# A CLOCKWORK WIMP

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# PLAN

- the clockwork mechanism illustrated
- construction of a clockwork wimp
- Majorana neutrino masses
- a bit of deconstruction
- conclusions

Based on arXiv:1612.06411, in collaboration with Thomas Hambye and Daniele Teresi

#### **ON SCALES AND MASSES**

new physics (local)  $\longrightarrow$  effective operators  $\longrightarrow$  new scale

Weinberg operator  $\longrightarrow$  L breaking scale Proton decay  $\longrightarrow$  GUT scale Axion  $\longrightarrow$  PQ symmetry breaking scale Quantum gravity  $\longrightarrow$  Planck scale

#### but MASS ≠ SCALE

. . .





# THE CLOCKWORK MECHANISM

Kaplan & Ratazzi (2015); Choi & Im (2015) ; Giuidice & McCullough (2016)

 $\Lambda \sim \frac{M}{y^x}$  large scale physics : large mass or tiny coupling?

We are often reluctant to introduce small parameters (even if natural in the sense of 't Hooft)

The clockwork is a possible mechanism to generate small numbers out of a theory with O(I) parameters or large effective scales out of dynamics at much lower energies

reading accessible energies...

# THE CLOCKWORK MECHANISM

#### a framework for model building....



# THE CLOCKWORK MECHANISM

- hierarchy problem Guidice & McCullough (2016)
- low scale invisible axion Guidice & McCullough (2016); Farina et al (2016)
- inflation Kehagias & Riotto (2016)
- neutrino physics Hambye, Teresi & MT (2016); Carena et al (2017); Ibarra et al (2017)
- Sugra Antoniadis, Delgado, Markou & Pokorski (2016)
- dark matter Hambye, Teresi & MT (2016) (this talk)





#### The Scalar Clockwork



one massless + N massive modes

$$\begin{split} \Phi_{0} & \Phi_{1} & \cdots & \Phi_{k} \sim e^{i\phi_{k}/f} & \cdots & \Phi_{N} \\ \mathbf{x} & \mathbf{x} \\ -\frac{m^{2}}{2} \sum_{0}^{N-1} (\phi_{k} - q\phi_{k+1})^{2} + \cdots & \longrightarrow & M^{2} = m^{2} \begin{pmatrix} 1 & -q & 0 & \cdots & 0 \\ -q & 1 + q^{2} & -q & \cdots & 0 \\ 0 & -q & 1 + q^{2} & -q & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ & & & -q & 1 + q^{2} & -q \\ 0 & 0 & 0 & \cdots & -q & q^{2} \end{pmatrix} \end{split}$$

 $q > 1 \rightarrow q^N \gg 1$ 

massless mode localized towards site k=0



plot from Farina et al (2016)



 $N = 15, \quad q = 3 \quad \longrightarrow q^N F \sim 10^{10} \left(\frac{F}{\text{TeV}}\right) \quad \text{GeV}$ 

large effective scale

# **A CLOCKWORK WIMP ?**

# WIMP abundance from thermal freeze-out



of GeV-TeV range

# WIMP stability protected by a symmetry

SUSY, SO(10), exact gauge symmetry, accidental symmetry,...

e.g. a Majorana SU(2) 5-plet (aka Minimal Dark Matter)

possible decay through dim-6 op.

$$\tau \sim \frac{\Lambda^4}{M^5}$$
 -

 $\longrightarrow \Lambda \gtrsim 10^{16} \mathrm{GeV}$ a large, difficult to test scale

#### CAN WE STABILIZE A WIMP THROUGH THE CLOCKWORK MECHANISM ?



chiral chain

$$\begin{cases} S_i \sim (1, -1) \text{ under } U(1)_{L_{i+1}} \times U(1)_{R_i} \\ C_i \sim (1, -1) \text{ under } U(1)_{L_i} \times U(1)_{R_i} \end{cases}$$

complex scalars

# a chiral chain (Weyl spinors) and complex scalars (spurions or dynamical)

construction goes through 4 steps



 $\mathcal{L} \supset -m \sum_{i=1}^{N} \left( \bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c.$ 





#### i.e. pretty much like the scalar clockwork



# step 2: break the residual chiral symmetry

i.e.

give a mass to the N state







$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left( \bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$

clockwork mechanism unaffected provided  $m_N \lesssim qm$ 

one light & localized mode N gears with O(1) couplings

(pseudo-Dirac if  $q \gg m_N/m$ )



Hambye, Teresi & MT (2016)



$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left( \bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$
$$-y (\bar{L}_{SM} \tilde{H} R_N + h.c.)$$

step 3: couple the chain to the SM



$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left( \bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$
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step 3: couple the chain to the SM

$$N \approx R_0 + \frac{1}{q}R_1 + \frac{1}{q^2}R_2 + \ldots + \frac{1}{q^N}R_N \rightarrow \mathcal{L} \supset -\frac{y}{q^N}\bar{L}_{SM}\tilde{H}N + h.c.$$

$$\lim_{k \to \infty} \underset{\text{coupling}}{\text{tiny}} \underset{\text{coupling}}{\text{tiny}}$$



The N mode is unstable but lifetime  $\gtrsim 10^{26} {
m sec}$  (gamma's, etc.)

if 
$$q^{2N} \gtrsim 10^{52} \left(\frac{m_N}{100 \text{GeV}}\right) y^2$$
 e.g.  $(q = 10, N = 26, y = 1)$ 



$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left( \bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$
$$-y (\bar{L}_{SM} \tilde{H} R_N + h.c.)$$

# last (but not the least) step: abundance from thermal freeze-out?

$$NN 
ightarrow hh, hZ, l\overline{l}, \dots$$
 but rate  $\propto rac{y^4}{q^{4N}}$ 

i.e. quite suppressed...

#### **A CLOCKWORK WIMP - AT LAST**



# A CLOCKWORK WIMP - AT LAST

with 
$$\xi_j^S \approx \sqrt{\frac{2}{N+1}} \sin\left(\frac{j\pi}{N+1}\right) y_S$$
 and  $\xi_j \approx \theta_S \xi_j^S$   
 $NN \leftrightarrow S_1 S_1, S_1 h, hh, \dots \rightarrow \propto \xi_S^2 \ (\xi^2) = \sum_{j=1}^N |\xi_j^S|^2 \approx y_S^2 \ (\theta_S)^2$ 



O(I) couplings, the N could be a WIMP!

### **A CLOCKWORK WIMP - TO RECAP**



# **A CLOCKWORK WIMP - TO RECAP**



# A CLOCKWORK WIMP — SCENARIO A

 $m_{S_1} = 150 \text{ GeV}$ 





<u>black solid</u>:  $y_S$  required for  $\Omega_{dm} \approx 0.25$ 

<u>blue solid</u>: LUX16 exclusions (below lines)

<u>red dashed:</u> Xenon I T reach

Higgs mediated  $\propto (\theta_S/q)^2$ 



#### A CLOCKWORK WIMP — SCENARIO B





black solid:  $y_S \theta_S$  required for  $\Omega_{dm} \approx 0.25$ 

<u>blue solid</u>: LUX16 exclusions (below lines)

# **A CLOCKWORK WIMP — BASIC PHENO**

Indirect detection p-wave annihilation, but decay into monochromatic neutrinos  $N \to h \nu$ 

e.g. El Aisati, Gustafsson & Hambye (2015)

# A CLOCKWORK WIMP — BASIC PHENO

**Indirect detection** p-wave annihilation, but decay into monochromatic neutrinos  $N \rightarrow h \nu$ 

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 $\psi_i$  in the hundred GeV/TeV range, coupled via Collider searches  $yL_{SM}HR_N$  with y sizable and  $\psi_j \supset R_N$ 



approximation assuming single state

observable pseudo-Dirac sterile RH neutrinos! **EWPT:**  $|\theta_{l\psi}|^2 \equiv y^2 v^2 / (2m_{\psi}^2) \lesssim 10^{-3}$ <u>LFV:</u>  $BR(\mu \to e\gamma) \sim 10^{-3} |\theta_{e\psi}|^2 |\theta_{\mu\psi}|^2 \lesssim 10^{-13}$ <u>L conserving searches:</u>  $m_{\psi} \lesssim 200 \text{ GeV}$  with  $300 \text{ fb}^{-1}$ Das, Dev & Okada (2014) <u>L violating searches:</u>  $m_{\psi} \lesssim 300 \text{GeV}$  with  $300 \text{ fb}^{-1}$  $m_N \lesssim m_\psi$ 

Deppisch, Dev & Pilaftsis (2015)

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assuming single state approximation observable pseudo-Dirac sterile RH neutrinos! <u>EWPT:</u>  $|\theta_{l\psi}|^2 \equiv y^2 v^2 / (2m_{\psi}^2) \lesssim 10^{-3}$ <u>LFV:</u>  $BR(\mu \to e\gamma) \sim 10^{-3} |\theta_{e\psi}|^2 |\theta_{\mu\psi}|^2 \lesssim 10^{-13}$ <u>L conserving searches:</u>  $m_{\psi} \lesssim 200 \text{ GeV}$  with 300 fb<sup>-1</sup> Das, Dev & Okada (2014) <u>L violating searches:</u>  $m_{\psi} \lesssim 300 \text{GeV}$  with 300 fb<sup>-1</sup>  $m_N \lesssim m_{\psi}$ Deppisch, Dev & Pilaftsis (2015)



### A CLOCKWORK WIMP IS POSSIBLE



# **CLOCKWORK MAJORANA NEUTRINO MASS**

Hambye, Teresi & MT (2016)

another chain...





has to go through the whole chain  $m_{\nu} \sim y^2 \frac{v^2}{q^N M}$ smaller chain than for DM stability  $(q = 10, N = 7, M = 1 \,\text{TeV})$ but one/SM neutrino mass!

(i.e. at least 2 chains for neutrino, one for the DM...)

See also Carena, Li, Machado, Machado & Wagner (2017); Ibarra, Kushwaha & Vempati (2017)

#### A LOT OF FIELDS, WITH VERY SPECIFIC COUPLINGS...



# **A CLOCKWORK WIMP FROM DECONSTRUCTION**



agrangian can be related to a discretized flat 5th dimension 
$$(Z)$$

$$\mathbf{L}_5 \supset \bar{\psi}(i\overleftrightarrow{\partial} - M)\psi = i\bar{\psi}\gamma^{\mu}\partial_{\mu}\psi + \left[\frac{1}{2}\left(\bar{L}\partial_Z R - \partial_Z \bar{L}R\right) - M\bar{L}R + h.c.\right]$$

- 2) naive discretization  $\Rightarrow$  fermion doubling
- $\Rightarrow$  add a Wilson term  $-\frac{a}{2}\partial_Z \bar{\psi} \partial_Z \psi$  with lattice spacing  $a = \frac{\pi R}{N} \to 0$
- 3) Dirichlet condition  $L(0) = 0 \Rightarrow$  surviving chiral mode  $R_0$
- 4) Standard Model degrees of freedom at  $Z = \pi R$  (~ braneworld)

$$\mathcal{L} \supset \sum_{i=0}^{N-1} \frac{1}{a} \bar{L}_{i+1} R_i - \sum_{i=1}^{N} \left( \frac{1}{a} + M \right) \bar{L}_i R_i \qquad \begin{array}{c} \text{well-defined} \\ \text{continuum limit} \\ m & qm & q^N = \left( 1 + \frac{\pi RM}{N} \right)^N & \rightarrow e^{\pi RM} \\ \end{array}$$
Hambye, Teresi & MT (2016) 
$$1/M \ll \pi R$$

### CONCLUSIONS

We have build a clockwork WIMP It is accidentally stable (not protected by a symmetry)

Its decay is mediated by the very same degrees of freedom that determine its abundance

These many degrees of freedom lie in the 100 GeV-TeV range

They could be seen at the LHC as RHN

There are also (possibly) plenty of scalars

Notice that it's also a framework for low scale SM neutrino Majorana mass

If you care, the clockwork could be seen as an extra spatial dimension



# [literotica.com]

Loving Wives > A Clockwork Wimp

### A Clockwork Wimp

by Lord\_Gino\_X ©

A Clockwork Wimp (Book 1 of 2)

#### Chapter 1

I slowly slide the large window open and climbed inside without making any noise. But there were noises. Noises that were tearing my ear drums apart. Noises that were directly coming from my chest. I knew I have to control my breathing in order to slow down my heart. It was becoming a distraction. A distraction I cannot afford. Luckily, I need less than a minute to do what I am here to do.

## **ON SCALES AND MASSES**

new physics (local)  $\longrightarrow$  effective operators  $\longrightarrow$  new scale



$$m_{(k)}^2 = m^2 \left( 1 + q^2 - 2q \cos\left(\frac{k\pi}{N+1}\right) \right)$$



plot from Farina et al (2016)



clockwork gears have **unsuppressed** overlap at all sites



plot from Farina, Pappadopulo, Rompineve & Tesi (2016)