notin SO(6)$, in general **not respected** by higher order terms in chiral Lagrangian

E.g. Wess-Zumino-Witten (see $\pi_0 \rightarrow \gamma\gamma$ in QCD, breaks $\pi \rightarrow -\pi$)

$$\frac{\eta}{16\pi^2} (n_W g^2 W_{\mu\nu}^a \tilde{W}^{a\mu\nu} + n_B g'^2 B_{\mu\nu} \tilde{B}^{\mu\nu})$$

Need to **assume the UV respects the full $O(6)$**

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see also:

Chala et al. 2016

for $SO(7)/SO(6)$

Ballesteros et al. 2017

for $SO(7)/G_2$

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Charged Composite Scalar Dark Matter

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$$SO(\textcolor{red}{7})/SO(\textcolor{red}{6}) \rightarrow (H, \chi) \sim \mathbf{4} + \mathbf{1}_{\pm}$$

Balkin, Ruhdorfer, ES, Weiler,
to appear

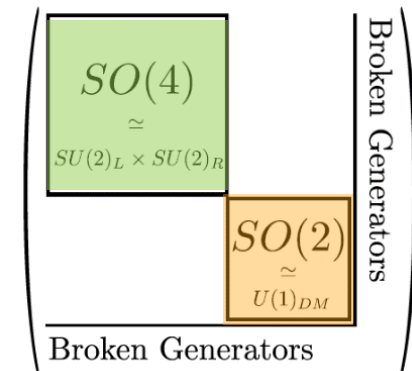
DM candidate is **complex scalar**,

charged under conserved $U(1)_{\text{DM}} \subset SO(6)$

$$\chi \rightarrow e^{i\alpha} \chi$$

Furthermore, no anomalies (no complex reps.)

UV automatically safe



Coupling to elementary fields

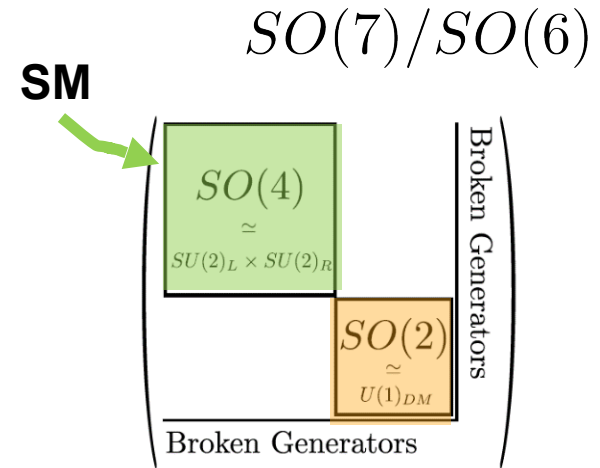
- SM weak gauging preserves $U(1)_{\text{DM}}$
- Fermion partial compositeness

$$\mathcal{L}_{\text{mix}} \sim \epsilon_q \bar{q}_L \mathcal{O}_q + \epsilon_t \bar{t}_R \mathcal{O}_t$$

- If SM fermions embedded in **7** (fundamental),
can leave $U(1)_{\text{DM}}$ intact

$$q_L \rightarrow \begin{pmatrix} ib_L \\ b_L \\ it_L \\ -t_L \\ 0 \\ 0 \\ 0 \end{pmatrix}, \quad t_R \rightarrow \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ t_R \end{pmatrix}$$

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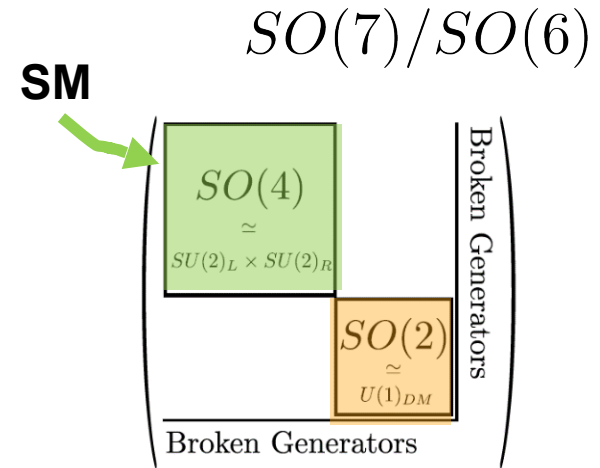
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[for other choices, shift sym. can also be preserved $\Rightarrow m_\chi \ll m_h$

Mass controlled by light fermions + gauging of $U(1)_{\text{DM}}$. E.g. $t_R \sim \mathbf{21}, \mathbf{27}$]



Radiative scalar potential

$$V(h, \chi) = \frac{1}{2}\mu_h^2 h^2 + \frac{\lambda_h}{4}h^4 + \underbrace{\mu_{DM}^2 \chi^* \chi}_{\text{red bracket}} + \lambda_{DM}(\chi^* \chi)^2 + \lambda h^2 \chi^* \chi$$



$$\mu_{DM}^2 \approx \frac{N_c}{4\pi^2 f^2} \int_0^\infty dp^2 p^2 \left(\underbrace{\sum_{i=1}^{N_S} \frac{|\epsilon_{tS}^i|^2}{p^2 + m_{S_i}^2} - \sum_{i=1}^{N_Q} \frac{|\epsilon_{tQ}^i|^2}{p^2 + m_{Q_i}^2}}_{\text{black bracket}} \right)$$

Marzocca et al. 2012
Pomarol, Riva 2012

obtain calculability through
Generalized **Weinberg Sum Rules**,
that give relations between parameters



UV-finite if $\sim \frac{1}{p^6}$ or faster

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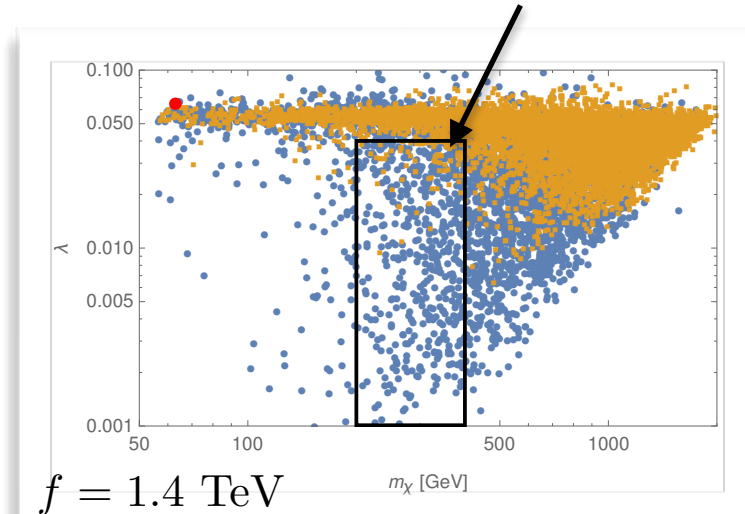
Leading order prediction is $\lambda \sim \frac{\lambda_h}{2} \sim 0.06$,
ruled out by direct detection

λ reduction is correlated with lighter top partners

blue: $M_{\text{lightest}} < 1 \text{ TeV}$, excluded by LHC

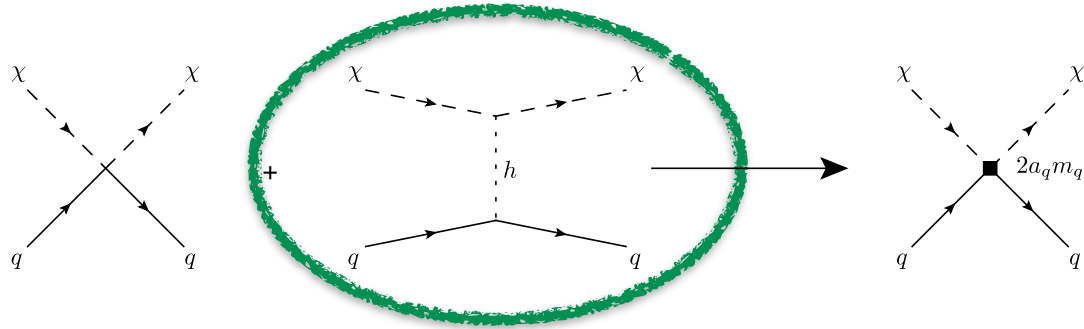
orange: LHC ok

pheno-viable region



Pheno: direct detection

- Higgs exchange in t-channel dominates




$$\frac{1}{2f^2} \ll \frac{\lambda}{m_h^2}$$

$$\mathcal{L} \sim \frac{1}{f^2} \cancel{\partial(h^2) \partial(\chi^* \chi)} - \lambda h^2 \chi^* \chi \quad \rightarrow \quad \sigma_{\text{SI}}^{\chi N} \simeq \frac{0.3^2}{\pi} \frac{m_N^4 \lambda^2}{m_\chi^2 m_h^4}$$

~ vanilla Higgs portal, with minor corrections

Pheno: relic abundance

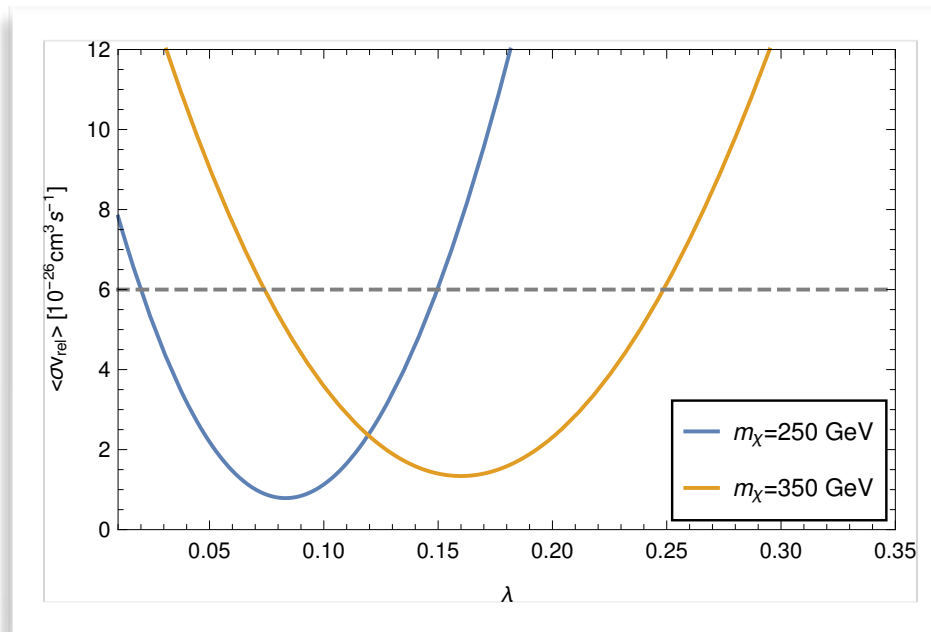
- Annihilation into $t\bar{t}$, WW , ZZ , hh
- Interplay of derivative and portal couplings


$$\mathcal{L} \sim \frac{1}{f^2} \partial(h^2) \partial(\chi^* \chi) - \lambda h^2 \chi^* \chi \quad \rightarrow \quad \sigma \propto \left(\frac{2m_\chi^2}{f^2} - \lambda \right)^2$$

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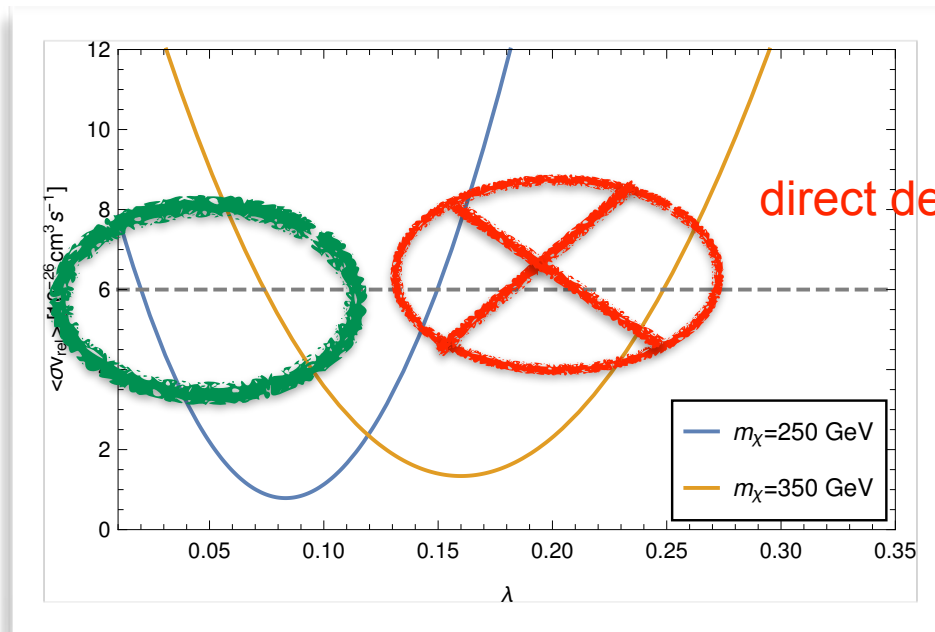


for fixed DM mass, **two** values of λ reproduce relic abundance

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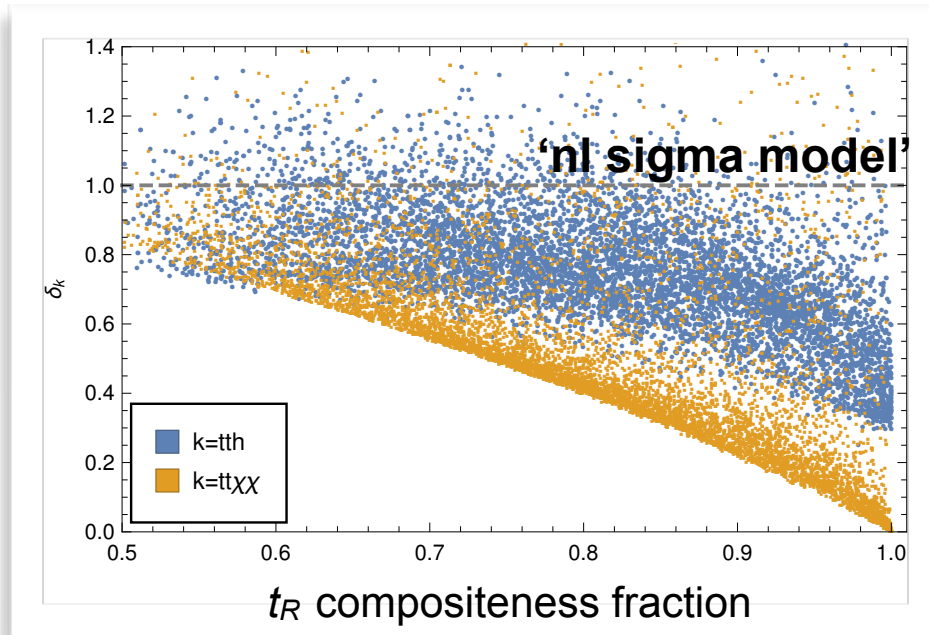
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Top partner mixing matters

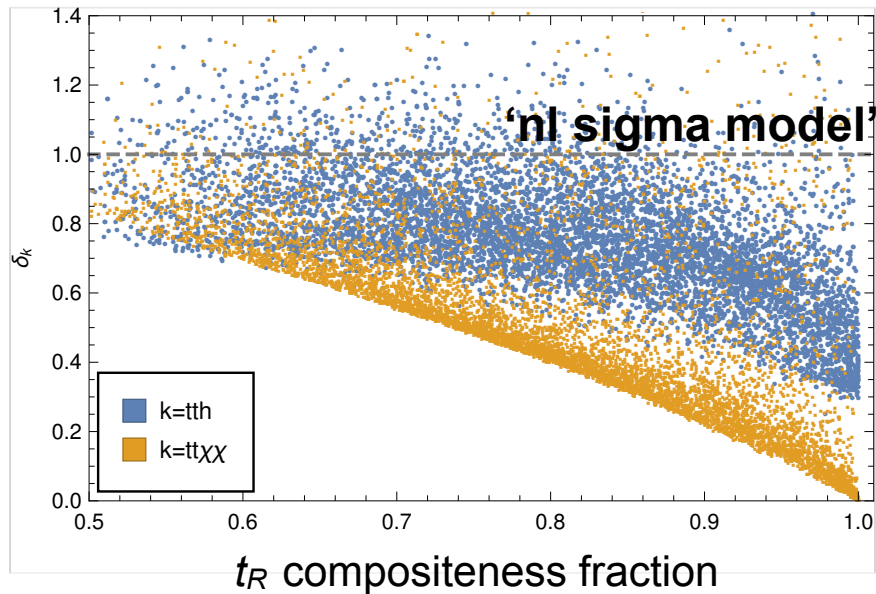


when t_R is fully composite,
it respects χ shift symmetry

➡ non-derivative couplings
vanish

orange: $tt\chi\chi$ coupling

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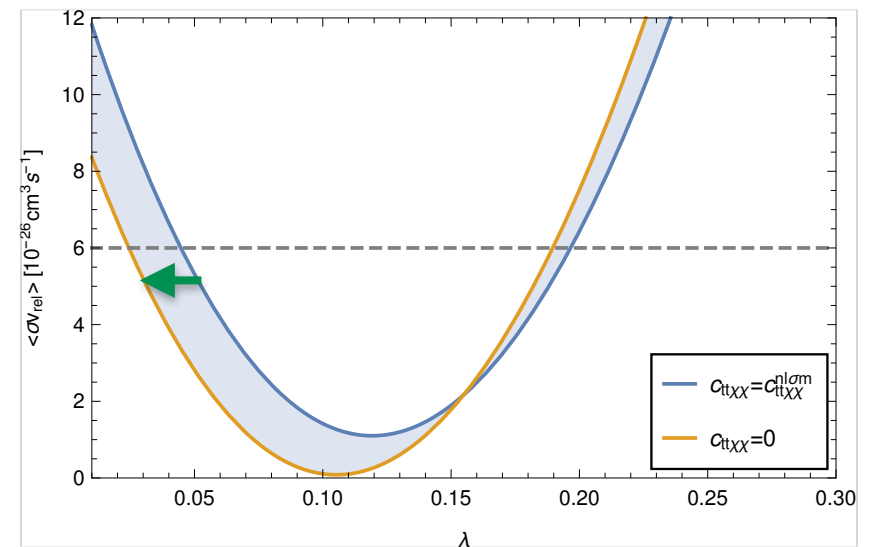


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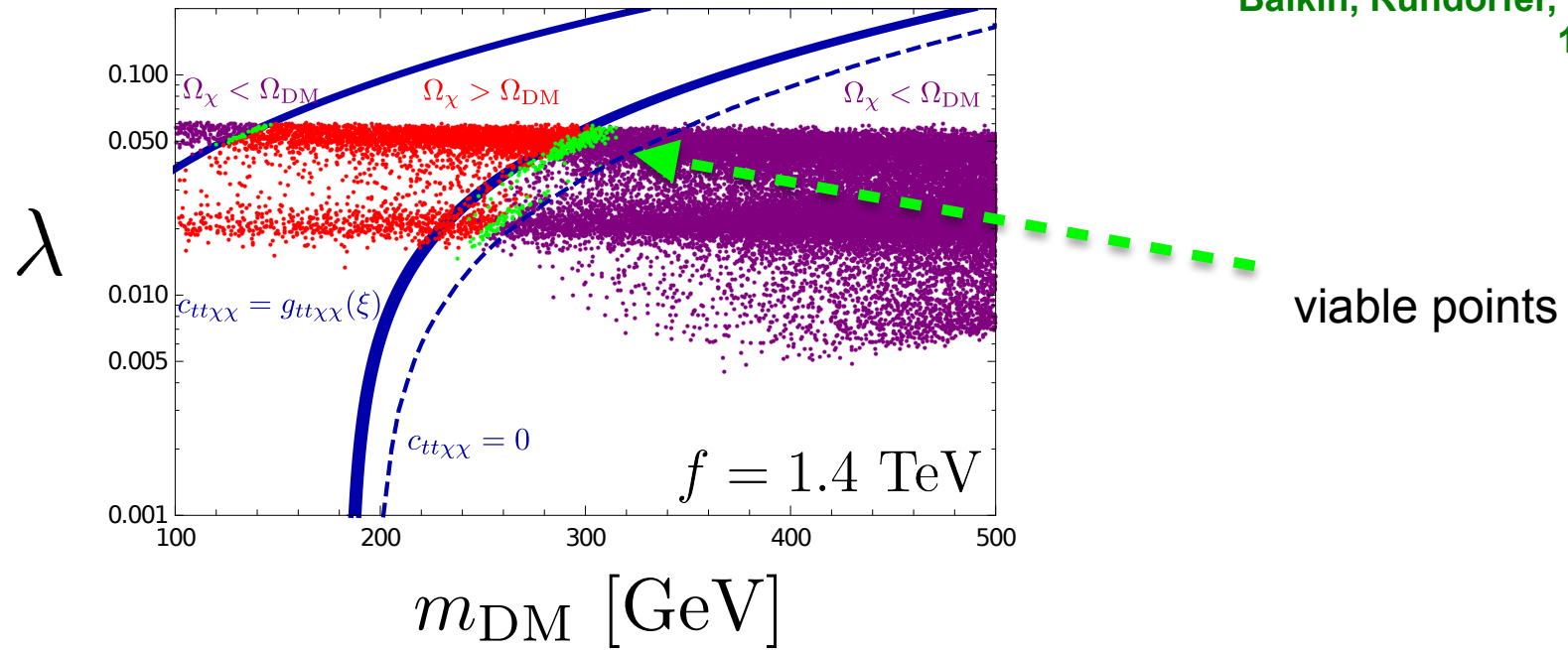
orange: $t\bar{t}\chi\chi$

top partner mixing lowers λ ,
helps with direct detection



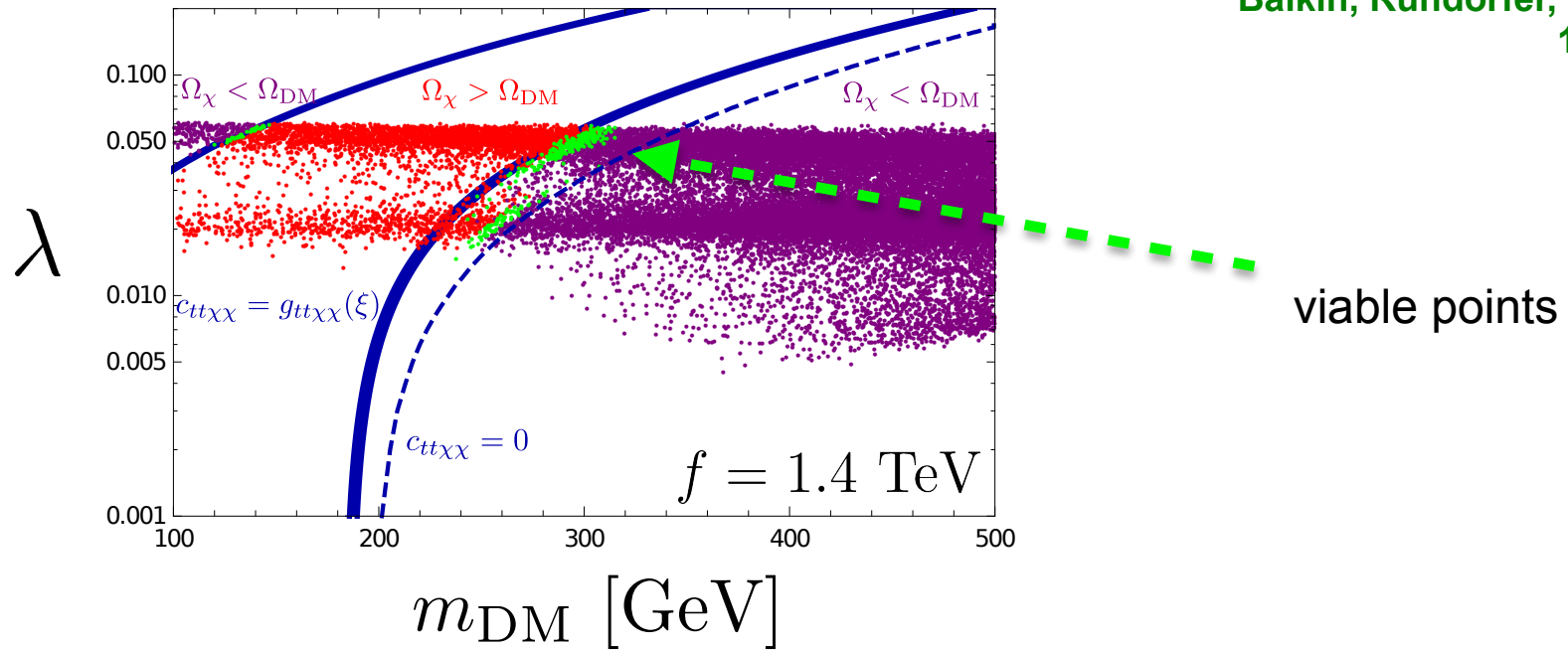
Relic Abundance

Balkin, Ruhdorfer, ES, Weiler
1705.xxxxx



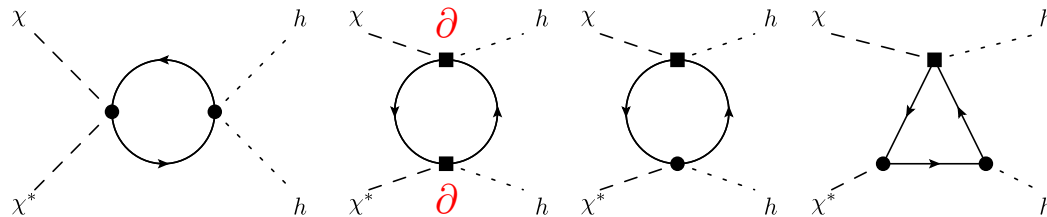
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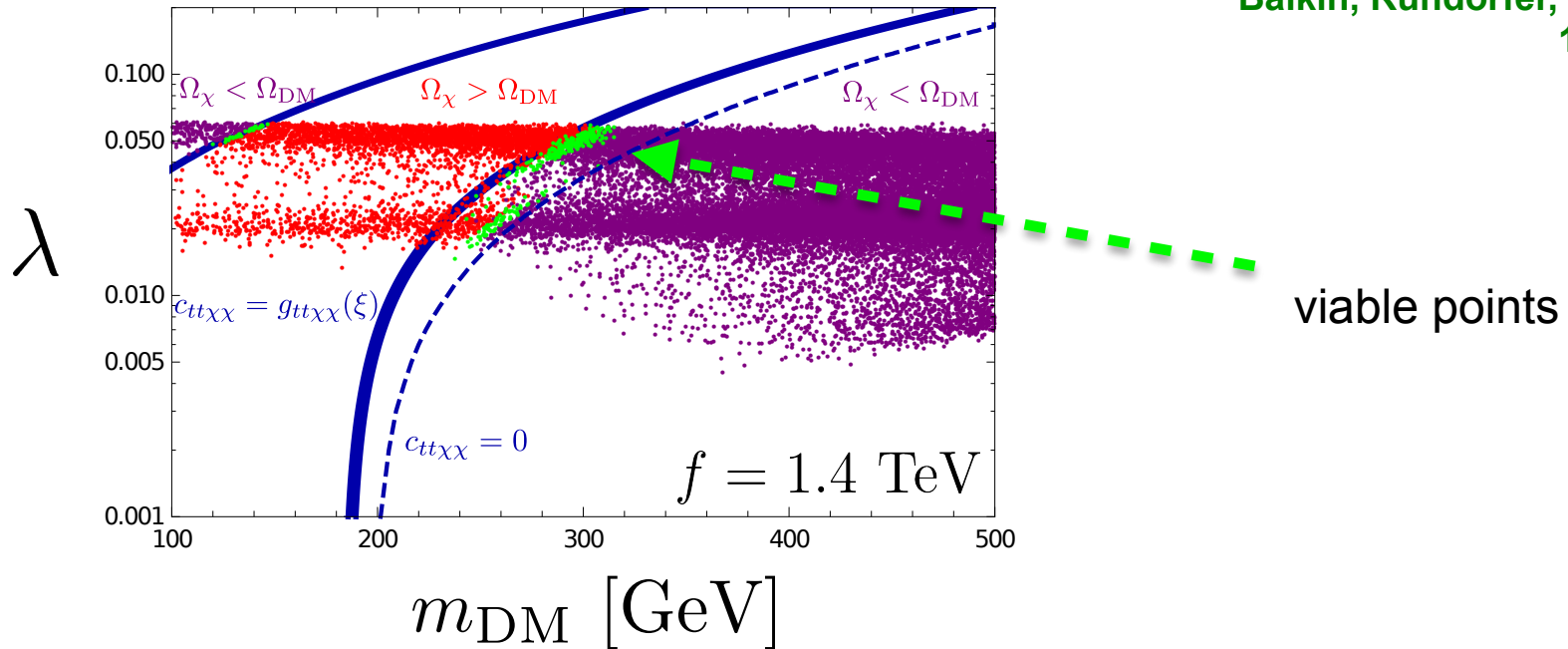
But [tree-level + CW] calculation
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$$c \frac{1}{f^2} \partial_\mu |H|^2 \partial^\mu |\chi|^2$$



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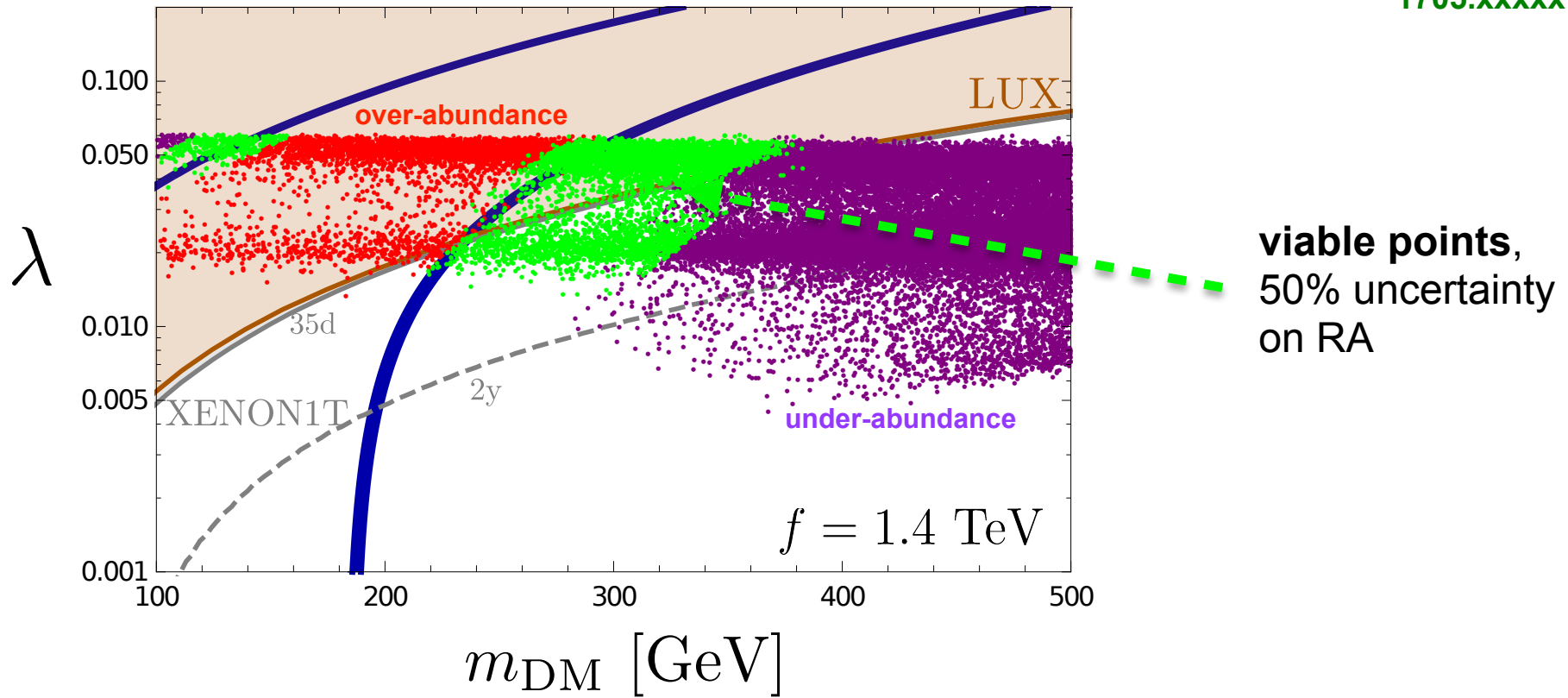
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$$c_{\text{tree}} = 1, \quad c_{1\text{-loop}} = \frac{N_c}{2\pi^2 f^2} \left(\epsilon_t^2 - \frac{\epsilon_q^2}{8} \right) \log \frac{\Lambda^2}{m_\psi^2}$$

irreducible uncertainty of $\sim 50\%$ on cross section

Composite DM pheno

Balkin, Ruhdorfer, ES, Weiler
1705.xxxxx



**Effect of top partner mixing pushes λ down,
relaxes direct detection constraint**

Tuning

- Simple estimate of tuning is $\frac{1}{\Delta} \sim 2\xi \sim 6\%$ for $f = 1.4$ TeV
- However, most natural value of portal coupling is $\lambda \sim \frac{\lambda_h}{2} \approx 0.06$

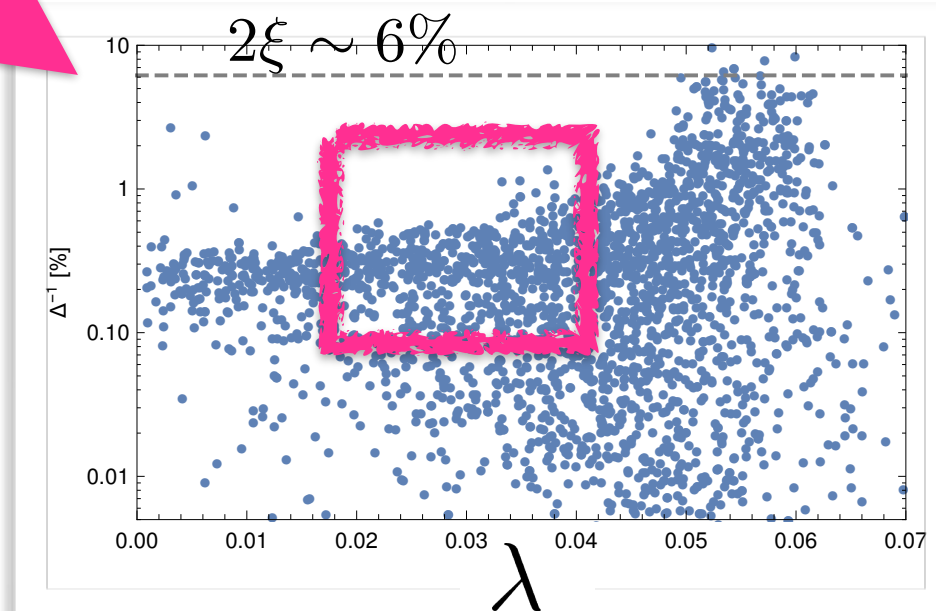
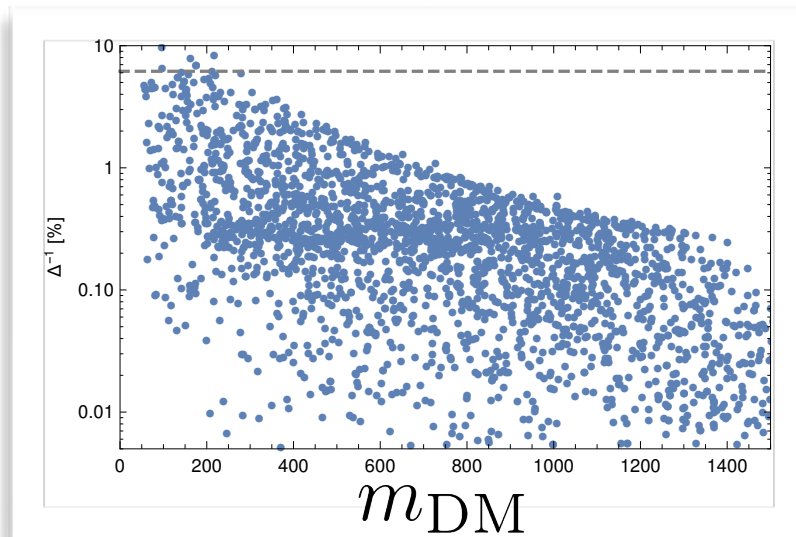
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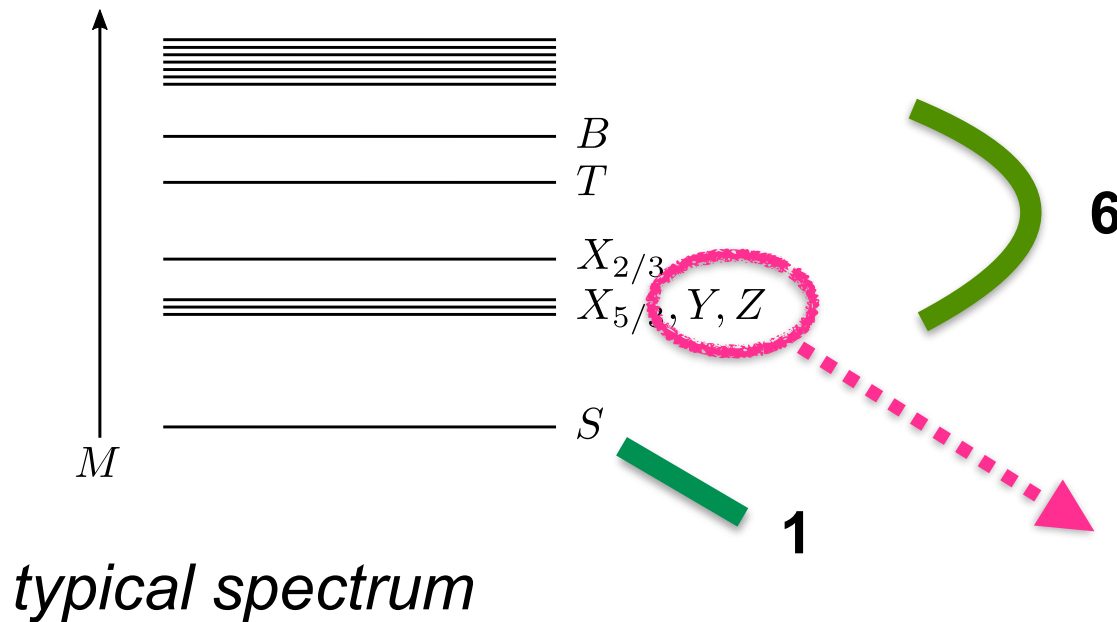
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Too large for direct detection.

Suppressing it costs extra tuning



Collider pheno, sketch



$$SO(7)/SO(6)$$

$$\mathbf{7} = \mathbf{6} \oplus \mathbf{1}$$

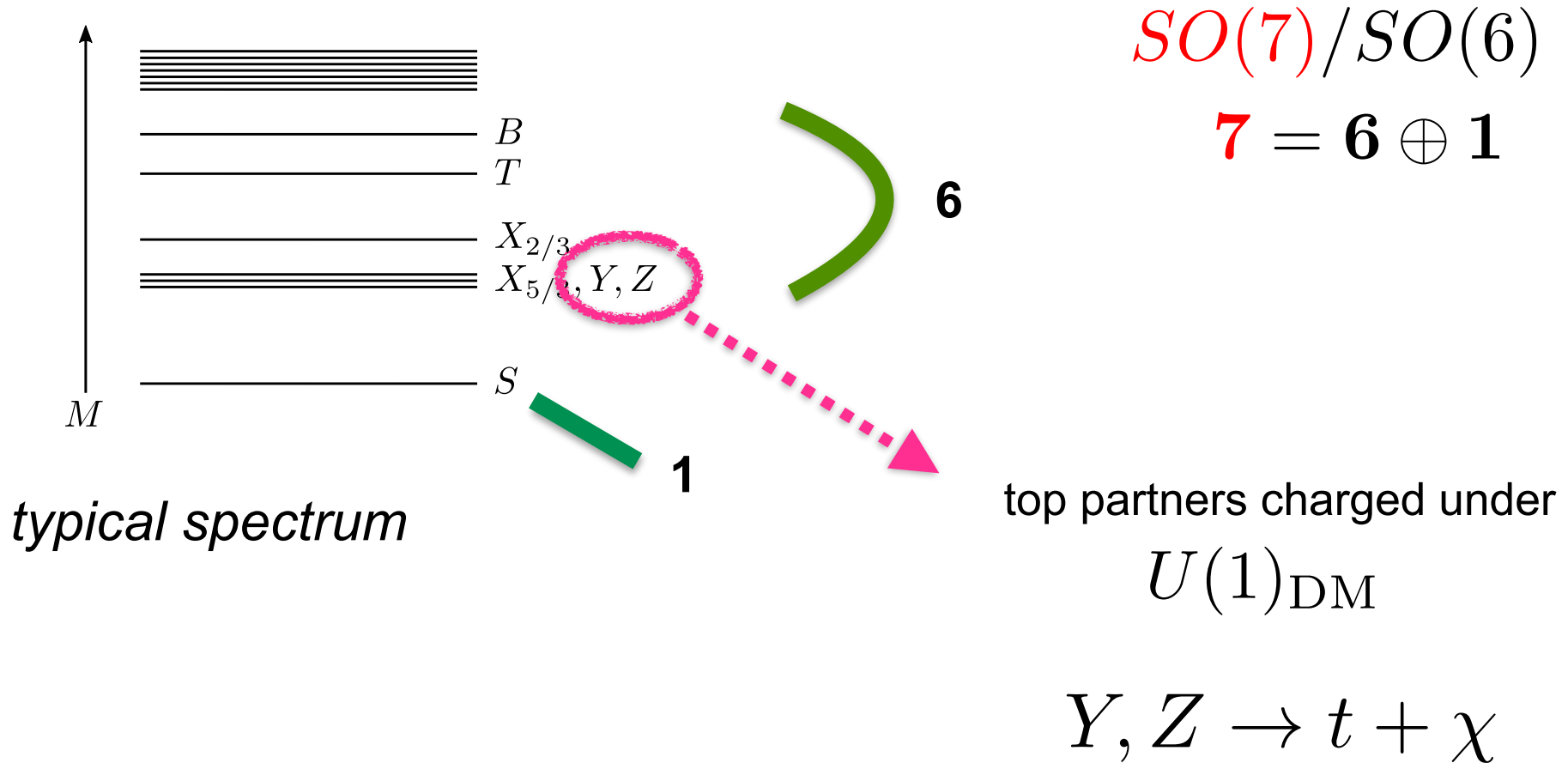
top partners charged under

$$U(1)_{\text{DM}}$$

$$Y, Z \rightarrow t + \chi$$

see e.g. Serra 2015

Collider pheno, sketch



current bounds:

$$M_{\text{singlet}} > 1 \text{ TeV}, \quad M_{\text{doublets}} > 1.2 \text{ TeV}, \quad M_{Y,Z} > 1.4 \text{ TeV}$$

Summary & Outlook

- Composite Higgs model with **UV-safe DM stabilization**
- DM is pGB scalar with 200 - 400 GeV mass
Will be fully tested by XENON1T
- Typically t_R is very composite, mixing with top partners has important effects in annihilation
 - + large radiative corrections to derivative operators
- New LHC signals from $U(1)_{\text{DM}}$ - charged top partners

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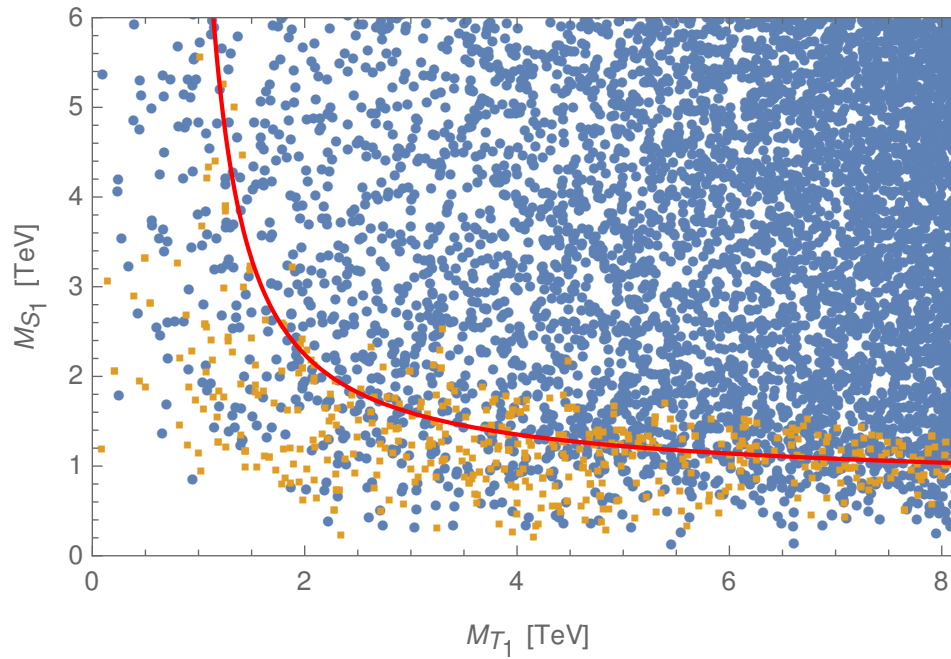
- New LHC signals from $U(1)_{\text{DM}}$ - charged top partners

- $U(1)_{\text{DM}}$ can be weakly gauged **Agrawal, Cyr-Racine, Randall, Scholtz 2016**

- Indirect detection: antiproton constraints? (AMS-02)

Backup

Light top partners

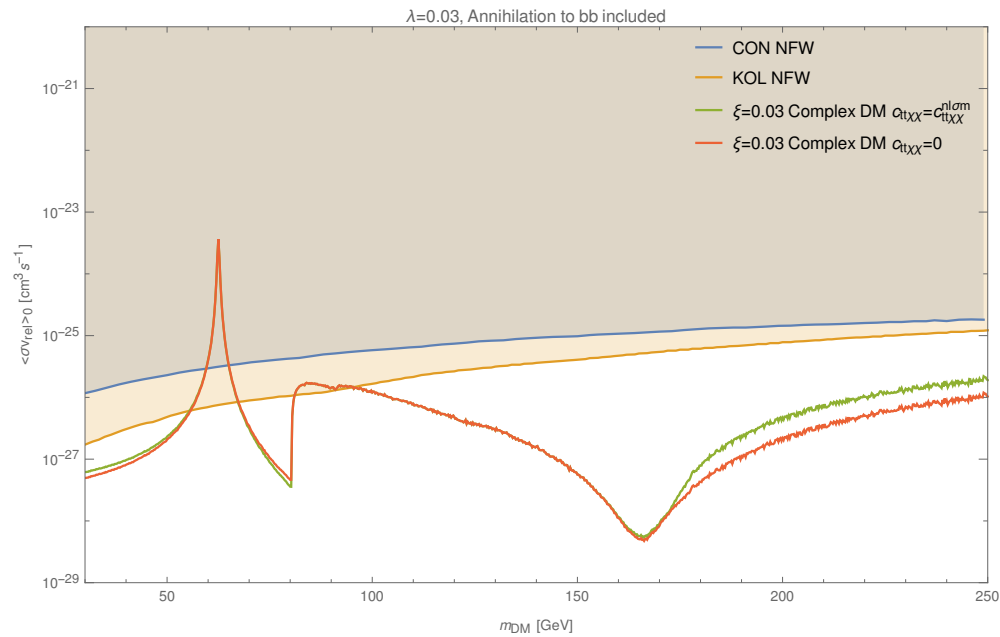


$$f = 1.4 \text{ TeV}$$

orange: $120 \text{ GeV} < m_h < 130 \text{ GeV}$

red:
$$\frac{m_h^2}{m_t^2} \approx \frac{N_c}{\pi^2 f^2} \frac{M_T^2 M_S^2}{M_T^2 - M_S^2} \log(M_T^2 / M_S^2)$$

Indirect detection: antiprotons



Bounds from PAMELA: recast from complete analysis
in **Marzocca and Urbano, 2014**