

Harmonia meeting

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R-symmetry

- additional symmetry of the SUSY algebra allowed by the Haag - Łopuszański - Sohnius theorem
- for N=1 it is a global $U_R(1)$ symmetry under which the SUSY generators are charged
- implies that the spinorial coordinates are also charged $Q_R(\theta) = 1, \theta \rightarrow e^{i\alpha}\theta$
- Superpotential example

$$\mathcal{L} \ni \int d^2\theta W$$

- Superpotential is polynomial in fields. For W to transform homogeneously superfields must have definite R-charges

$$e^{i\alpha Q_R} \Phi = e^{i\alpha Q_R} \phi(y) + \sqrt{2}\theta\psi(y) + \theta\theta F(y)$$

- Similarly one can work out other parts of the Lagrangian

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(we want it to be)
R-invariant $\longrightarrow$ $\mathcal{L} \ni \int d^2\theta \quad W$

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(we want it to be)
R-invariant

$\longrightarrow$
 $\mathcal{L} \ni \int d^2\theta \quad W$

transforms as $e^{-2i\alpha}$

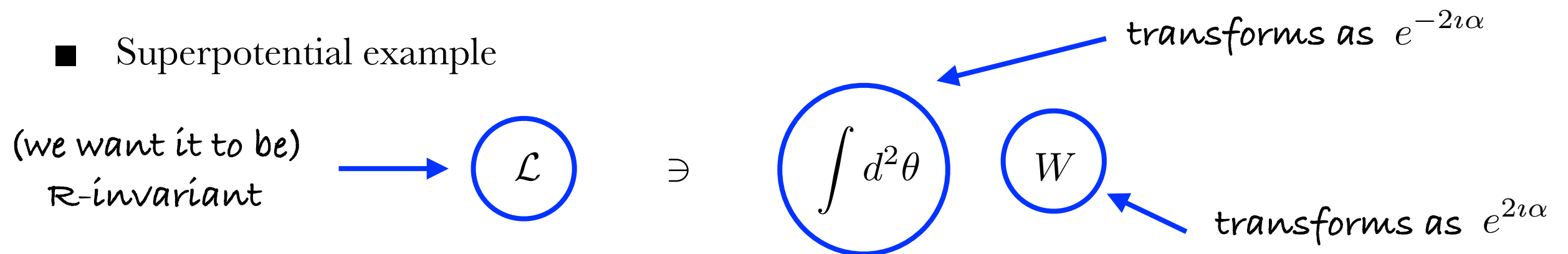
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$$\begin{array}{ccccccc}
 e^{i\alpha Q_R} & & e^{i\alpha Q_R} & & e^{i\alpha(Q_R-1)} & & \\
 \Phi & = & \phi(y) & + & \sqrt{2}\theta\psi(y) & + & \theta\theta F(y)
 \end{array}$$

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- Similarly one can work out other parts of the Lagrangian

Low-energy R-symmetry realization

- Different possible models that one can construct

- “Natural” choice

$$e^{i\alpha Q_R} \Phi = e^{i\alpha Q_R} \phi(y) + e^{i\alpha(Q_R-1)} \sqrt{2\theta} \psi(y) + \theta F(y)$$

leptons and quarks	$Q_R = 1$	$Q_R = 1$	$Q_R = 0$
Higgs	$Q_R = 0$	$Q_R = 0$	$Q_R = -1$

- Good: no barion and lepton number violating terms
- Bad: No Majorana masses for higgsinos and gauginos

One way to fix it: [Dirac masses](#)

Minimal R-Symmetric Supersymmetric Standardmodel (MRSSM)

Kribs et.al. arXiv:0712.2039

		$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	$U(1)_R$
Additional fields:	Singlet $\hat{S}$	1	1	0	0
	Triplet $\hat{T}$	1	3	0	0
	Octet $\hat{O}$	8	1	0	0
	R-Higgses $\hat{R}_u$	1	2	-1/2	2
	$\hat{R}_d$	1	2	1/2	2

$$W = \mu_d \hat{R}_d \hat{H}_d + \mu_u \hat{R}_u \hat{H}_u$$

$$+ \Lambda_d \hat{R}_d \hat{T} \hat{H}_d + \Lambda_u \hat{R}_u \hat{T} \hat{H}_u + \lambda_d \hat{S} \hat{R}_d \hat{H}_d + \lambda_u \hat{S} \hat{R}_u \hat{H}_u$$

$$- Y_d \hat{d} \hat{q} \hat{H}_d - Y_e \hat{e} \hat{l} \hat{H}_d + Y_u \hat{u} \hat{q} \hat{H}_u$$

MSSM vs. MRSSM

■ superpotencial

$$\mu \hat{H}_u \hat{H}_d \quad \text{!}$$

$$-Y_d \hat{d} \hat{q} \hat{H}_d - Y_e \hat{e} \hat{l} \hat{H}_d + Y_u \hat{u} \hat{q} \hat{H}_u \quad \text{✓}$$

■ soft-SUSY breaking terms

$$\square B_\mu \text{-term} \quad \text{✓}$$

$$\square \text{soft scalar masses} \quad \text{✓}$$

$$\square \text{Majorana gaugino masses} \quad \text{!}$$

$$\square \text{A - terms} \quad \text{!}$$

■ superpotencial

$$\mu_d \hat{R}_d \hat{H}_d + \mu_u \hat{R}_u \hat{H}_u$$

$$-Y_d \hat{d} \hat{q} \hat{H}_d - Y_e \hat{e} \hat{l} \hat{H}_d + Y_u \hat{u} \hat{q} \hat{H}_u$$

$$\Lambda_d \hat{R}_d \hat{T} \hat{H}_d + \Lambda_u \hat{R}_u \hat{T} \hat{H}_u + \lambda_d \hat{S} \hat{R}_d \hat{H}_d + \lambda_u \hat{S} \hat{R}_u \hat{H}_u$$

■ soft-SUSY breaking terms

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$$\square \text{Dirac gaugino masses}$$

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Kribs, Popitz, Weiter (2008)

MSSM vs. MRSSM

■ superpotencial

$$\mu \hat{H}_u \hat{H}_d$$

$$-Y_d \hat{d} \hat{q} \hat{H}_d - Y_e \hat{e} \hat{l} \hat{H}_d + Y_u \hat{u} \hat{q} \hat{H}_u$$



■ soft-SUSY breaking terms

□ B_μ - term



□ soft scalar masses



□ Majorana gaugino masses



□ A - terms



■ superpotencial

$$\mu_d \hat{R}_d \hat{H}_d + \mu_u \hat{R}_u \hat{H}_u$$

$$-Y_d \hat{d} \hat{q} \hat{H}_d - Y_e \hat{e} \hat{l} \hat{H}_d + Y_u \hat{u} \hat{q} \hat{H}_u$$

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■ soft-SUSY breaking terms

□ B_μ -term

□ soft scalar masses

□ Dirac gaugino masses

□ no A-terms

One way to fix it: [Dirac masses](#)

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		$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	$U(1)_R$	
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Kribs, Popitz, Weiter (2008)

Particle content summary: MSSM vs. MRSSM

different number of physical states completely new states

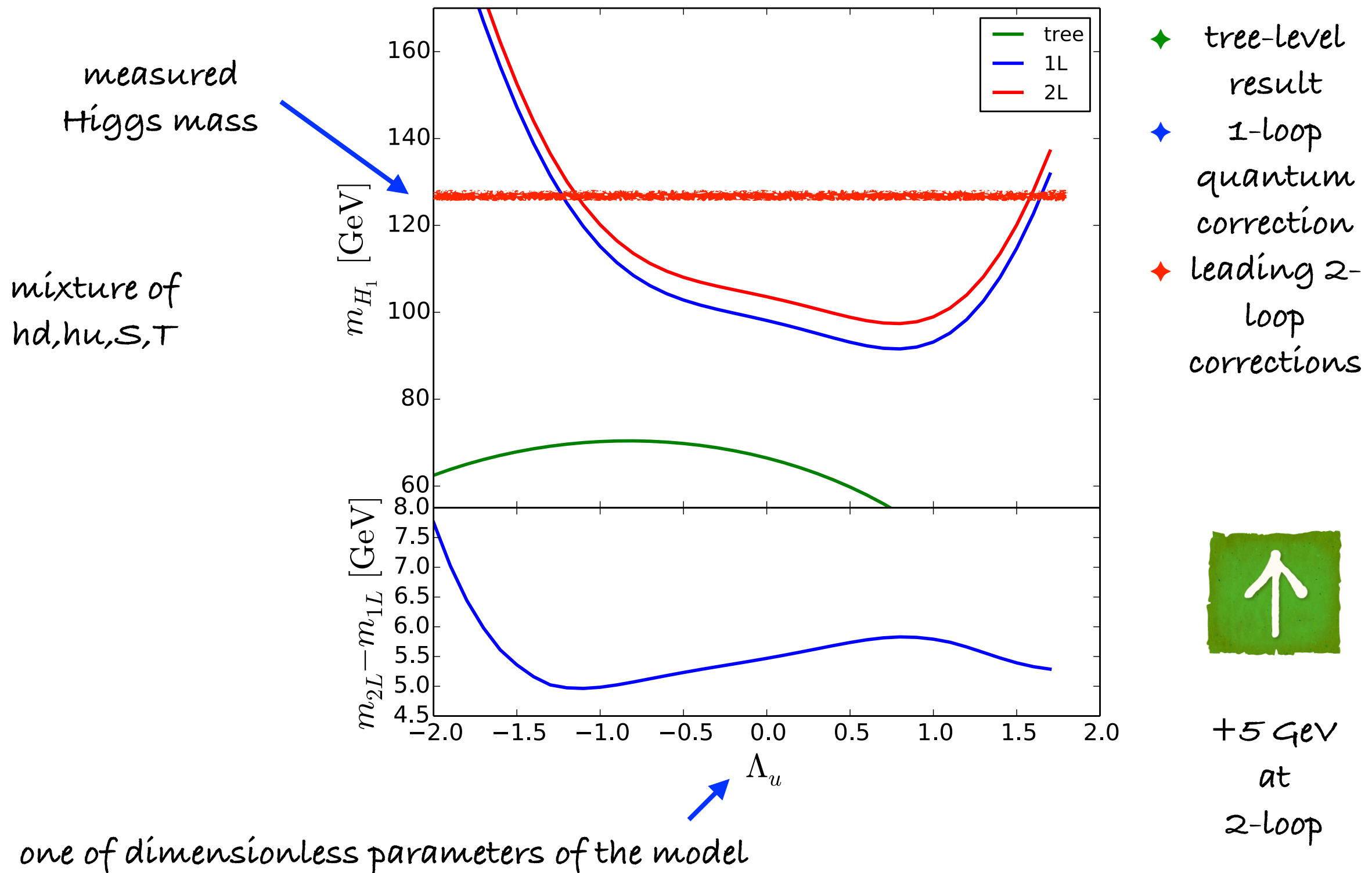
	Higgs			charginos	R-Higgs		sgluon
	CP-even	CP-odd	charged		neutral	charged	
MSSM	2	1	1	2	0	0	0
MRSSM	4	3	3	2+2	2	2	1

	neutralino	gluino
MSSM	4	1
MRSSM	4	1

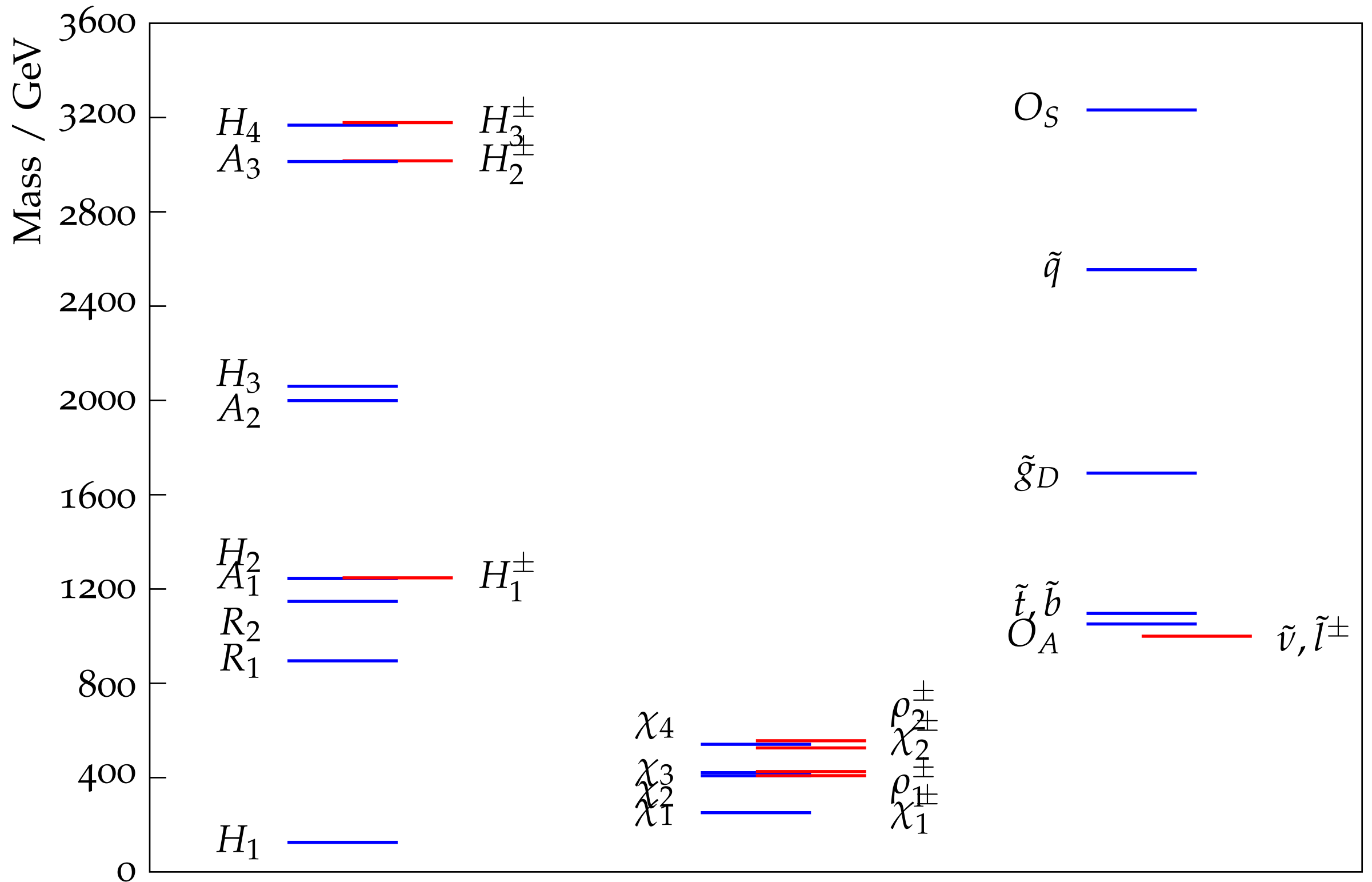
Majorana fermions

Dirac fermions

Lightest Higgs boson mass in the MRSSM



Exemplary mass spectrum



What has been done

- Higgs masses at 1- and 2-loop order
- analysis of constraints from m_W and STU parameters
- light singlet scenario
- (1 component) dark matter
- analysis of pseudoscalar sgluon production at the LHC

What is to be done

- NLO SQCD corrections to pair production
 - NLO R-symmetric SQCD corrections to squark pair production
 - pheno analysis
 - remaining processes
- 2-component dark matter
- $g-2$

MRSSM signatures at the LHC

■ sgluon pair production $pp \rightarrow OO$

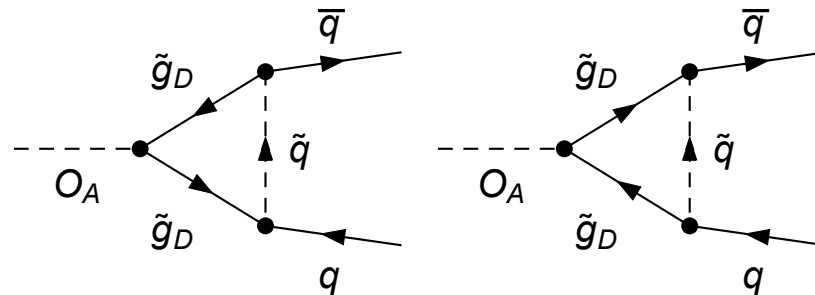
- complex fields O split by D-term contribution into scalar (S) and pseudoscalar (A) parts

$$m_{O_A}^2 = m_O^2$$

$$m_{O_S}^2 = m_O^2 + 4(M_O^D)^2$$

- O_S naturally heavy, O_A might decay into quarks through loop-induced coupling

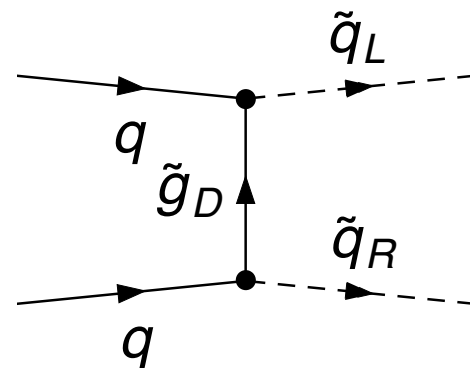
□



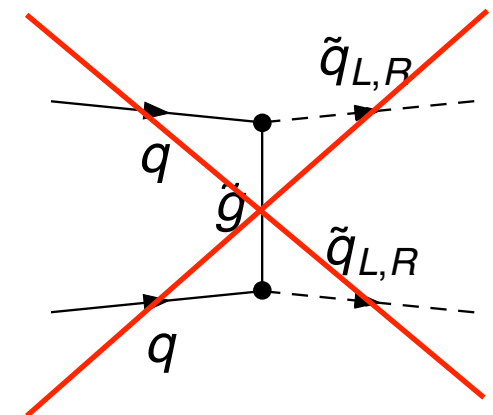
- Coupling proportional to m_q . If O_A is lighter than other SUSY particles but $m_{O_A} > 2m_t$ O_A decays exclusively to top quarks

■ same sign squark pair production $pp \rightarrow \tilde{q}_L \tilde{q}_R$

- diagram present in the MRSSM

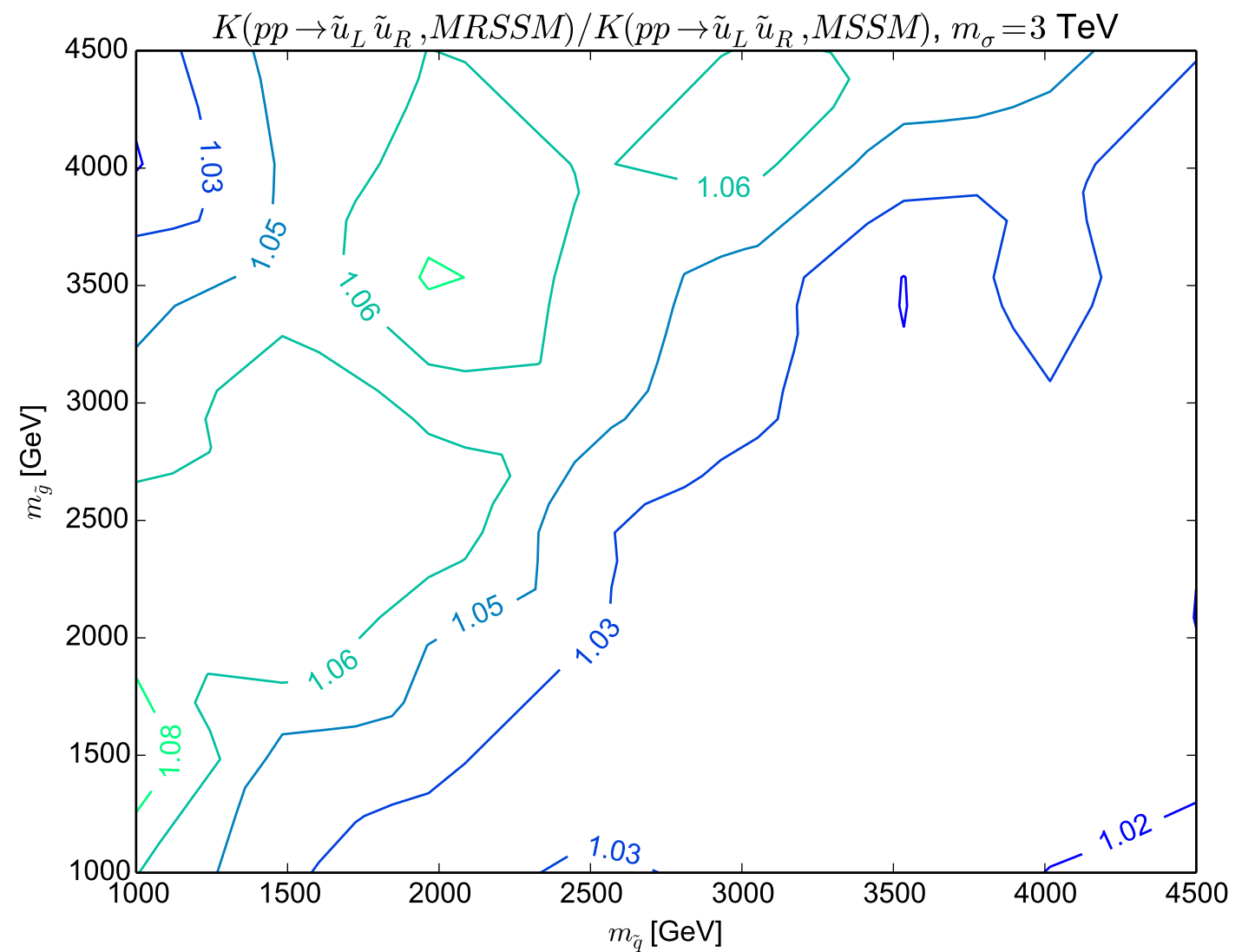
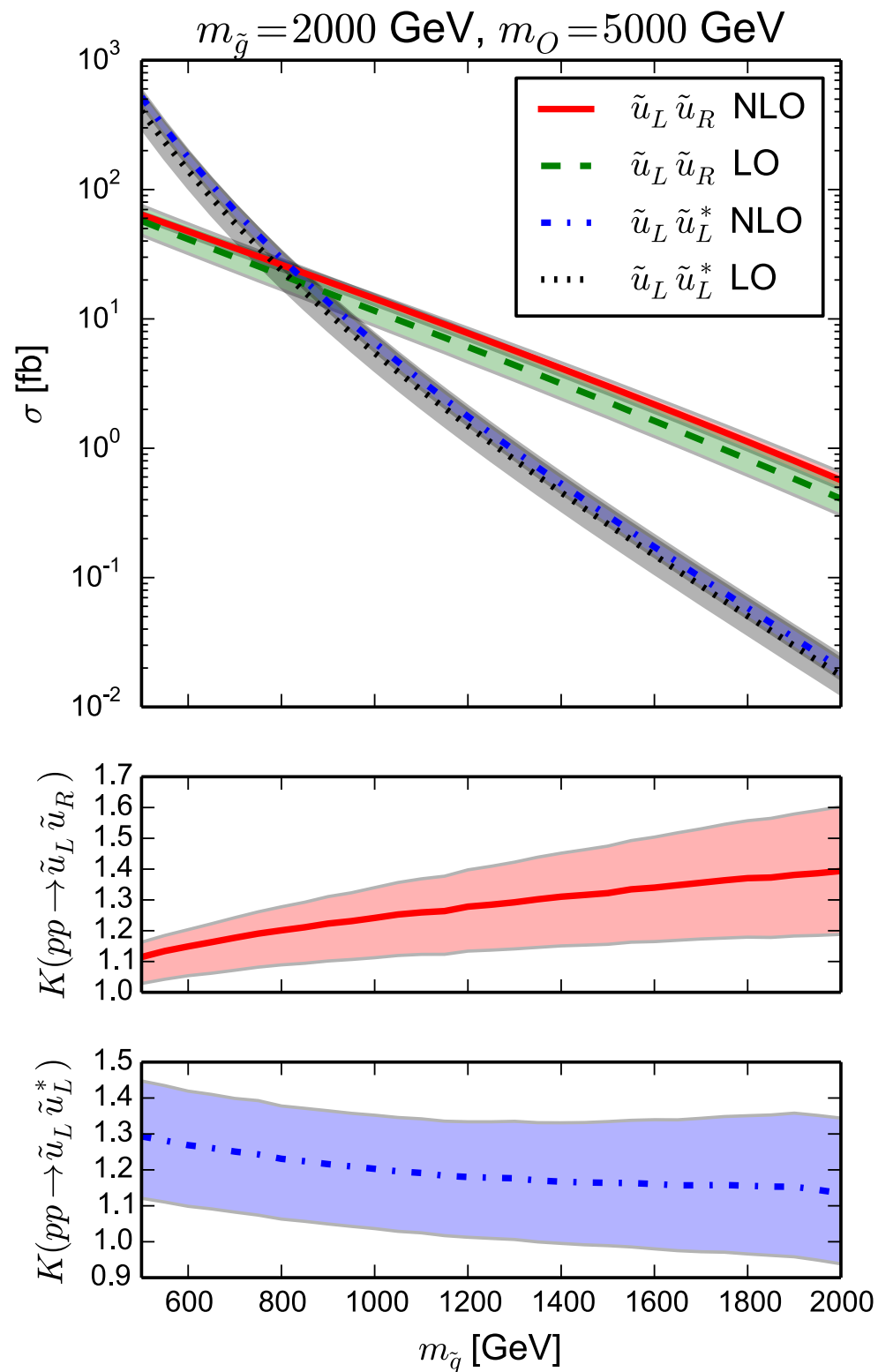


MSSM diagram with Majorana mass insertion doesn't exist



■ Dirac gluino $\tilde{g}_D$ pair production, with cross section roughly twice as large as in the MSSM

NLO R-symmetric SQCD corrections

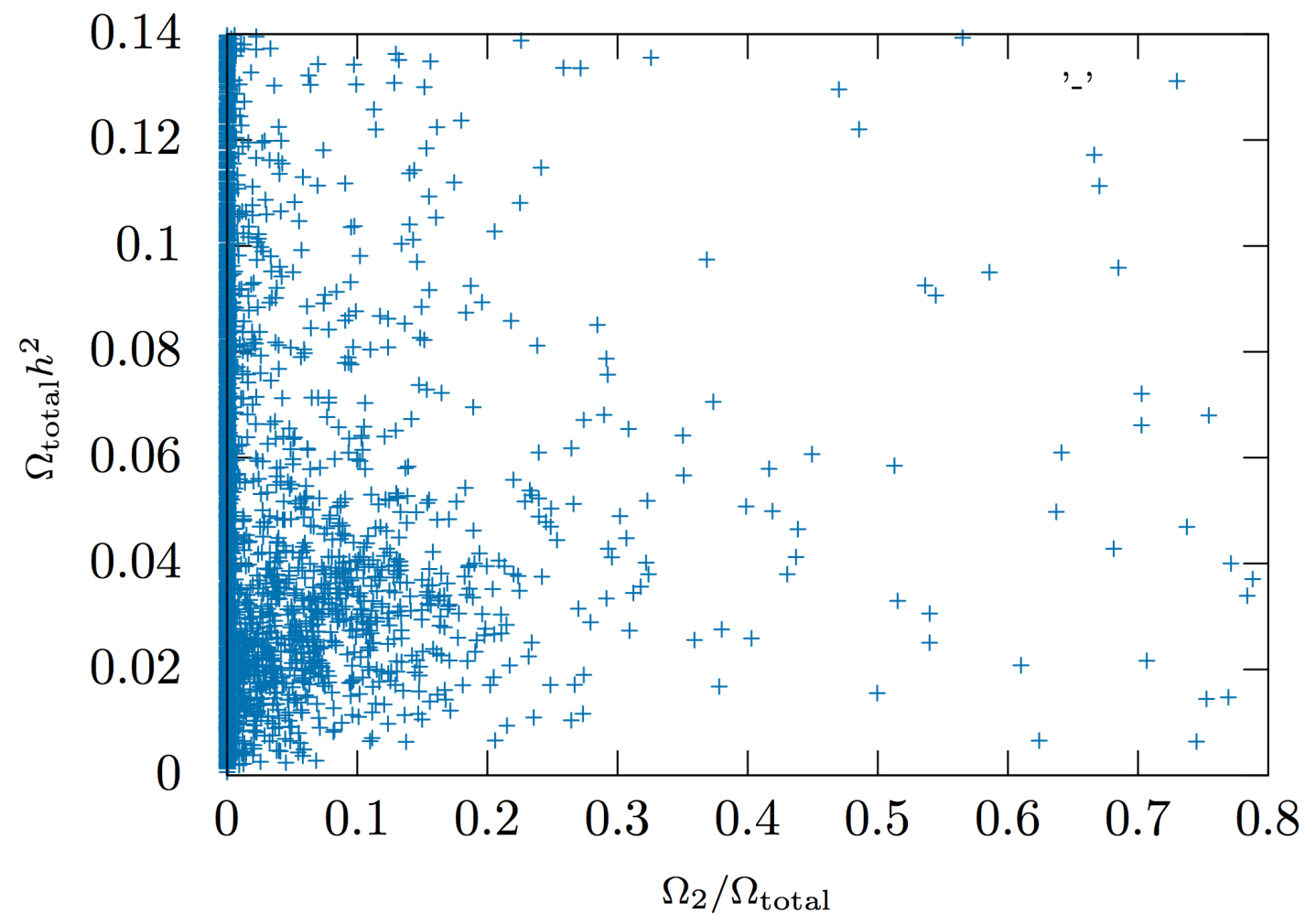


2 component dark matter

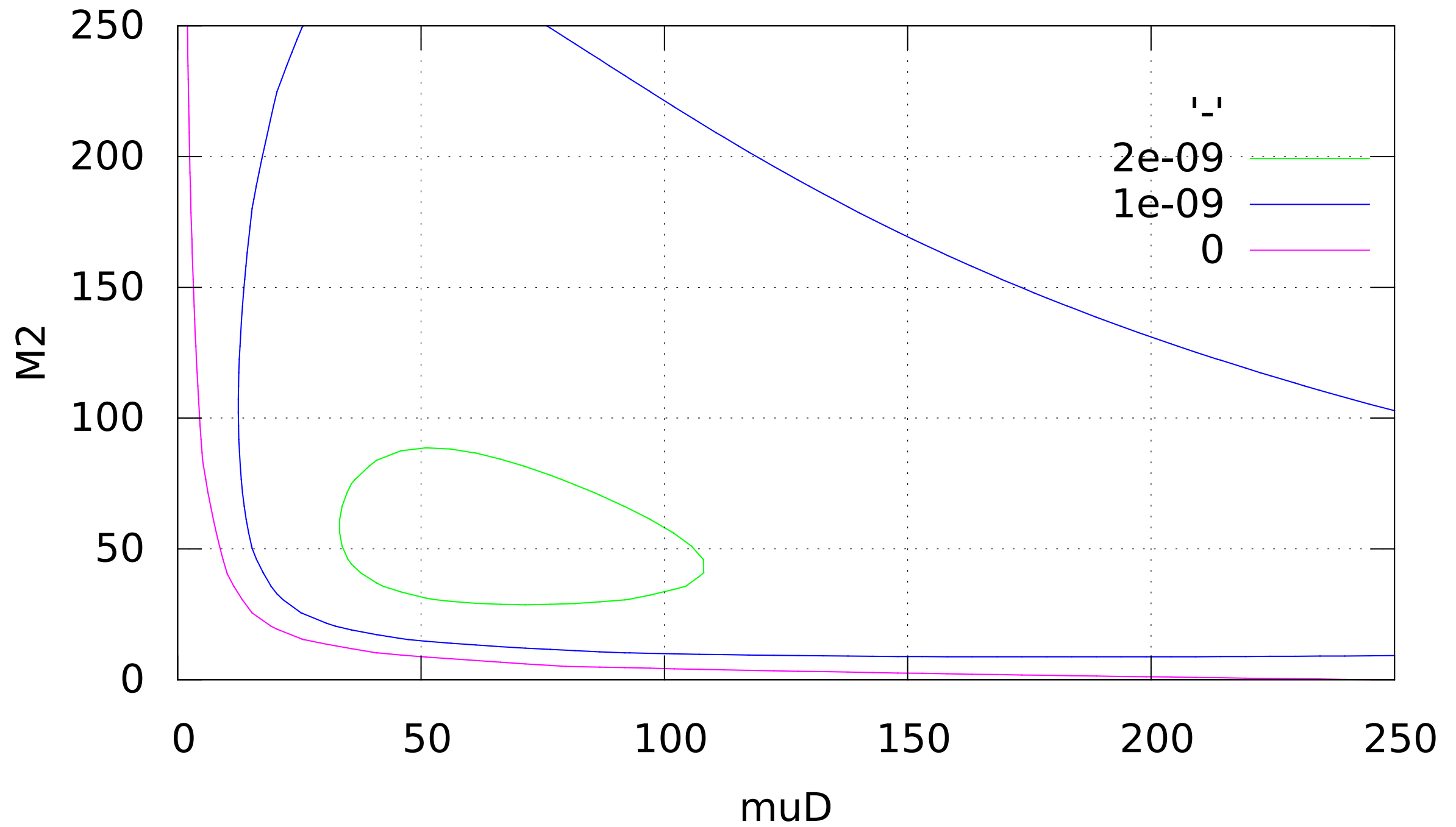
- consider scenarios where the lightest particle with $R=1$ is neutralino or sneutrino with mass $m_{\text{LSP}1}$
- if $m_{R_1^0} < 2m_{\text{LSP}1}$, lightest neutral R -Higgs is also stable
- two SUSY dark matter candidates with relic densities Ω_1 and Ω_2
- requirements
 - $\Omega_{\text{total}}h^2 \equiv (\Omega_1 + \Omega_2)h^2 \approx 0.11$
 - substantial fraction $\Omega_2/\Omega_{\text{total}}$
- (for now) best points are not collinear friendly:

$$m_{\tilde{\chi}_1^0} = 367 \text{ GeV}$$

$$m_{R_1^0} = 571 \text{ GeV}$$



$g - 2$

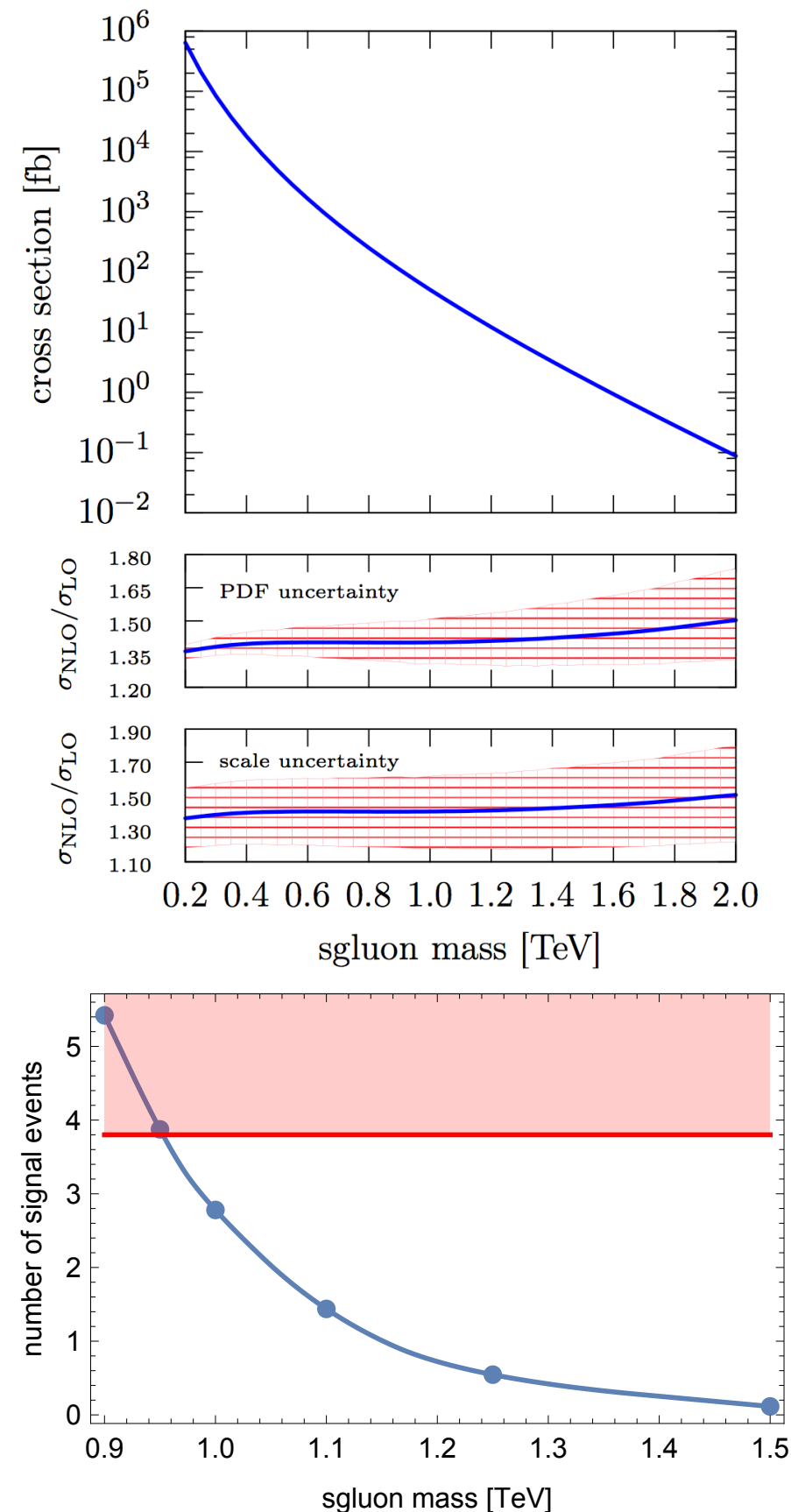


Backup

Sgluon pair production at 13 TeV LHC

W. Kotlarski [arXiv:1608.00915]

- Analysis of the sgluon pair production with subsequent decay into $t\bar{t}$ pairs. Recasting ATLAS search in the same-sign lepton channel using 3.2/fb of integrated luminosity
- Signal simulated at NLO using **MadGraph5_aMC@NLO** + **FeynRules** + **NLOCT** and matched to parton shower in the **MC@NLO** scheme
- Detector response parametrized using **Delphes3**
- Analysis validated on background processes $t\bar{t}l^+l^-$, $t\bar{t}l^\pm\nu$
- Mass of pair produced real sgluons decaying with $\text{BR}(O \rightarrow t\bar{t}) = 1$ excluded up to 950 GeV



Leading order analysis

LO cross-sections for sparticle production at the LHC at $\sqrt{s}=13\text{TeV}$

