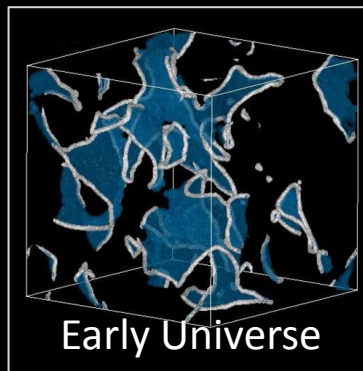




Solar Axions



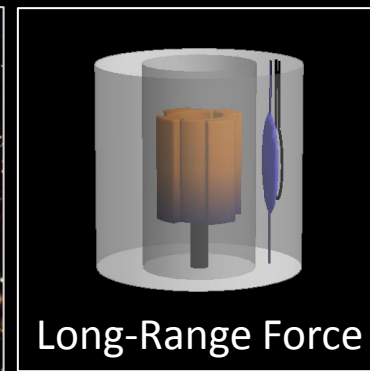
Supernova 1987A



Early Universe



Dark Matter



Long-Range Force

Axion Landscape



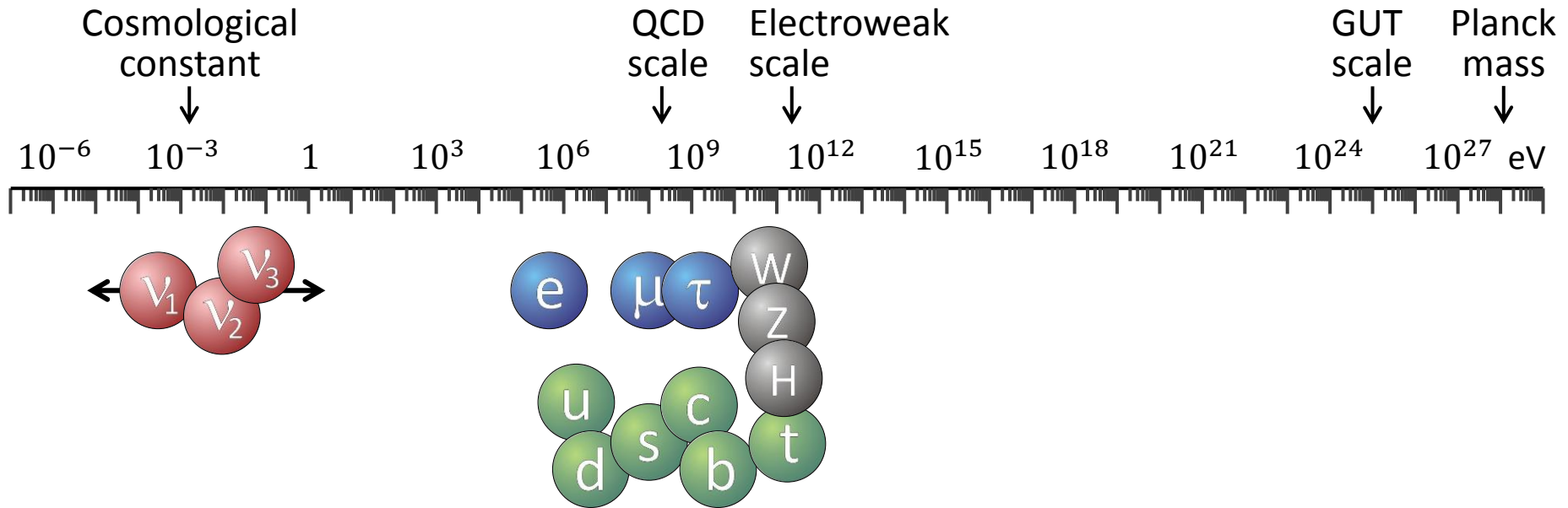
Georg Raffelt

Max Planck Institute for Physics, Munich

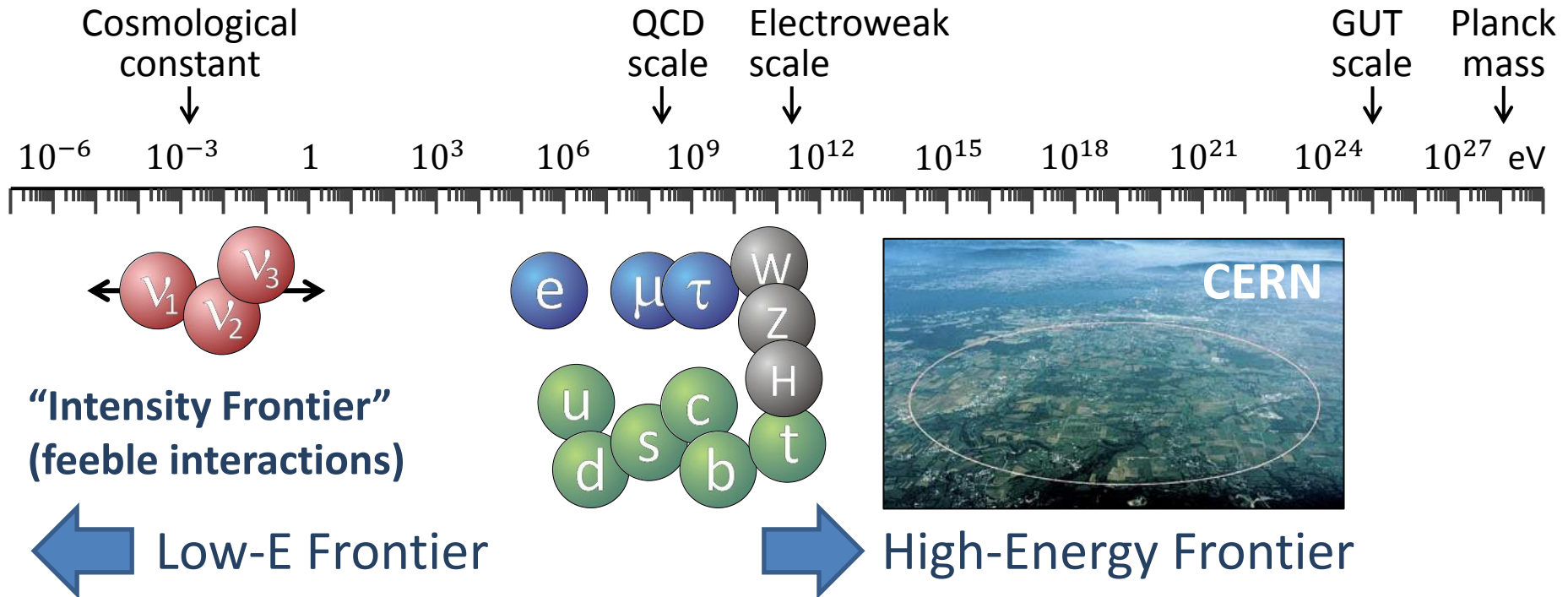


Max-Planck-Institut für Physik
(Werner-Heisenberg-Institut)

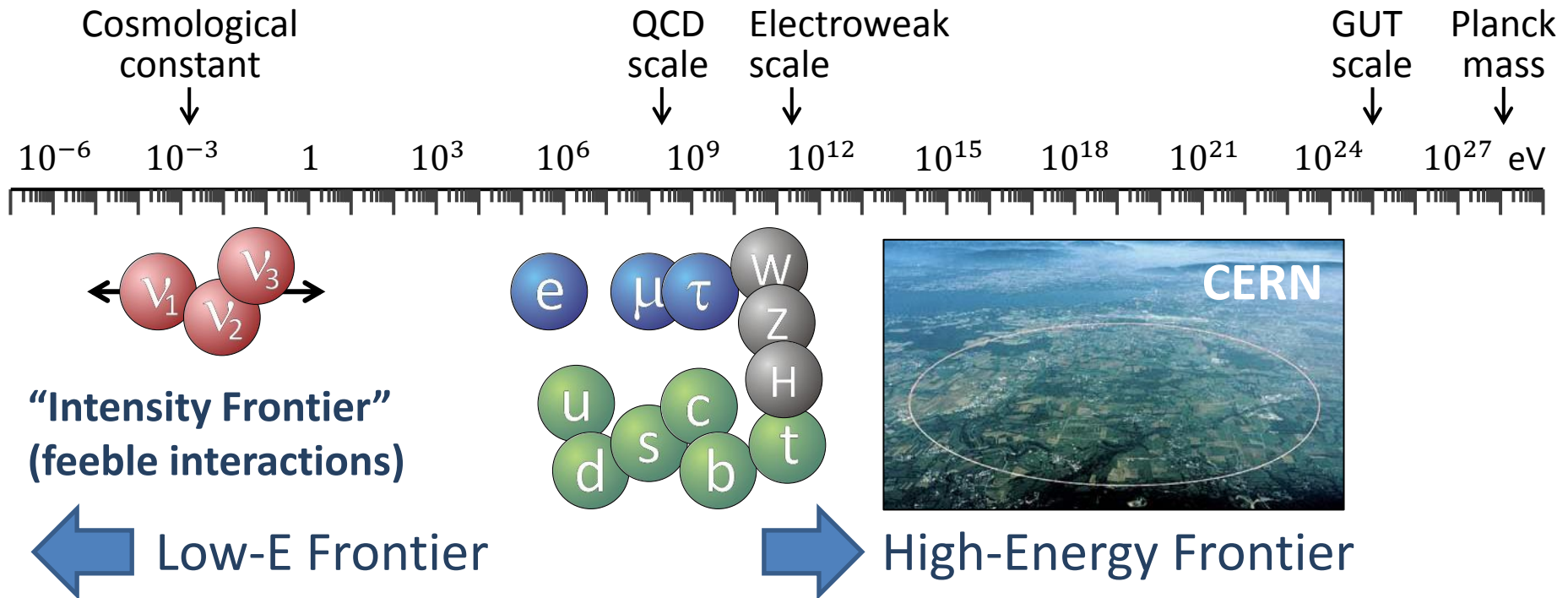
High- and Low-Energy Frontiers in Particle Physics



High- and Low-Energy Frontiers in Particle Physics

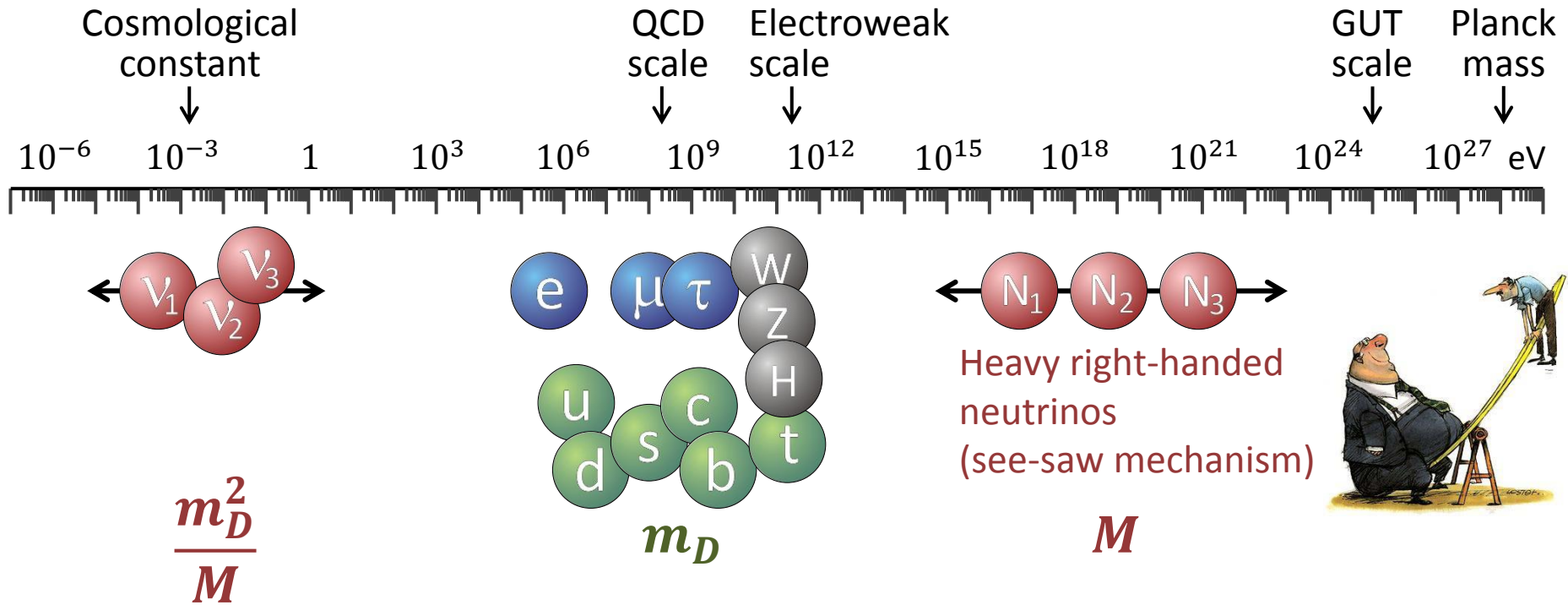


High- and Low-Energy Frontiers in Particle Physics

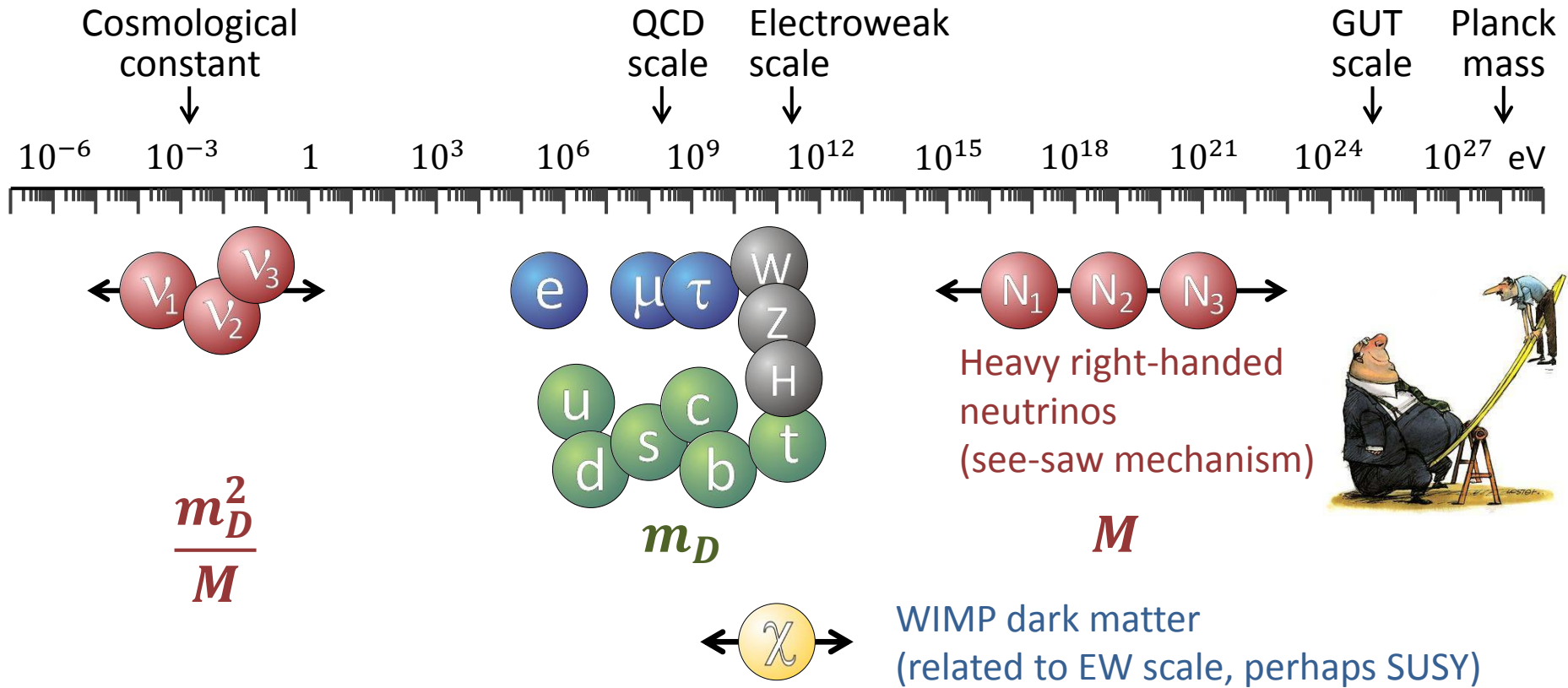


New 750 GeV boson?
few-sigma hint @ CMS & ATLAS
many theory papers since Dec 2015

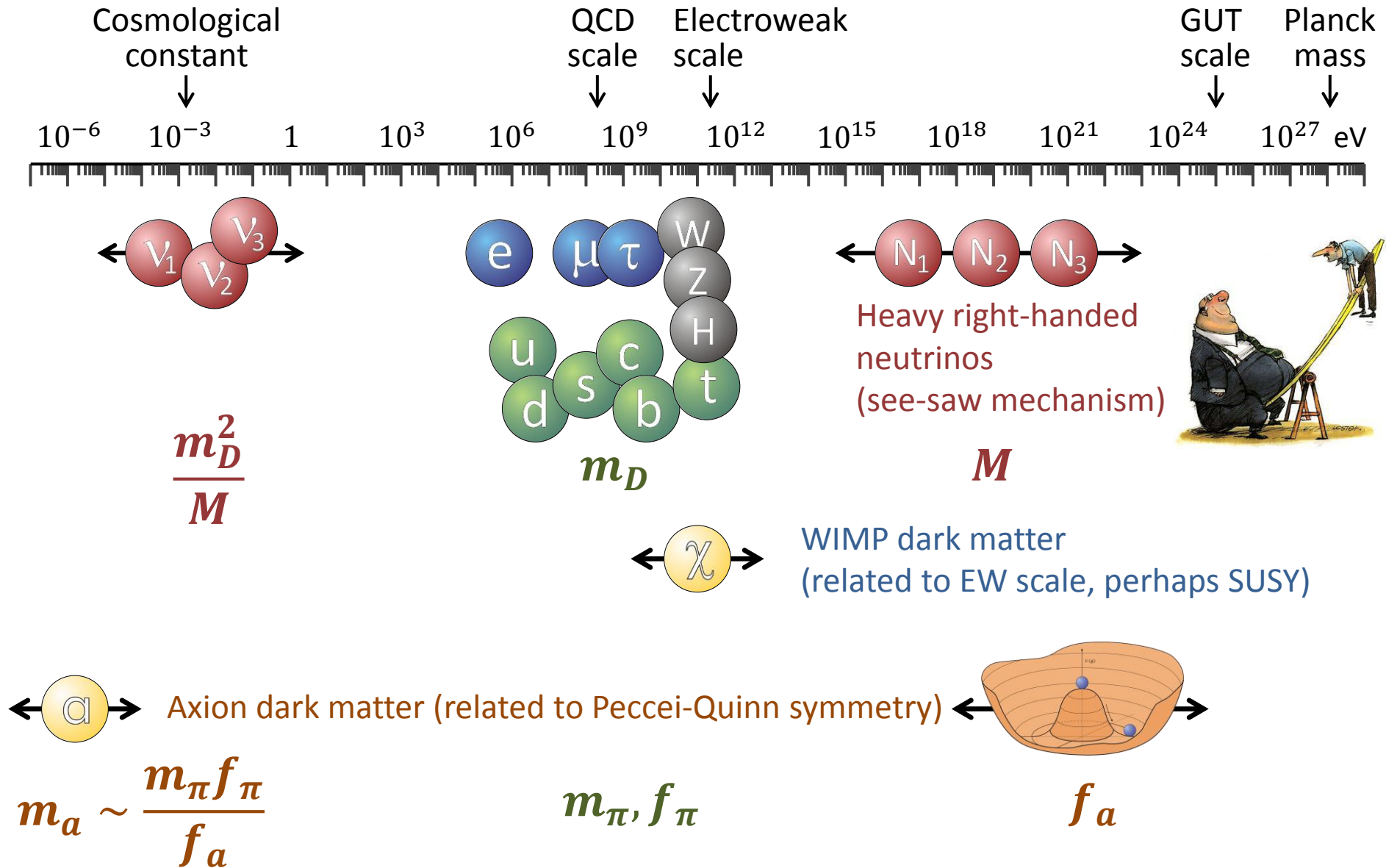
High- and Low-Energy Frontiers in Particle Physics



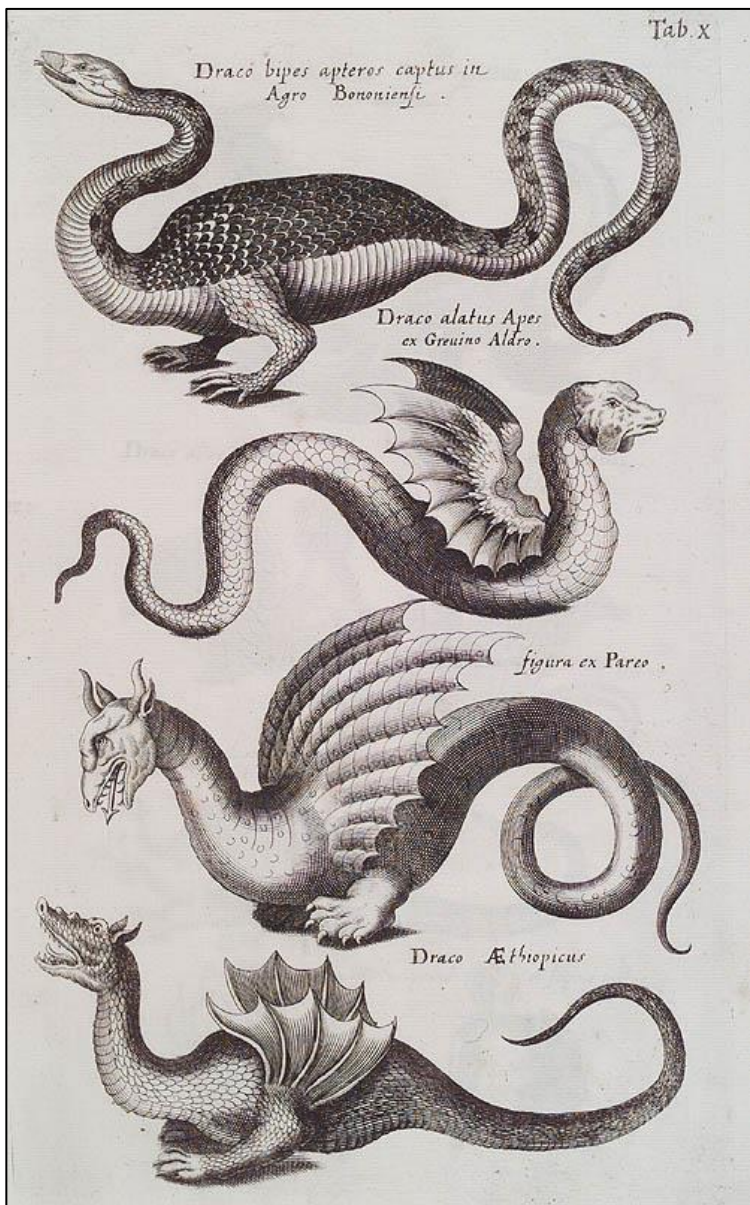
High- and Low-Energy Frontiers in Particle Physics



High- and Low-Energy Frontiers in Particle Physics



Bestiary of Low-Mass Bosons



Weakly Interacting Sub-eV Particles (WISPs)

- **Axions** (1 parameter family $m_a f_a \sim m_\pi f_\pi$)
Solves strong CP problem
Could be dark matter
- **Axion-like particles (ALPs)**
Generic two-photon vertex, could be dark matter
(2 parameters m_a and $g_{a\gamma}$)
- **String axions**
(almost massless pseudoscalars in string theory)
One of them may solve CP problem
- **Hidden photons**
Low-mass gauge bosons from $U'(1)$
(kinetic mixing parameter χ and mass $m_{\gamma'}$)
- **Chameleons**
Scalars in certain models of scalar-tensor gravity
Motivated by dark energy
Environment-dependent properties

The CP Problem of Strong Interactions

$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{\psi}_q (iD - \underbrace{m_q}_{\text{Real quark mass}} e^{i\theta_q} \underbrace{e^{i\theta_q}}_{\text{Phase from Yukawa coupling}}) \psi_q - \frac{1}{4} G_{\mu\nu a} G_a^{\mu\nu} - \underbrace{\Theta}_{\text{Angle variable}} \frac{\alpha_s}{8\pi} \underbrace{G_{\mu\nu a} \tilde{G}_a^{\mu\nu}}_{\text{CP-odd quantity} \sim \mathbf{E} \cdot \mathbf{B}}$$

Remove phase of mass term by chiral transformation of quark fields

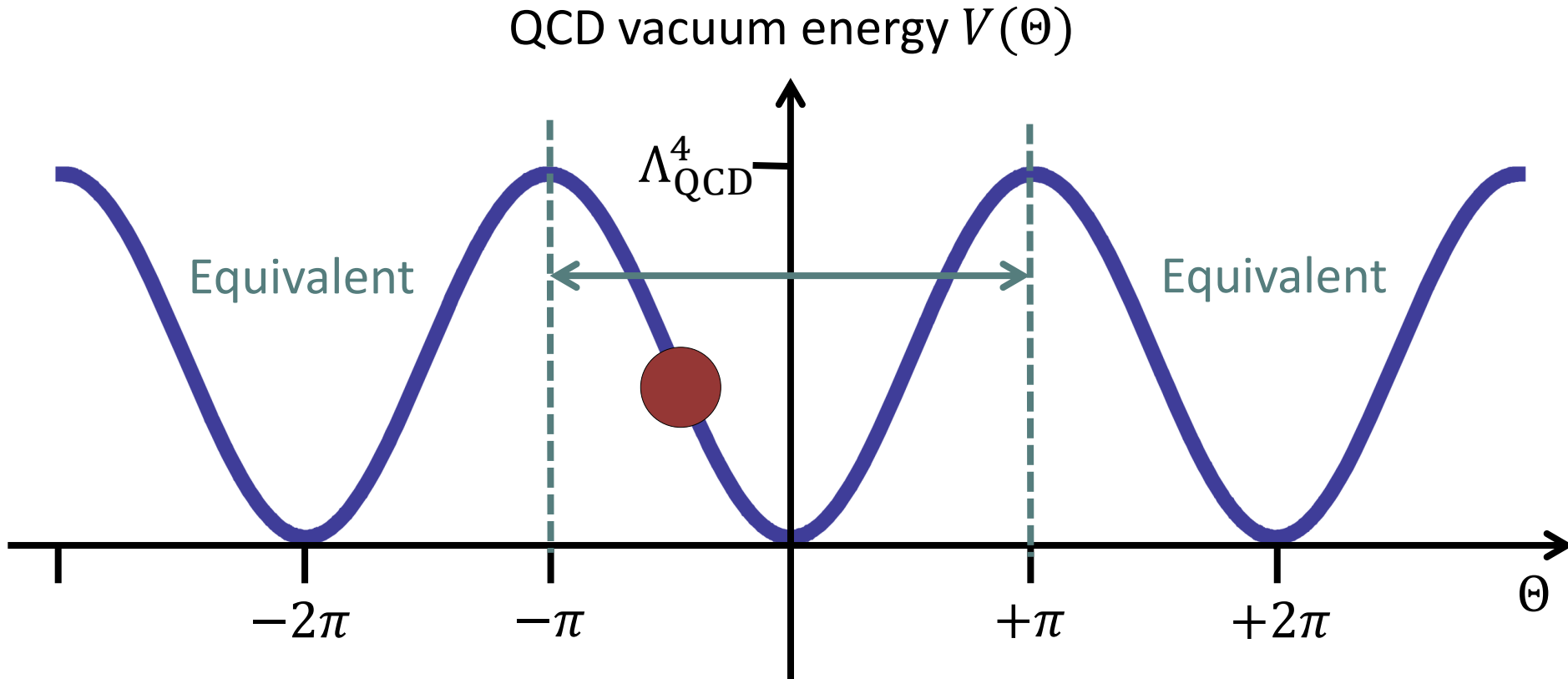
$$\psi_q \rightarrow e^{-i\gamma_5 \theta_q / 2} \psi_q$$

$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{\psi}_q (iD - m_q) \psi_q - \frac{1}{4} GG - \underbrace{(\Theta - \arg \det M_q)}_{-\pi \leq \bar{\Theta} \leq +\pi} \frac{\alpha_s}{8\pi} G \tilde{G}$$

- ❖ $\bar{\Theta}$ can be traded between quark phases and $G\tilde{G}$ term
- ❖ No physical impact if at least one $m_q = 0$
- ❖ Induces a large neutron electric dipole moment (a T-violating quantity)

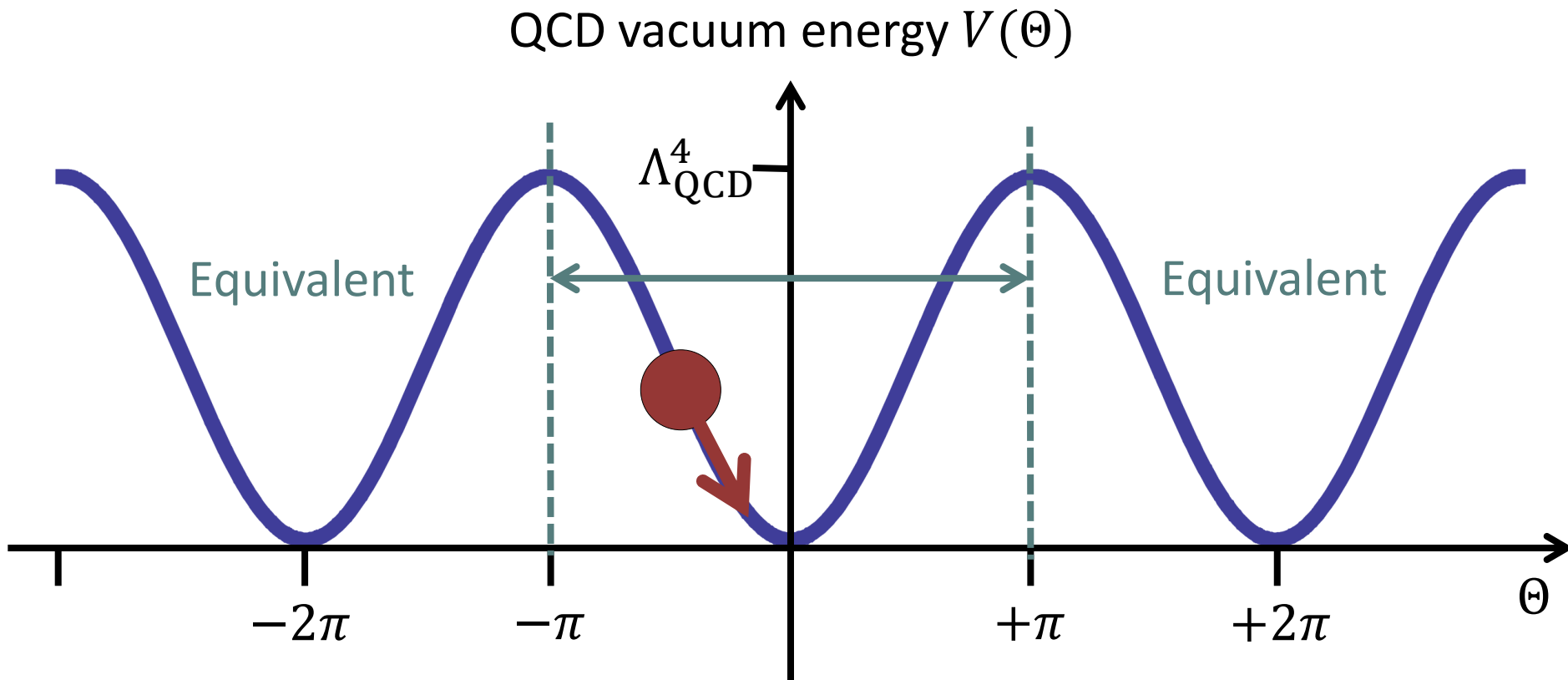
Experimental limits: $|\bar{\Theta}| < 10^{-11}$ Why so small?

Strong CP Problem



- CP conserving vacuum has $\Theta = 0$ (Vafa and Witten 1984)
- QCD could have any $-\pi \leq \Theta \leq +\pi$, is “constant of nature”
- Energy can not be minimized: Θ not dynamical

Strong CP Problem



- CP conserving vacuum has $\Theta = 0$ (Vafa and Witten 1984)
- QCD could have any $-\pi \leq \Theta \leq +\pi$, is “constant of nature”
- Energy can not be minimized: Θ not dynamical

Peccei-Quinn solution:

Make Θ dynamical, let system relax to lowest energy

38 Years of Axions

VOLUME 40, NUMBER 4

PHYSICAL REVIEW LETTERS

23 JANUARY 1978

A New Light Boson?

Steven Weinberg

Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138

(Received 6 December 1977)

It is pointed out that a global $U(1)$ symmetry, that has been introduced in order to preserve the parity and time-reversal invariance of strong interactions despite the effects of instantons, would lead to a neutral pseudoscalar boson, the "axion," with mass roughly of order 100 keV to 1 MeV. Experimental implications are discussed.

VOLUME 40, NUMBER 5

PHYSICAL REVIEW LETTERS

30 JANUARY 1978

Problem of Strong P and T Invariance in the Presence of Instantons

F. Wilczek^(a)

Columbia University, New York, New York 10027, and The Institute for Advanced Studies, Princeton, New Jersey 08540^(b)

(Received 29 November 1977)

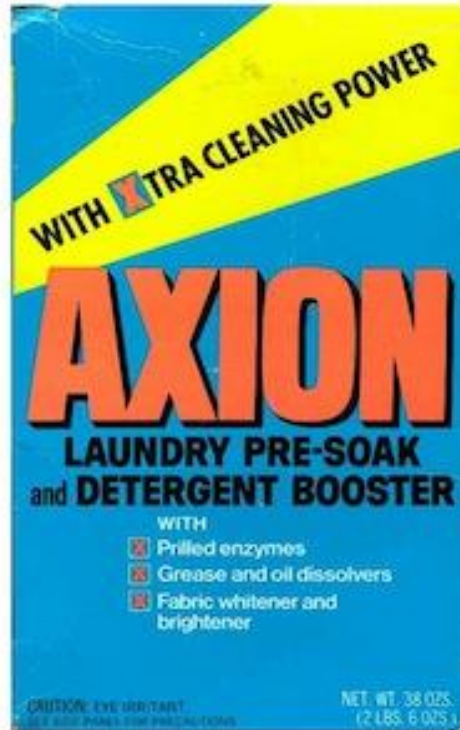
The requirement that P and T be approximately conserved in the color gauge theory of strong interactions without arbitrary adjustment of parameters is analyzed. Several possibilities are identified, including one which would give a remarkable new kind of very light, long-lived pseudoscalar boson.

One of the main advantages of the color gauge theory of strong interactions is that so many of the observed symmetries of strong interactions seem to follow automatically as a consequence of the gauge principle and renormalizability— P , T , C , flavor conservation, the $3 \oplus 3^*$ structure of chi-

a certain class of theories^{4,5,7} the parameter θ is physically meaningless,^{4,5} or dynamically determined.⁷ In this case, if the strong interaction conserves P and T , we shall say the conservation is *automatic*.

I regard a theory of type (i) as very unattrac-

The Cleansing Axion



Frank Wilczek

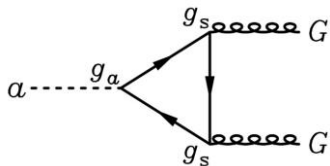


**“I named them after a laundry detergent, since they clean up a problem with an axial current.”
(Nobel lecture 2004)**

Phenomenological Axion Properties

Gluon coupling (generic), defines normalization of axion scale f_a

$$\mathcal{L}_{aG} = \frac{\alpha_s}{8\pi} \frac{a}{f_a} G \tilde{G}$$

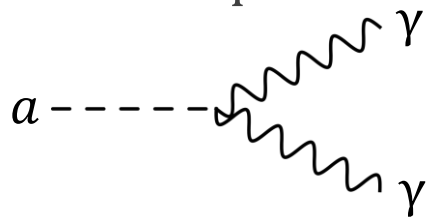


Mass (generic) depends on up/down quark masses

$$m_a = \frac{\sqrt{m_u m_d}}{m_u + m_d} \frac{m_\pi}{f_\pi f_a} \approx \frac{6 \mu\text{eV}}{f_a / 10^{12} \text{ GeV}}$$

Axion-photon coupling (model dependent)

$$\mathcal{L}_{a\gamma} = -\frac{g_{a\gamma}}{4} F \tilde{F} a = g_{a\gamma} \mathbf{E} \cdot \mathbf{B} a = \frac{\alpha}{2\pi} \left(\frac{E}{N} - 1.92 \right) \frac{a}{f_a} \mathbf{E} \cdot \mathbf{B}$$



Generic from a - π - η mixing

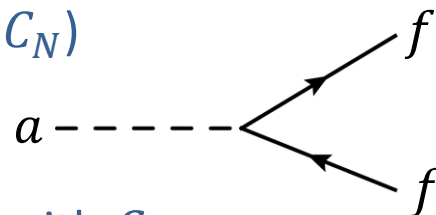


Model-dependent,
 $E/N = 0$ (KSVZ), $8/3$ (DFSZ), many others

Axion-nucleon coupling (model-dependent numerical factors C_N)

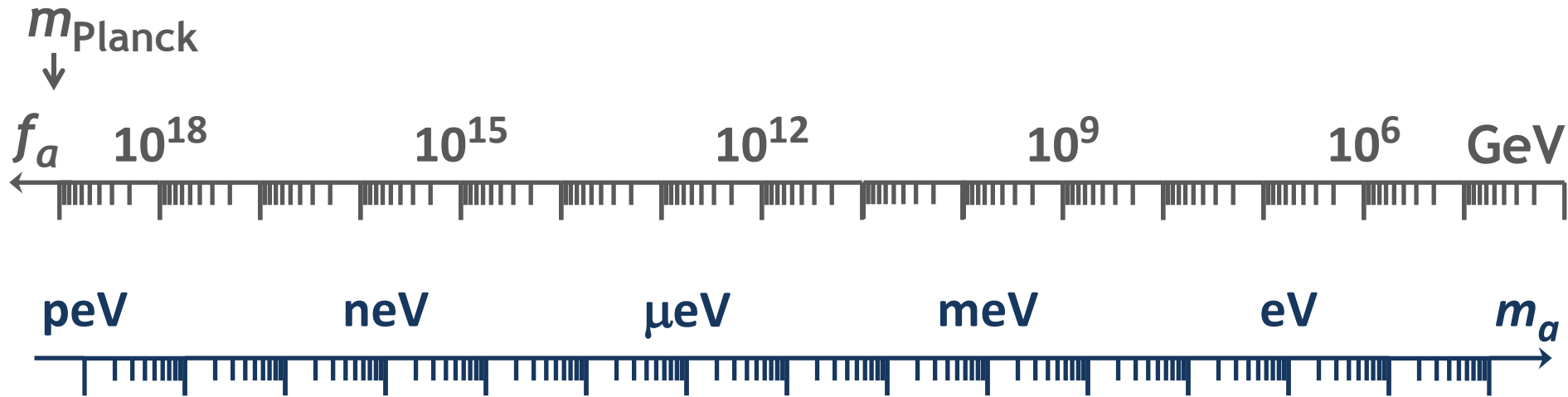
$$\mathcal{L}_{aN} = C_N \bar{\Psi}_N \gamma^\mu \gamma_5 \Psi_N \frac{\partial_\mu a}{2f_a}$$

- Axial-vector current
- Spin-dependent int'n



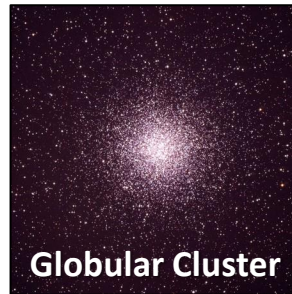
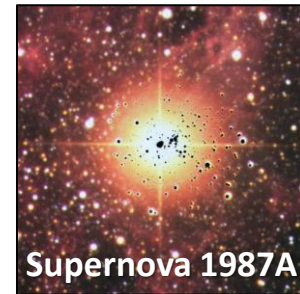
Axion-electron coupling in non-hadronic models is analogous with C_e

Axion Bounds

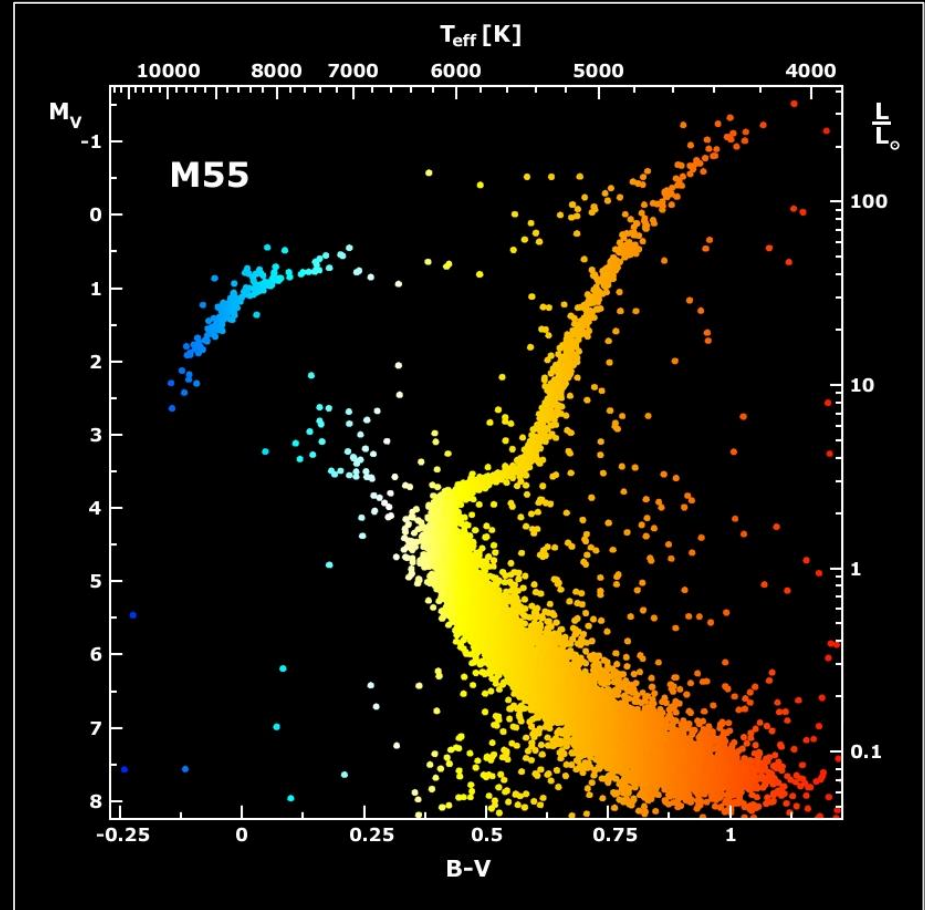


Opportunities for detection

Astrophysical Bounds
(Energy loss of stars)



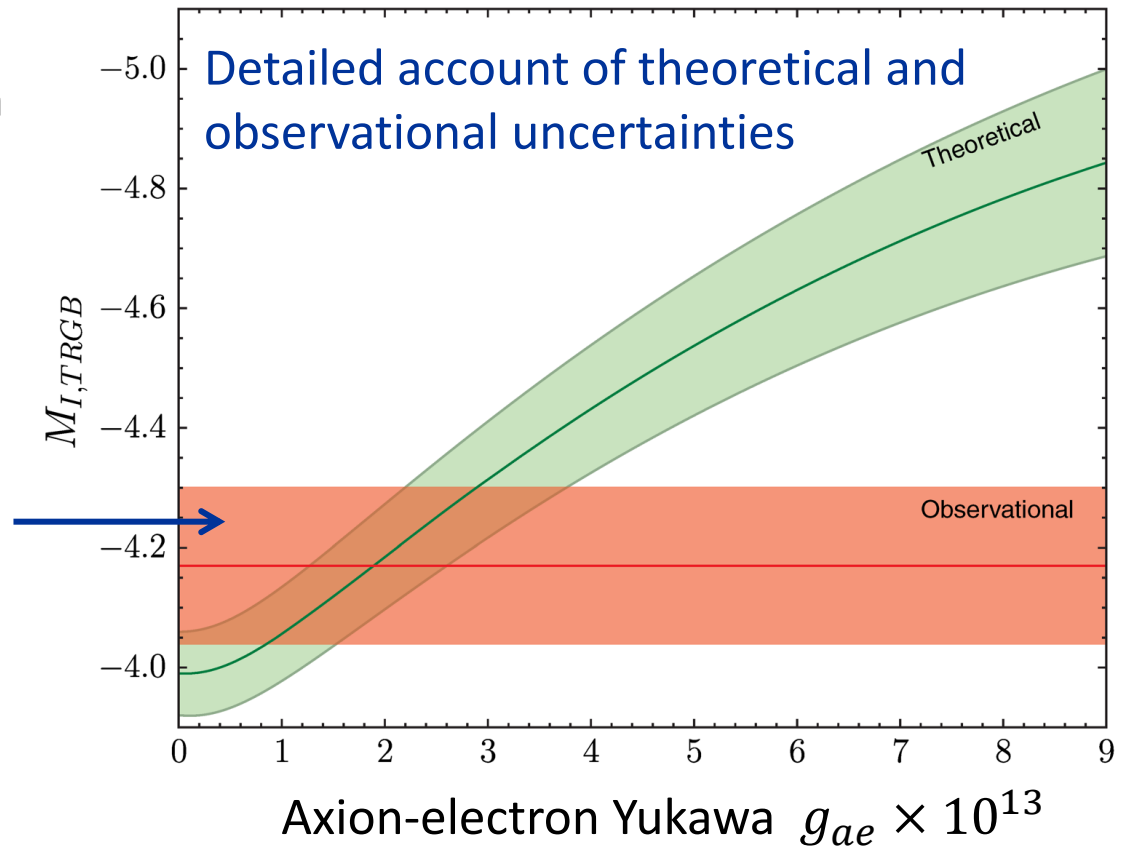
Galactic Globular Clusters



Limits on Axion-Electron Coupling from GC M5

I-band brightness
of tip of red-giant branch
[magnitudes]

- Uncertainty dominated by distance
- Can be improved in future (GAIA mission)



Limit on axion-electron Yukawa

$$g_{ae} < \begin{cases} 2.6 \times 10^{-13} & (68\% \text{ CL}) \\ 4.3 \times 10^{-13} & (95\% \text{ CL}) \end{cases}$$

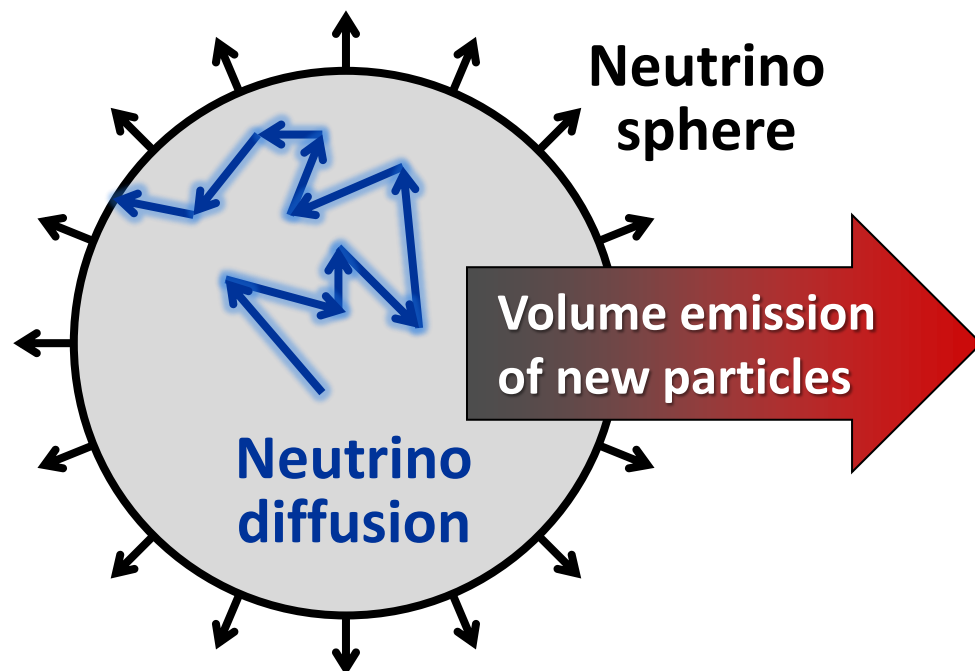
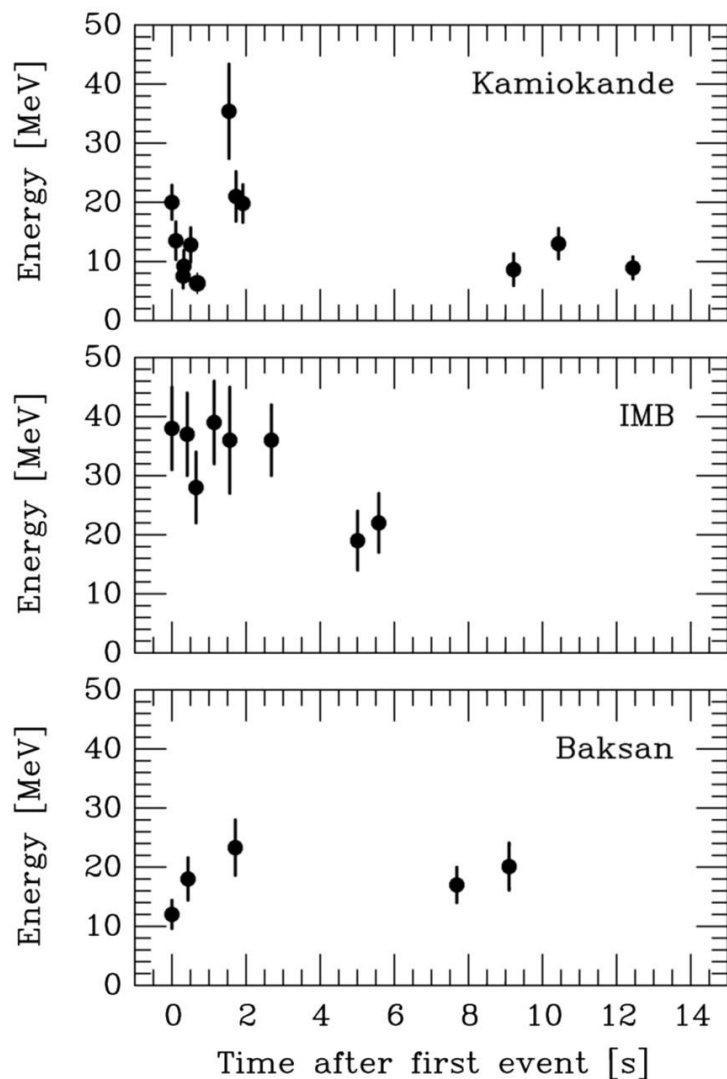
Mass limit in DFSZ model

$$m_a \cos^2 \beta < \begin{cases} 9.3 \text{ meV} & (68\% \text{ CL}) \\ 15.4 \text{ meV} & (95\% \text{ CL}) \end{cases}$$

Viaux, Catelan, Stetson, Raffelt, Redondo, Valcarce & Weiss, arXiv:1311.1669

Supernova 1987A Energy-Loss Argument

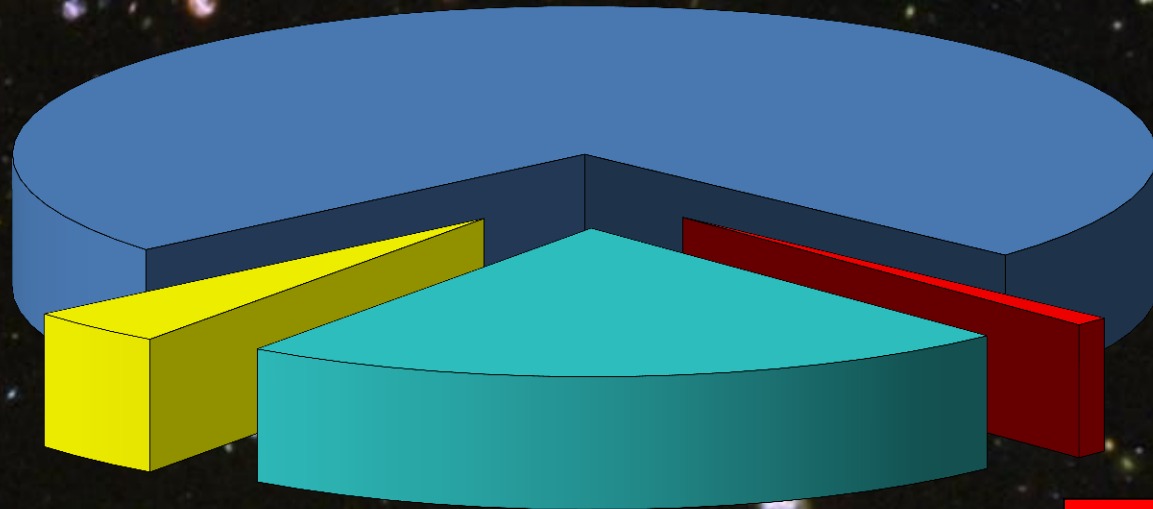
SN 1987A neutrino signal



Emission of very weakly interacting particles would “steal” energy from the neutrino burst and shorten it.
(Early neutrino burst powered by accretion, not sensitive to volume energy loss.)

Late-time signal most sensitive observable

Dark Energy 70%
(Cosmological Constant)

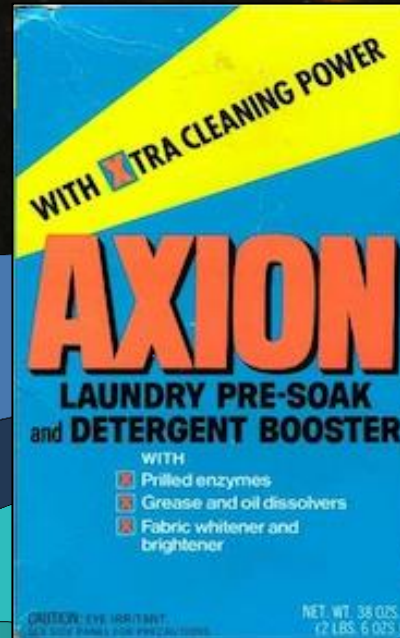


Ordinary Matter 5%
(of this only about
10% luminous)

Dark Matter
25%

Neutrinos
0.1–0.4%

Dark Energy 70%
(Cosmological Constant)



Ordinary Matter 5%
(of this only about
10% luminous)

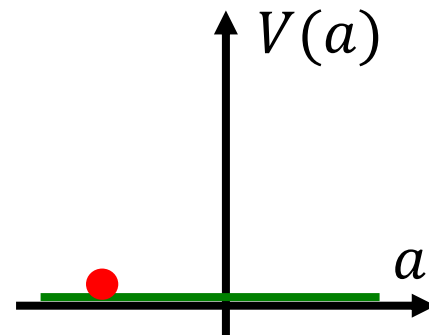
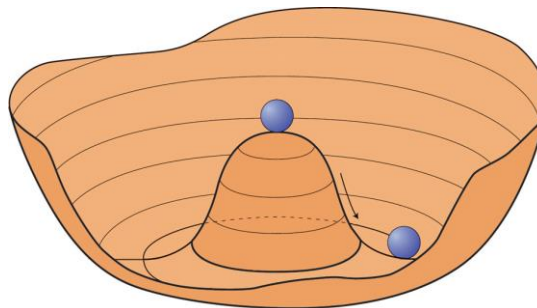
Dark Matter
25%

Neutrinos
0.1–0.4%

Creation of Cosmological Axions by Re-alignment

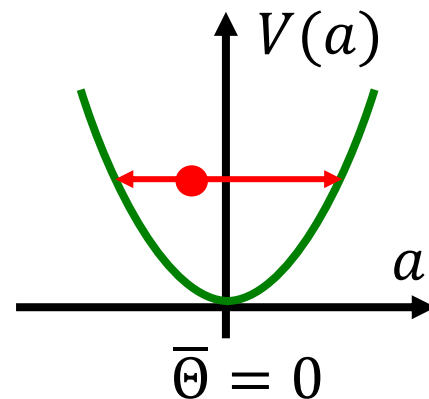
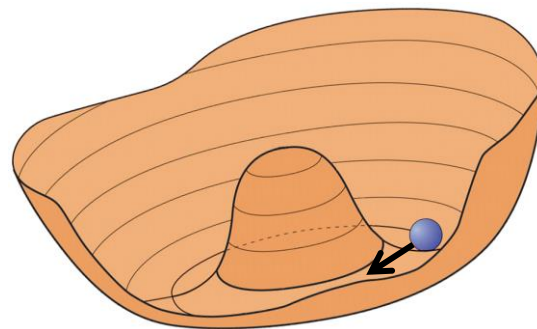
$T \sim f_a$ (very early universe)

- $U_{\text{PQ}}(1)$ spontaneously broken
- Higgs field settles in “Mexican hat”
- Axion field sits fixed at $a_i = \Theta_i f_a$



$T \sim 1 \text{ GeV}$ ($H \sim 10^{-9} \text{ eV}$)

- Axion mass turns on quickly by thermal instanton gas
- Field starts oscillating when $m_a \gtrsim 3H$
- Classical field oscillations (axions at rest)



Axions are born as nonrelativistic, classical field oscillations
Very small mass, yet cold dark matter

Axion Cosmology in PLB 120 (1983)

THE NOT-SO-HARMLESS AXION

Michael DINE

The Institute for Advanced Study, Princeton, NJ 08540, USA

and

Willy FISCHLER

Department of Physics

Received 17 September 1982

Received manuscript

Cosmological aspects discussed by Sikivie is not to give an upper bound

A COSMOLOGICAL BOUND ON THE INVISIBLE AXION

L.F. ABBOTT ¹

Physics Department, Brandeis University, Waltham, MA 02254, USA

and

P. SIKIVIE ²

Particle Theory

Received 14 September 1982

The product of f_a and m_a are found to be

COSMOLOGY OF THE INVISIBLE AXION

John PRESKILL ¹, Mark B. WISE ²

Lyman Laboratory of Physics, Harvard University, Cambridge, MA 02138, USA

and

Frank WILCZEK

Institute for Theoretical Physics, University of California, Santa Barbara, CA 93106, USA

Received 10 September 1982

We identify a new cosmological problem for models which solve the strong CP puzzle with an invisible axion, unrelated to the domain wall problem. Because the axion is very weakly coupled, the energy density stored in the oscillations of the classical axion field does not dissipate rapidly; it exceeds the critical density needed to close the universe unless $f_a \leq 10^{12}$ GeV, where f_a is the axion decay constant. If this bound is saturated, axions may comprise the dark matter of the universe.

Galactic WIMP vs Axion Dark Matter

Local dark matter density around 0.3 GeV/cm^3

WIMPs: Typical mass 100 GeV

- Local number density:

$$3000 \text{ m}^{-3} \frac{100 \text{ GeV}}{m_{\text{WIMP}}}$$

Dilute gas of particles

Axions: Typical mass $10 \mu\text{eV}$

- Local number density:

$$3 \times 10^{19} \text{ m}^{-3} \frac{10 \mu\text{eV}}{m_a}$$

- Typical galactic virial velocity:

$$v_{\text{gal}} \sim 10^{-3} c$$

- Typical de Broglie wavelength (“localisation”)

$$\lambda \sim \frac{2\pi}{m_a v_{\text{gal}}} \sim 100 \text{ m} \frac{10 \mu\text{eV}}{m_a}$$

- Minimal occupation number

$$f \sim 10^{25} \left(\frac{10 \mu\text{eV}}{m_a} \right)^4$$

Highly degenerate Bose gas!

Bose-Einstein Condensation of Dark Matter Axions

P. Sikivie and Q. Yang

Department of Physics, University of Florida, Gainesville, Florida 32611, USA

(Received 8 January 2009; revised manuscript received 23 July 2009; published 9 September 2009)

We show that cold dark matter axions thermalize and form a Bose-Einstein condensate (BEC). We obtain the axion state in a homogeneous and isotropic universe, and derive the equations governing small axion perturbations. Because they form a BEC, axions differ from ordinary cold dark matter in the nonlinear regime of structure formation and upon entering the horizon. Axion BEC provides a mechanism for the production of net overall rotation in dark matter halos, and for the alignment of cosmic microwave anisotropy multipoles.

- Axions $\sim$ WIMP dark matter on scales $\gg$ axion de Broglie wavelength?
- Larger-range correlation established by BEC dynamics? (Observable?)
- BEC formation caused by gravitational interactions possible?

See also • Erken, Sikivie, Tam & Yang, arXiv:1111.1157

• Saikawa & Yamaguchi, arXiv:1210.7080

• Noumi, Saikawa, Sato & Yamaguchi, arXiv:1310.0167

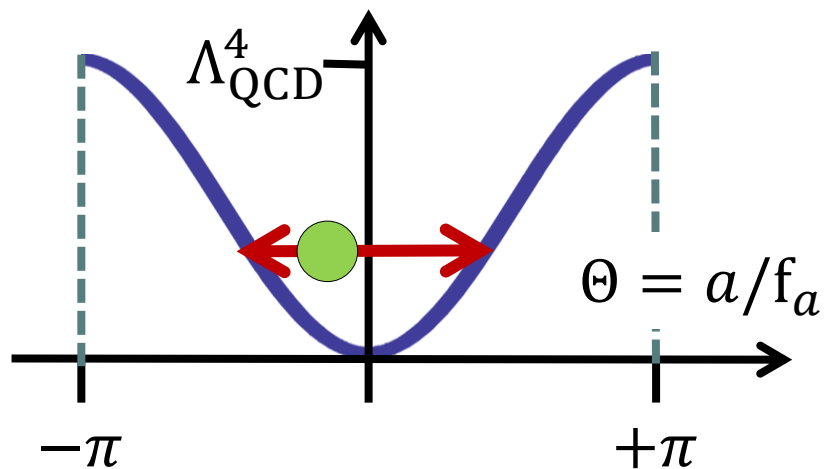
• Davidson & Elmer, arXiv:1307.8024, Davidson, arXiv:1405.1139

• Berges & Jaeckel, arXiv:1402.4776

• Guth, Hertzberg & Prescod-Weinstein, arXiv:1412.5930

~170 citations

Axion Dark Matter Driving Oscillators



Oscillating axion field (DM)

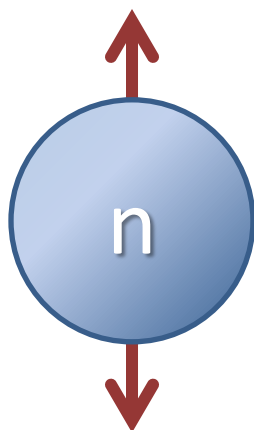
→ Oscillating Θ term

- Drives oscillating neutron EDM
- Drives oscillating E-field in microwave cavity w/ B-field

Assume axions are the galactic dark matter: $\rho_a \sim 300 \text{ MeV}/\text{cm}^3$

$$\rho_a = \frac{1}{2} m_a^2 \Phi_a^2 = \frac{1}{2} m_a^2 (\Theta f_a)^2 \sim \Theta^2 (m_\pi f_\pi)^2 \sim \Theta^2 \Lambda_{\text{QCD}}^4$$

$$\Theta(t) = a(t)/f_a \sim 3 \times 10^{-19} \cos(m_a t)$$



Neutron EDM
oscillating with
 $\omega = m_a$

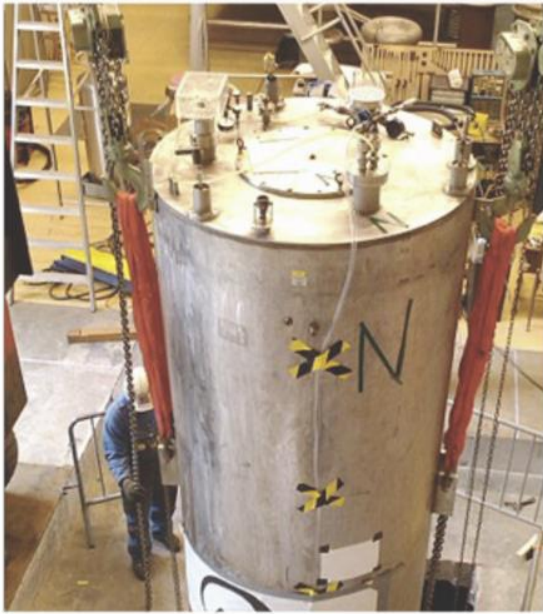
**External
static
B-field**



**Axion-induced
electric field
oscillating with
 $\omega = m_a$**

Landscape of Cavity Experiments

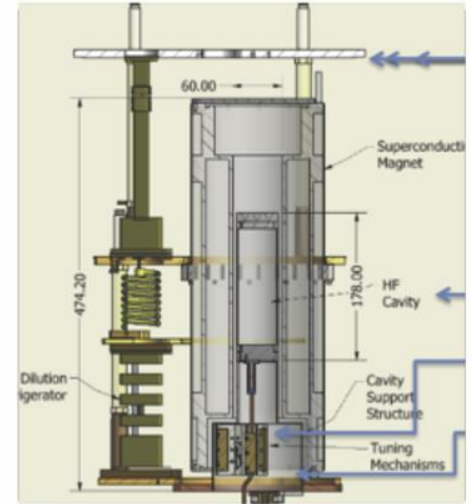
ADMX



CARRACK (discontinued)



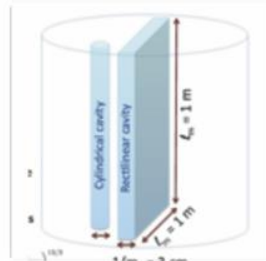
CULTASK - CAPP - Korea



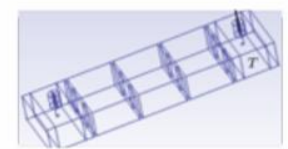
ADMX-HF



ADMX-Fermilab

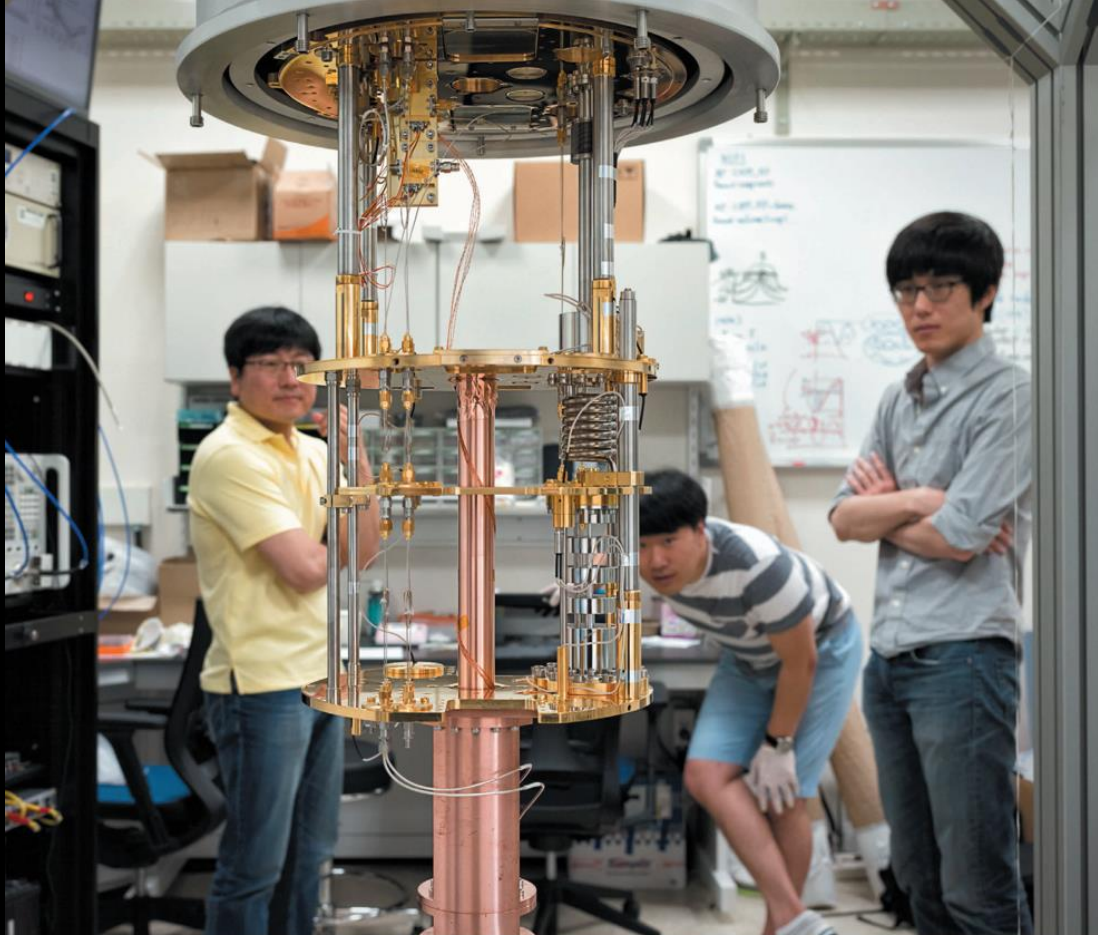


RADES



CAST-CAPP





South Korea's Nobel dream

The Asian nation spends more of its economic output on research than anywhere else in the world. But it will need more than cash to realize its ambitions.

BY MARK ZASTROW

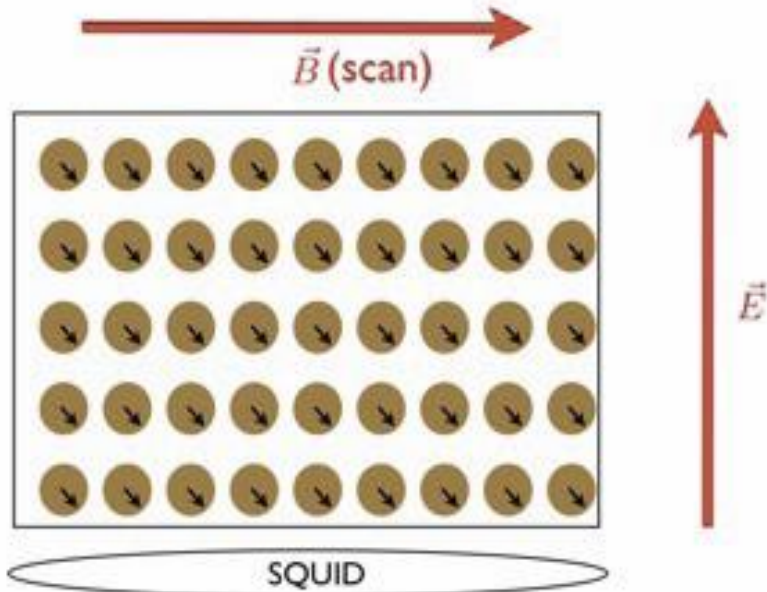
Behind the doors of a drab brick building in Daejeon, South Korea, a major experiment is slowly taking shape. Much of the first-floor lab space is under construction, and one glass door, taped shut, leads directly to a pit in the ground. But at the end of the hall, in a pristine lab, sits a gleaming cylindrical apparatus of copper and gold. It's a prototype of a device that might one day answer a major mystery about the Universe by detecting a particle called the axion — a possible component of dark matter.

If it succeeds, this apparatus has the potential to rewrite physics and win its designers a Nobel prize. "It will transform Korea, there's no question about it," says physicist Yannis Semertzidis, who leads the US\$7.6-million-per-year centre at South Korea's premier technical university, KAIST. But there's a catch: no one knows whether axions even exist. It's the kind of high-risk, high-reward project

SHIN WONG-AE

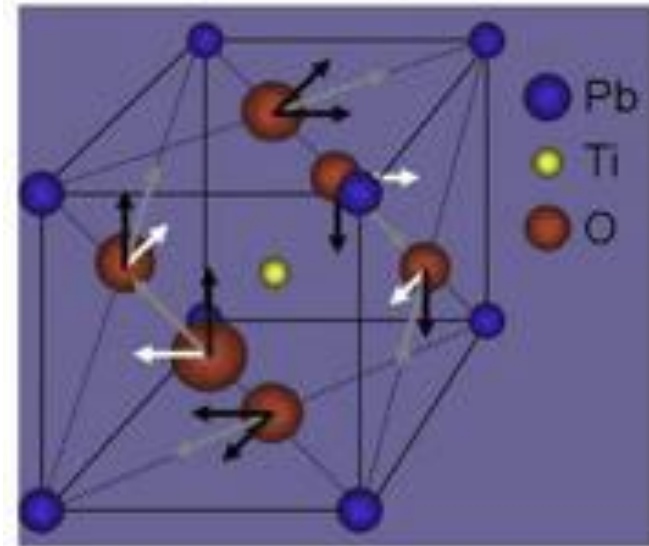
Searching for Axions in the Anthropic Range

Solid State Precision Magnetometry



$$\delta B \sim n\mu_N \frac{d_N E}{2\mu_N B - m_a} \sin((2\mu_N B - m_a)t) \sin(2\mu_N B t)$$

Polar Crystal



Lead Titanate

CASPER experiment

**Precise magnetometry to measure
tiny deviations from Larmor frequency**

Graham & Rajendran, arXiv:1101.2691

Budker, Graham, Ledbetter, Rajendran & Sushkov, arXiv:1306.6089

Resonantly Detecting Axion-Mediated Forces with Nuclear Magnetic Resonance

Asimina Arvanitaki¹ and Andrew A. Geraci^{2,*}

¹*Perimeter Institute for Theoretical Physics, Waterloo, Ontario N2L 2Y5, Canada*

²*Department of Physics, University of Nevada, Reno, Nevada 89557, USA*

(Received 5 March 2014; revised manuscript received 27 August 2014; published 14 October 2014)

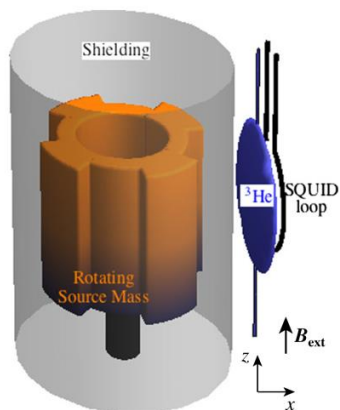


FIG. 1 (color online). A source mass consisting of a segmented cylinder with n sections is rotated around its axis of symmetry at frequency ω_{rot} , which results in a resonance between the frequency $\omega = n\omega_{\text{rot}}$ at which the segments pass near the sample and the resonant frequency $2\vec{\mu}_N \cdot \vec{B}_{\text{ext}}/\hbar$ of the NMR sample. Superconducting cylinders screen the NMR sample from the source mass and (not shown) the setup from the environment.

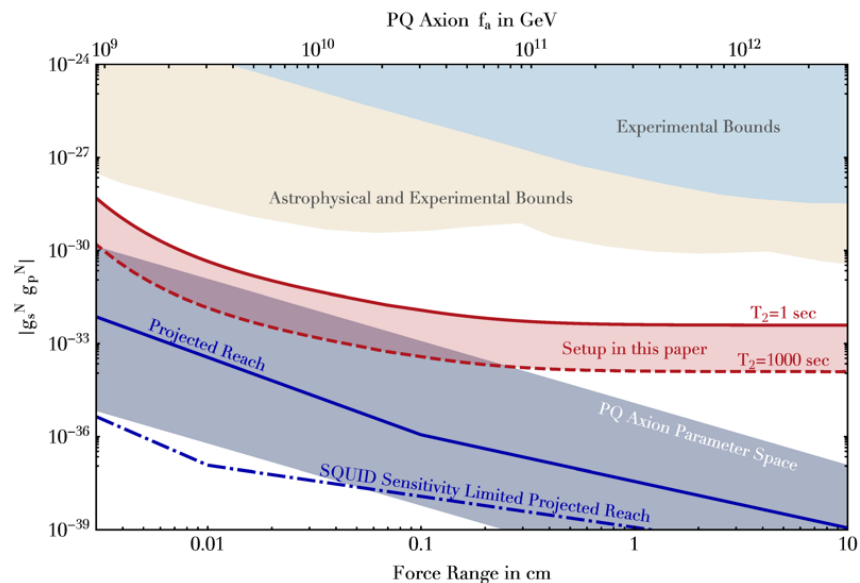
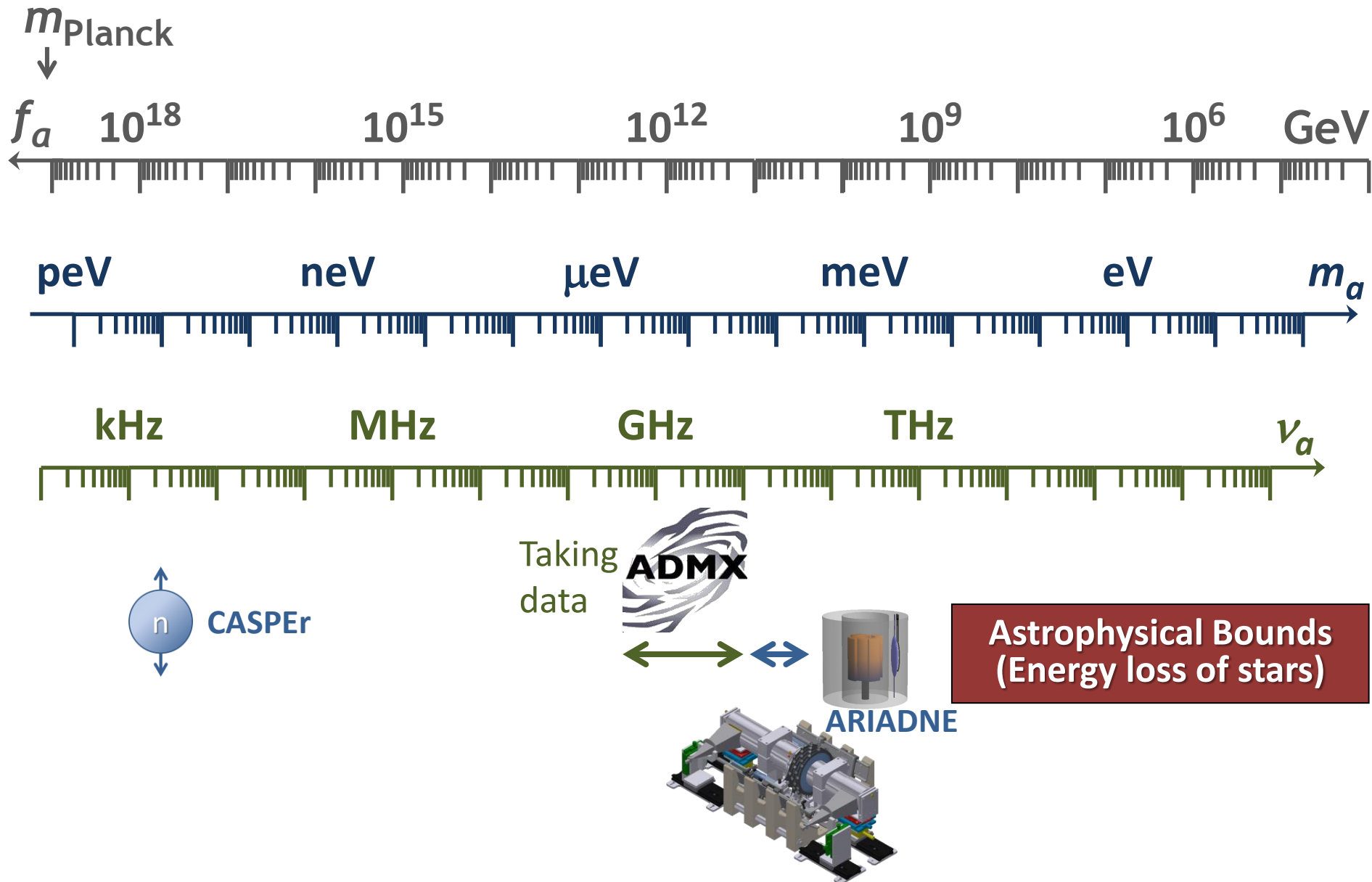


FIG. 2 (color online). Projected reach for monopole-dipole axion mediated interactions.

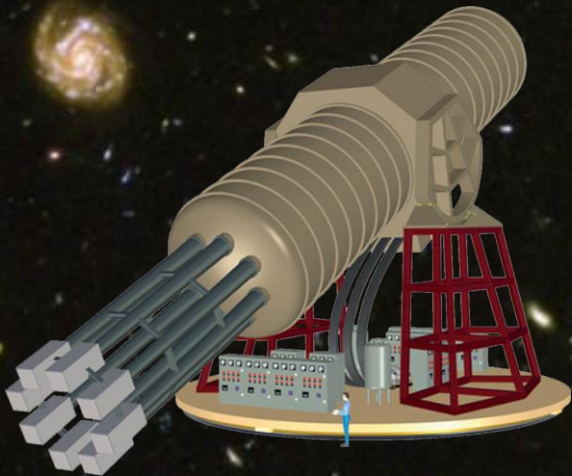
ARIADNE: Axion Resonant InterAction DetectionN Experiment

A.Geraci, A.Arvanitaki, A.Kapitulnik, Chen-Yu Liu, J.Long, Y.Semertzidis, M.Snow
(to be supported by NSF and/or DoE?)

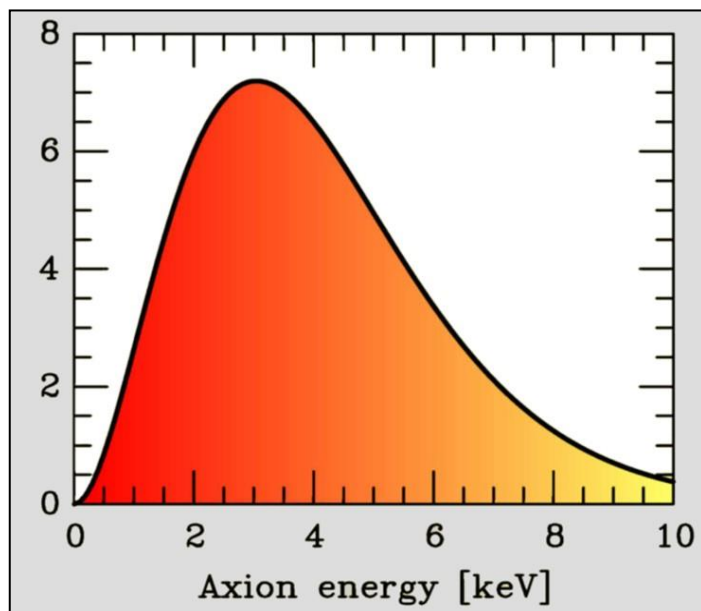
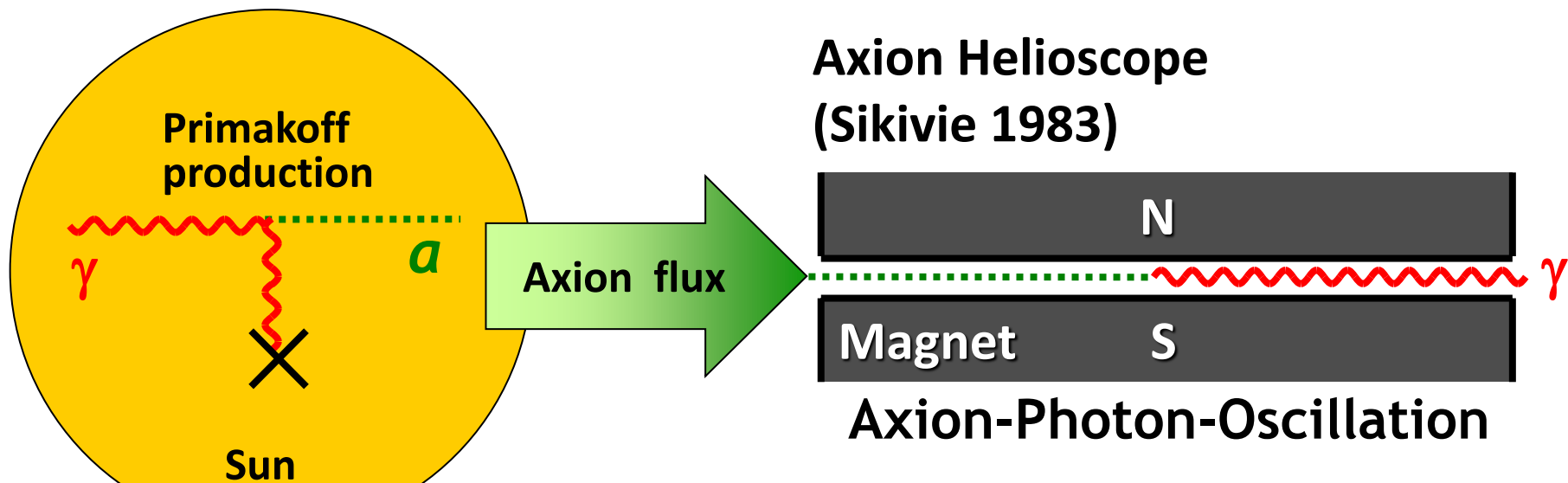
Axion Searches



Looking for axion-like particles (ALPs)



Search for Solar Axions



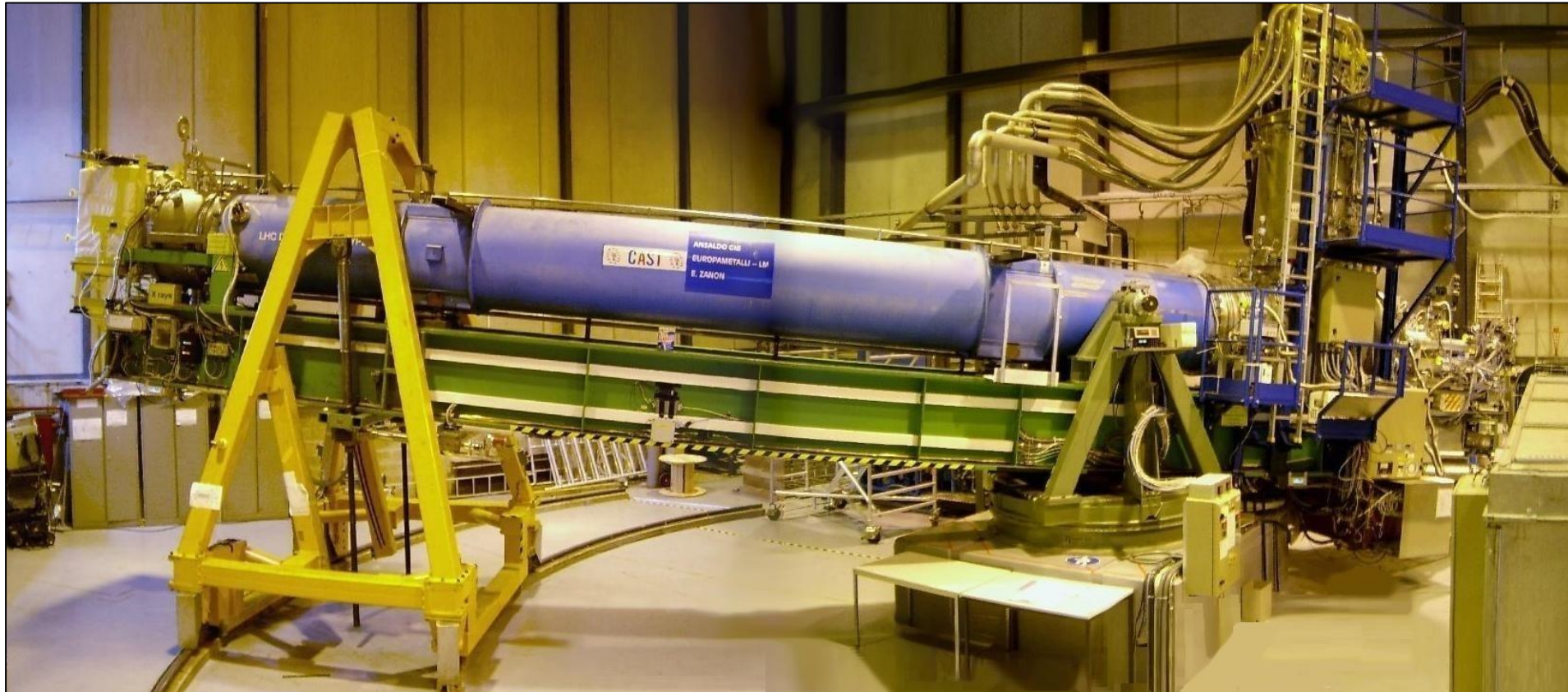
- Tokyo Axion Helioscope ("Sumico") (Results since 1998, up again 2008)
- CERN Axion Solar Telescope (CAST) (Data taking 2003–2015)

Alternative technique:

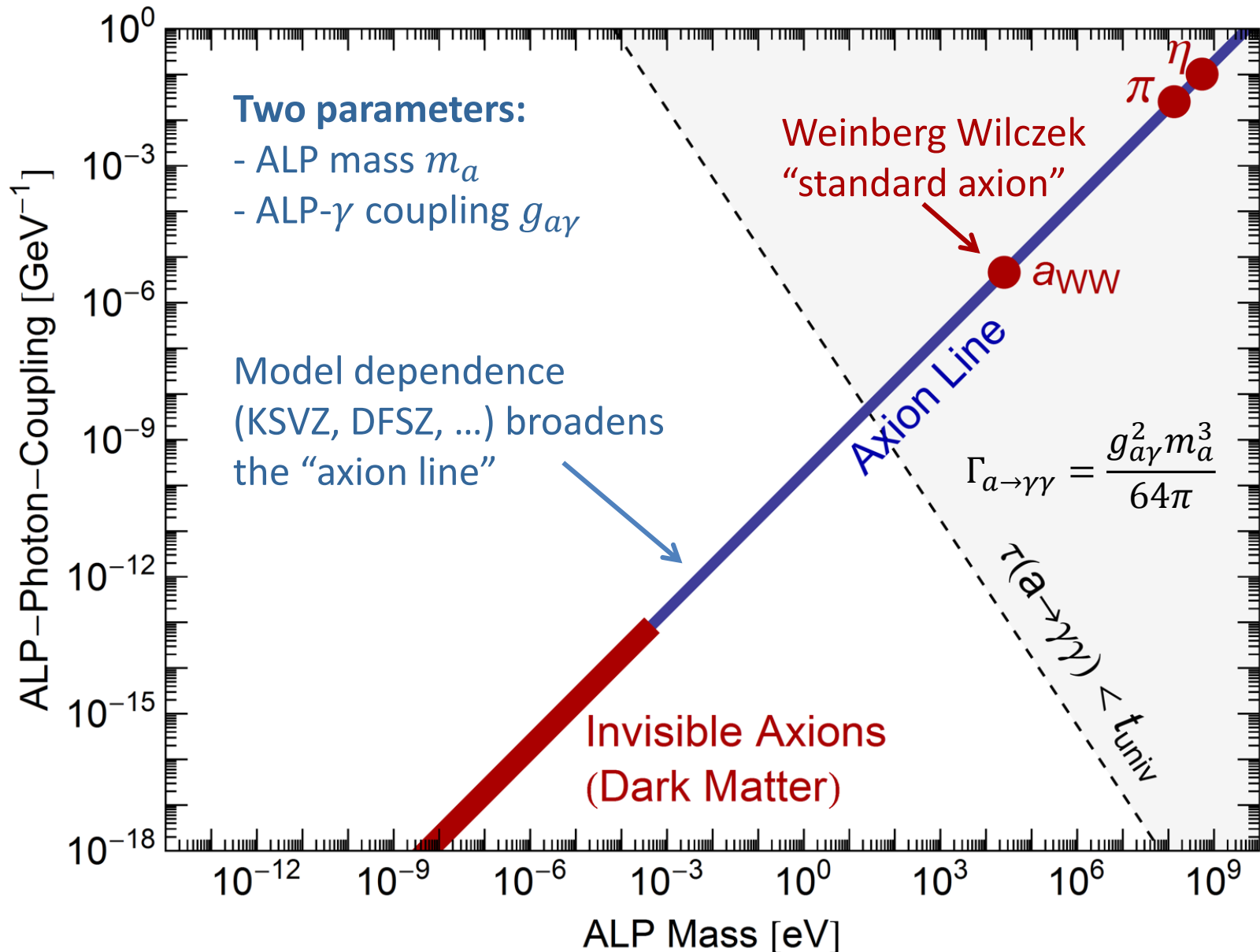
Bragg conversion in crystal

Experimental limits on solar axion flux from dark-matter experiments (SOLAX, COSME, DAMA, CDMS ...)

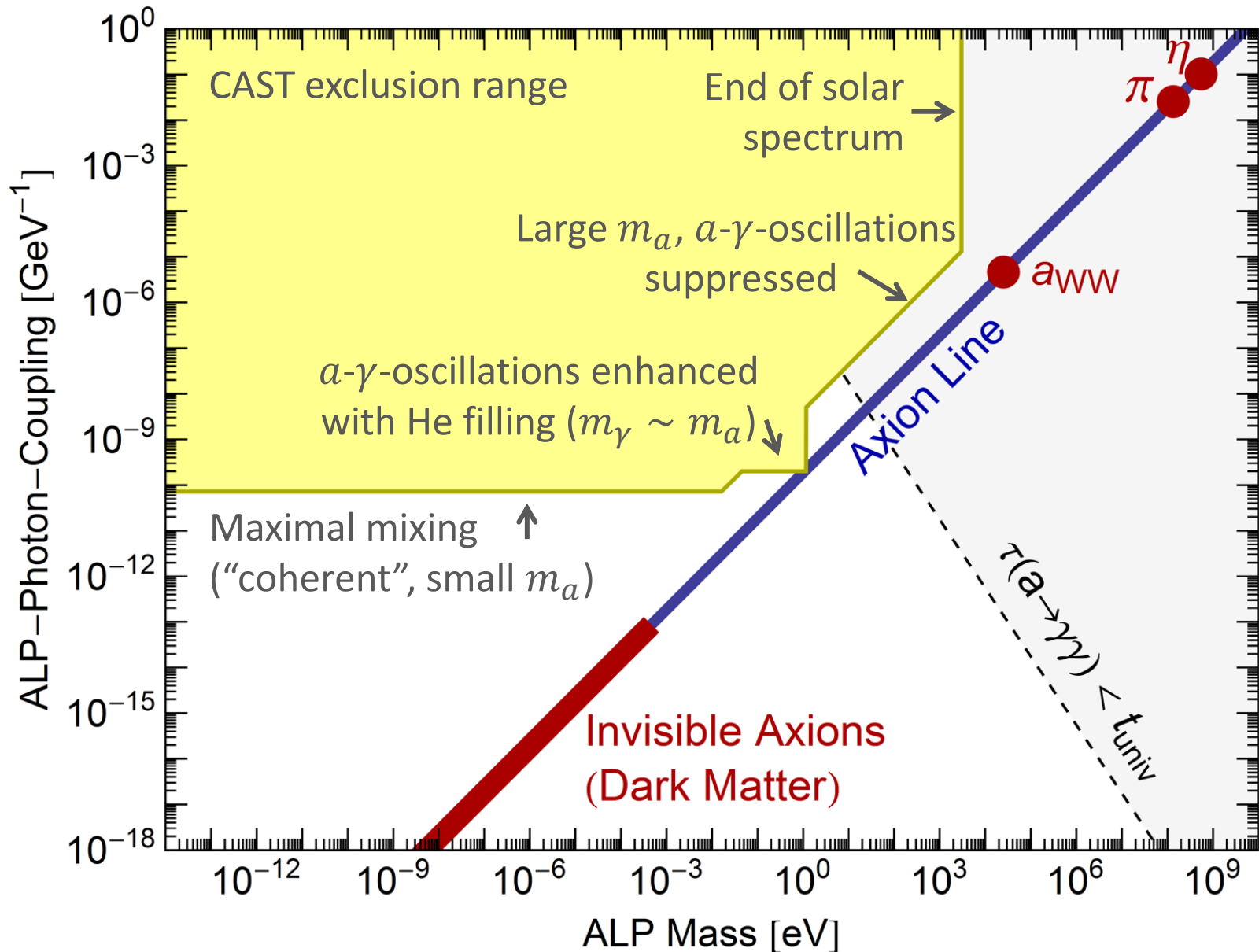
CERN Axion Solar Telescope (CAST)



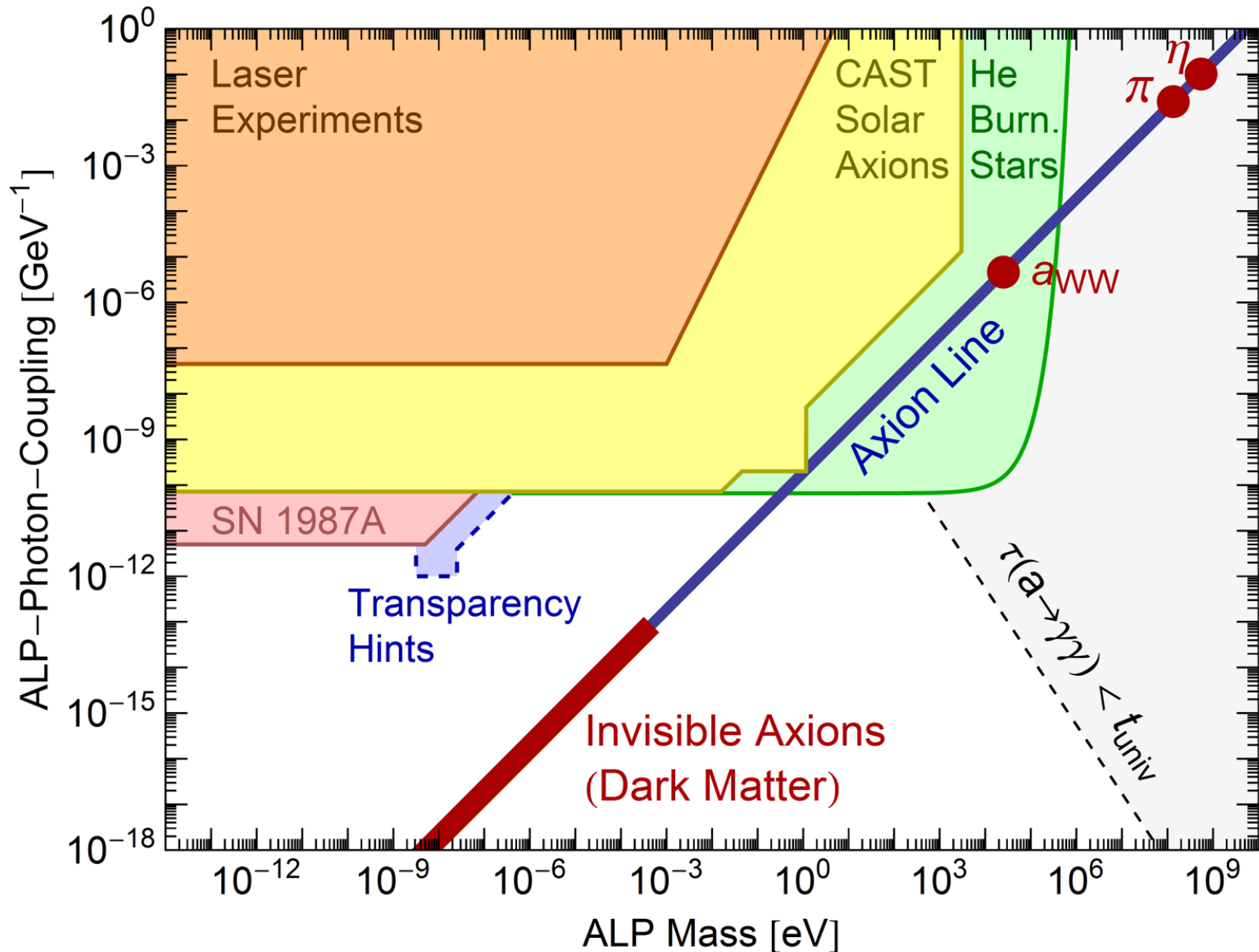
Parameter Space for Axion-Like Particles (ALPs)



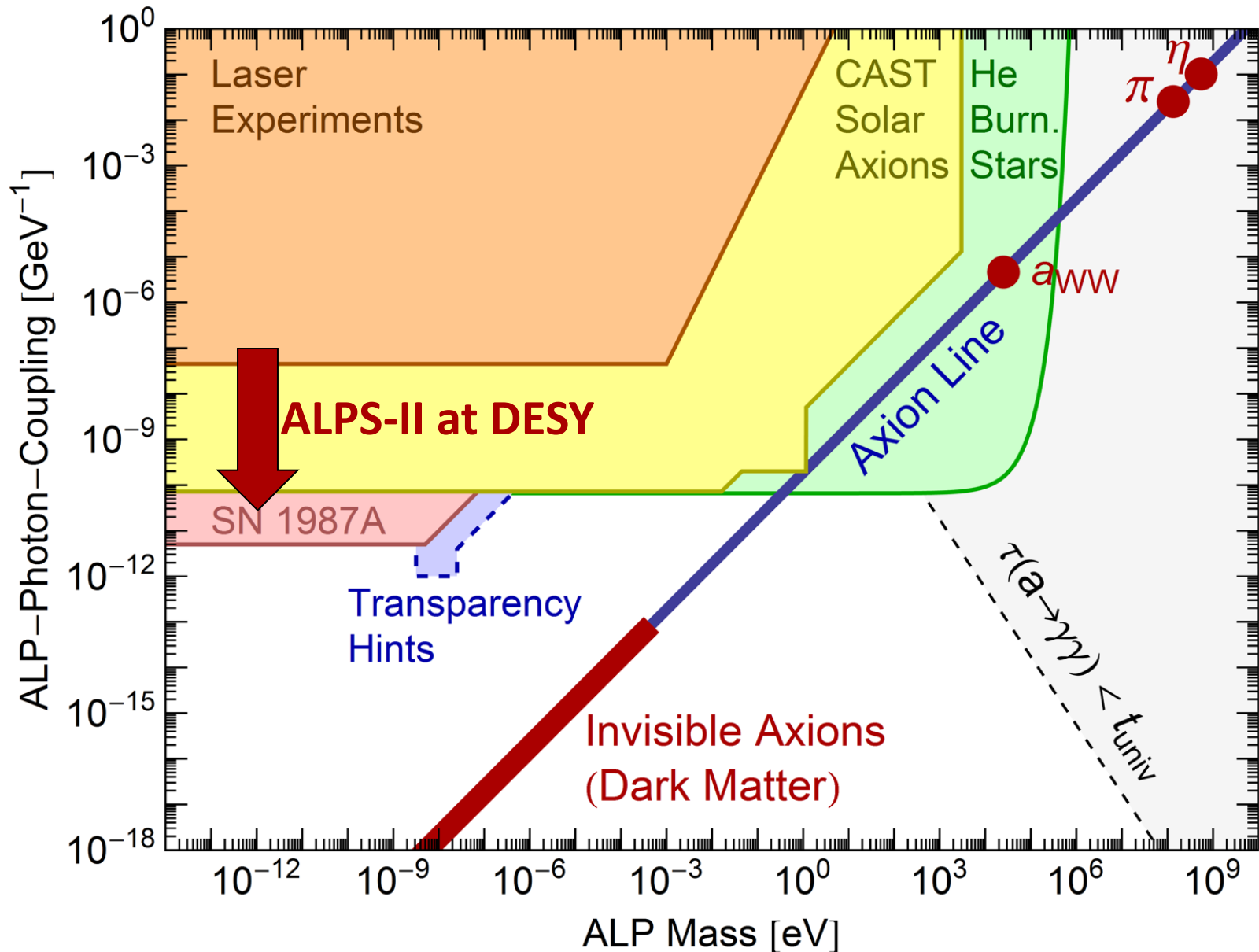
Parameter Space for Axion-Like Particles (ALPs)



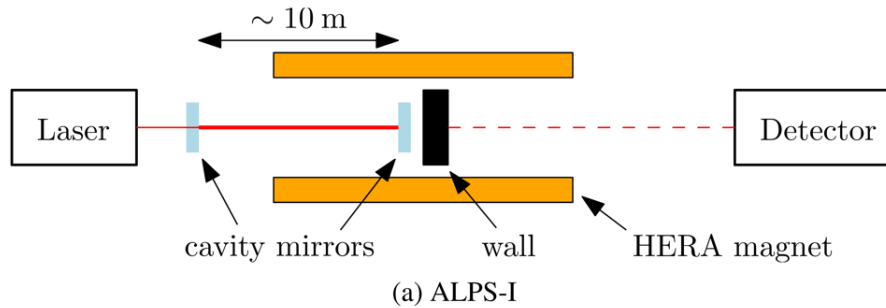
Parameter Space for Axion-Like Particles (ALPs)



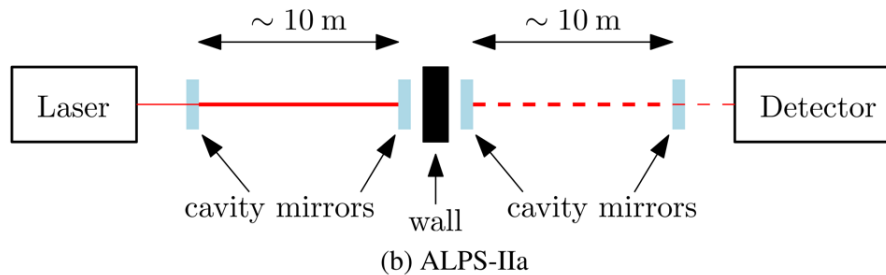
Parameter Space for Axion-Like Particles (ALPs)



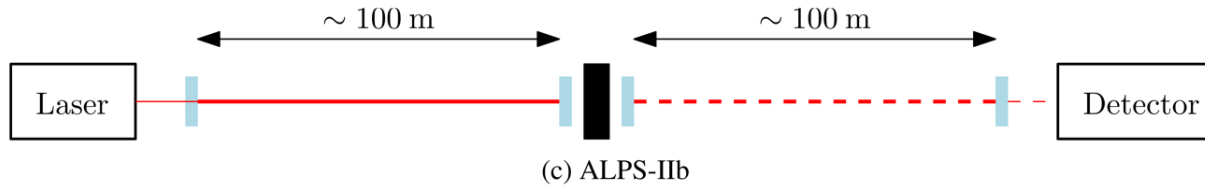
Any Light Particle Search II (ALPS-II) at DESY



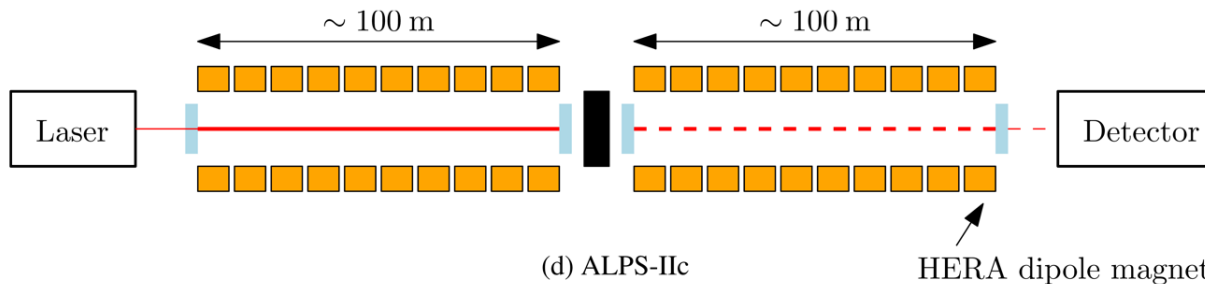
ALPS-I (finished)



ALPS-IIa (2014)



ALPS-IIb (2015)



ALPS-IIc (2017)

ALPS-II Technical Design Report, [arXiv:1302.5647](https://arxiv.org/abs/1302.5647)

Shining TeV Gamma Rays through the Universe

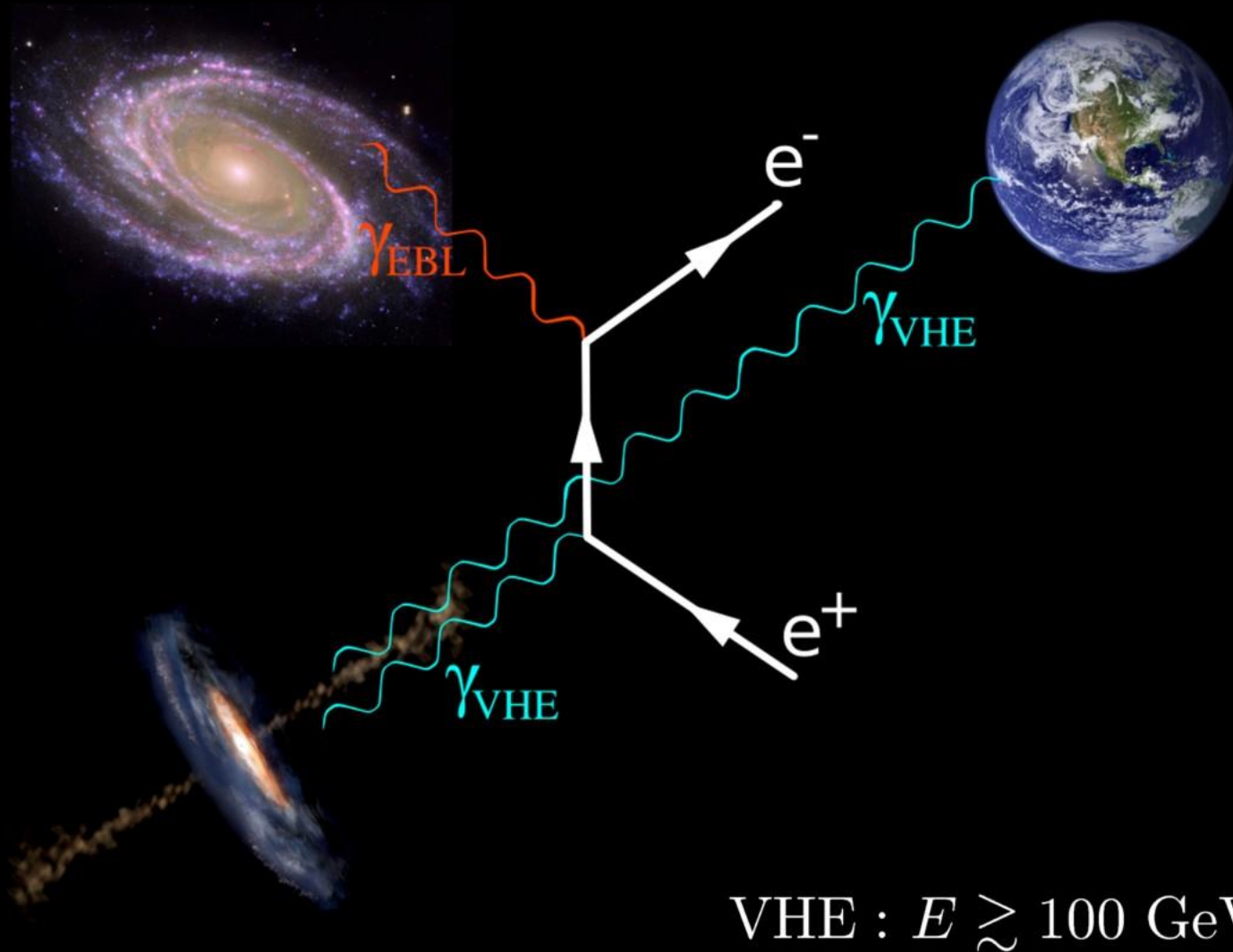


Figure from a talk by Manuel Meyer (Univ. Hamburg)

Shining TeV Gamma Rays through the Universe

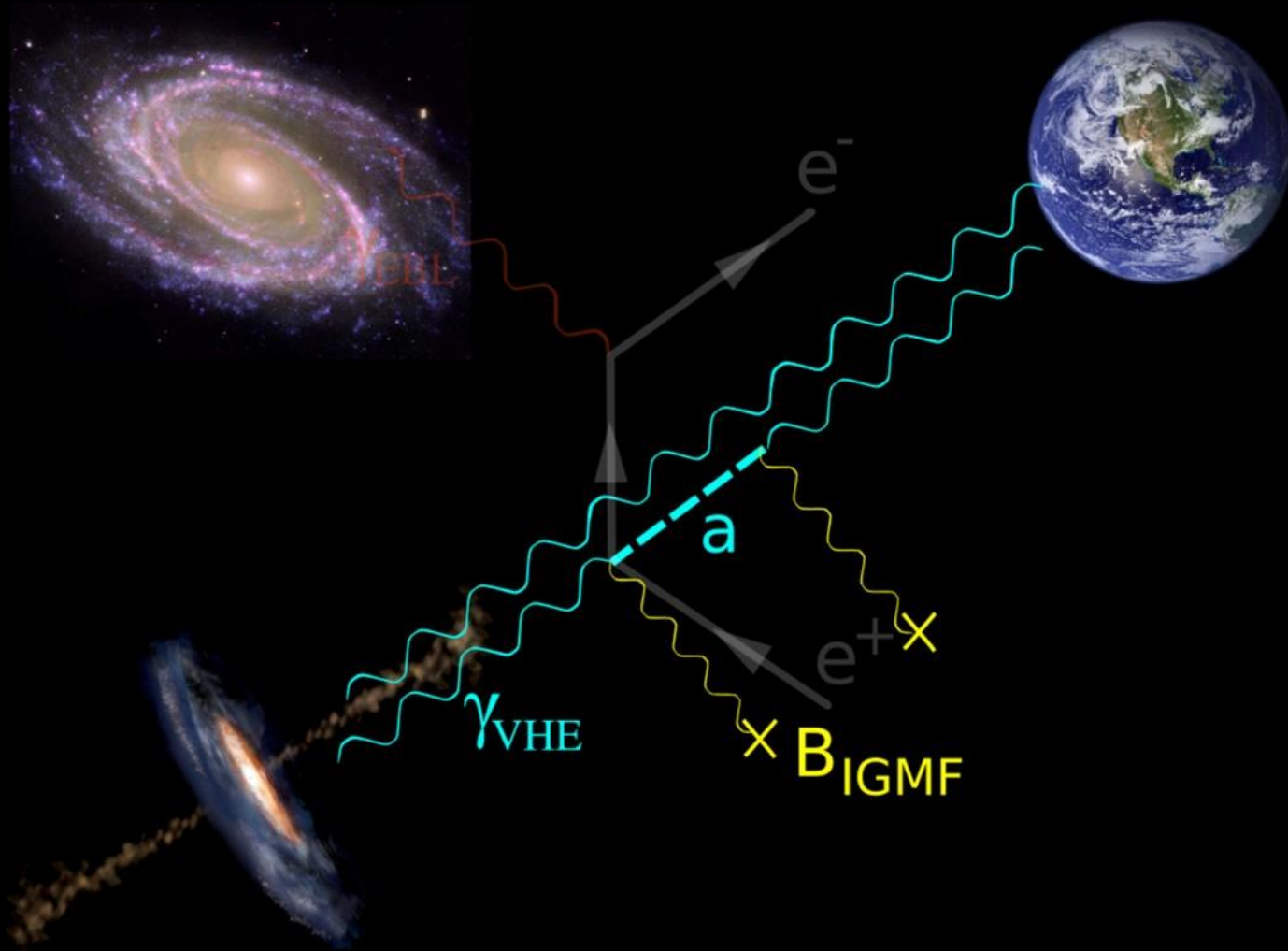
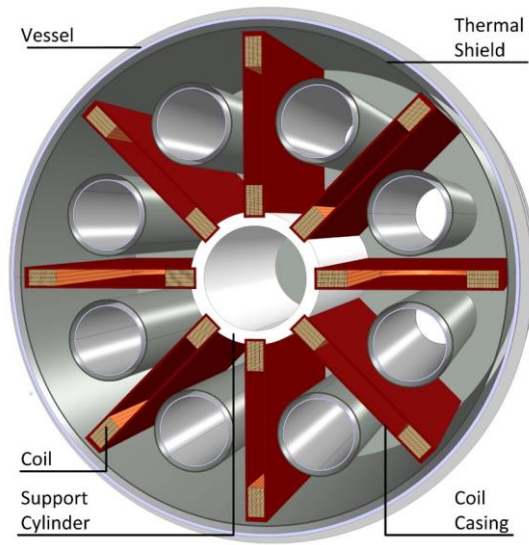
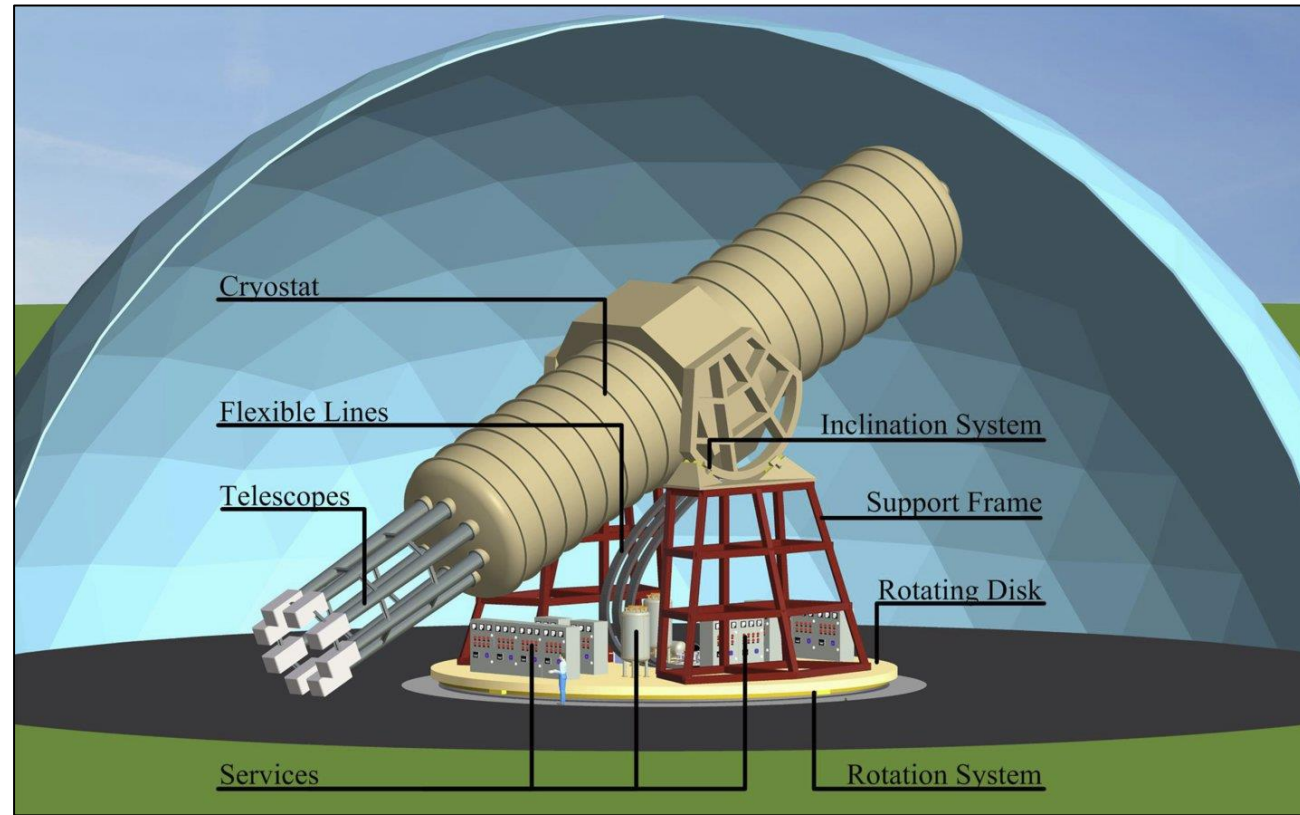


Figure from a talk by Manuel Meyer (Univ. Hamburg)

Next Generation Axion Helioscope (IAXO) at CERN

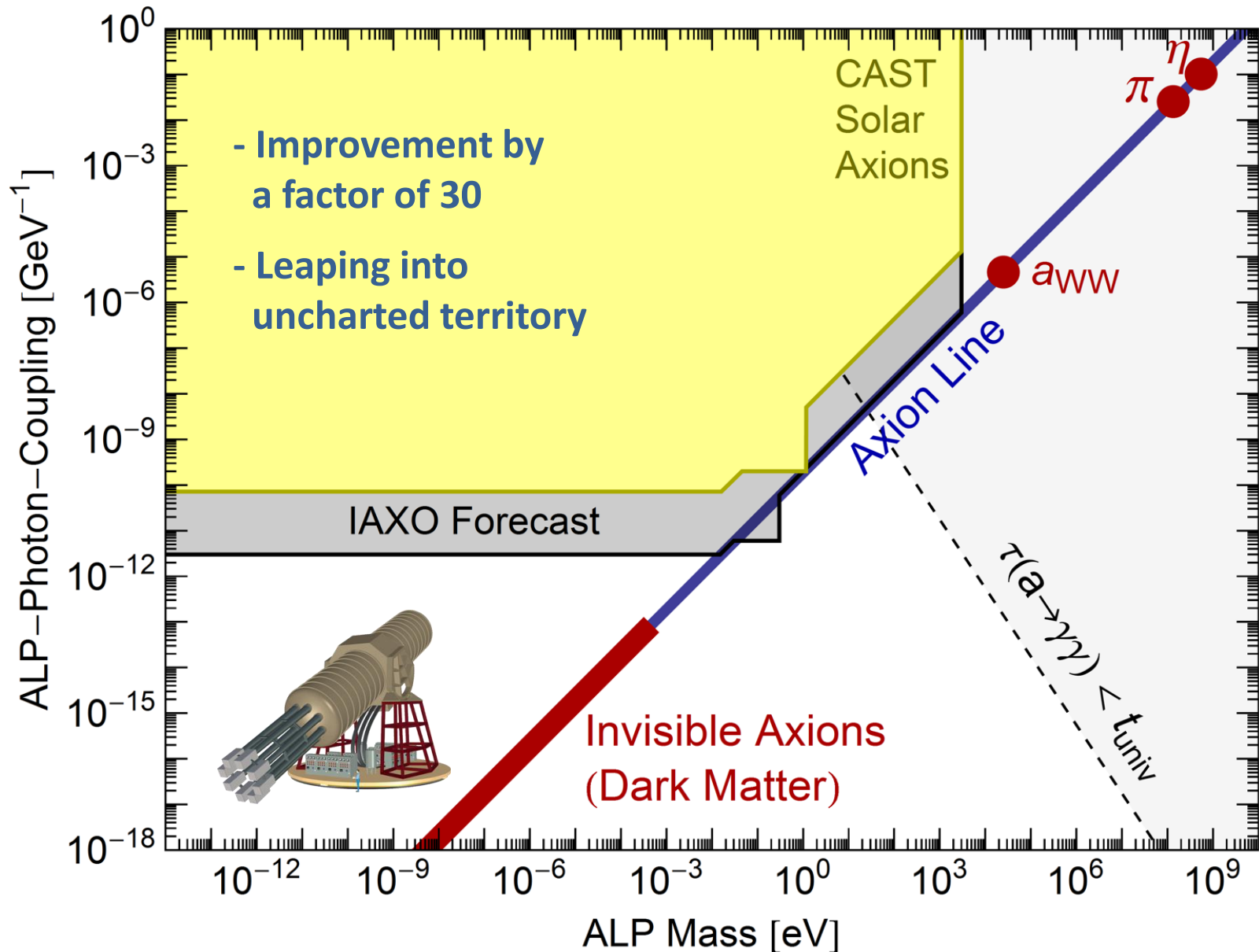


Need new magnet w/
– Much bigger aperture:
 $\sim 1 \text{ m}^2$ per bore
– Lighter (no iron yoke)
– Bores at T_{room}

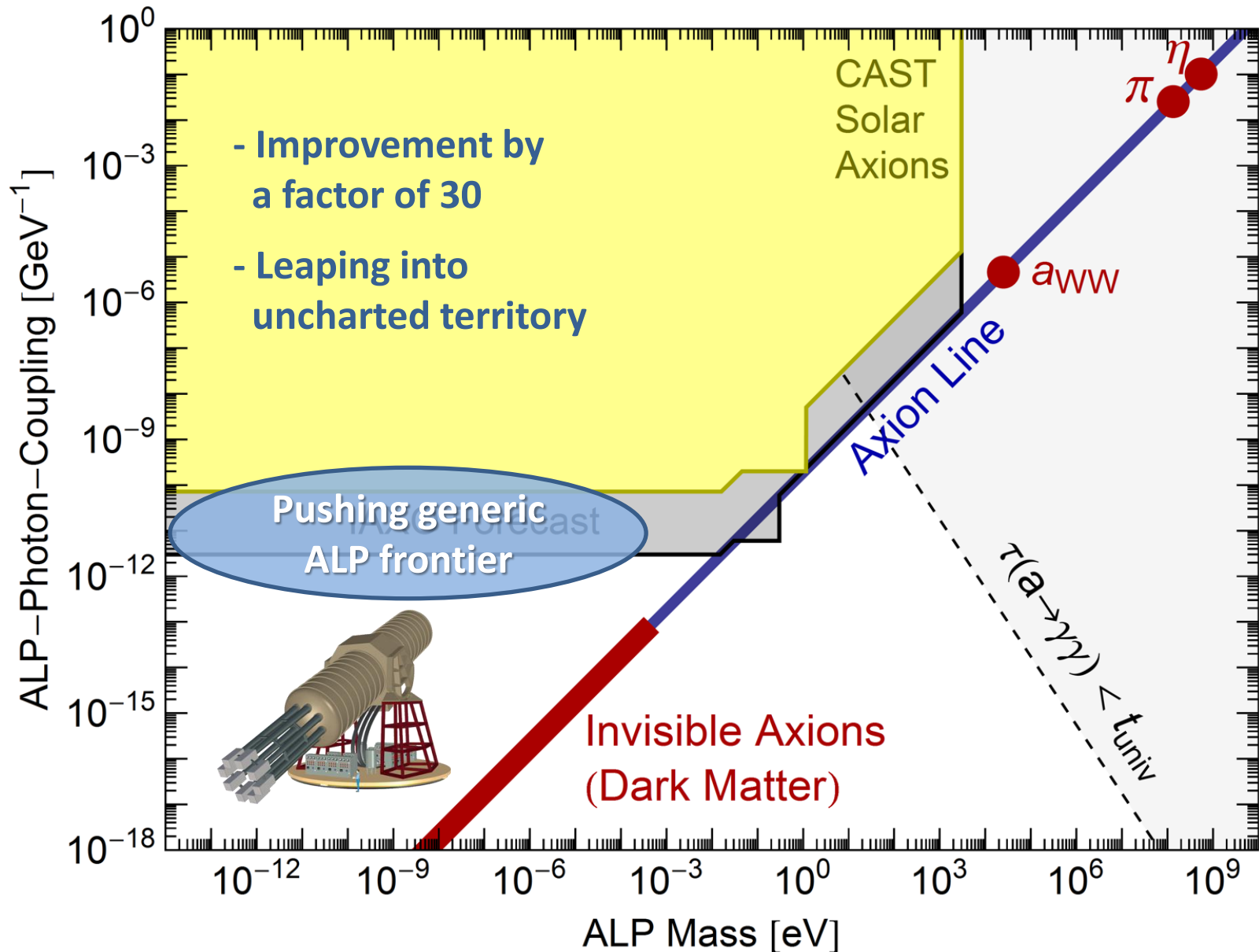


- Irastorza et al.: Towards a new generation axion helioscope, arXiv:1103.5334
- Armengaud et al.:
 Conceptual Design of the International Axion Observatory (IAXO), arXiv:1401.3233

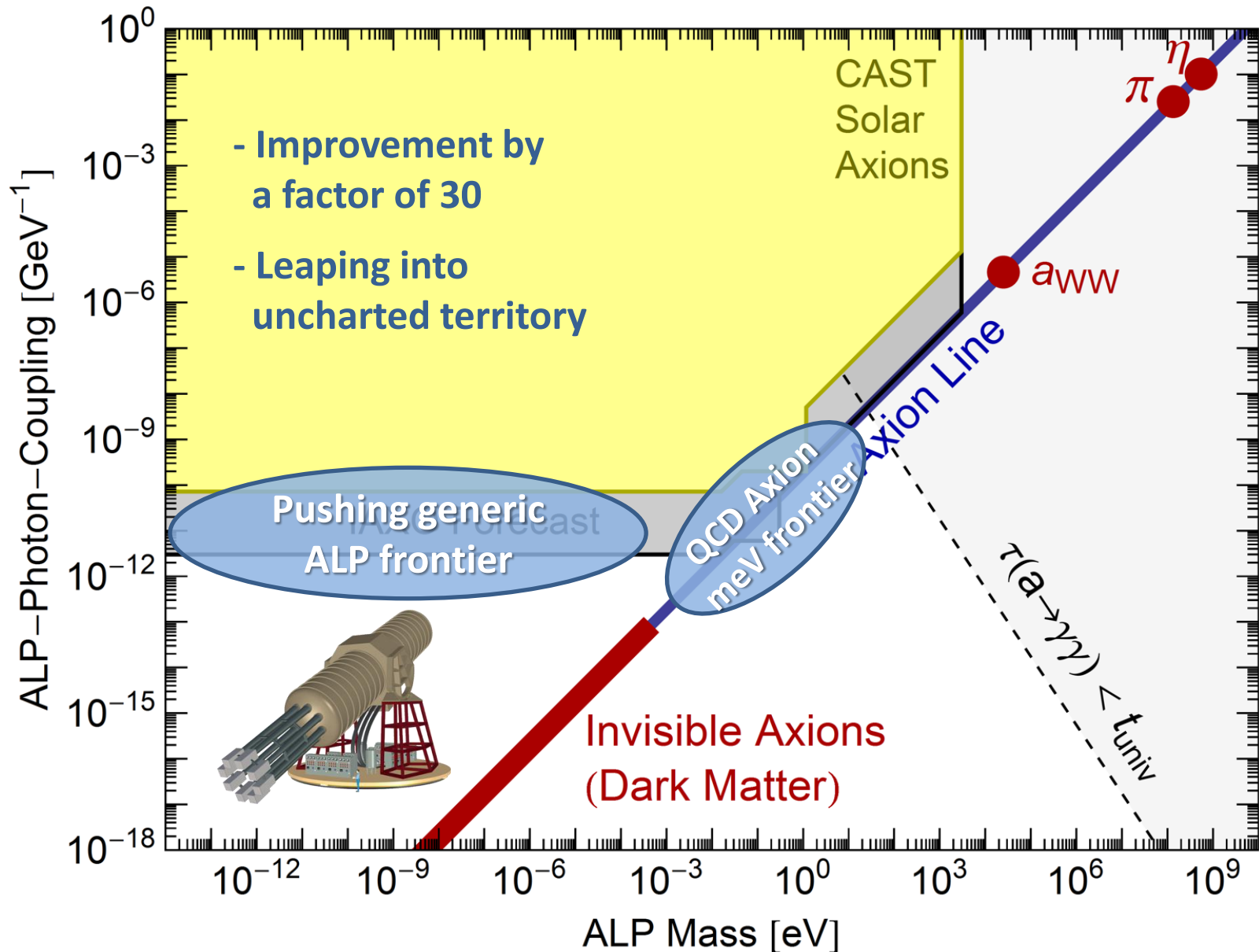
Parameter Space for Axion-Like Particles (ALPs)



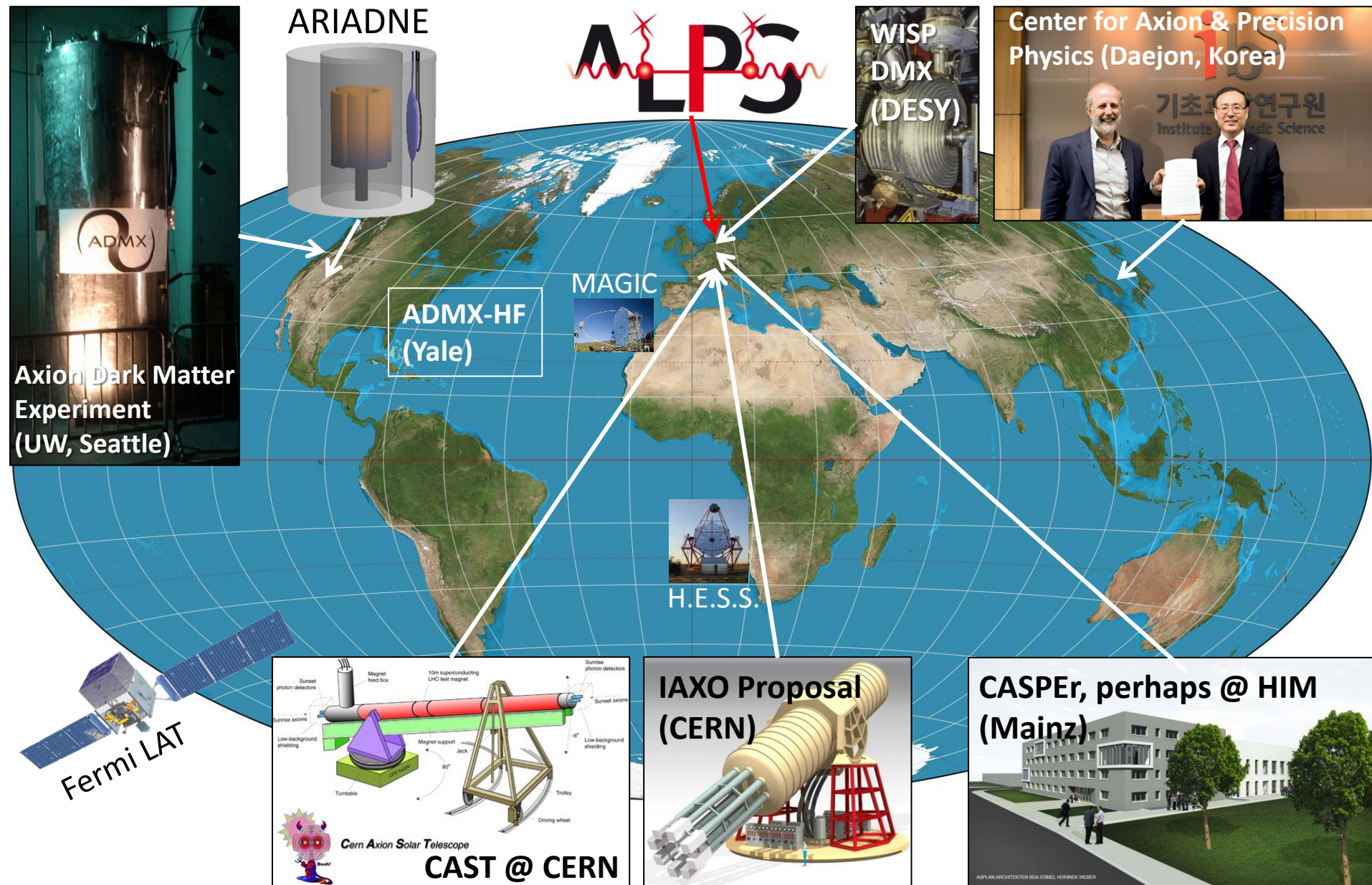
Parameter Space for Axion-Like Particles (ALPs)



Parameter Space for Axion-Like Particles (ALPs)



Axion and Axion-Like Particle Searches



Dow Jones Index for Axions: The Rally Continues

inSPIRE: Citation of Peccei-Quinn papers or title axion or axino

