

# Back to Weinberg's 3HDM: boundedness from below

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Harmonia VI, Warsaw, September 9-11, 2019

based on: [I. P. Ivanov, F. Faro, PRD100 \(2019\) 035038](#)



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# 3HDM with $U(1) \times U(1)$

The Higgs potential of the  $U(1) \times U(1)$ -symmetric 3HDM:

$$V = V_2 + V_N + V_{CB},$$

where  $V_2 = m_{11}^2(\phi_1^\dagger\phi_1) + m_{22}^2(\phi_2^\dagger\phi_2) + m_{33}^2(\phi_3^\dagger\phi_3)$  and

$$\begin{aligned} V_N &= \frac{\lambda_{11}}{2}(\phi_1^\dagger\phi_1)^2 + \frac{\lambda_{22}}{2}(\phi_2^\dagger\phi_2)^2 + \frac{\lambda_{33}}{2}(\phi_3^\dagger\phi_3)^2 \\ &\quad + \lambda_{12}(\phi_1^\dagger\phi_1)(\phi_2^\dagger\phi_2) + \lambda_{13}(\phi_1^\dagger\phi_1)(\phi_3^\dagger\phi_3) + \lambda_{23}(\phi_2^\dagger\phi_2)(\phi_3^\dagger\phi_3), \\ V_{CB} &= \lambda'_{12}z_{12} + \lambda'_{13}z_{13} + \lambda'_{23}z_{23}. \end{aligned}$$

Here,  $z_{ij} = (\phi_i^\dagger\phi_i)(\phi_j^\dagger\phi_j) - (\phi_i^\dagger\phi_j)(\phi_j^\dagger\phi_i)$ .

"Neutral" directions:  $\phi_1 \propto \phi_2 \propto \phi_3 \rightarrow z_{ij} = 0$ .

Other directions in the Higgs space are called "charge-breaking" directions.

# 3HDM with $\mathbb{Z}_2 \times \mathbb{Z}_2$

Weinberg's 3HDM [PRL37 (1976) 657] with  $\approx 1000$  citations:

$$V = V_2 + V_N + V_{CB} + V_{\mathbb{Z}_2 \times \mathbb{Z}_2},$$

where

$$V_{\mathbb{Z}_2 \times \mathbb{Z}_2} = \bar{\lambda}_{12}(\phi_1^\dagger \phi_2)^2 + \bar{\lambda}_{13}(\phi_1^\dagger \phi_3)^2 + \bar{\lambda}_{23}(\phi_2^\dagger \phi_3)^2 + H.c.$$

with  $\bar{\lambda}_{ij}$  in general complex (explicit CPV + natural flavor conservation).

- When exploring Weinberg's model, it is usually **assumed** that the potential is bounded from below (BFB);
- **Necessary and sufficient BFB conditions** are still not found!
- If we want to explore all pheno situations offered by Weinberg's model, we need to establish the exact BFB conditions.

# 3HDM with $\mathbb{Z}_2 \times \mathbb{Z}_2$

A recent twist:

- Recent years saw a revived interest in 3HDMs, in particular when combining active and inert doublets ( $\rightarrow$  DM candidates), including cases with "dark democracy": [0904.2173, 1012.4680, 1302.3713, 1504.06432, 1608.01673].
- Although  $\mathbb{Z}_2 \times \mathbb{Z}_2$  is not always exact in these models, they are equivalent to Weinberg's model in what concerns the BFB conditions.
- In 2009, [Grzadkowski, Ogreid, Osland, 0904.2173] proposed a set of the exact BFB conditions.
- Very recently, in [Faro, Ivanov, 1907.01963], it was shown that these conditions are **sufficient but not necessary**.
- Sufficient conditions are easy to establish, but they **miss viable, potentially intriguing parameter space regions** offer by the model.

# Towards the BFB conditions in Weinberg's model

We tried to derive the exact BFB conditions in Weinberg's model, but have not succeeded so far. Here is what we do have:

- Exact BFB conditions in  $U(1) \times U(1)$  3HDM [Faro, Ivanov, 1907.01963], which can be a viable model (if unbroken by vevs or with soft breaking). The method is to apply a couple of cute tricks to render the problem treatable by the [copositivity conditions](#).
- A couple of ideas how to attack the  $\mathbb{Z}_2 \times \mathbb{Z}_2$ -symmetric case, not yet finished!
- The [suspicion](#) that it may be [impossible](#) to present these conditions as familiar inequalities on coefs!
- The [hope](#) that this is not the end of the story: a different way to parametrize the model may help!

# Conclusions

- Our main goal, when building bSM models, is to look for **new interesting pheno consequences**. But all these examples must be technically self-consistent.
- A basic task needed for a systematic exploration of Weinberg's model (and similar 3HDMs) has not yet been solved in a satisfactory way. Previous attempts were too simplistic and the difficulty of the problem was not widely recognized.
- Finding the exact BFB conditions for Weinberg's model **remains an intriguing and challenging problem**. It may teach us something useful in general, applicable to other multi-Higgs models.