

# Well-tempered n-plet dark matter

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# Motivation: WIMPs

## Observation:

EW-scale mass + EW-scale cross section  
⇒ thermal relic density  $\Omega h^2 \approx 0.1$  = observed value

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“WIMP miracle”

# Motivation: WIMPs

## A closer look:

To match the observed relic density

- a thermal higgsino needs  $m = 1.1 \text{ TeV} \gg m_Z$
- a thermal wino needs  $m = 2.5 \text{ TeV} \gg m_Z$
- minimal DM (fermionic 5-plet or scalar 7-plet) needs  $m \approx 10 \text{ TeV} \gg \gg m_Z$
- ...

A WIMP with generic electroweak quantum numbers and a mass of  $\sim 100 \text{ GeV}$  has a **too large annihilation cross section**.

A WIMP which is within kinematic reach of LHC is necessarily **mostly  $SU(2) \times U(1)$  singlet**.

# Motivation: WIMPs

## Assumptions for this talk:

- DM = mostly  $SU(2) \times U(1)$  singlet
- Some admixture of non-singlet state for right relic density from coannihilation
- Mass of order 100 GeV to be within LHC reach

**Q:** Why should nature care about LHC discovery potential?

**A:** She might not — but we do!



# Motivation: WIMPs

## Further assumptions:

- DM is a fermion
- Stabilized by  $\mathbb{Z}_2$
- Sub-TeV particle content is minimal
- We don't consider EW doublets
  - extensive literature on well-tempered bino-higgsino and its non-SUSY version

Dark matter is a singlet fermion  $\chi$  mixing with an  $n$ -plet fermion  $\psi$  ( $n \geq 3$ ) through higher-dimensional operators.

States inducing the mixing live at scales  $\gtrsim$  TeV  $\Rightarrow$  mostly irrelevant for LHC if carrying only EW charges.

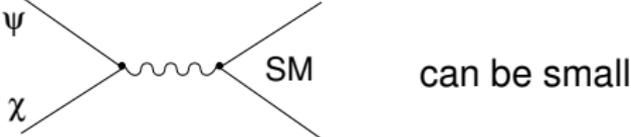
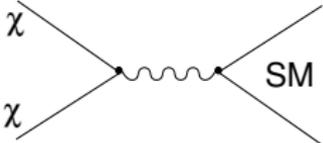
Familiar example: Split SUSY with somewhat heavy higgsinos and  $M_1 < M_2$ . DM is mostly bino, mixing with wino through dimension-5 operator

$$\mathcal{L}_{\text{mix}} = \frac{\tilde{g}_u \tilde{g}'_d + \tilde{g}_d \tilde{g}'_u}{\mu} \phi^\dagger \tau^a \phi \widetilde{W}^a \widetilde{B}$$

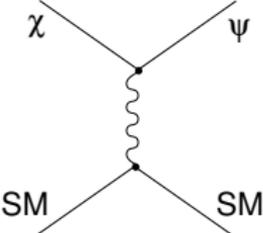
→ e.g. Arkani-Hamed/Delgado/Giudice '06

# Relic density

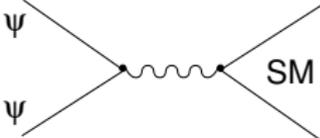
**Somewhat under-appreciated fact:** Coannihilation can be **very efficient** in reducing the DM relic density **even when DM coupling to SM is tiny**.



provided that  $\chi$  remains in equilibrium via



and that

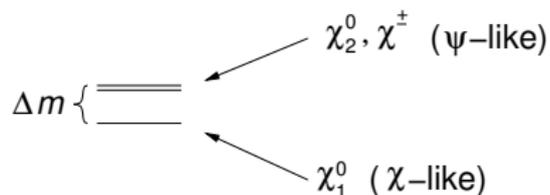


is efficient.

**Annihilation of  $\psi$  efficiently depletes  $\chi$  density,**  
even if  $\psi - \chi$  mixing angle is  $\ll 1$ .

# First case study: SU(2) triplets

Particle content: one charged and two neutral fermions



Interactions of  $\chi_1^0$  with SM mainly through two operators:

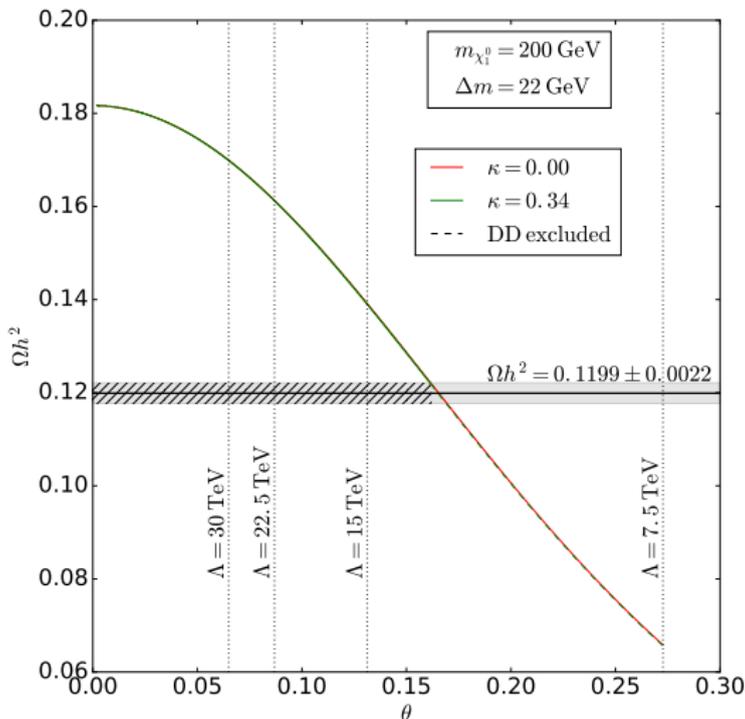
$$\mathcal{L} = \frac{1}{2} \frac{\kappa}{\Lambda} \phi^\dagger \phi \chi \chi + \frac{\lambda}{\Lambda} \phi^\dagger \tau^a \phi \psi^a \chi + \text{h.c.}$$

$\phi$  = SM Higgs;  $\Lambda$  = cutoff scale.

Wilson coefficients  $\kappa, \lambda$  both contribute to **DM annihilation** ( $\Rightarrow$  thermal relic density) and **DM-nucleus scattering** ( $\Rightarrow$  direct detection).

Trade  $\lambda$  for **mixing angle**  $\theta = \frac{\lambda}{\Lambda} \frac{v^2}{\Delta m}$

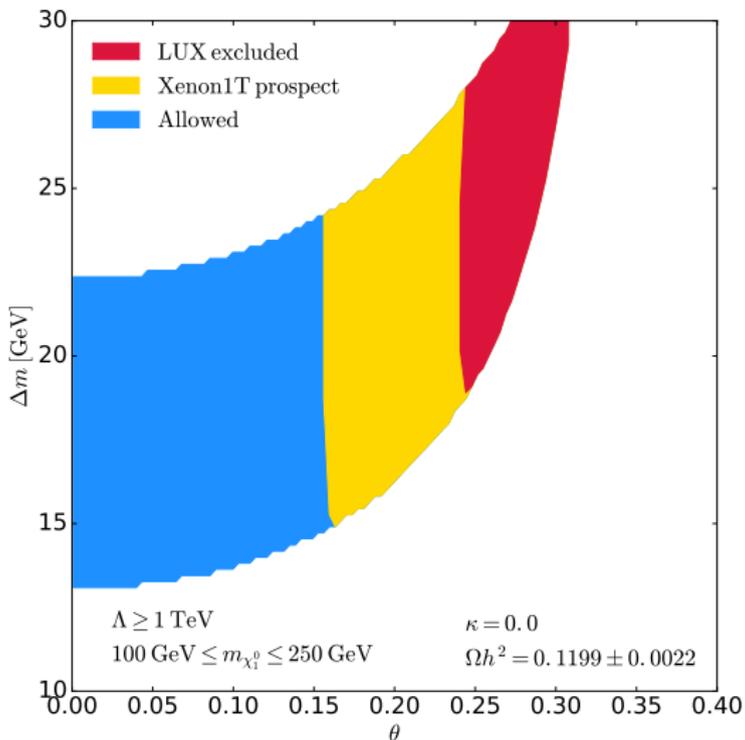
# First case study: SU(2) triplets



Any  $\kappa$  large enough to influence relic density significantly is ruled out by direct detection.

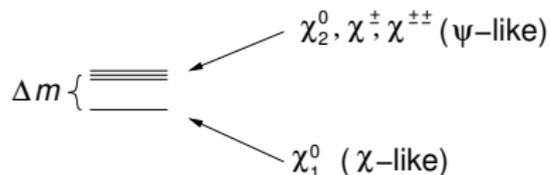
# First case study: SU(2) triplets

For  $\kappa = 0$ :



## Second case study: SU(2) quintuplets

Particle content: one doubly charged, one singly charged and two neutral fermions

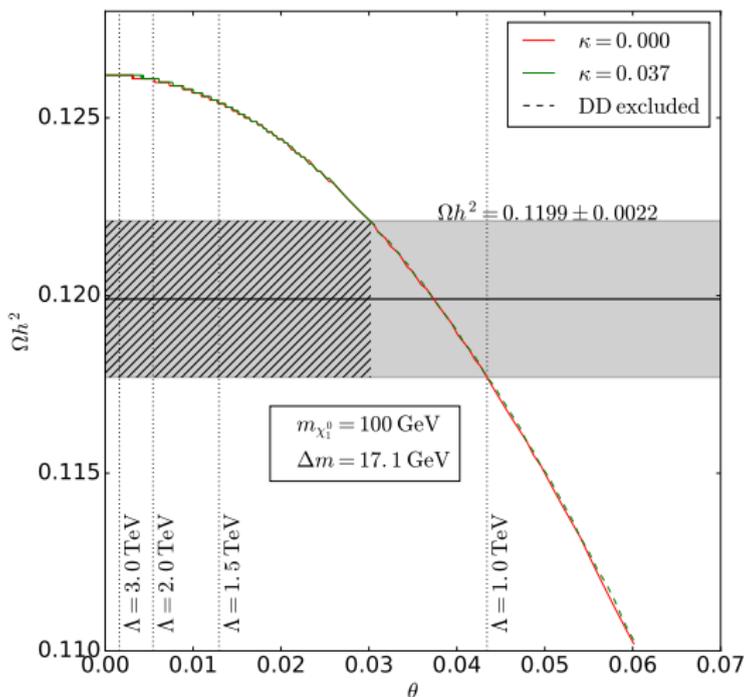


- Doubly charged state  $\Rightarrow$  potentially characteristic signatures at LHC (long-lived)
- Mixing operator is now **dimension 7**:

$$\frac{\lambda}{\Lambda^3} \phi^{\dagger i} \phi_j \phi^{\dagger k} \phi_\ell C_{A ik}^{j\ell} \psi^A \chi + \text{h.c.}$$

$\Rightarrow$  mixing angles **small** for reasonable cutoff scale ( $\Lambda \gtrsim \text{TeV}$ )

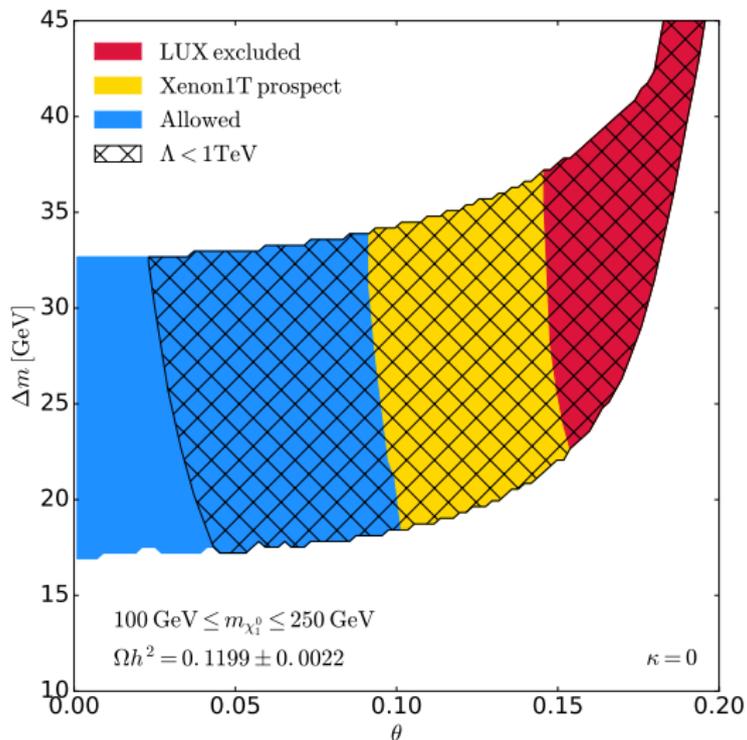
## Second case study: SU(2) quintuplets



Similar to triplet case. Mass parameters need some tuning. Wilson coefficient  $\kappa$  for dimension-5 operator  $\frac{\kappa}{\Lambda} \phi^\dagger \phi \chi \chi$  already **tightly constrained** by direct detection.

# First case study: SU(2) quintuplets

For  $\kappa = 0$ :



# On quadruplets

The quadruplet (isospin 3/2) case is more complicated because

- we now need a Dirac fermion ( $\psi, \bar{\psi}$ )
- one doubly charged, two singly charged mass eigenstates
- **three** neutral mass eigenstates  $\Rightarrow$  **two** relevant neutral mixing angles
- spectrum depends on additional Wilson coefficients inducing non-universal mass splittings in the  $\psi$  sector
- **proliferation of parameters**

**Result of our analysis:** Qualitatively similar conclusions as for triplet/quintuplet case  $\rightarrow$  [FB/Bharucha/Ruffault '17](#)

# Conclusions

- Dark matter could be a mixed singlet- $n$ -plet with an EW-scale mass.
- Effective theory. Mixing induced by higher-dimensional operators.
- Simplest example: Well-tempered bino-wino in split SUSY.
- More interesting for colliders: Singlet-quintuplet (doubly charged states).
- LUX already very constraining, even more so with Xenon1T.
- Collider phenomenology: under investigation.