EMERGING JETS

DANIEL STOLARSKI WITH PEDRO SCHWALLER AND ANDREAS WEILER

arXiv:141?.????

Warsaw University November 24, 2014

MOTIVATION 1

Getting away from the lamp post

MOTIVATION 2

We have seen dark matter in the sky

But not in the label and Plots

ASYMMETRIC DARK MATTER

$\Omega_{DM} \simeq 5\Omega_B$

 $\Omega_B = m_p n_B$ Controlled by complicated (known) QCD dynamics Unknown dynamics of baryogenesis $\Omega_{DM} = m_{DM} n_{DM}$ **?**

MANY PAPERS

S. Nussinov, Phys.Lett.B.165 (1985) 55.

D. B. Kaplan, Phys.Rev.Lett.B.68 (1992) 741-3.

D. E. Kaplan, M. A. Luty, K. M. Zurek, Phys.Rev.D **79** 115016 (2009) [arXiv:0901.4117 [hep-ph]].

K. K. Boddy, J. L. Feng, M. Kaplinghat, and T. M. P. Tait, Phys. Rev. D. **89** 11, 115017 (2014) [arXiv:1402.3629 [hep-ph]].

For a review see K. Petraki and R. R. Volkas, Int.J.Mod.Phys.A 28, 1330028 (2013) [arXiv:1305.4939 [hep-ph]].

...

ASYMMETRIC DARK MATTER

$\Omega_{DM} \simeq 5\Omega_B$

 \mathbf{Can} get $n_{DM} \sim n_B$, usually have to assume $m_{DM} \sim m_B$

Can we get both?

GETTING THE MASS $\Omega_{DM} \simeq 5\Omega_B$

DARK QCD

Propose new $SU(N_d)$ "dark QCD," dark quarks **Bai, Schwaller, PRD 13.**

Dark matter is dark sector baryons with mass ~ Λ_dQCD

Massive bifundamental fields decouple at mass $M \gg \Lambda_\mathrm{dQCD}$ ecouple at il

Search for model with perturbative fixed point search for model with perturbative fixed *dg dt* $\mathbf{\hat{g}} = g^*$!"#\$%&'(")*+&

$$
\frac{dg}{dt} = \beta(g) = 0 \text{ for } g = g^*
$$

$$
\begin{array}{c}\n\swarrow \\
\hline\n\end{array}
$$

β(*g*)

$\sum_{c,d} \left(g_c, g_d \right) = \frac{1}{2^{(2\alpha_0 - \frac{1}{2} \log D)} \cdot 1}$ \sum_{s} fixed points: $\beta^{(2)}_{c. d.}$ $(g_c, g_d) = 0$

• Search perturbative

 $\mu_{DM} \sim 0.5$ der DM mass: $M_{DM} \approx 3.5 \; GeV$

*Y*1*,*²

Xⁱ

```
\alpha_d (\Lambdag_c, g_d) = 0\alpha_d(\Lambda_{dQCD}) \equiv\pi4
                         M_{DM} \approx 1.5*\Lambda_{dQCD}
```


 α_s

 Λ_{QCD} Λ_{dQCD} M Λ_{U}

 α_s

DARK MATTER

Can co-generate DM and baryon asymmetry

$$
\overline{Q} \oplus d_i \longrightarrow \text{SM quark}
$$

Subifundamental scalar

Dark matter is strongly self interacting — potentially solves various problems of cold dark matter

Rocha et. al. '12. Peter et. al. '12. Vogelsberger, Zavala, Loeb, '12. Zavala, Vogelsberger, Walker '12.

- Cusp vs core Missing satellites
	- Too big to fail

PHENOMENOLOGY

DARK QCD

Confining SU(*Nc*) gauge group with *Nf* flavors Q_i \overline{Q}_j $G_d^{\mu\nu}$

This sector is QCD like, and it confines at a scale $\Lambda_d \sim 1 - 10 \text{ GeV}$

At the confining scale we have all the usual states

$$
\begin{array}{c}\n\boxed{Pd} & \boxed{\pi_d} & \text{Zoo}_d \\
\text{Stable} & \text{Decays} \\
\hline\n\text{to SM}\n\end{array}
$$

DANIEL STOLARSKI November 24, 2014 University of Warsaw 13

MEDIATORS

Motivated by getting comparable asymmetries, put in heavy mediator which couples to SM and dark sector

$M_{\Phi} \gg \Lambda_d$

<u>Example 1:</u> Φ is a scalar charged under both color and dark color

MEDIATORS

Motivated by getting comparable asymmetries, put in heavy mediator which couples to SM and dark sector

$M_{\Phi} \gg \Lambda_{d}$

<u>Example 2:</u> Z_d is a vector that couples to quarks and dark quarks **Strassler, Zurek, PLB 07.** *q* Z_d *Q^d*

 Q_d

q

QCD JETS

ALEPH event

PION DECAY

Operator used to generate asymmetry mediates decay

 $Q \Phi d_i$ Φ *q* \overline{q} *Q*^{*d*} Q_d $q\!\!\!q$ \overline{q} \overline{Q}_d $\bar{\overline{\eta}}$ Integrate out Φ **Notify** π_d Dark pion decays to quarks

PION DECAY

Same story for Z_d model:

DECAY LENGTH

$$
\frac{1}{M_X^2}\overline{Q}\gamma_\mu Q\,\bar{d}_R\gamma^\mu d_R
$$

Can use (dark) chiral Lagrangian to estimate:

$$
\Gamma(\pi_d \to \bar{d}d) \approx \frac{f_{\pi_d}^2 m_\text{d}^2}{32\pi M_{X_d}^4} m_{\pi_d}
$$
\n
$$
c\tau \approx 5 \,\text{cm} \times \left(\frac{1 \text{ GeV}}{f_{\pi_d}}\right)^2 \left(\frac{100 \text{ MeV}}{m_\text{d}}\right)^2 \left(\frac{1 \text{ GeV}}{m_{\pi_d}}\right) \left(\frac{M_{X_d}}{1 \text{ TeV}}\right)^4
$$

DECAY LENGTH

 \sim 2 m (CMS)

Look for jets with no/few tracks in the circle

pp → *QQ*

BACKGROUND?

QCD 4-jet production in PYTHIA 8 *p^T >* 200 GeV

BACKGROUND COMPOSITION

10

 1

1500

 $\mathbf{1}$

 $\mathbf{1}$

 $\mathbf{1}$

Γ neutron strange Flavor of earliest decaying track

jet $p_T > 200 \text{ GeV}$ track $p_T > 1$ GeV

TRACKLESS BACKGROUND

Composition of completely trackless background

jet $p_T > 200 \text{ GeV}$ track $p_T > 1$ GeV

DARK SECTOR Λ*^d* 10 GeV 4 GeV **n** α 20 α 20 α 20 α

Choose two benchmarks:

$$
N_c = 3 \text{ and } n_f = 7
$$

Dark QCD alrea Dark QCD already in PYTHIA! is the mass of the dark vector mesons, and *m*π*^d* is the pseudo-scalar mass. *c* τπ*^d* is the rest frame decay length of the pseudo-scalars. We take *N^c* = 3 and *n^f* = 7 in both benchmarks.

Carloni, Sjorstrand, 2010.

Carloni, Rathsman, Sjorstrand, 2011. Carloni, and the mediator, and we will describe the mediator and we will describe the mediator.

Run modified version with running. Run modified version with running IN THIS SECTION WE WE WE WE WILL DESCRIPTION WE WILL DESCRIPTION OF THE MEDIATOR, AND THE MEDIATOR, AND WE WILL DESCRIPTION OF THE MEDIATOR, AND IN THE MEDIATOR, AND IN THE MEDIATOR, AND IN THE MEDIATOR, AND IN THE MEDIATO

COUPLING \bullet 0 \bullet 0 \bullet 400 \bullet 600 \bullet 800 \bullet $0.$ p_T (GeV) 0° 0.00

Modify PYTHIA to include gauge comparision of the Modify PYTHIA to include gauge coupling running in the dark sector implemented. The left plot is the orphan *p^T* : the scalar sum of the *p^T* of visible particles which are not clustered into a jet of *p^T >* 200 GeV. The right plot is the girth distribution

SIMULATION \blacklozenge 6 \blacklozenge 6 0 " " " \blacksquare " " " "

"

"

"

 β II SIIII WAGALION II2GAKES SE β AGC DY 4000 Check to see if simulation makes sense by *s* !GeV" looking at average particle multiplicity

"

30

MESON MULTIPLICITY

Number of dark mesons in a jet

BENCHMARK MEDIATOR 1 $pp \rightarrow \Phi \Phi^{\dagger} \rightarrow \bar{q} Q_d \overline{Q}_d q$ \mathbf{F} Pair \mathbf{F}

BENCHMARK MEDIATOR 1 $pp \rightarrow \Phi \Phi^{\dagger} \rightarrow \bar{q} Q_d \overline{Q}_d q$

Final state is

- 2 QCD jets
- 2 emerging jets

Cross section is stop-like $\sigma \approx \text{few} \times \sigma(pp \to \tilde{t}_1 \tilde{t}_1)$ $\sigma(M_{\Phi} = 1 \,\text{TeV}) \approx 10 \,\text{fb}$ @ LHC14

BENCHMARK MEDIATOR 2 $pp \rightarrow Z_d \rightarrow Q_d Q_d$

Final state is

• 2 emerging jets

Cross section depends on couplings

Work in progress

JET MOMENTA

Hardest jet p_T

Four hard jets is enough to pass trigger κ inmatic cuts described in the text and also have at left plot is for κ Four hard jets is enough to pass trigger

JET SHAPES

JET SHAPES

SEARCH STRATEGY 2 tracks $\overline{}$ fraction

$pp \rightarrow \Phi \Phi^{\dagger} \rightarrow \bar{q} Q_d \bar{Q}_d q$ PP I I $\alpha \propto a \vee a$ <u>.</u>

 $m_{\pi_d} = 2 \text{ GeV}$ *c* $\tau_{\pi_d} = 5 \text{ mm}$ $\int \mathbf{r} \cdot \mathbf{r}_{d} = 2 \text{ GeV}$

fraction

 $c\tau_{\pi_d} = 5$ mm

DIFFERENT MODEL POINTS

Model **A'** Model **A''**

CUT FLOW

Cross sections in fb at LHC14:

Paired di-jet resonance search very difficult! For the signal we take the mass of the bifundamental *M^X* = 1 TeV. The two right most columns are Requiring emerging jets changes the game.

ALTERNATIVE STRATEGY

Fraction of jet energy reconstructing outside of circle

Neutrals (photon, neutron) do not contribute, hard to get *F*=1

Much more robust to pile-up

F DISTRIBUTIONS ———————————————————— 1.!10"⁶ $\sqrt{10}$ 0.0 0.2 0.4 0.6 0.8 1.0 0.00 0.01

0.0 0.2 0.4 0.6 0.8 1.0

0.0 0.2 0.4 0.6 0.8 1.0

ALTERNATIVE CUT FLOW

Cross sections in fb:

b-jet background too large at *r*=10 mm mm than for 10 mm. This is a consequence of *b* hadrons; in Fig. 5 we see that *b* hadrons tend to decay

Works pretty well at *r*=100 mm even for short lifetime model $b = \frac{10}{2}$ protty $m = 100$ mm, such for contribute to *F*, but for *r* = 100 mm, only strange mesons contribute. Looking at the third row of $\frac{1}{2}$ short lifetime model

CAMPLES SEARCH CMS Physics Analysis Summary

Search for long-lived neutral particles decaying to dijets

The CMS Collaboration

Abstract

A search is performed for long-lived massive neutral particles decaying to quarkantiquark pairs. The experimental signature is a distinctive topology of a pair of jets originating at a secondary vertex. Events were collected by the CMS detector at the LHC during pp collisions at $\sqrt{s} = 8$ TeV, and selected from data samples corresponding to 18.6 fb−¹ of integrated luminosity. No significant excess is observed above standard model expectations and an upper limit is set with 95% confidence level on the production cross section of a heavy scalar particle, H^0 , in the mass range 200 to 1000 GeV, decaying into a pair of long-lived neutral X^0 particles in the mass range 50 to 350 GeV, which each decay to quark-antiquark pairs. For X^0 mean proper lifetimes of 0.1 to 200 cm the upper limits are typically 0.3−300 fb.

Available on the CERN CDS information server **CMS PAS EXO-12-038**

sis Summary

CMS SEARCH

Available on the CERN CDS information server **CMS PAS EXO-12-038**

$ICERN$ is important to evaluate the performance of the performance of the LJ search criteria by setting limits on I models that predict LJs in the final state \mathcal{A} relevance are models which predicted which predicted which predicted which predicted with predicted values of particular relevance are models which predicted with predic non-Sm Higgs boson decays to LJs. Indeed, the phenomenology of the phenomenology of the phenomenology of the H **ATLAS SEARCH** radiated photons, however, varies on an event-by-event basis and depends on unknown model-dependent parameters such as the hidden gauge coupling α*d*. possibility is not considered here.

possibility is not considered here.

HLSP

Search for long-lived neutral particles decaying into lepton jets

in grater proton, collisions of $\sqrt{2}$, a TeV with the ATLAC **in proton–proton collisions at** \sqrt{s} = 8 TeV with the ATLAS **detector**

The ATLAS Collaboration

Abstract

H lepton jets are derived as a function of the particle's proper decay length. sions at the centre-or-mass energy or \sqrt{s} = 6 TeV in a sample or 20.3 ib $^{-1}$ collected during 2012 with
the ATLAS detector at the LHC. Limits on models predicting Higgs boson decays to neutral long-lived sions at the centre-of-mass energy of \sqrt{s} = 8 TeV in a sample of 20.3 fb⁻¹ collected during 2012 with states consisting or commated jets or light reptons and hadrons (so-called repton jets). These parti-
cles can also be long-lived with decay length comparable to, or even larger than, the LHC detectors' Several models of physics beyond the Standard Model predict neutral particles that decay into final states consisting of collimated jets of light leptons and hadrons (so-called "lepton jets"). These partilinear dimensions. This paper presents the results of a search for lepton jets in proton–proton colli-

 $arXiv:1409.0746v2$ [hep-ex]

ATLAS SEARCH

neutral collisionsossibly for long ..
M presents des l'al
Locs S
length
of the space of th extremely low efficiency except **Figure 15. The 95 and 77 (16)** possi possibly for long lifetimes **Against 19** (*right)*, and *a* function of the *possibly* for tong experies the dashed curve and the almost identical solid curve shown as the observed curve shows the observed curve

r:
T **at d** .
Tr $\overline{}$ 014 **Abstract** l. iee al e
An leptons and $\overline{\mathbf{C}}$ in \bullet data photons. See also ATLAS trigger paper: arXiv:1305.2204 [hep-ex].

LHC b

21 stations

LHCb has excellent tracking

Limited coverage of event

13.9.2012 1:48:03 Run 128262 Event 100499354 bld 1886

LHC b

Still get a good fraction of events with dark pions in LHCb

Discussions with collaboration ongoing

POWER OF EMERGING JET HLSP and to a hidden scalar, *s*d¹ that in turn decays to pairs of dark photons. For the γ^d decays, only electron, muon and pion final states are considered. In general, radiation in the hidden sector may occur, resulting in additional hidden photons. The number of such

Emerging jet sear sensitive to other scenarios

Lepton jets

• …

RPV neutralinos

Work in progress

CONCLUSIONS

- Important to explore different ways LHC can search for NP
- DM exists, exhaustively search for different classes of models
- Emerging jets are novel and motivated, no current searches are sensitive
- Strategies presented here can reach very low cross sections, sensitive to broad class of displaced models
- ATLAS and CMS exotics groups are investigating

THANK YOU

DIFFERENT MODEL POINTS

Model **A** Model **B**

SIGNAL VS BG

$$
m_{\pi_d} = 5 \,\text{GeV}
$$

$$
c\tau_{\pi_d} = 50 \,\text{mm}
$$

MESON MOMENTUM FRACTION 10 400

Fraction of let momentum carried by the control of **Contract Contract C** In[70]:= **GraphicsRow!"FragA, FragB#\$** Fraction of jet momentum carried by any individual dark meson

