

# Light Dark Higgs boson in Minimal Sub-GeV Dark Matter Scenarios

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# Sub-GeV Dark Matter

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- ▶ Sub-GeV dark matter has been studied in a variety of scenarios particularly vector portals
- ▶ A sub-GeV dark matter with a mediator of the same mass scale behaves like a WIMP requiring  $\langle\sigma v\rangle \sim 10^{-26}\text{cm}^3\text{s}^{-1}$  for the correct relic density
- ▶ We focus here on vector portal with a dark Higgs mechanism for providing mass to the vector boson
- ▶ Sub-GeV range searches for visible and invisible decays of the mediator/dark Higgs provide the most stringent constraints on such models
- ▶ For long lived scalars, BBN and CMB constraints also play an important role in constraining the parameter space

# Minimal Dark Sector Scenarios

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- ▶ Minimal dark sector refers to extension of SM by adding an extra  $U(1)_D$  gauge group and corresponding fields
- ▶ Minimal dark sector consists of
  - ▶  $U(1)_D$  gauge boson - Dark photon  $V$
  - ▶ Complex scalar breaking  $U(1)_D$  - dark Higgs  $S$
  - ▶ Dark matter candidate charged under  $U(1)_D$  - Complex scalar or pseudo Dirac fermion  $\chi$
- ▶  $U(1)_D$  breaking leads to a massive dark photon, dark Higgs and dark matter particles
- ▶ Dark photon interacts with SM fields through kinetic mixing of  $U(1)_D$  and  $U(1)_Y$
- ▶ Dark Higgs interacts with the SM fields through dark photon, e.g.  $SV_\mu^2$

$$\mathcal{L}_V = -\frac{1}{4}F'^{\mu\nu}F'_{\mu\nu} - \frac{1}{2}\frac{\epsilon}{\cos\theta_w}B_{\mu\nu}F'^{\mu\nu} ,$$

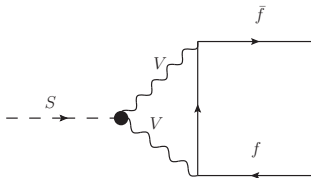
$$\mathcal{L}_S = (D^\mu S)^*(D_\mu S) + \mu_S^2|S|^2 - \frac{\lambda_S}{2}|S|^4 - \frac{\lambda_{SH}}{2}|S|^2|H|^2 ,$$

- ▶ Focus on effect of including dark Higgs in the dark sector in addition to a dark photon
- ▶ Two models of dark matter studied -  $CS$ (complex scalar) and  $pDF$ (pseudo Dirac fermion)

$$V_{pDF} = y_{SL} S \chi_L \chi_L + y_{SR} S \chi_R^c \chi_R^c + \text{h.c.}$$

$$V_{CS} = \lambda_\chi |\chi|^4 + \lambda_{\chi S} |\chi|^2 |S|^2 + \lambda_{\chi H} |\chi|^2 |H|^2$$

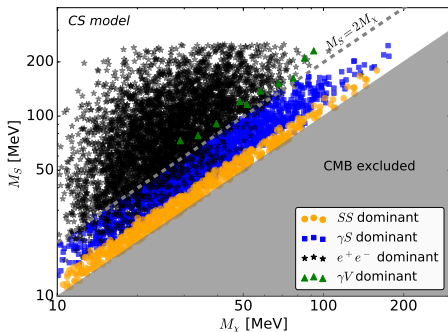
- ▶ Parameters :  $g_V, \epsilon, M_S, M_V, m_\chi, y_{SL} (\lambda_{S\chi})$  and  $y_{SR} (\lambda_{\chi})$
- ▶ Restrict analysis to  $M_S \leq 200$  MeV and  $M_V \leq 500$  MeV and thereby avoid hadronic final states so as to simplify the analysis
- ▶ Only decay mode for dark Higgs is to leptons



## Relic Density and CMB constraints

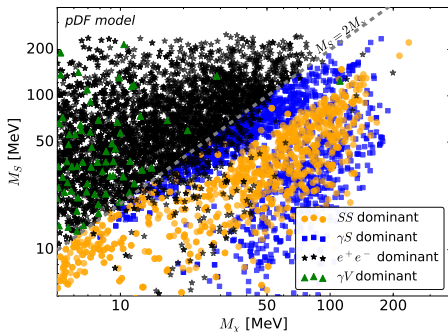
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- ▶ Relic density achieved through the freeze-out mechanism in a two-component dark matter scenario due to presence of long lived dark Higgs
- ▶ Both mediator and DM particles in the MeV mass scale requiring a annihilation cross section of  $10^{-26}\text{cm}^3\text{s}^{-1}$  as in the case of a traditional WIMP
- ▶ Presence of dark Higgs provides additional annihilation channels thus modifying relic density evaluation
- ▶ Injection of charged particles and photons in the inter-galactic medium (IGM) can significantly alter the recombination history of the universe by ionizing and heating the IGM gas
- ▶ Injections from DM annihilation parametrized as
$$p_{ann} = f \langle \sigma v \rangle / M_\chi$$
- ▶ Stringent constraints from  $s$ -wave DM annihilations but much more relaxed for velocity suppressed  $p$ -wave annihilations



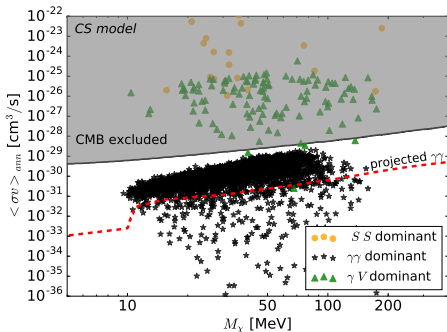
## CS

- ▶  $M_S > 2M_\chi$  :  $e^+e^-$  no  $S$
- ▶  $M_S \simeq M_\chi$  :  $SS$
- ▶  $M_S \sim 2M_\chi$  :  $S\gamma$
- ▶  $M_S < M_\chi$  : excluded



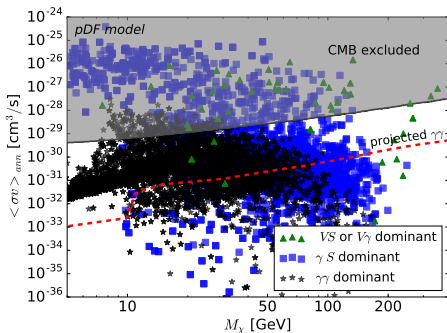
## pDF

- ▶  $M_S < M_\chi$  :  $S\gamma$
- ▶ CMB bounds relaxed due to coannihilation channels



CS

- $M_\chi < M_V$ ,  $M_S$  only  
 $\gamma\gamma$  available,  
suppressed by  $\varepsilon^4$
- $M_\chi > M_S$ ,  $S\gamma$   
dominates, only  
suppressed by  $\varepsilon^2$



pDF

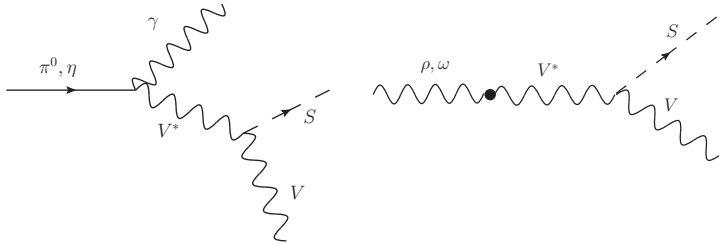
- $M_\chi > M_S$  completely  
excluded, no  $\varepsilon$   
suppression
- $M_\chi > M_V$ , channel  
with  $V$  in final state  
open

# Search for Dark Higgs in Beam Dump Experiments

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- ▶ Fixed target experiments well suited for detection of light dark sector particles
- ▶ High-intensity, but relatively low-energy proton or electron beam impacting target produces “dark matter beam” e.g. dark photon decay
- ▶ Proton beam-dump experiments could be practically seen as light meson factories, with around one neutral pion created for each proton on the target.
- ▶ Electron beam-dump experiments can also lead to dark sector beams through dark photon production by bremsstrahlung
- ▶ Bounds on the kinetic mixing parameter  $\varepsilon$  from electron beam dumps were found to be always significantly weaker than the current missing energy bound from BABAR,  $\varepsilon < 10^{-3}$



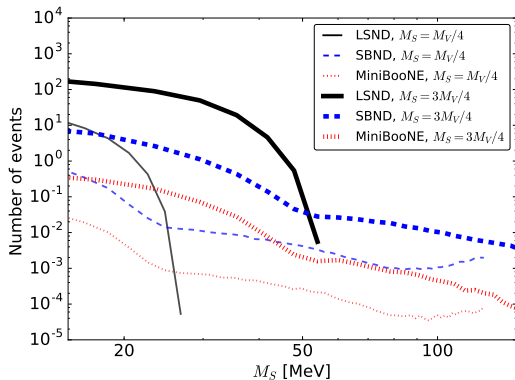


- ▶ Meson produced in proton beam dump decays into dark photon with a small probability
- ▶ Dark Higgs can be produced from an excited dark photon through a “dark” Higgstrahlung mechanism
- ▶ Scalar meson decay

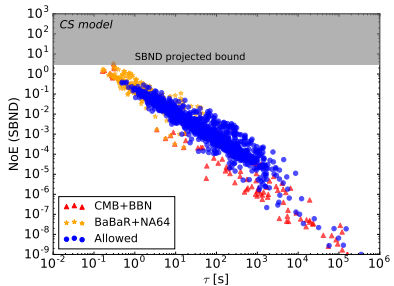
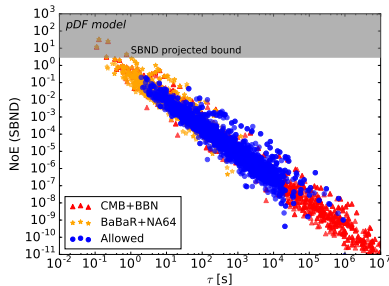
$$\pi^0, \eta \rightarrow \gamma V^*, V^* \rightarrow SV,$$

- ▶ Vector meson decay

$$\rho, \omega \rightarrow V^*, V^* \rightarrow SV.$$

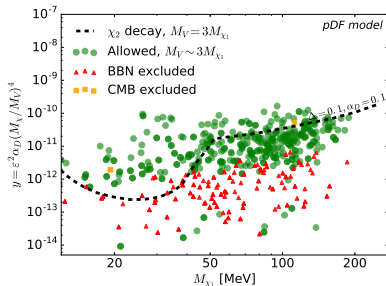
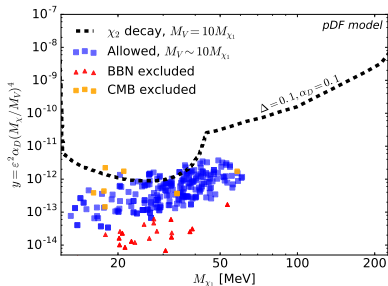


- ▶ We have chosen  $\varepsilon = 0.001$  and  $\alpha_D = \alpha_{\text{em}}$
- ▶ SBND will improve on the miniBooNE bound by one order of magnitude



- ▶ No. of dark Higgs detected goes as  $\varepsilon^6$  and so more sensitive to SBND compared to DM searches
- ▶ Thermal value target out of reach of beam dump experiments

# Probing CSmodel through DM decay



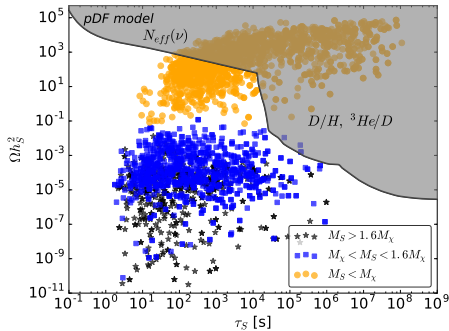
- Search at LSND recast to probe  $\chi_2 \rightarrow \chi_1 e^+ e^-$   
*Izaguirre et al, PRD (2017)*
- Very sensitive to mass splitting,  $\Delta = M_{\chi_2} - M_{\chi_1}$

$$\Gamma_{\chi_2} \propto \alpha_D \epsilon^2 \Delta^5 M_V^{-4}$$

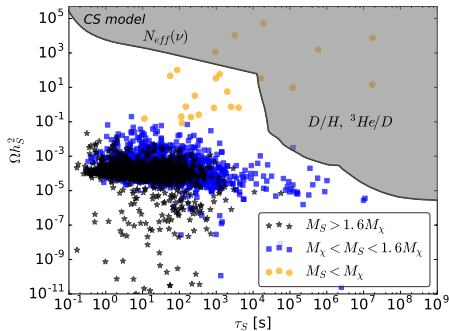
# BBN Constraints

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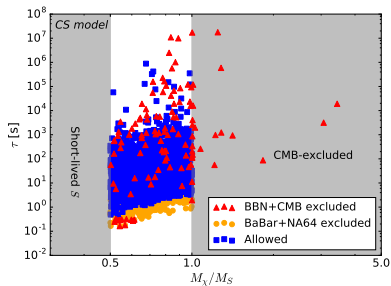
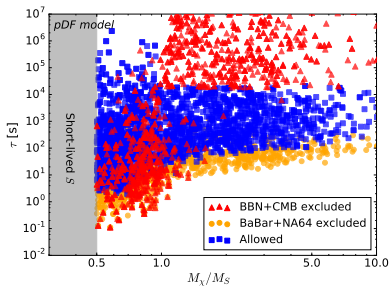
- ▶ BBN bounds arise from either early time ( $t \lesssim 10$  s) injection or late time ( $t > 100$  s) injection of energy from decays of long lived particles like the dark Higgs
- ▶ Early time injection constraints are obtained from  $n/p$  ratio or the effective number of neutrino species  $N_{\text{eff}}$
- ▶ Late time injection constraints arise from primordial abundances of light nuclei ( $^3\text{He}$  and  $\text{D}$ )
- ▶ BBN bounds on Higgs portal models from  $N_{\text{eff}}$  limit dark Higgs lifetime to be smaller than 0.1 s
- ▶ Bounds from primordial abundances affect  $\tau_S > 10^4$  s, and can be quite stringent
- ▶ These bounds are mitigated due to lepton-only decay modes of dark Higgs and efficient annihilation mechanism



- ▶ When  $M_S > M_\chi$   
 $SS \rightarrow \chi\chi$  decreases abundance
- ▶ For  $M_S < M_\chi$ ,  $N_{\text{eff}}$  bound rules out part of parameter space above 100 s
- ▶ Large part of parameter space ruled out above  $10^4$  s



# Summary of Constraints



## Conclusions

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- ▶ We have argued that models with a massive, but light, dark vector mediator, the spectrum should naturally contain a light dark Higgs, which can substantially modify predictions
- ▶ Cosmological constraints are considerably modified due to the presence of dark Higgs and additional modes of annihilation for DM
- ▶ We particularly study a long lived, light dark Higgs boson
- ▶ The low abundance region,  $M_\chi < M_S < 2M_\chi$ , where bounds from BBN are weakened due to presence of dark Higgs while relic density is realized mainly from  $\chi\chi \rightarrow S\gamma$  leading to a lower bound on  $\varepsilon$
- ▶ The high abundance region,  $M_S < M_\chi$ , where there are relatively strong bounds from BBN-related observables and for CSmodel, this region is ruled out by CMB constraints