

Does Quantum Gravity constrain Large-Field Inflation?

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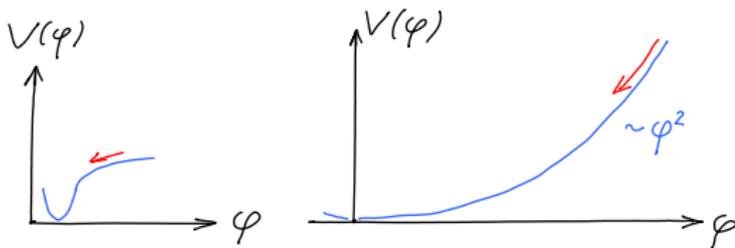
(including work with P. Mangat, F. Rompineve, L. Witkowski)

Outline

- The interest in **large-field models** of inflation
- Fundamental obstructions to large-field inflation
- Problems with large-field inflation in string theory
- Axion alignment and Axion monodromy:
Early models and recent progress

Small vs. Large Field Models

- The inflaton φ roll some distance $\Delta\varphi$ in a potential $V(\varphi)$



- Small field: If $V(\varphi)$ has some **very flat** region, we can have $\Delta\varphi \ll 1$ ($M_P \equiv 1$ here and below)
 - Large field: To get inflation in '**generic**' potentials (e.g. $V(\varphi) \sim \varphi^2$), we must have $\Delta\varphi \gg 1$
- (We will see that this may be a problem in quantum gravity)

I will now focus on large-field models for two reasons....

1) Observations

- The amount of primordial gravity waves is measured by the tensor-to-scalar ratio:

$$r = \frac{\Delta_T^2}{\Delta_R^2} \simeq 8 \left| \frac{d\varphi}{dN} \right|^2 \Rightarrow \Delta\varphi \simeq 20\sqrt{r}$$

- Thus, even though the BICEP 'discovery' of $r \simeq 0.15$ went away, the need to consider large-field models may return
- Note: The Planck/BICEP analysis still sees a ($\sim 1.8\sigma$) hint for $r \simeq 0.05$
- Much better values/bounds are expected soon

...reasons for interest in large-field models...

2) Fundamental

- On the one hand, large-field models are more 'robust'
- On the other hand, there are generic arguments against large-field models in consistent quantum gravity theories

see e.g. Arkani-Hamed/Motl/Nicolis/Vafa '06 Conlon '12

.....

Kaloper/Kleban/Lawrence/Sloth '15

- This goes hand in hand with certain problems in constructing large-field models in (the known part of) the string theory landscape

'Fundamental reasoning' continued...

- **However**, triggered by BICEP, new promising classes of stringy large-field have been constructed (e.g. F -term axion monodromy)

Kim/Nilles/Peloso '07

McAllister, Silverstein, Westphal '08

.....

Marchesano/Shiu/Uranga '14

Blumenhagen/Plauschinn '14

AH/Kraus/Witkowski '14

- At the same time, there are ongoing efforts to sharpen the 'no-go arguments' as well as to refute them

Rudelius '14...'15

Montero, Uranga, Valenzuela '15

Brown, Cottrell, Shiu, Soler '15

AH/Mangat/Rompineve/Witkowski '15

...

- I will try to explain some aspects of this debate....

No-go argument I: (Gravitational) instantons

- One of the leading inflaton candidates is a shift-symmetric, periodic scalar (axion)

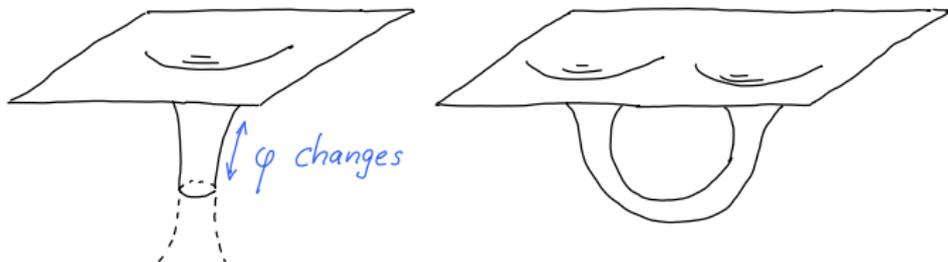
Freese/Frieman/Olinto '90

Kawasaki/Yamaguchi/Yanagida '00

- In Euclidean Einstein gravity, supplemented with an axionic scalar φ ($\varphi \equiv \varphi + f$), instantonic solutions exist:

Giddings/Strominger '88

...



- The 'throat' is supported by the kinetic energy of φ , hence the large field range is essential

Caveats:

- a) Euclidean quantum gravity has its own fundamental problems
- b) It is not completely clear 'where the throat should connect' (our world, another world, 'crunch', 'baby universe')
- Hence the interpretation of these instanton solutions still has issues...

Gravitational instantons (continued)

- Their Euclidean action is

$$S \sim n/f \quad (\text{with } n \text{ the instanton number})$$

- Their maximal curvature scale is f/n , which should not exceed the UV cutoff:

$$f/n < \Lambda$$

- This fixes the lowest n that we can trust and hence the minimal size of the instanton correction to the potential $V(\varphi)$:

$$\delta V \sim e^{-S} \sim e^{-n/f} \sim e^{-1/\Lambda}$$

Gravitational instantons (continued)

- For gravitational instantons **not** to prevent inflation, the **relative** correction must remain small:

$$\frac{\delta V}{V} \sim \frac{e^{-1/\Lambda}}{H^2} \ll 1$$

- For a Planck-scale cutoff, $\Lambda \sim 1$, this is never possible
- However, the UV cutoff can in principle be as low as H
- Then, if also $H \ll 1$, everything might be fine....

$$\frac{\delta V}{V} \sim \frac{e^{-1/H}}{H^2}$$

Gravitational instantons (continued)

- Now, most string models of inflation do indeed have a low cutoff (e.g. compactification scale)
- **However**, it may be too naive to assume that 'uncalculable' gravitational instantons can simply be ignored
- They may find their 'continuation' in the gauge or D-brane instantons of the concrete string model
- The closely related issue of (singular) 'core instantons' has been brought up
Heidenreich, Reece, Rudelius '15
- UV completion and moduli stabilization are crucial open issues
...our own ongoing work

No-go argument II: Weak gravity conjecture

Arkani-Hamed/Motl/Nicolis/Vafa '06

- Roughly speaking: 'Gravity is always the weakest force.'
- More concretely (mild form):
For any U(1) gauge theory there exists a charged particle with

$$q/m > 1.$$

- Strong form:
The above relation holds for the lightest charged particle.

Weak gravity conjecture (continued)

- One supporting argument:

Quantum gravity forbids **global symmetries**. We should not be able to take the limit of small gauge couplings.

The WGC quantifies this on the basis of stringy examples.

- Another supporting argument:

In the absence of **sufficiently light**, charged particles, extremal BHs are stable. Such **remnants** are believed to cause inconsistencies.

see e.g. Susskind '95

The boundary of stability of extremal black holes is precisely $q/m = 1$ for the decay products

Generalizations of the weak gravity conjecture

- The basic lagrangian underlying the above is

$$S \sim \int (F_2)^2 + m \int_{1-dim.} d\ell + q \int_{1-dim.} A_1$$

- This generalizes to charged **strings, domain walls etc.** Crucially, the degree of the corresponding form-field (gauge-field) changes:

$$S \sim \int (F_{p+1})^2 + m \int_{p-dim.} dV + q \int_{p-dim.} A_p$$

with

$$F_{p+1} = dA_p$$

Generalizations to instantons

- One can also **lower** the dimension of the charged object, making it a point in space-time:

$$S \sim \int (d\varphi)^2 + m + q\varphi(x_{inst.})$$

- One easily recognizes that this is just a more general way of talking about instantons and axions:

$$m \Leftrightarrow S_{inst.} \quad , \quad q\varphi(x_{inst.}) \Leftrightarrow \frac{1}{f} \int \varphi F\tilde{F}$$

WGC for instantons and inflation

- The consequences for inflation are easy to derive
- First, recall that the instantons induce a potential

$$V(\varphi) \sim e^{-m} \cos(\varphi/f)$$

- Since, for instantons, $q \equiv 1/f$, we have

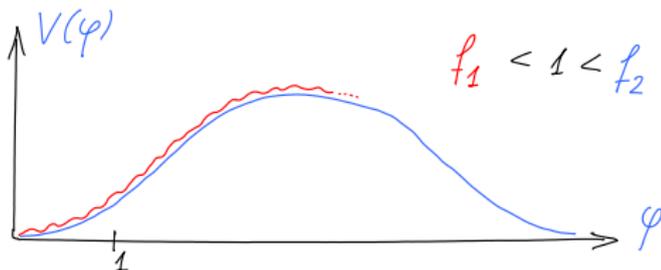
$$q/m > 1 \quad \Rightarrow \quad mf < 1$$

- Theoretical control (dilute instanton gas) requires $m > 1$
- This implies $f < 1$ and hence
large-field 'natural' inflation is in trouble

A Loophole

Rudelius '15

- Suppose that **only the mild form** of the WGC holds
- In this case, we can have one **'sub-planckian'** instanton maintaining the WGC, together with a lighter **'super-planckian'** instanton realizing inflation:

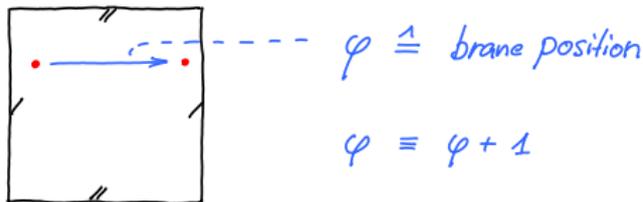


For other arguments and loopholes see e.g.
de la Fuente, Saraswat, Sundrum '14
Bachlechner, Long, McAllister '15

What do explicit string constructions have to say about $\Delta\varphi \gg 1$?

- The problem is that (more or less) all 4d fields φ (moduli) have a small field range.
- An obvious example arises if φ is a brane position. Clearly, this field is periodic and the field space is hence limited:

Dvali/Tye '98



- Note: Thus, we naturally get the axionic scalars discussed earlier. But their periodicity is always too short.

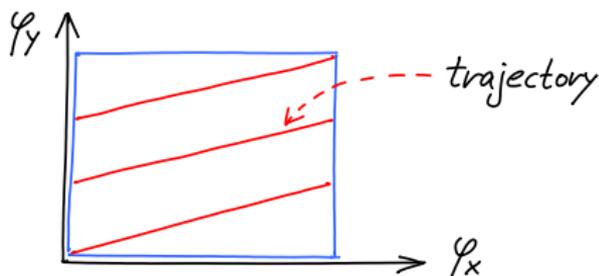
Banks, Dine, Fox, Gorbatov '03

One needs ideas!

(I) Winding inflation / KNP

Kim/Nilles/Peloso '04; Berg/Pajer/Sjors '09; Ben-Dayan/Pedro/Westphal '14

- One such idea is to realize a 'winding' trajectory on a 2d periodic field space:



- Clearly, such a trajectory can be much longer than the (naive) field range
- The technical challenge is the realization of the required potential in concrete string models

Winding inflation (continued)

- The fields φ_x and φ_y are two 'string theory axions', both with $f < 1$ (obeying the WGC)
- They are also moduli. Hence, fluxes (e.g. $\langle F_3 \rangle \neq 0$ on the compact space) can be used to stabilize them
- A judicious choice of fluxes allows for stabilizing just one linear combination, forcing the remaining light field on the winding trajectory:

$$V \supset (\varphi_x - N\varphi_y)^2 + e^{-M} \cos(\varphi_x/f) + e^{-m} \cos(\varphi_y/F)$$

with $N \gg 1$

- This realizes inflation and avoids the WGC!

Winding inflation (continued)

- To be more precise, let's change variables:

$$\varphi \equiv \varphi_x, \quad \psi \equiv \varphi_x - N\varphi_y$$

- While ψ is 'frozen', our inflaton φ 'sees' both the instantons belonging to φ_x as well as those belonging to φ_y :

$$V \supset \psi^2 + e^{-M} \cos(\varphi/f) + e^{-m} \cos[(\varphi - \psi)/NF]$$

- Crucially, in our proposal the quantities M and m are precisely the type of variables that can be tuned in the landscape (like the vacuum energy)

....thus, getting a largish M is not a problem

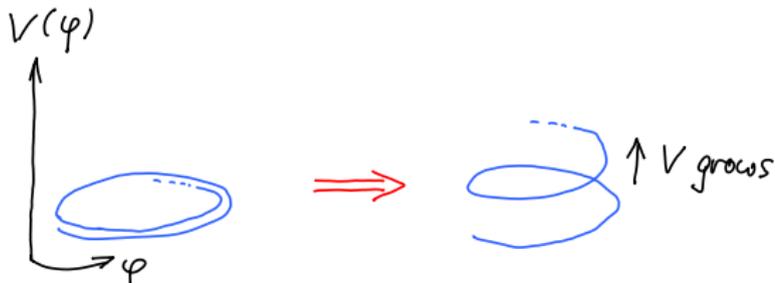
- For further issues and recent developments see e.g. ...

Buchmüller, Dudas, Heurtier, Westphal, Wieck, Winkler '15
Kobayashi, Oikawa, Otsuka '15

(II) Monodromy inflation

Silverstein/Westphal/McAllister '08

- We start with a single, periodic inflaton φ
- The periodicity is then **weakly** broken by the scalar potential



- The status of constraints from the Weak Gravity Conjecture is less clear. See however...

Ibanez, Montero, Uranga, Valenzuela '15

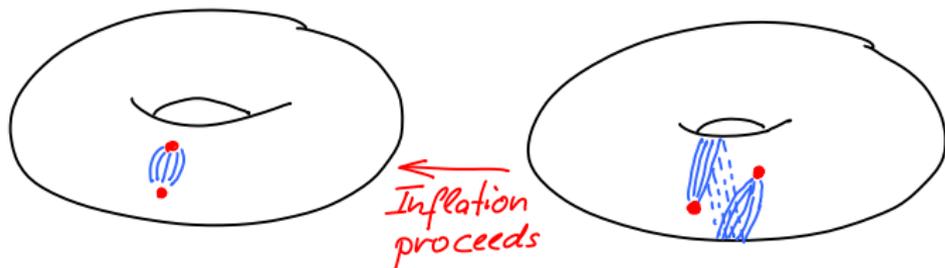
F-term axion monodromy

- First suggestions have emerged how this could be realized in a quantitatively controlled way

(i.e. in a 4d supergravity description, with a stabilized compact space)

Marchesano/Shiu/Uranga '14
Blumenhagen/Plauschinn '14
AH/Kraus/Witkowski '14

- In particular, in our suggestion inflation corresponds to **brane-motion**
- The monodromy arises from a flux sourced by the brane



Summary/Conclusions

- Quantum gravity (Instantons / Weak gravity conjecture) may be constraining large-field inflation at a very fundamental level
- Concrete problems with large-field inflation in string theory reflect these fundamental 'issues'
- Progress is being made both in understanding the generic constraints as well as in constructing counterexamples (i.e. models)

In primordial gravity waves / large-field inflation, fundamental quantum gravity problems may meet reality!

- Some of this discussion may also be relevant for relaxions