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Scattering of massive electroweak gauge bosons - Definition and significance

Measurement of electroweak production of two like-charge W bosons at ATLAS Signal definition

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Vector Boson Scattering and Measurement of Electroweak Production of two Like-Charge W Bosons and two Jets at the ATLAS detector

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Peferences

LHC Run 1:

 Start to investigate the nature of electroweak symmetry breaking after Higgs discovery



• Great experimental support for the Standard Model

After the first run of the LHC





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Many issues remain:

- · properties of the Higgs
- is there new physics beyond the Standard Model?
 - in the electroweak symmetry breaking sector
 - elsewhere

... to explain hierarchy problem, dark matter, ...

ATLAS Prelim - σ(stat.) Total uncertainty mu = 125.36 GeV o(sys inc.) t1σ on μ $H \rightarrow \gamma \gamma$ + 0.23 $\mu = 1.17^{+0.27}_{-0.27}$ $H \rightarrow 77^* \rightarrow 4I$ + 0.34 $\mu = 1.44^{+0.40}_{-0.33}$ $H \rightarrow WW^* \rightarrow IvIv$ $\mu = 1.08^{+0.22}_{-0.20} + 0.16_{-0.13}$ W.Z $H \rightarrow b\overline{b}$ $\mu = 0.5^{+0.4}_{-0.4}$ ATLAS CONF-2014-06 $H \rightarrow \tau \tau$ $\mu = 1.4^{+0.4}_{-0.4}$ 0.5 1.5 2 (s = 7 TeV (1 dt = 4.5-4.7 fb Signal strength (µ) s = 8 TeV Ldt = 20.3 fb⁻¹

One process to address remaining questions on the electroweak symmetry breaking mechanism and the nature of gauge interactions:

The more we know, the more we don't know

Electroweak vector boson scattering



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Scattering of massive electroweak gauge bosons - Definition and significance

Outline of this talk

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Peferences

Scattering of electroweak gauge bosons

Process definition of electroweak gauge boson scattering at the LHC:



V = W, Z



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 $VV \rightarrow VV$ contains:



Scattering of electroweak gauge bosons

Diagrams with triple gauge couplings

Constrained by diboson production

Quartic gauge boson vertex First observation via $W^{\pm}W^{\pm}jj$ -EW production

Observed in other channels

Higgs exchange

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Vector Boson Scattering (VBS)

- Contains VVVV vertex (predicted by the SM)
 ⇒ probe self-interactions of the electroweak gauge bosons
- Contains scattering of longitudinal components of the gauge bosons
 ⇒ Probe electroweak symmetry breaking mechanism



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Longitudinal VBS

The scattering of longitudinal vector bosons violates unitarity in the absence of a SM Higgs boson:



Gauge bosons:

$$\mathcal{M}^{gauge} = -\frac{g_w^2}{4m_W^2}u + \mathcal{O}\left(\frac{E}{m_W}\right)^0$$

 $\Rightarrow \sim E^2$, violates unitarity

Higgs exchange: (s, t, u $>> m_W^2, m_H^2$)



 \Rightarrow terms cancel



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Longitudinal VBS

The scattering of longitudinal vector bosons violates unitarity in the absence of a SM Higgs boson:

 $W^+Z \rightarrow W^+Z$, m_H = 125 GeV · $W^+Z \rightarrow W^+Z$, my = 10¹⁰ GeV 3×10^{6} $ZZ \rightarrow W^+W^-$, m_H = 125 GeV $ZZ \rightarrow W^+W^-$, m_H = 10¹⁰ GeV ----- $W^+W^+ \rightarrow W^+W^+$, m_H = 125 GeV $W^+W^+ \rightarrow W^+W^+$, m_H = 10¹⁰ GeV 2.5×10^{6} $W^+W^- \rightarrow W^+W^-$, m_H = 125 GeV $W^+W^- \rightarrow W^+W^-$, m_H = 10¹⁰ GeV 2×10^{6} 1.5×10^{6} 1×10^{6} 5×10^{5} 0 0 500 1000 1500 2000 2500 CME [GeV]

Gauge bosons:

$$\mathcal{M}^{gauge} = -\frac{g_{w}^{2}}{4m_{W}^{2}}u + \mathcal{O}\left(\frac{E}{m_{W}}\right)^{0}$$

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Higgs exchange: (s, t, u $>> m_W^2, m_H^2$)

$$\mathcal{M}^{\mathrm{Higgs}} = rac{g_w^2}{4m_W^2}u$$

 \Rightarrow terms cancel



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Higgs exchange: (s, t, u $>> m_W^2, m_H^2$)

$$\mathcal{M}^{\mathrm{Higgs}} = rac{g_{W}^{2}}{4m_{W}^{2}}u$$

 \Rightarrow terms cancel

 \Rightarrow Probing the Higgs properties through the scattering of electroweak gauge bosons



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Scattering of two like-charge W bosons

$W^{\pm}W^{\pm}$ scattering analysis at 8 TeV published by the ATLAS collaboration in Phys. Rev. Lett. 113, 141803 [1]



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VBS signal definition



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VBS signal definition

VVjj-EW: $\mathcal{O}(EW) = 6$ (in	ncl. decays)		
VBS	irreducible EW	vvv	
	+ +		~~ *~ *
$\frac{\nabla \nabla jj}{\partial CD} + \frac{\partial C}{\partial C} + \frac{\partial C}{\partial C} + \frac{\partial C}{\partial C}$	$\mathcal{O}(QCD) = 2 \text{ (incl. dec})$	ays)	x not included in W [±] W [±] jj

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VBS in $W^{\pm}W^{\pm}jj$ channel

C <u>ross</u>	sections	in	various	VV	channels at 8 TeV:	

Final state	Process	VVjj-EW	VVjj-QCD	Ratio
$\ell^{\pm}\nu\ell'^{\pm}\nu'jj$	$W^{\pm}W^{\pm}$	19.5 fb	18.8 fb	1:1
$\ell^{\pm} \nu \ell^{\mp} \nu j j$	$W^{\pm}W^{\mp} + ZZ$	93.7 fb	3192 fb	1:30
$\ell^{\pm}\ell^{\mp}\ell'^{\pm} u'jj$	$W^{\pm}Z$	30.2 fb	687 fb	1:20
lllljj	ZZ	1.5 fb	106 fb	1:70

(cuts: 2 leptons: $p_T >$ 5 GeV, $m_{\ell\ell} > 4$ GeV, 2 jets: $p_T >$ 10 GeV)

VBS by channel

- W[±] W[±] jj most promising: high EW/QCD ratio (no gluons in initial state)
- $W^{\pm}Zjj$ has larger QCD contribution, but profits from clean 3 lepton signature
- $W^{\pm}W^{\mp}$ overwhelmed by ttbar background \rightarrow use MVA methods



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VVjj-EW vs. VVjj-QCD

Comparison of $W^{\pm}W^{\pm}jj$ -EW and $W^{\pm}W^{\pm}jj$ -QCD

Tagging the VBS process according to dijet kinematics





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VVjj-EW vs. VVjj-QCD

Comparison of $W^{\pm}W^{\pm}jj$ -EW and $W^{\pm}W^{\pm}jj$ -QCD

Kinematics of the leptons, i.e. the decay products of the *W* bosons:





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Tile calorimeters LAr hadronic end-cap and forward calorimeters Pixel detector Toroid magnets LAr electromagnetic calorimeters Solenoid magnet Transition radiation tracker Muon chambers

44m

ATLAS detector

Semiconductor tracker



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VBS topology

- two forward energetic tagging jets (initial quarks radiating off *Ws*)
- two central leptons (decay products of Ws)
- Missing E_T from Ws

Inclusive analysis phase space

- 2 same-sign high-p_T leptons
- at least 2 high- $p_{\rm T}$ jets
- $E_{\rm T}^{\rm miss} > 40 \, {\rm GeV}$
- leading jets: $m_{jj} > 500 \text{ GeV}$

VBS analysis phase space

inclusive analysis phase space +

• $|\Delta y(jj)| > 2.4$

Measure $W^{\pm}W^{\pm}jj$ (incl. EW, QCD)

Measure $W^{\pm}W^{\pm}jj$ EW and set aQGC limits

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 $p_{\rm T}({\rm lep 1}) = 180.2 \,{\rm GeV}$ $p_{\rm T}({\rm lep 2}) = 37.5 \,{\rm GeV}$ $E_{\rm T}^{\rm miss} = 74.8 \,{\rm GeV}$

Phys. Rev. Lett. 113, 141803

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Backgrounds to $W^{\pm}W^{\pm}$ jj

Backgrounds arise due to

- · Events with two prompt like-charge leptons in the final state
- Events with two opposite-charge leptons in the final state, if one charge is mis-measured
- Events with one charged lepton and one lepton from jet mis-identification or hadronic decays
- · Double parton interaction (negligible)



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Containing two real prompt same-charge leptons

Prompt background

- WZ/γ^* +jets
- ZZ+jets
- $t\overline{t}+W/Z$, tZj

- Main contribution: WZ/γ^* +jets
- · Suppressed by a veto on additional, looser leptons
- Estimated from MC and cross checked in a control region where additional lepton is required



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Photon conversions

Prompt photon conversion: $W\gamma$

 \rightarrow estimated from MC

Charge mis-identification due to bremsstrahlung with conversion: Z/γ^* +jets or $t\bar{t}$

ightarrow estimated from a control-region with two oppositely charged electrons by applying charge-flip rates determined in Z
ightarrow ee events

Cross check in \leq 1 jet control region:



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Other non-prompt leptons

Jets mis-reconstructed as leptons

or

W+jets

leptons from hadron decays in jets

- · Suppressed by requiring high quality, isolated leptons
- Contribution is estimated by applying fake rates to a dijet control region

Cross check in low M_{jj} control region:



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Event yields

Process	Inclusive Region	VBS Region
Prompt backgrounds Conversions Fake leptons W [±] W [±] jj-QCD	$11.7 \pm 1.6 \\ 5.6 \pm 1.1 \\ 2.3 \pm 0.9 \\ (ext{is signal})$	8.3 ± 1.3 4.0 ± 0.9 2.3 ± 0.7 1.3 ± 0.1
Background total Expected signal	$\begin{array}{c} 21\pm2.4\\ 22\pm1.5\end{array}$	$\begin{array}{c} 16\pm2\\ 14\pm1 \end{array}$
Data	50	34

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Tagging jets invariant mass



Kinematics of jets

(Inclusive region before m_{ii} cut)

Phys. Rev. Lett. 113, 141803



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Background estimates

Invariant mass split by channel Inclusive region before M_{ii} cut

 $\mu\mu$

all

m_i [GeV]

m, [GeV]



Phys. Rev. Lett. 113, 141803



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Kinematics of leptons



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Results

Extract cross section by fitting a likelihood function to the observed data

Total cross sections

	Observed	Significance
EWQCD (incl. Region)	2.1 ± 0.5 (stat.) \pm 0.3 (syst.) fb	4.5 σ
EW (VBS Region)	1.3 \pm 0.4 (stat.) \pm 0.2 (syst.) fb	3.6σ

Comparison to theory prediction:







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Model effects of new physics via higher-order operators

Effective field theory ansatz:

$$\mathcal{L}_{ ext{EFT}} = \mathcal{L}_{ ext{SM}} + \sum_{d>4} \sum_{i} rac{c_i^{(d)}}{\Lambda^{d-4}} \mathcal{O}_i^{(d)}.$$

 \Rightarrow Introduce additional higher-dimensional (*d*) operators \mathcal{O} which modify the electroweak gauge couplings, suppressed by scale of new physics Λ



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Electroweak chiral Lagrangian

In analogy to chiral perturbation theory in QCD: Break electroweak $SU(2)_L \otimes U(1)_Y$ to $U(1)_Q$ and custodial symmetry $SU(2)_C$

 \Rightarrow Electroweak chiral Lagrangian with the operators

$$\mathcal{L}_4 = lpha_4 (\mathrm{tr}[\mathbf{V}_\mu \mathbf{V}_
u])^2$$

$$\mathcal{L}_5 = \alpha_5 (\mathrm{tr}[\mathbf{V}_{\mu}\mathbf{V}^{\mu}])^2$$

with parameters α_4 and α_5



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In general, additional operator terms lead to unitarity violation of the amplitude:



⇒ apply unitarization procedure to retain physical events

Unitarity

unitarization with K-matrix method (projection of amplitude on Argand circle) as implemented in WHIZARD [2, 3]



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Influence on kinematics

Anomalous quartic gauge couplings modify the kinematics of the decay products of the gauge bosons:





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Limits on aQGC

Set limits on α_4, α_5 in the VBS region of the $W^{\pm}W^{\pm}jj$ -EW analysis



Use fiducial cross sections with anomalous couplings to derive confidence intervals



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с⁵

0.6

04

0.2

0

-0.2

-0.4

-0.6F

confidence intervals

× Standard Model

0 01 02

68% CL

95% CL – expected 95% CL

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Limits on aQGC

Set limits on α_4, α_5 in the VBS region of the $W^{\pm}W^{\pm}jj$ -EW analysis

 20.3 fb^{-1} , $\sqrt{s} = 8 \text{ TeV}$

K-matrix unitarization

ATLAS

 $pp \rightarrow W^{\pm}W^{\pm}ii$

Use fiducial cross sections with anomalous couplings to derive confidence intervals One-dimensional limits at 95 % C.L. :

$$\alpha_4 \in [-0.229, 0.244]$$

 $\alpha_5 \in [-0.229, 0.244]$

 \Rightarrow First limits on aQGC with massive electroweak vector bosons

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Interpretation in terms of resonances

Anomalous couplings can be interpreted in terms of models such as:

- Extended Higgs sectors (2HDM, etc)
- · Composite Higgs, Technicolor

or simplified models of generic resonances ordered by spin, isospin, with different masses, couplings, and widths [4]:

resonance	spin J	isospin l	Γ/Γ_0	$\Delta lpha_4$ / $lpha$	$\Delta lpha_{5}$ / $lpha$
ϕ	0	2	1	1/4	-1/12
t	2	2	1/30	-5/8	35/8



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Interpretation in terms of resonances



- ϕ resonance (spin = 0, isospin = 2): $\frac{M}{g}$ < 130 GeV excluded
- t resonance (spin = isospin = 2): $\frac{M}{g} < 90$ GeV excluded



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Prospects at 14 TeV

Estimate sensitivity to anomalous couplings in $W^{\pm}W^{\pm}jj$ -EW at High-Luminosity LHC ($\sqrt{s} = 14$ TeV, $\mathcal{L} = 3000$ fb⁻¹) \Rightarrow Use performance projections for the future ATLAS detector in terms of

- efficiency (trigger, particle identification, flavor tagging)
- resolutions (momentum, energy)

Backgrounds: MC for WZ, WWss-QCD.

WZ background is scaled by a factor 2 to account for additional background contributions from charge-flip, jet fakes, photon conversion

Selection:

- 2 leptons with same charge, $p_T > 25$ GeV
 - \geq 2 jets, $p_T > 50$ GeV
 - m(jj) > 1000 GeV



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3000F

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L = 3000 fb⁻¹

VBS ssWW (SM)

SM VBS ssWW + f_{so} = 10 TeV⁴

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 m_{llii} with a signal of

 $f_{S,0} = 10 \text{ TeV}^{-4}$

Prospects at 14 TeV [5]



signal significance (in standard deviations) as a function of f_{S0}



Conclusions

W+W+ scattering at ATLAS

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- Measurement of electroweak production of two like-charge W bosons at ATLAS
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- Scattering of massive electroweak gauge bosons is a crucial process to study the electroweak gauge interactions and the mechanism of electroweak symmetry breaking
 - First evidence of $W^{\pm}W^{\pm}$ scattering and $W^{\pm}W^{\pm}jj$ production \Rightarrow Measurements in good agreement with the SM prediction
 - First limits on anomalous couplings of massive electroweak gauge bosons of parameters α_4, α_5



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Backup

Backup



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References



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- A. Alboteanu, W. Kilian, and J. Reuter, *Resonances and Unitarity in Weak Boson Scattering at the LHC*, JHEP **0811** (2008) 010, arXiv:0806.4145 [hep-ph].
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Scattering of massive electroweak gauge bosons - Definition and significance

- Measurement of electroweak production of two like-charge W bosons at ATLAS
- Signal definition Phase space definition Background estimates Results
- Anomalous quartic gauge couplings Effective field theory Limits on aQGC Interpretation in terms of resonances

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References

Available experimental constraints on 4 electroweak gauge boson couplings

- No observation of processes with quartic vertex of massive electroweak gauge bosons before the presented measurement
- · No constraints on massive gauge boson quartic couplings
- Constraints on anomalous quartic gauge couplings with photons exist:

Experiment	channel	Publication	year
CMS	WW γ , WZ γ	CMS-PAS-SMP-13-009	2013
CMS	WW $\rightarrow \gamma \gamma$	JHEP07 (2013) 116	2013
D0	$\gamma\gamma ightarrow WW$	Phys Rev D 88 012005	2013
L3	$WW\gamma$	Phys. Lett. B 527:29-38, 2002	2002
OPAL	$WW\gamma$	Phys. Lett. B580:17-36, 2004	2004
DELPHI	$WW\gamma$	Eur. Phys. J C31:139-147, 2003	2003



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References



- well isolated, tight quality cuts
- $p_{\rm T}$ > 25 GeV, $|\eta| < 2.5$
- $m_{\ell\ell} > 20 \, \text{GeV}$
- at least 2
 - $p_{\rm T}$ > 30 GeV, $|\eta| < 4.5$

 $E_{\rm T}^{\rm miss} > 40 \, {\rm GeV}$

Veto

Jets

 events with additional softer and less isolated leptons

Analysis event selection

events with b-tagged jets

VBS cuts

- $m_{ii} > 500 \, \text{GeV}$
- $|\Delta y_{ii}| > 2.4$ (in VBS region)

