

New signs of Supersymmetry: the R-axion

Bellazzini Mariotti Redigolo FS Serra, 1702.02152 + works in progress

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Take-home messages

Any other sign of SUSY to be looked for at colliders?

beyond gluino, stop, EWinos, extra Higgses,...

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Any other sign of SUSY to be looked for at colliders?

beyond gluino, stop, EWinos, extra Higgses,...

Spontaneously broken R-symmetry is quite generic

provides a naturally light state, the “R-axion” (whose pheno has been overlooked)

could be the first sign of SUSY at colliders

gives access to properties of the SUSY-breaking sector

opens new pheno & model-building avenues

Take-home messages

Spontaneously broken R-symmetry is quite generic
provides a naturally light state, the “R-axion”

R symmetry: one slide recap

N = 1 SUSY always accompanied by a continuous $U(1)_R = \text{“R-symmetry”}$

$$R : \theta_\alpha \rightarrow e^{i\epsilon} \theta_\alpha \quad [R, Q] = -Q$$

R-charge assignments:

$$\Phi = \phi + \sqrt{2}\theta \psi + \theta^2 F$$

$$r_\phi = r_\Phi$$

$$r_\psi = r_\Phi - 1$$

$$r_F = r_\Phi - 2$$

Vector superfields are real \Rightarrow **gauginos** have $r_\lambda = 1$

Lagrangian \mathcal{L} R-symmetric $\Rightarrow R(W) = 2$

(\Leftarrow if Kahler canonical)

$$\mathcal{L} \supset \int d^2\theta W + \text{c.c.}$$

W superpotential

Why bothering with R-symmetry?

Nelson Seiberg hep-ph/9309299

- i) SUSY broken in global minimum
- ii) W generic
(i.e. contains all terms not forbidden by symmetries)



Lagrangian respects a $U(1)_R$

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In MSSM + closer relatives $U(1)_R$ needs to be broken because of

1. gaugino masses

$$\mathcal{L} \supset m_\lambda \lambda\lambda$$

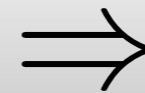
2. EW symmetry breaking & Higgsino masses

$$\mathcal{L} \supset B_\mu H_u H_d + \text{c.c.} \quad \& \quad W \supset \mu H_u H_d$$

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Ways out?

❖ Break $U(1)_R$ spontaneously \Rightarrow Massless Goldstone in the spectrum **R-axion a !**

❖ Violate i) or ii) (e.g. SUSY broken in metastable vacuum) [Intriligator Seiberg Shih 2006, 2007](#)

but then an analogous of Nelson Seiberg holds for an approximate $U(1)_R$

\rightarrow the R-axion is a pseudo-Goldstone, gets small mass m_a

✓ Motivation

Setup

Phenomenology

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A general description

Least possible ingredients in the IR:

SM 



Goldstone from $U(1)_R$ breaking

Goldstino from SUSY breaking

Free parameters

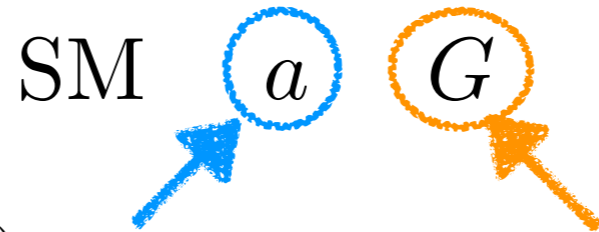
f_a
 F
 m_a

$U(1)_R$ breaking scale, controls a pheno

SUSY breaking scale, controls G pheno

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m_a From small R-symmetry breaking, quite generic

→ from cosmological constant [Bagger+ hep-ph/9405345](#)

$$m_a^2 \sim (10 \text{ MeV})^2 \times \frac{M_{\text{SUSY}}}{10 \text{ TeV}} \times \frac{m_{3/2}}{\text{eV}}$$

→ we live in a metastable SUSY-breaking vacuum

[Intriligator Seiberg Shih 2007](#)

→ happens in many models

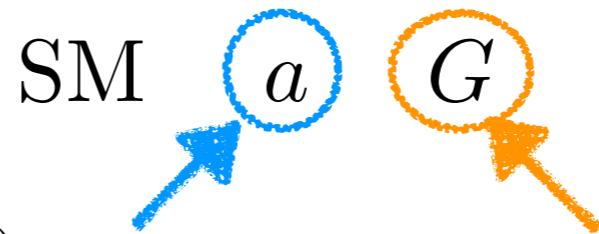
[Nelson Seiberg hep-ph/9309299](#)

[Affleck Dine Seiberg 1985,](#)

[Dine+ 1612.05770, Bellazzini+ in progress](#)

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[Goh Ibe 0810.5773](#)

Only mass range studied before our work

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[Intriligator Seiberg Shih 2007](#)

→ happens in many models

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The R-axion pheno Lagrangian-I

Komargodski Seiberg 0907.2441

Tool: constrained superfield formalism

$$X = \frac{G^2}{2F_X} + \sqrt{2}\theta G + \theta^2 F_X$$

$$\mathcal{R} = e^{i\mathcal{A}/f_a} = e^{ia/f_a} + O(aG, \dots)$$

satisfy the constraints
$$\begin{cases} X^2 = 0 \\ X(R^\dagger R - 1) = 0 \end{cases}$$

~ analogous to ordinary Goldstones $U^\dagger U = 1$ $U = e^{i\pi}$

Most general effective Lagrangian:

$$r_X = 2 \quad r_{\mathcal{R}} = 1$$

$$\mathcal{L}_{G+a} = \int d^4\theta (X^\dagger X + f_a^2 \mathcal{R}^\dagger \mathcal{R}) + \int d^2\theta (FX + w_R \mathcal{R}^2) + \text{c.c.}$$

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$$-\frac{w_R}{f_a F^2} \square a \bar{G} i \gamma_5 G$$

First pheno prediction (valid for any UV completion!):

R-axion decays to missing energy

$$w_R < \frac{1}{2} f_a F$$

$$\Gamma_{a \rightarrow GG} < \frac{1}{32\pi} \frac{m_a^5}{F^2}$$

Dine Festuccia Komargodski 0910.2527
see also Bellazzini 1605.06111

The R-axion pheno Lagrangian-II

$$\mathcal{L}_{\text{gauge}} = \int d^2\theta \left(\frac{1}{4} - ig_i^2 \frac{c_i^{\text{hid}}}{16\pi^2} \mathcal{A} \right) \mathcal{W}_i^2 - \int d^2\theta \frac{m_{\lambda_i}}{2F} X \mathcal{R}^{-2} \mathcal{W}_i^2 + \text{c.c.} \quad r_X = 2 \quad r_{\mathcal{W}} = 1$$

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between a and MSSM pseudoscalar A

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$$-i a R_H \left(c_\beta^2 \frac{m_u}{f_a} \bar{u} \gamma_5 u + s_\beta^2 \frac{m_d}{f_a} \bar{d} \gamma_5 d + s_\beta^2 \frac{m_\ell}{f_a} \bar{\ell} \gamma_5 \ell \right)$$

$$\frac{\delta^2}{v} (\partial_\mu a)^2 h$$

(Here and in the following $m_A, m_{\tilde{f}} \gg m_a$)

A strongly coupled “UV” completion

Very low energy SUSY breaking F

motivated by:

Naturalness + Higgs mass [Gherghetta Pomarol 1107.4697](#)

+ LHC exclusions [Buckley et al. 1610.08059](#)

Gravitino cosmology [Ibe Yanagida 1608.01610](#)

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needs a strongly coupled sector

so that $m_{\lambda_i} \sim \frac{g_i^2}{g_*^2} m_*$ OK with LHC bounds

m_* mass gap of the hidden sector
(e.g. mass of messengers in gauge mediation)

$g_* > 1$ coupling between hidden sector states

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$$m_*$$

$$m_{\text{soft}} \approx \frac{g^2}{g_*^2} m_*$$

$$m_a \approx \sqrt{\epsilon_R} m_*$$

$$m_G = \frac{F}{\sqrt{3}M_{Pl}}$$

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SUSY Naive Dimensional Analysis

$$M_{\text{SUSY}} \sim m_* \sim g_* f \quad f_a \sim f$$

$$F \sim g_* f^2 \quad w_R \sim g_* f^3$$

inspired by
[Cohen et al. 1997](#)
[Luty 1998](#)
[Giudice+ 2007](#)

$a \rightarrow GG$ saturates the upper bound

✓ Motivation

✓ Setup

Phenomenology

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Pheno overview

We will focus on

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just for definiteness

otherwise ~ not a pseudo Goldstone

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i.e. beyond reach of LHC14

$$m_{\tilde{B}} \simeq 500 \text{ GeV} \quad m_{\tilde{W}} \simeq 1 \text{ TeV} \quad m_{\tilde{g}} \simeq 2.5 \text{ TeV} \quad \text{would correspond to} \quad m_* \simeq 22 \text{ TeV} \left(\frac{g_*}{3}\right)^2$$

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Value of their masses do not matter for R-axion pheno (as long as they are decoupled)

Example:

$$\Gamma_{a \rightarrow gg} = k_{gg} \frac{\alpha_s^2}{16\pi^3} \frac{m_a^3}{f_a^2} \left| c_3^{\text{hid}} + 3 \text{Loop}_{\tilde{g}} - \sum_{q=t,b,\dots} r_q \text{Loop}_q \right|^2$$

$c_i^{\text{hid}} = -N_{\text{mess}}$
 $\text{Loop}_{\tilde{g}} \xrightarrow{m_{\tilde{g}} \gg m_a/2} 1$

(messengers in $5 + \bar{5}$ of $SU(5)$ with zero R-charge)
 i.e. an anomaly

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Could the R-axion be the first SUSY sign at colliders?

Value of their production cross-section for R-axion pheno (as long as they are decoupled)

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Resonant production @ LHC

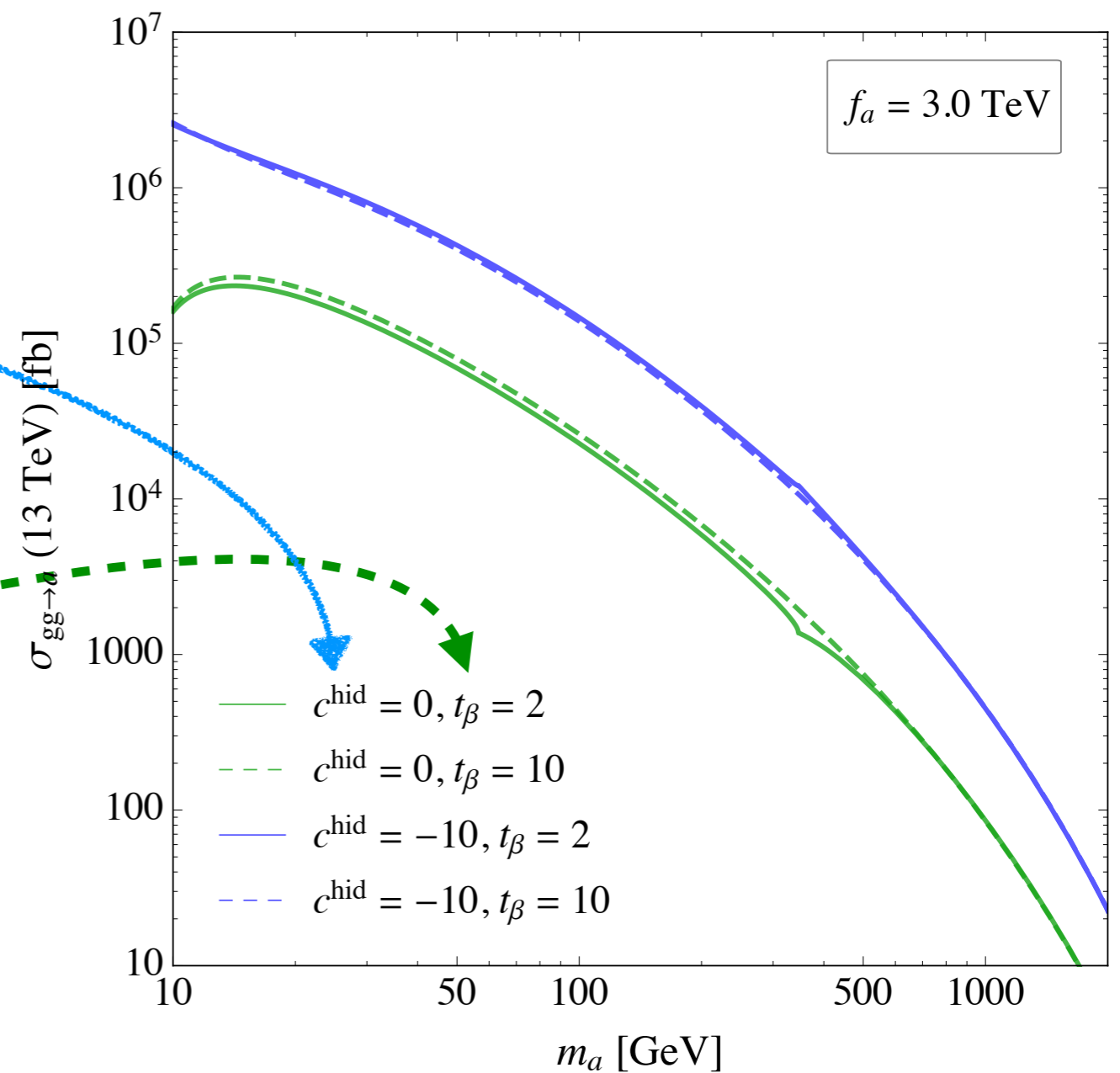
3 contributions to $a G_{\mu\nu}^a \tilde{G}^{a,\mu\nu}$ vertex:

UV anomaly

gluino loop

quark loops

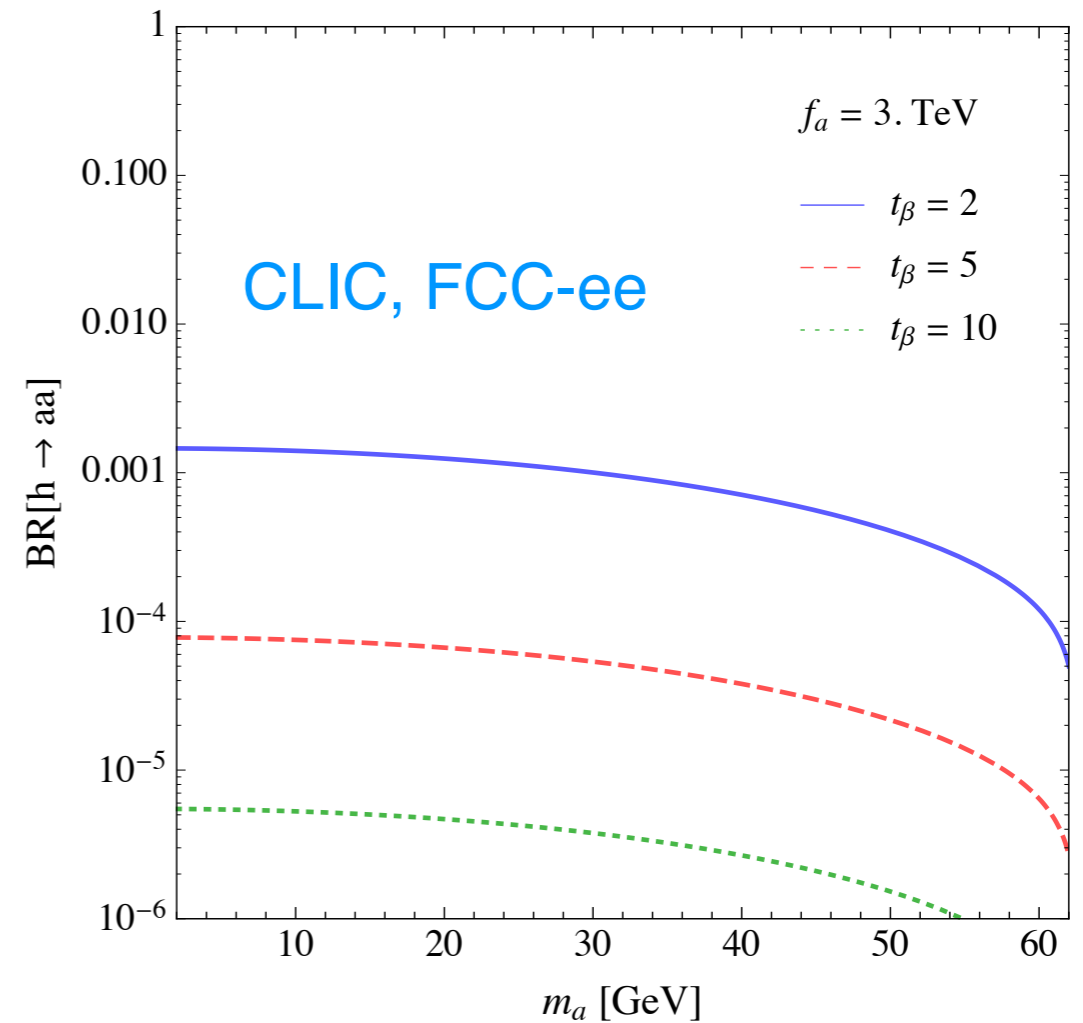
$$r_H c_\beta^2 \frac{m_t}{f_a} a \bar{t} \gamma_5 t$$



a from decays of h , Υ and B

$$\mathcal{L}_{ha^2} = \frac{\delta^2}{v} (\partial_\mu a)^2 h$$

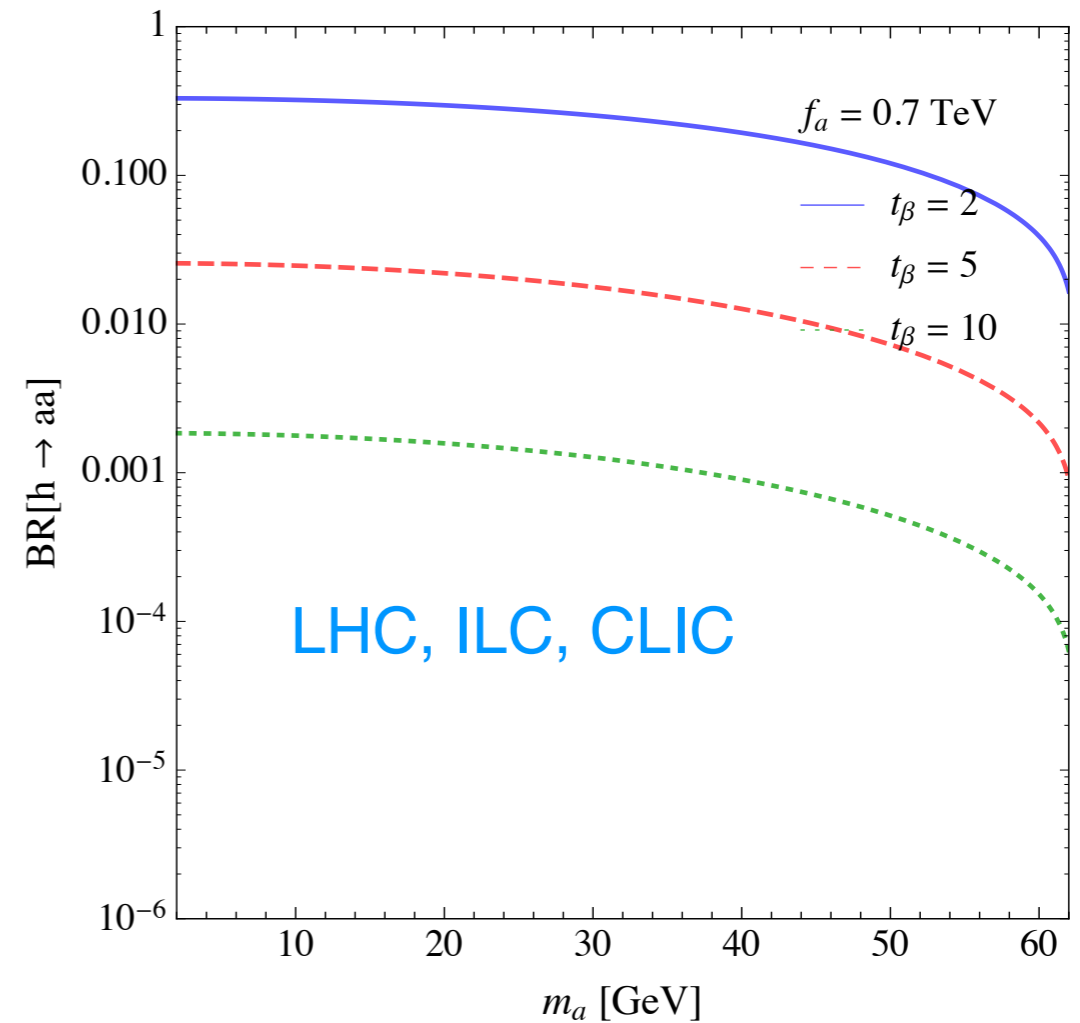
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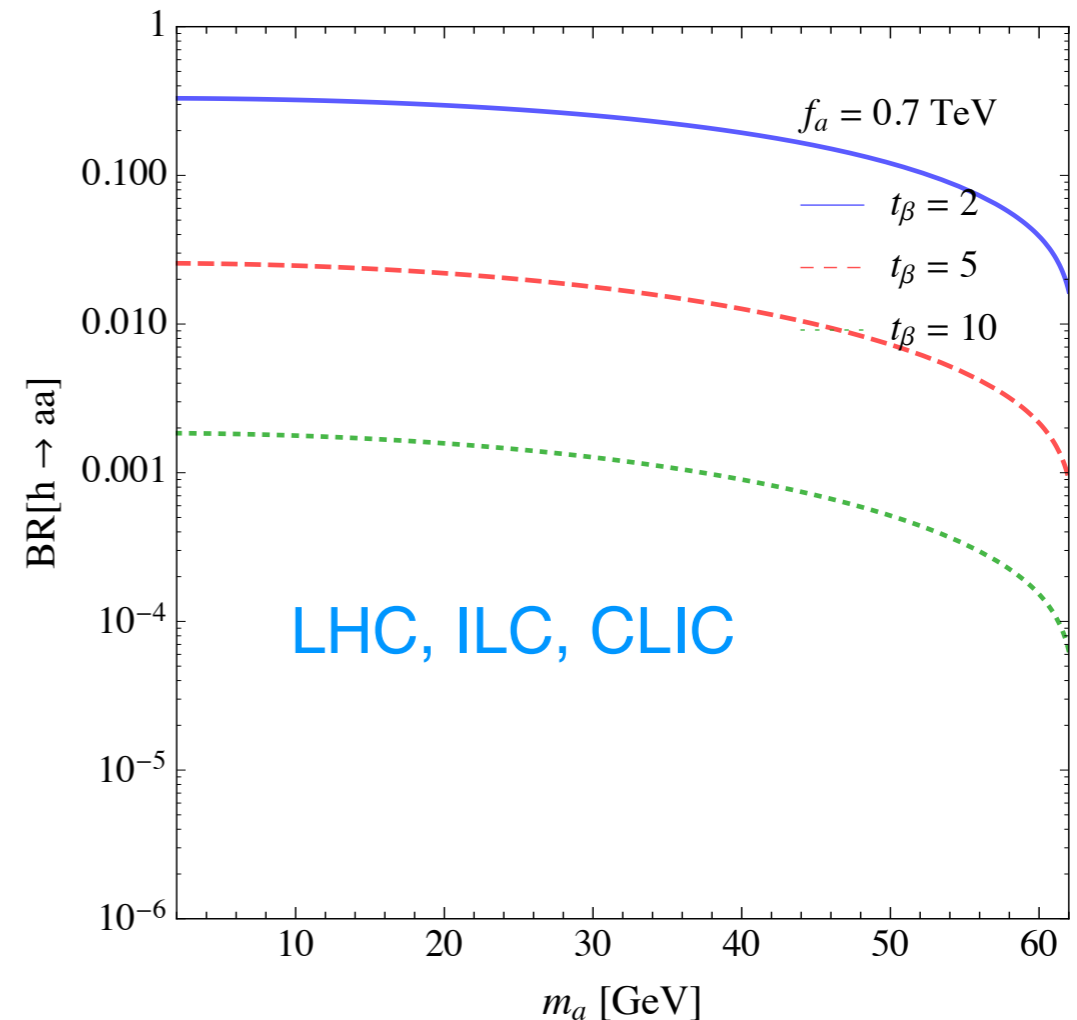
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$$\text{BR}_{\Upsilon \rightarrow \gamma a} \simeq 3 - 5 \times 10^{-5} \left(\frac{\text{TeV}}{f_a} \right)^2$$

since Wilczek PRL39 (1977)

experiments: **BABAR**
Belle-II

$$\text{BR}_{B \rightarrow K a, K^* a} \simeq 3 - 5 \times 10^{-4} \left(\frac{\text{TeV}}{f_a} \right)^2$$

see Hall Wise 1981, Freytsis Ligeti Thaler 0911.5355

LHCb
Belle, Belle-II

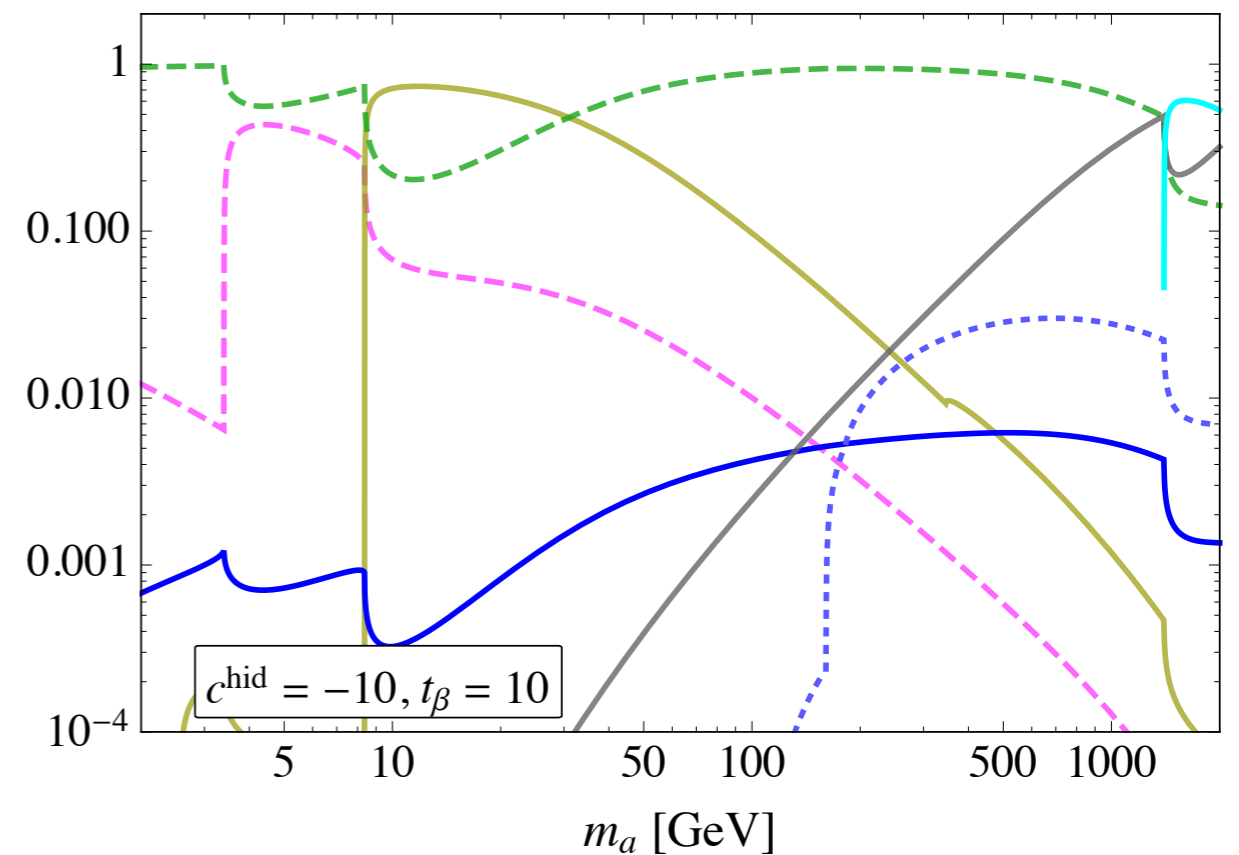
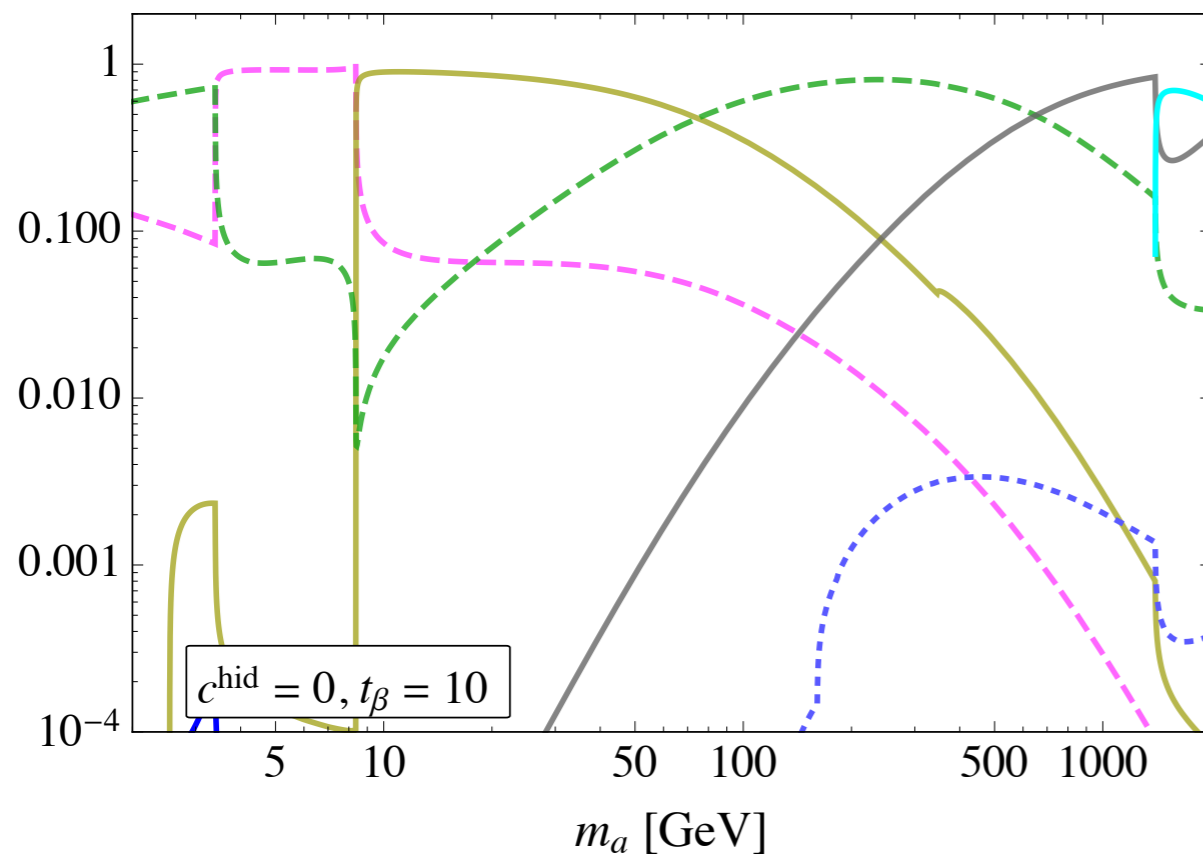
R axion branching ratios

Both plots: $t_\beta = 10$

No anomaly

Large anomalies

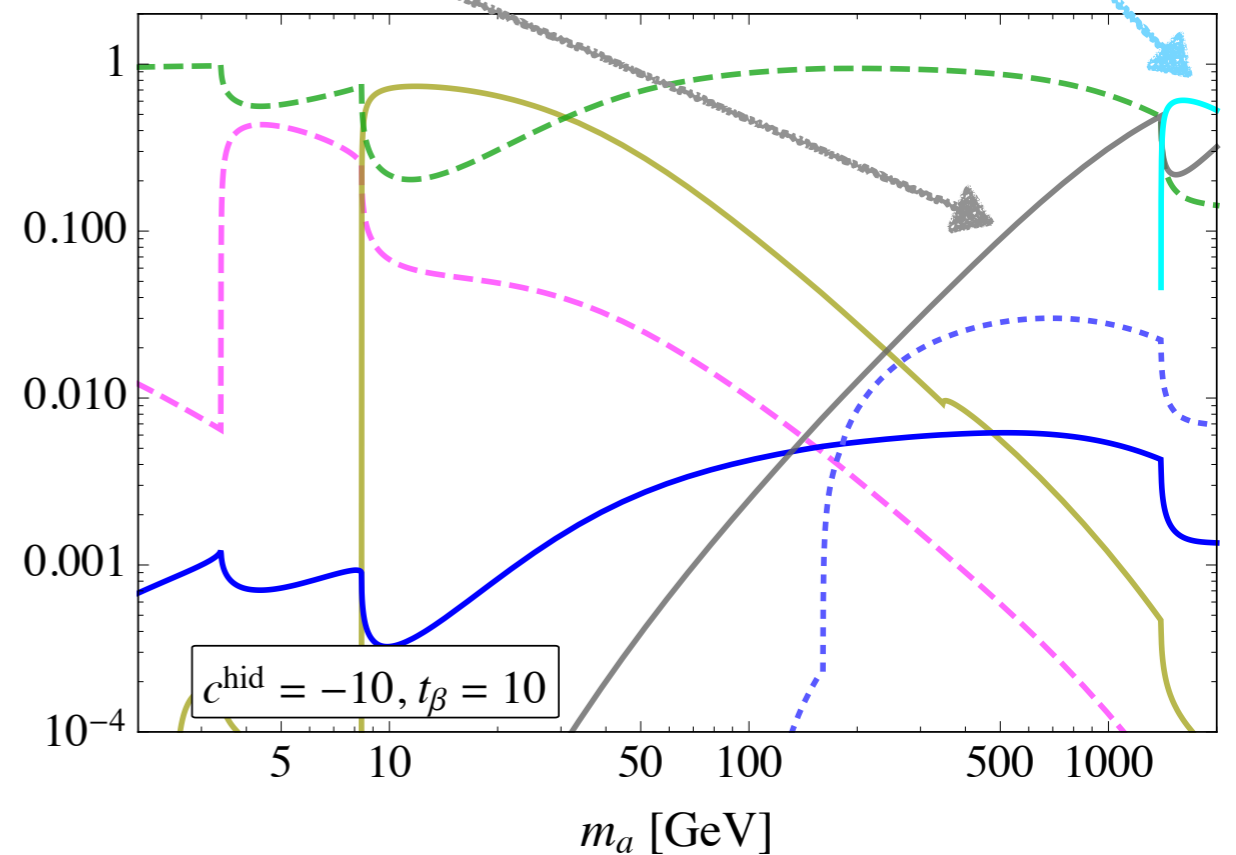
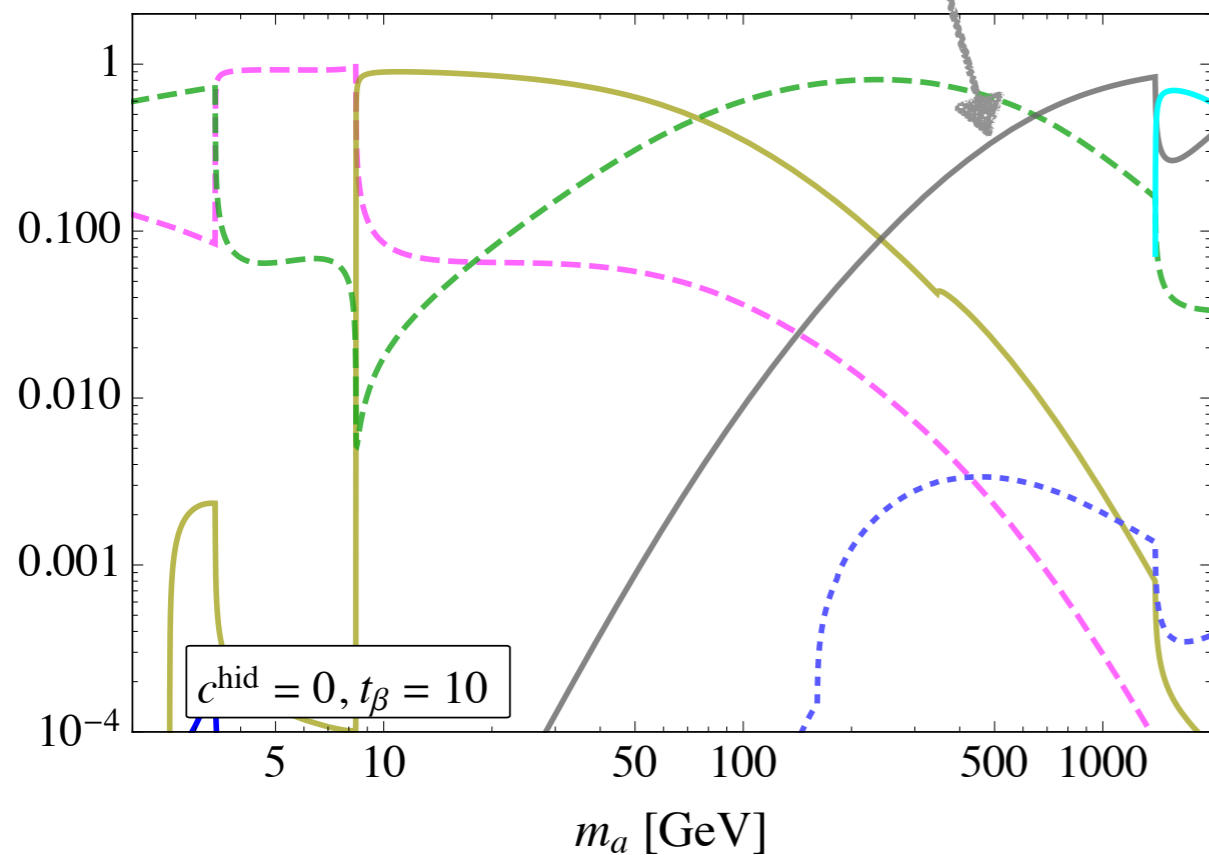
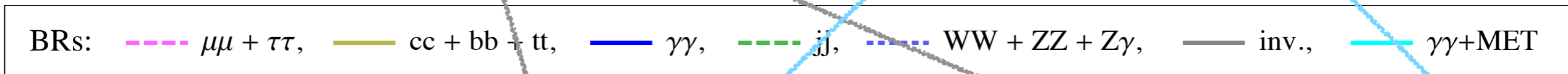
BRs: $\mu\mu + \tau\tau$, $cc + bb + tt$, $\gamma\gamma$, jj , $WW + ZZ + Z\gamma$, $inv.$, $\gamma\gamma+MET$



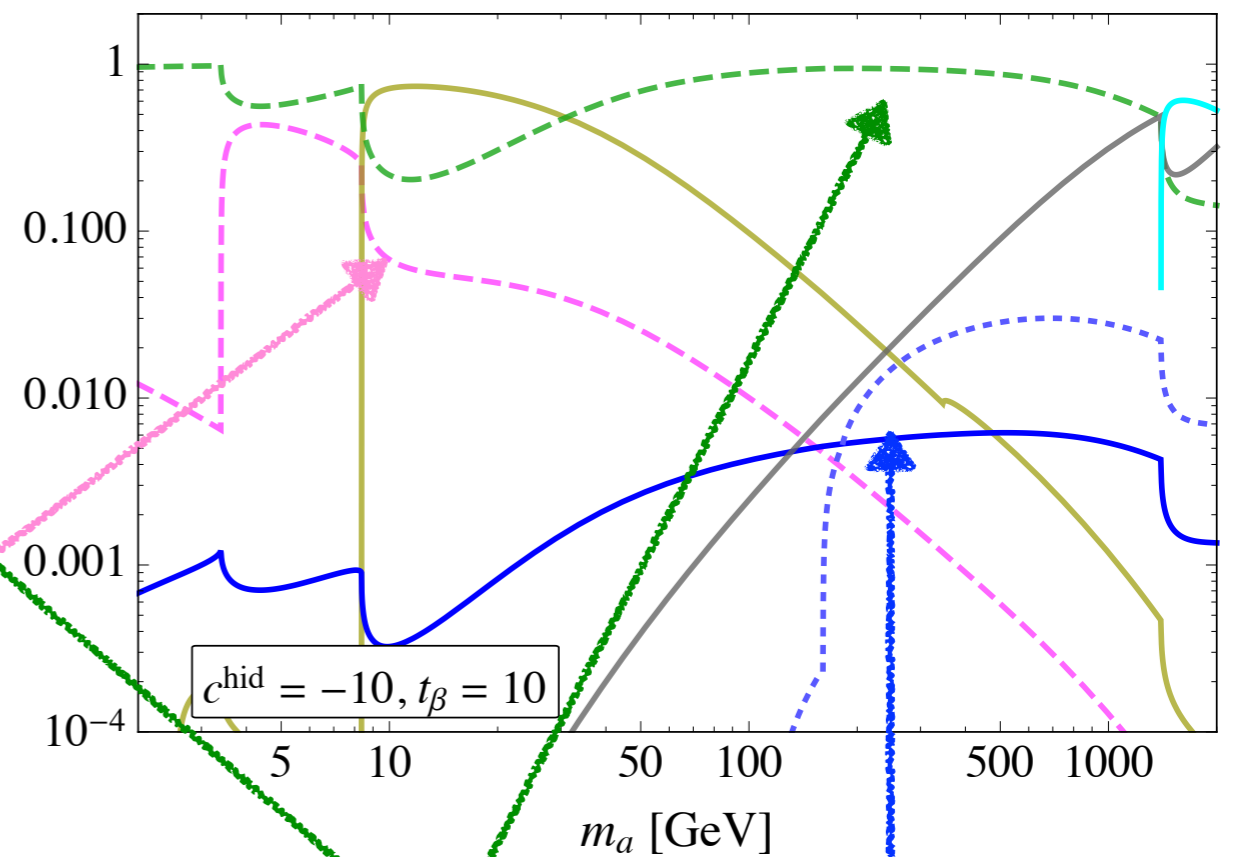
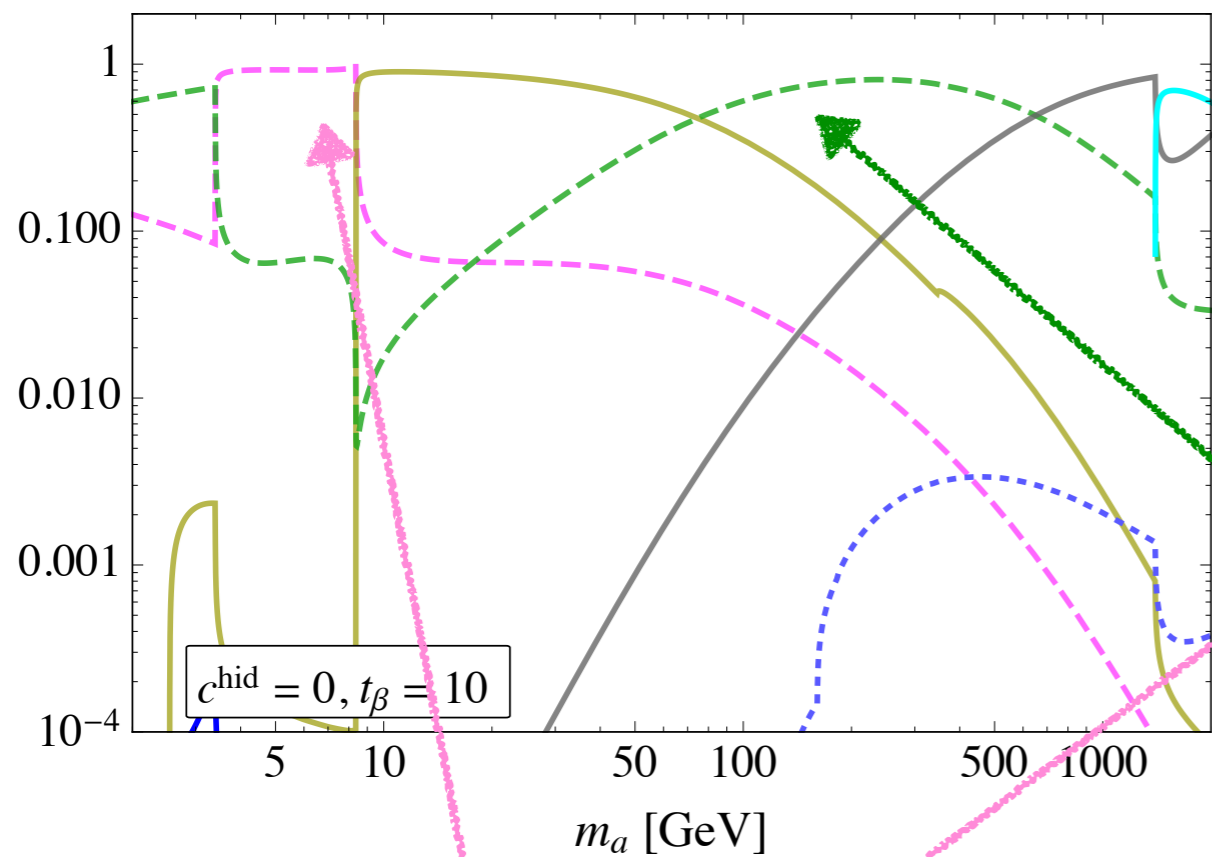
R axion branching ratios

MET from decays into gravitinos

more model dependent
 $a \rightarrow \tilde{B}\tilde{B} \rightarrow \gamma\gamma + \text{MET}$



R axion branching ratios



$\mu\mu$ $\tau\tau$ $b\bar{b}$

dijet diphoton

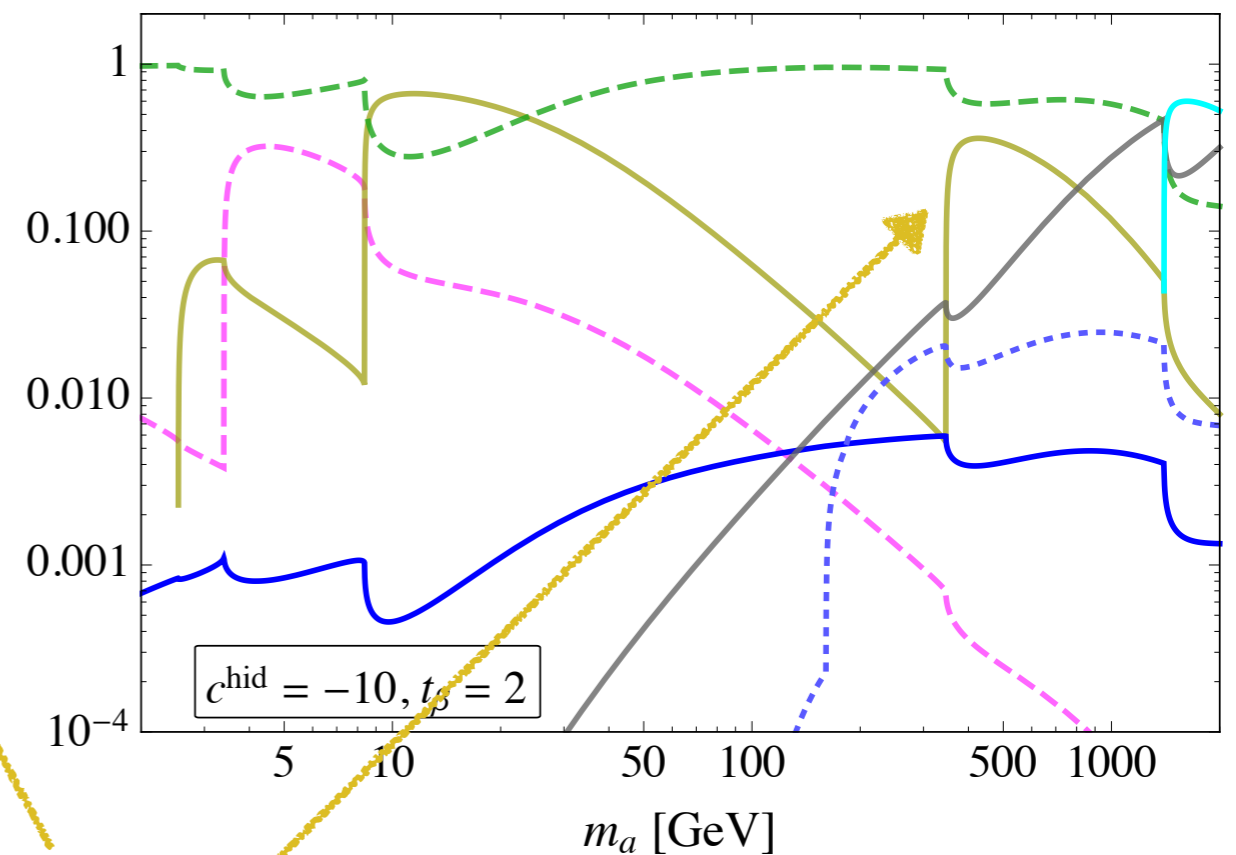
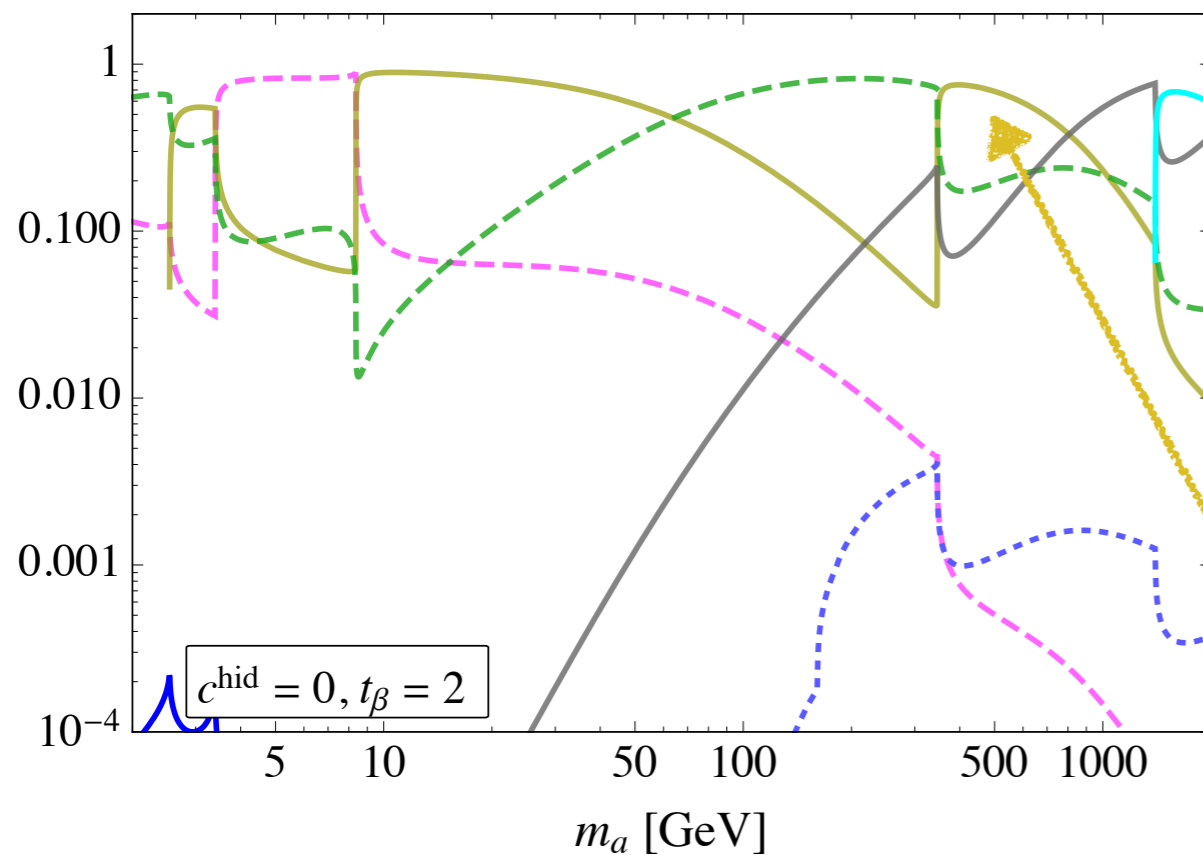
R axion branching ratios

Both plots: $t_\beta = 2$

No anomaly

Large anomalies

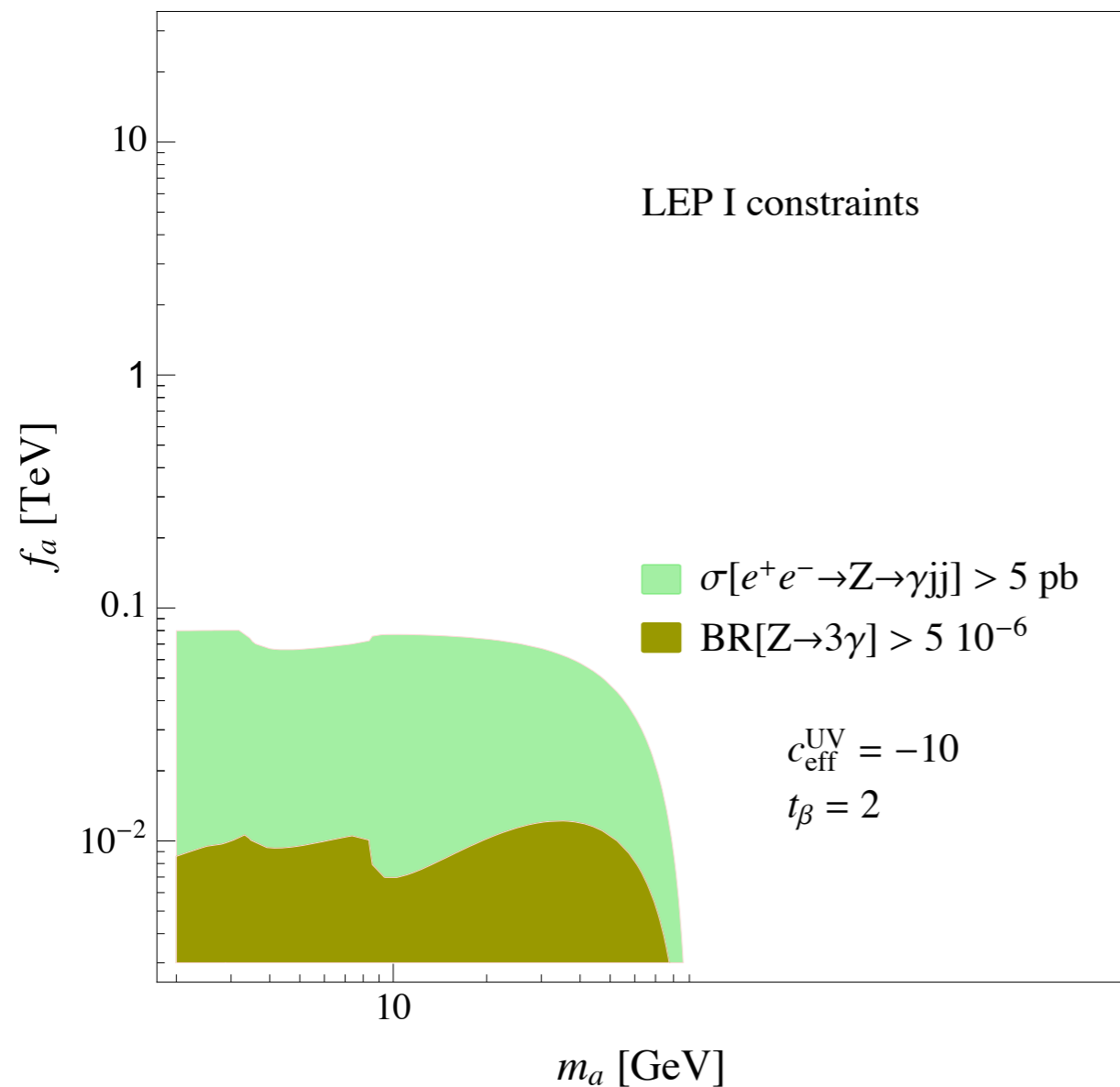
BRs: $\mu\mu + \tau\tau$, $cc + bb + tt$, $\gamma\gamma$, jj , $WW + ZZ + Z\gamma$, $inv.$, $\gamma\gamma+MET$



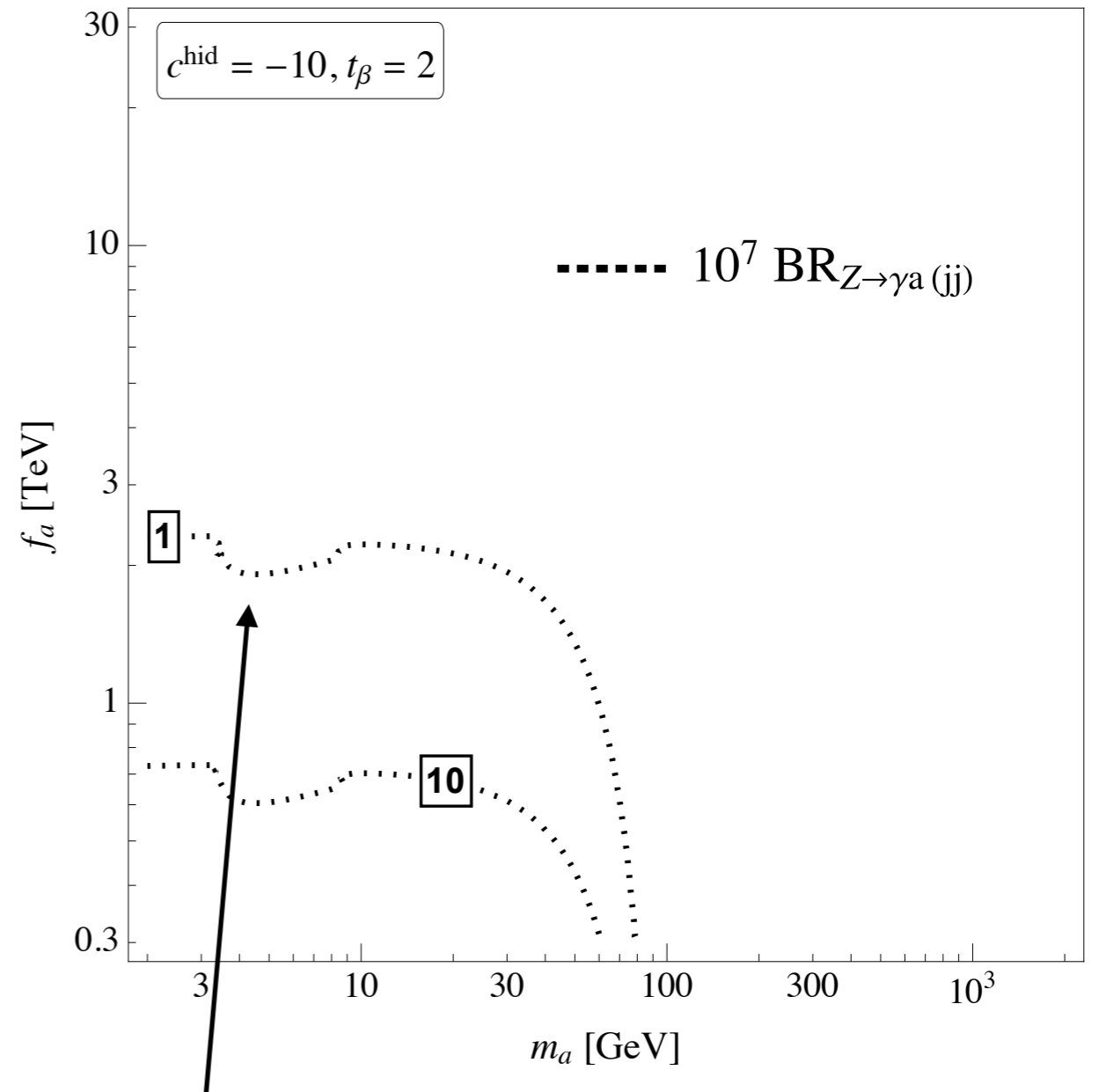
$t\bar{t}$ becomes more important!

LEP & future lepton colliders

$$\mathcal{L}_{aZ\gamma} = \frac{\alpha_w}{2\pi} c_{Z\gamma} \frac{a}{f_a} Z_{\mu\nu} \tilde{F}^{\mu\nu}$$



LEP: not really relevant

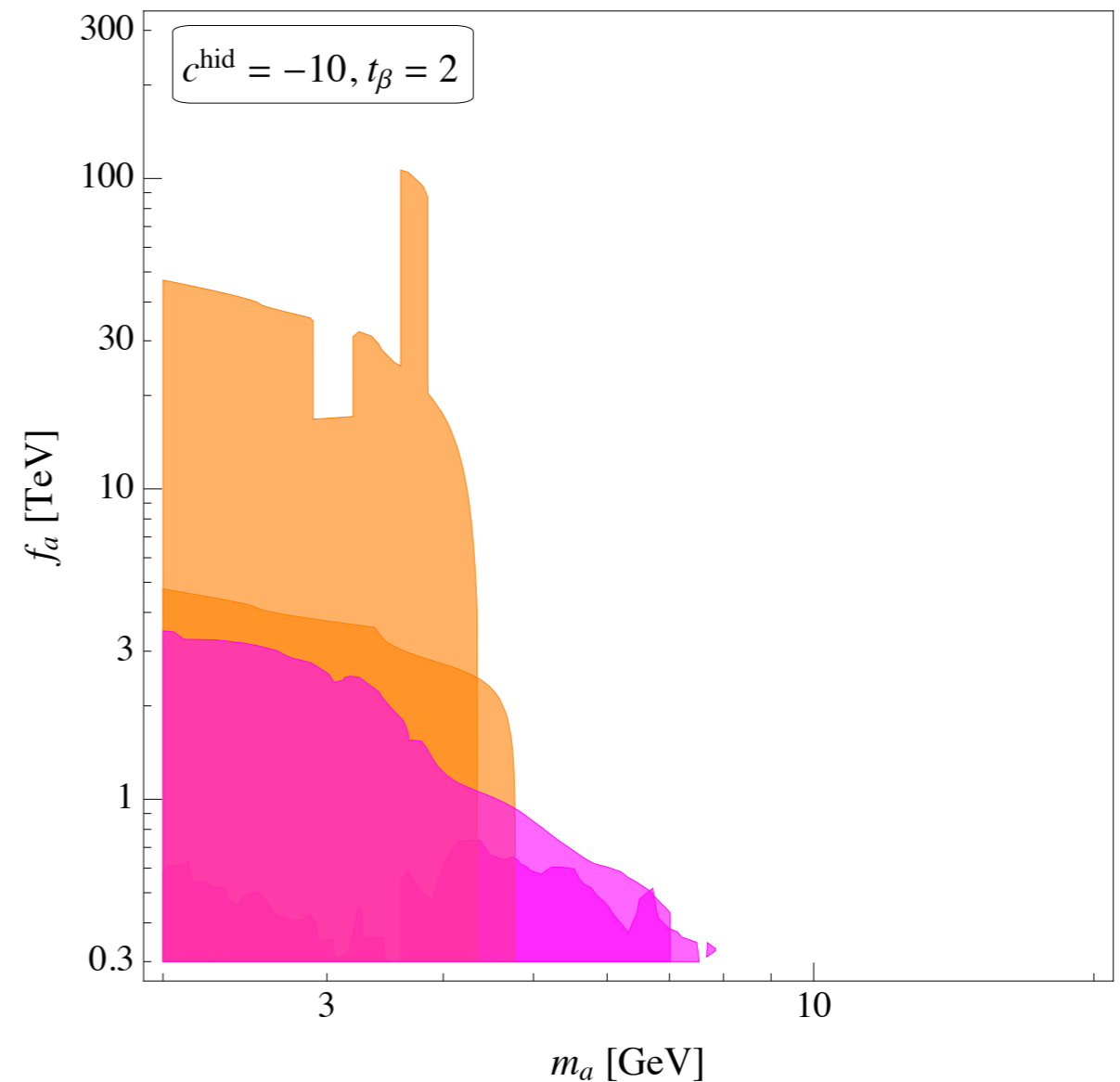
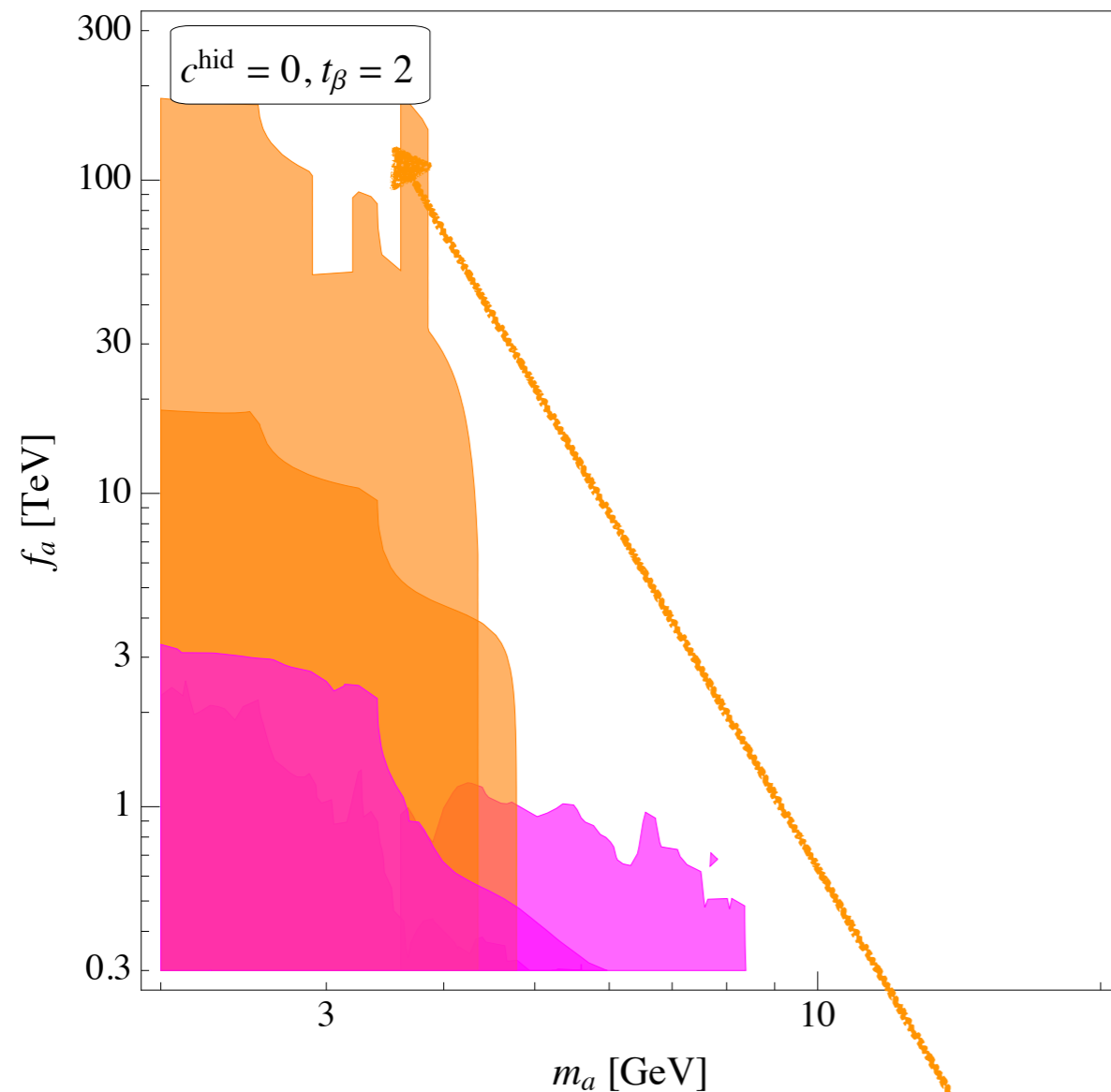


Reach of **FCC-ee** ?
(from naive luminosity rescaling of LEP I search)

Decays of B and Upsilon

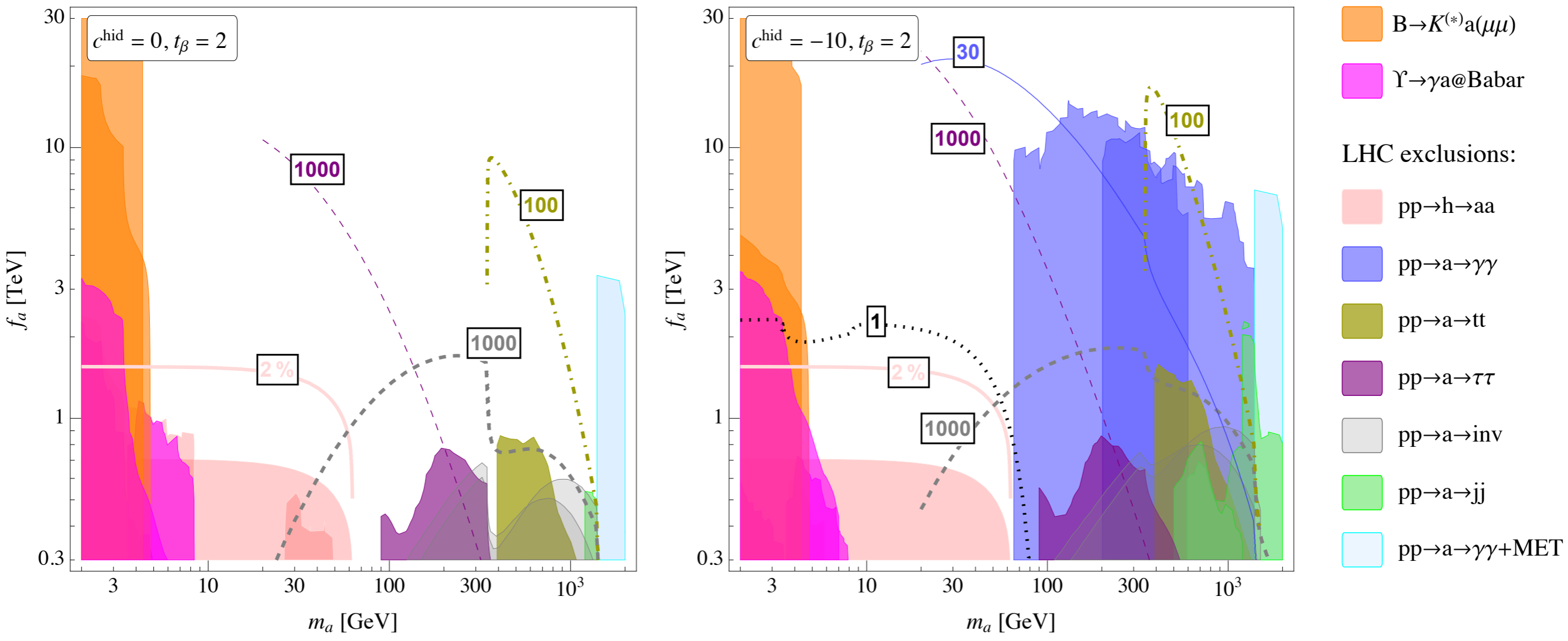
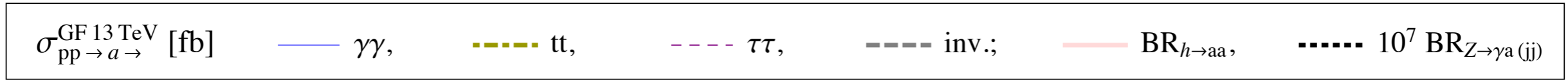
 $B \rightarrow K^{(*)} a(\mu\mu)$ LHCb 1508.04094 (+ Belle)

 $\Upsilon \rightarrow \gamma a @ \text{Babar}$ BABAR 1210.0287 (muons), 1210.5669 (taus), 1108.3549 (hadrons)

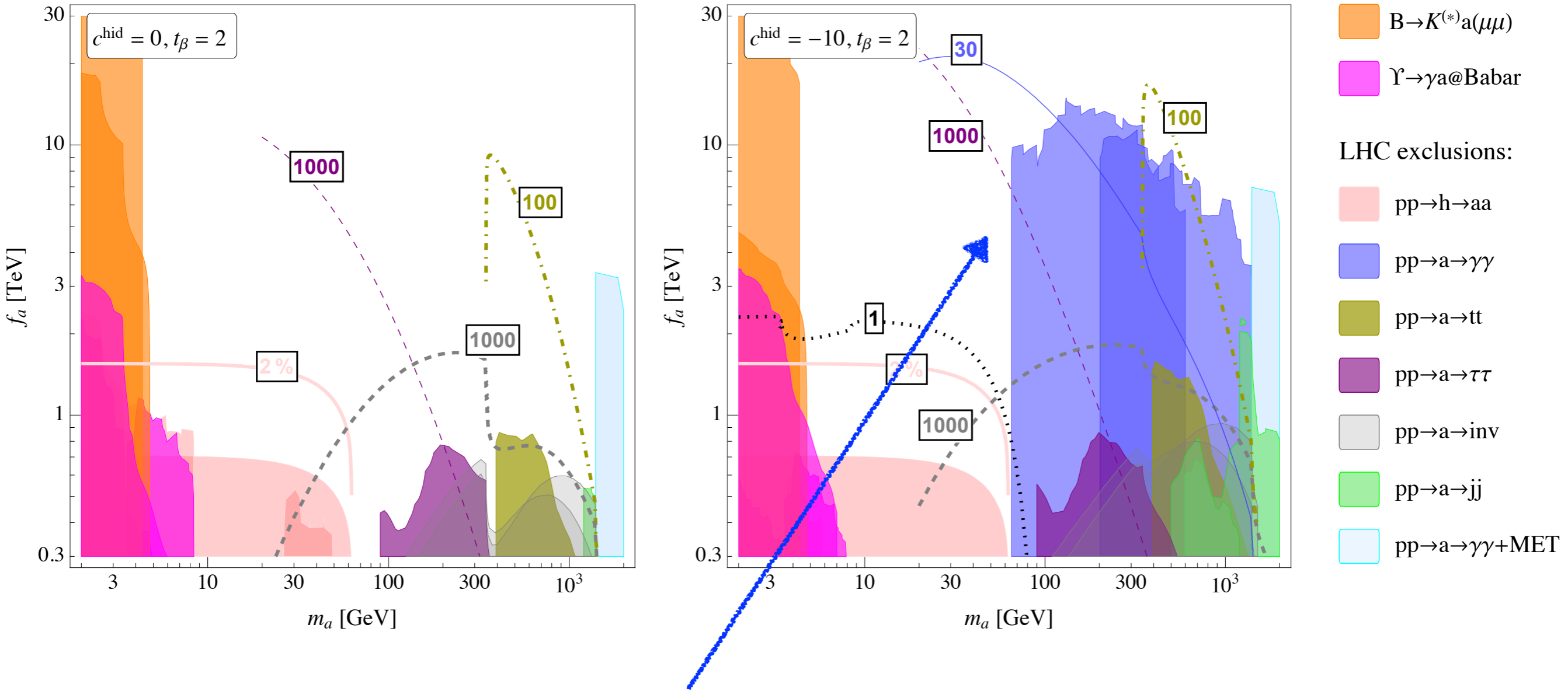
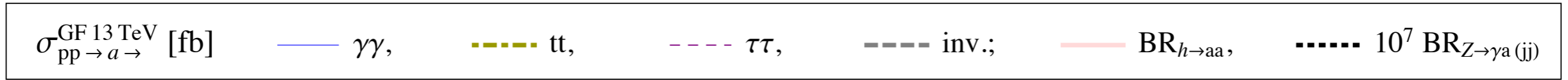


LHCb sensitive to beyond-100-TeV flavour-preserving SUSY!!

Pheno Summary



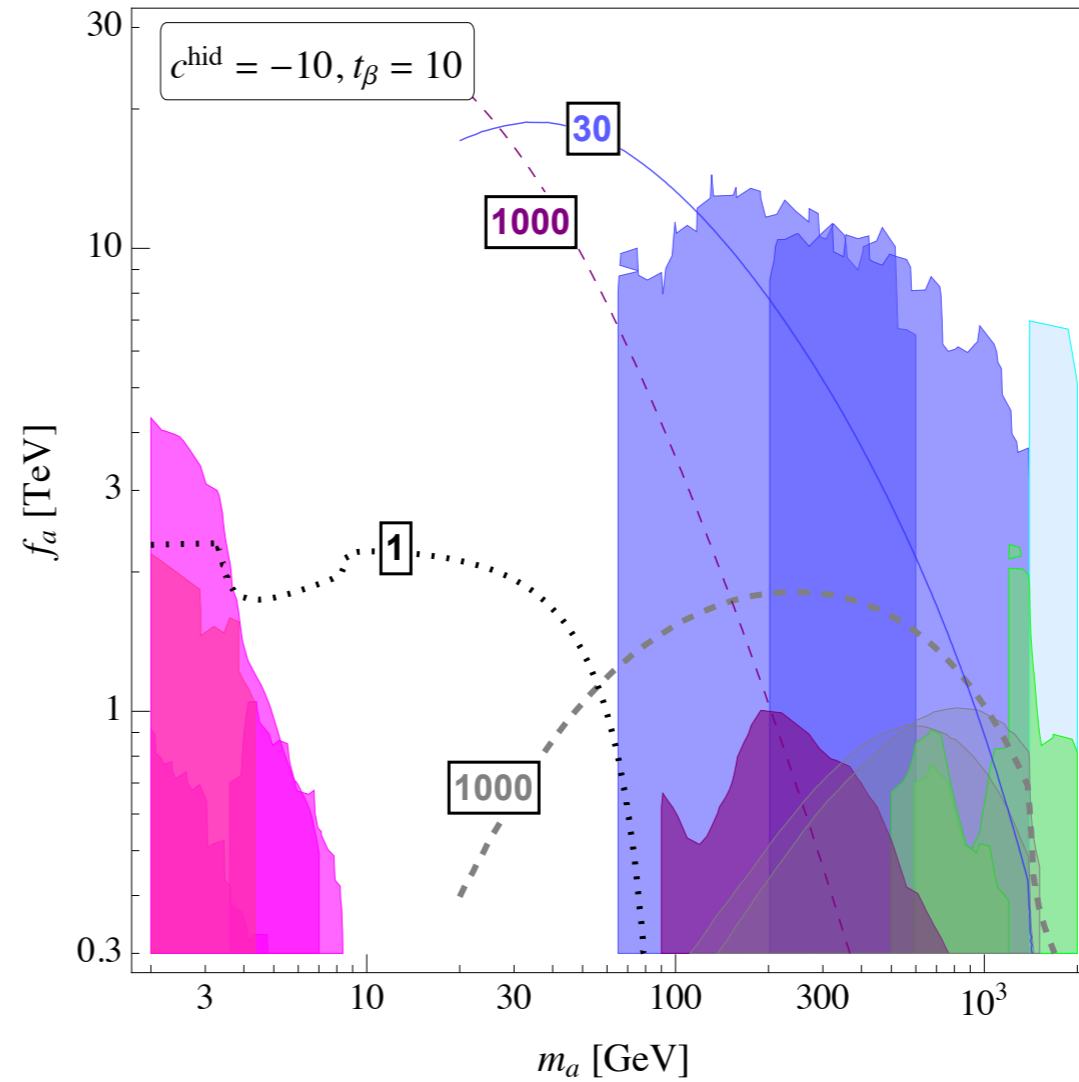
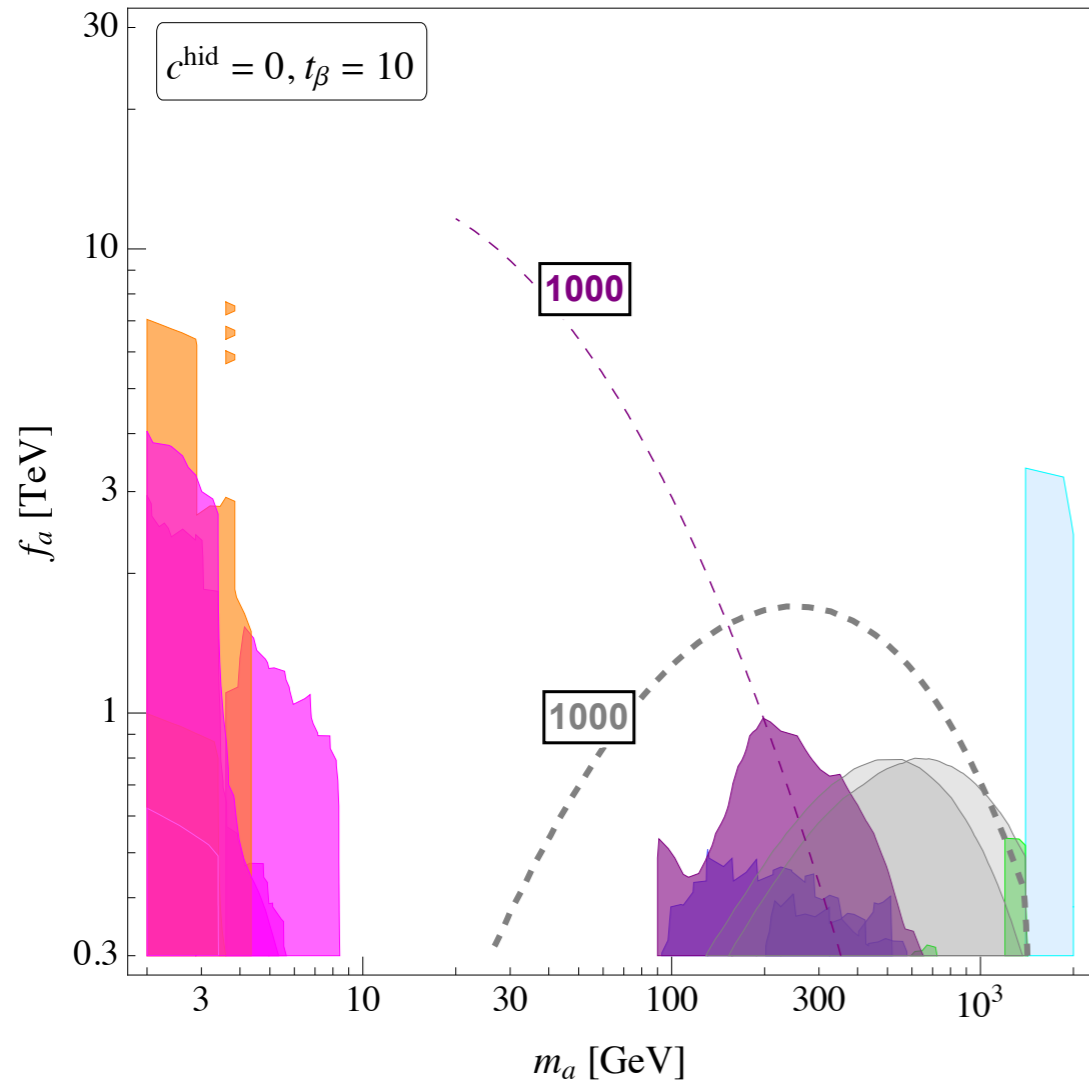
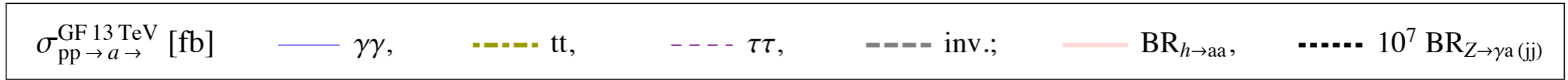
Pheno Summary



Why not diphoton below 65 GeV?

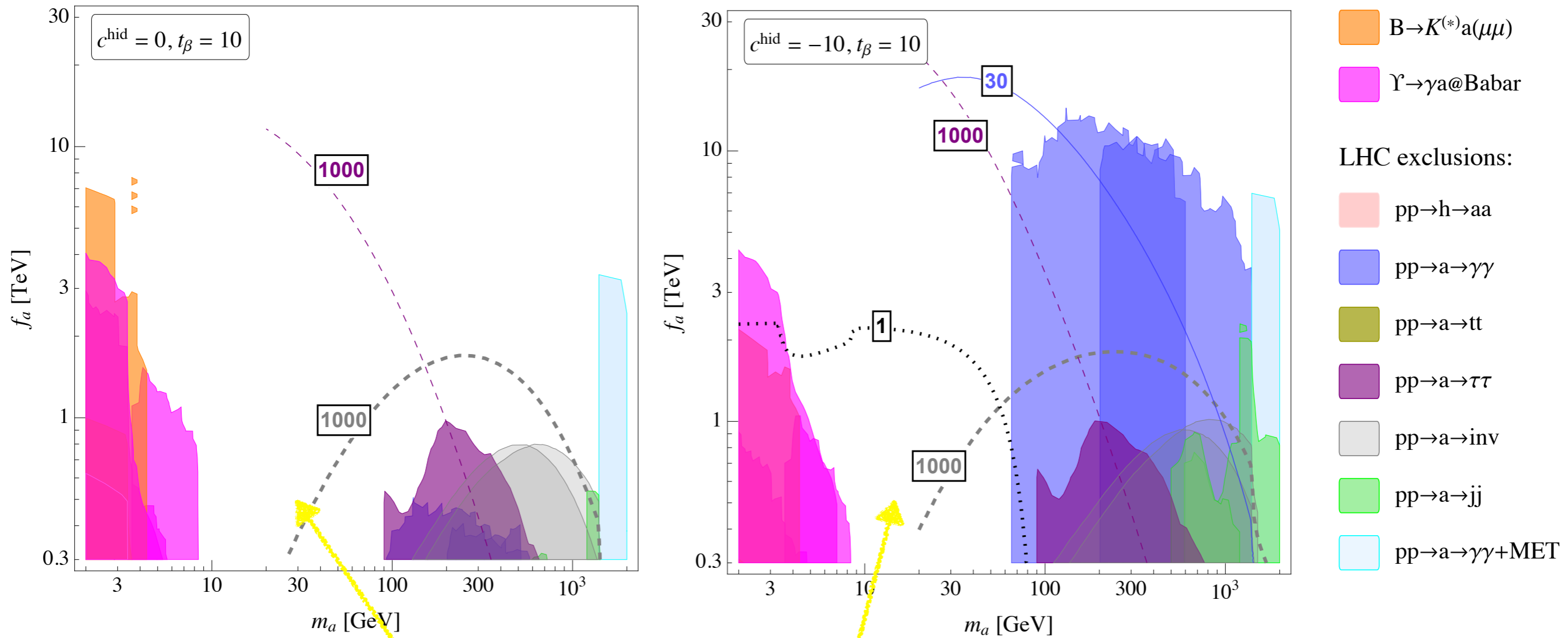
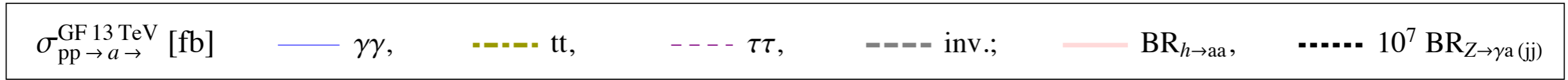
Pheno Summary

$$t_\beta = 10$$



Pheno Summary

$$t_\beta = 10$$



How to explore $10 \text{ GeV} \lesssim m_a \lesssim 60 \text{ GeV}$?

Crucial for any ALP that couples dominantly to gauge bosons!

Take-home messages

Any other sign of SUSY to be looked for at colliders?

beyond gluino, stop, EWinos, extra Higgses,...

Spontaneously broken R-symmetry is quite generic

provides a naturally light state, the “**R-axion**” (whose pheno had been overlooked)

Bellazzini Mariotti Redigolo FS Serra, 1702.02152

could be the first sign of SUSY at colliders

gives access to properties of the SUSY-breaking sector

opens new pheno & model-building avenues

Conclusions and outlook

How to know it is R-axion and not another scalar?

Hint: other peculiar decays like MET; pattern of signals

Observe other SUSY & correlate parameters

light diphoton resonances *work in progress...*

Rich pheno adds motivation to

LHC (heavy & light) scalar searches!

flavour factories (e.g. look for bump in B decays)

lepton colliders program

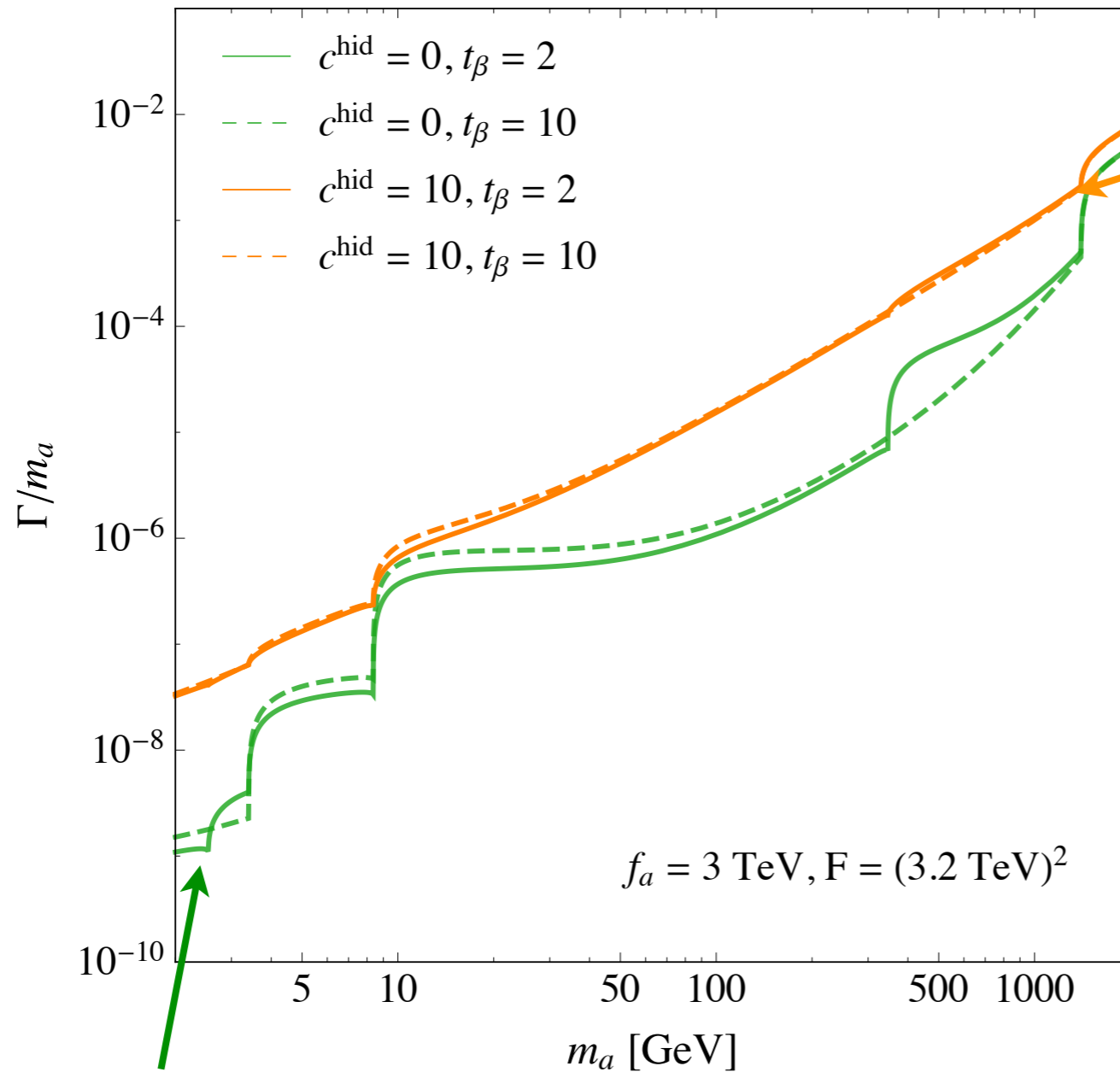
Future more model building (e.g. split SUSY? models for m_a ?)

pheno other than colliders (e.g. cosmo?)

....

Back-up

R axion total width



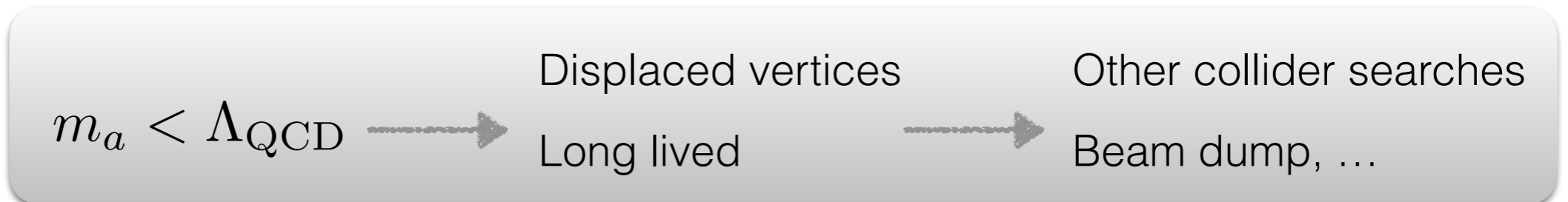
$$\Gamma_a/m_a < 10^{-3}$$

so interference with SM in $t\bar{t}$ should not give problems...

see e.g. [Craig et al. 1504.04630](#)

(unlike usual targets for $t\bar{t}$, like MSSM Higgses)

$$c\tau \approx \mu m$$



More on R-axion production

$$\sigma_{pp \rightarrow a} = \frac{K}{m_a s} C_{gg} \Gamma_{a \rightarrow gg}^{\text{tree}} \quad K \simeq 2.4 \quad \text{see e.g. Ahmed et al. 1606.00837}$$

$$C_{gg} = \frac{\pi^2}{8} \int_{m_a^2/s}^1 \frac{dx}{x} f_g(x) f_g\left(\frac{m_a^2}{s x}\right) \quad \text{gluon pdfs } f_g \text{ from mstw2008}$$

$$\mathcal{L}_{aZ\gamma} = \frac{\alpha_w}{2\pi} c_{Z\gamma} \frac{a}{f_a} Z_{\mu\nu} \tilde{F}^{\mu\nu} \quad \text{too small for LEP, not for FCC-ee!}$$

$$\text{BR}(Z \rightarrow a\gamma) \approx 3 \cdot 10^{-7} \left(\frac{\text{TeV}}{f_a}\right)^2 \quad \text{for } N_{\text{mess}} \approx 10$$

All single productions scale as

$$\sim 1/f_a^2$$

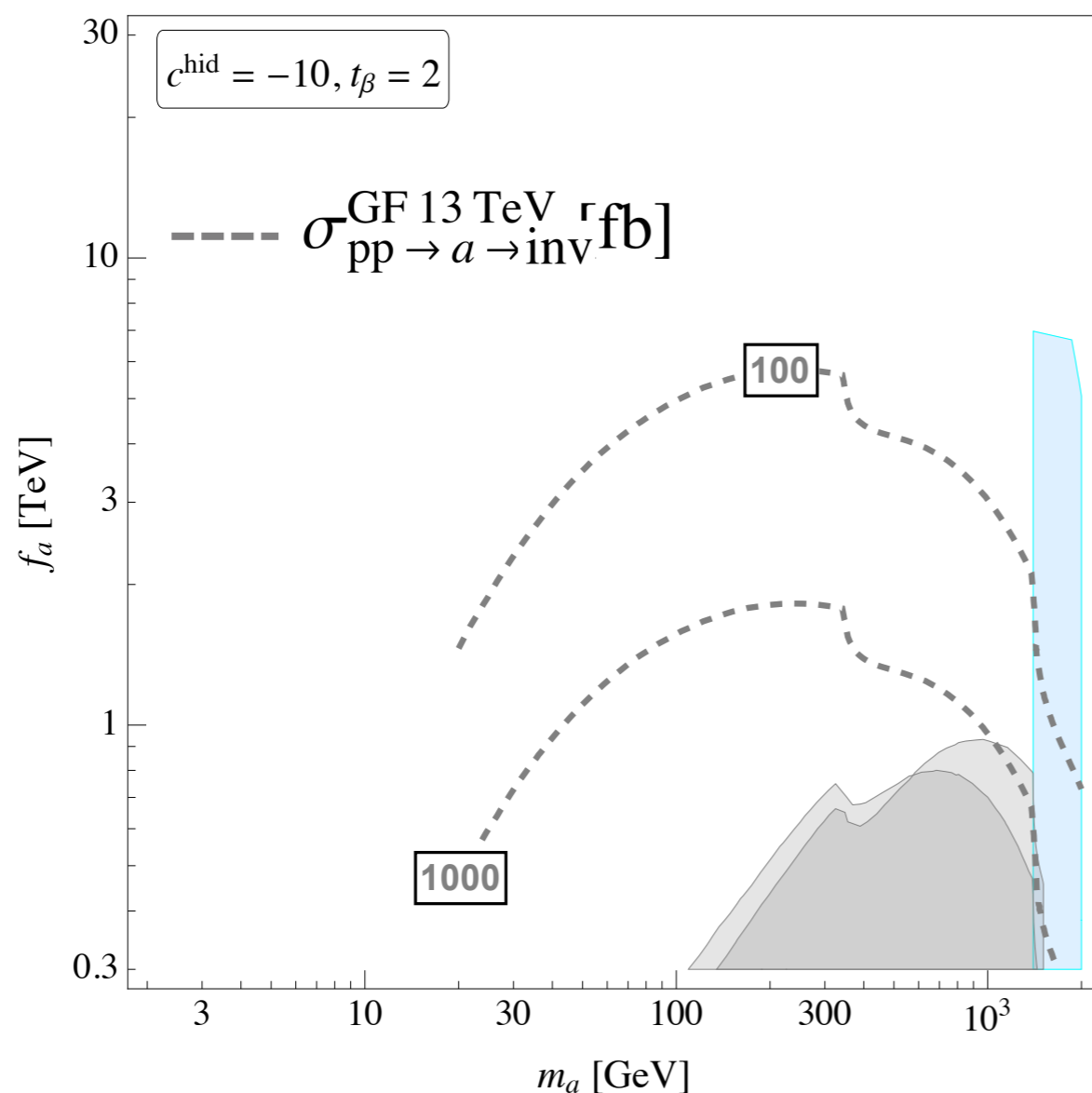
double productions (here only $h \rightarrow aa$) as

$$\sim 1/f_a^4$$

LHC: MET + monojet, MET + diphoton

	$p_T > 250$	$p_T > 500$	$p_T > 700$	
σ_{95} 8 TeV	90 fb	7.2 fb	3.4 fb	ATLAS 1502.01518
σ_{95} 13 TeV [3.2 fb^{-1}]	553 fb	61 fb	19 fb	ATLAS 1604.07773

Rough procedure: gluon gluon resonance w/ and w/o extra jet simulated with Madgraph
 ratio used to rescale $\sigma_{pp \rightarrow a \rightarrow GG}$



$a \rightarrow \tilde{B}\tilde{B} \rightarrow \gamma\gamma + \text{MET}$
 excludes signals
 down to **0.3 fb** @LHC8

These signals have little dependence on anomalies, and in particular on t_β, r_H

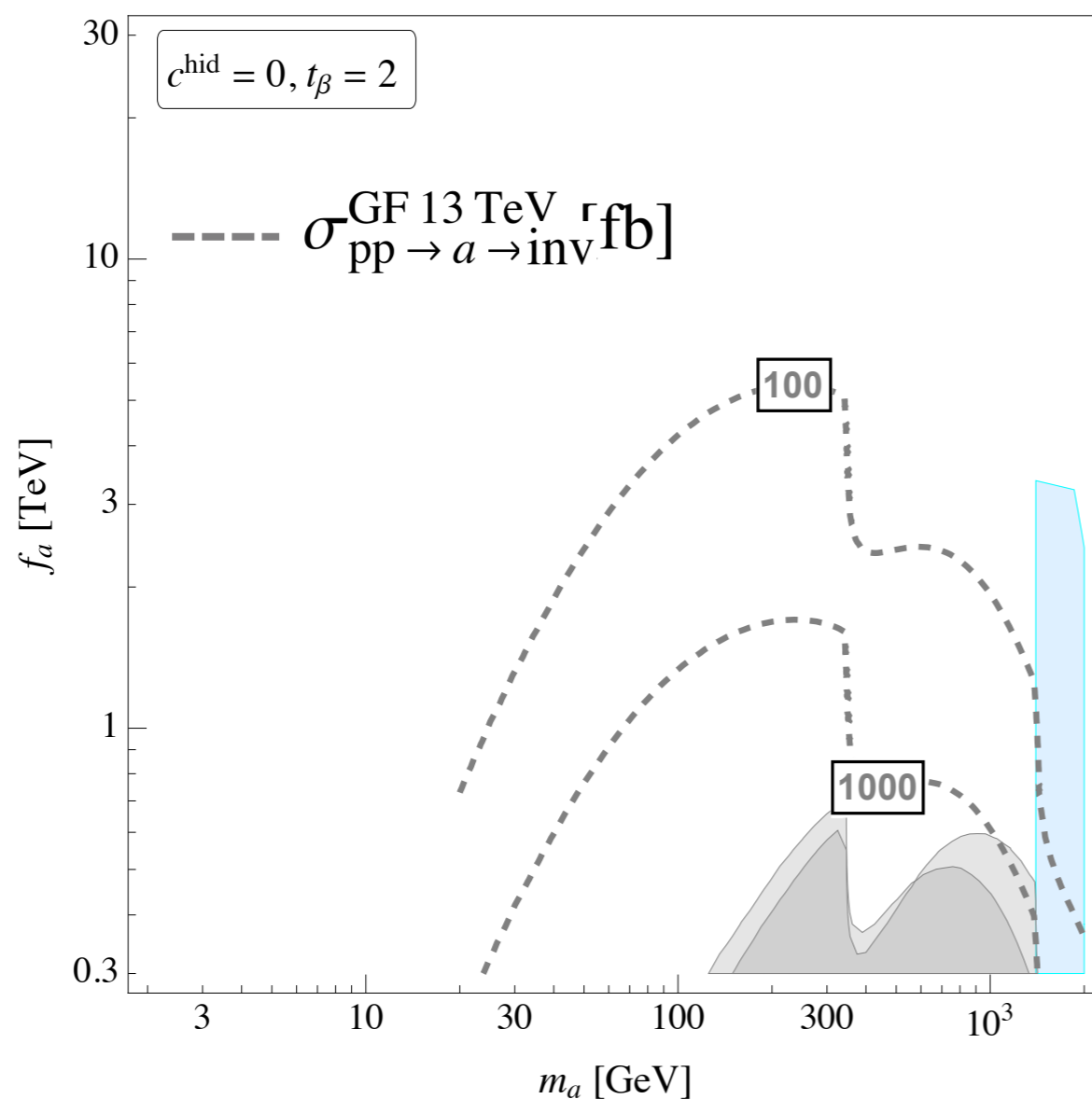
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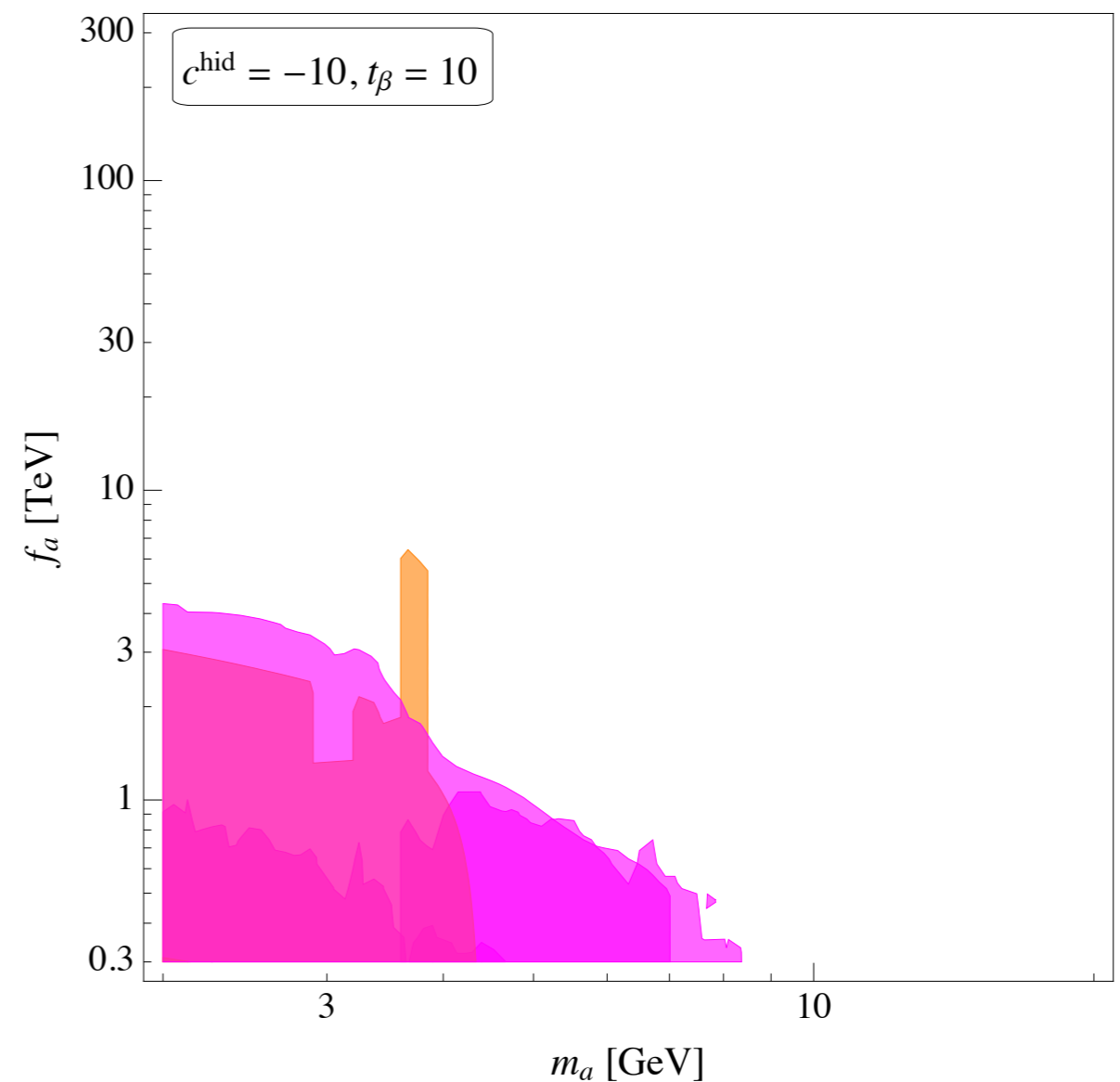
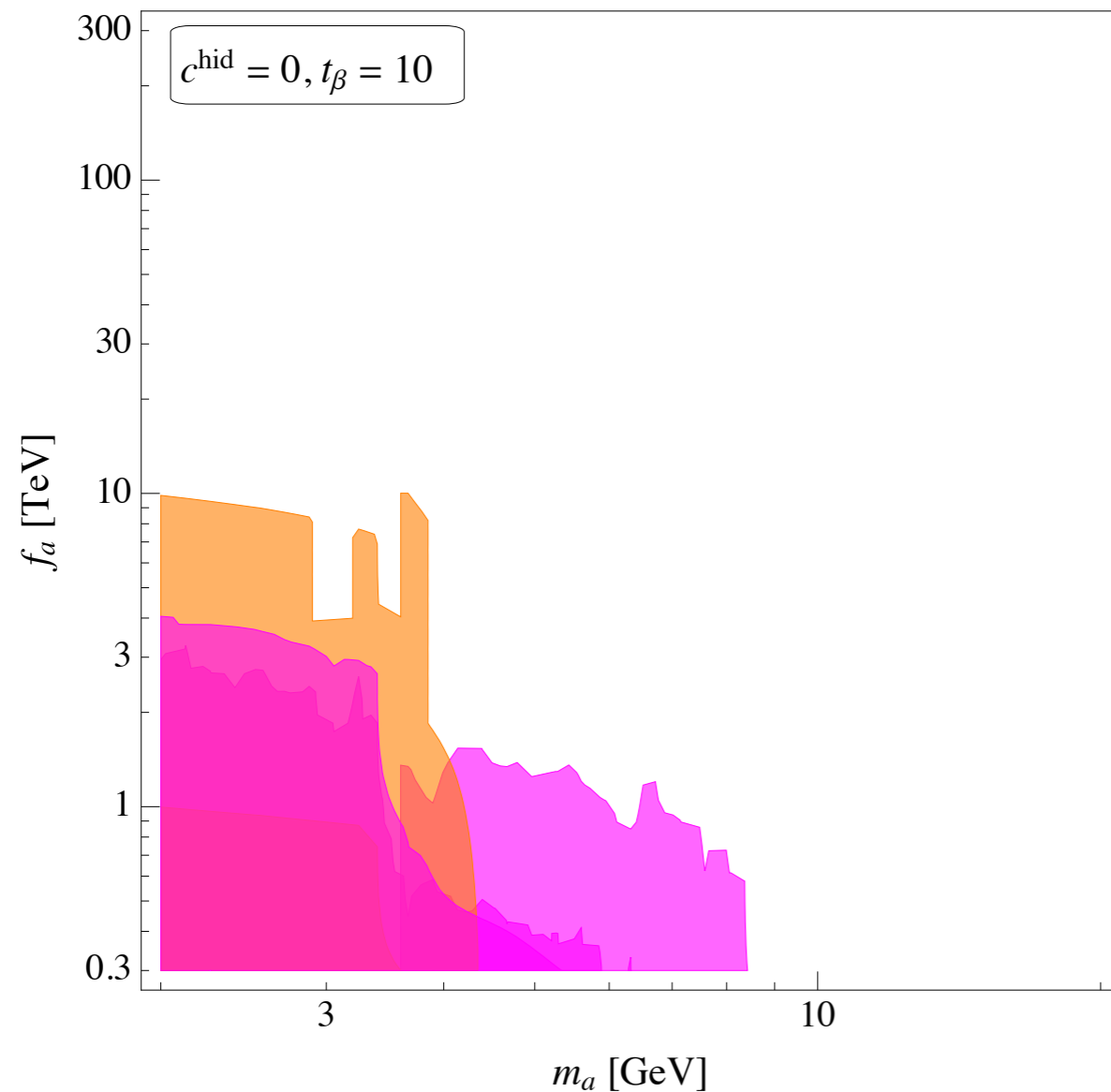
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$B \rightarrow K^{(*)} a(\mu\mu) \sim$ sensitive to value of t_β

Both disappear for $r_H = 0$

maybe not, work in progress...

R-axion from decays of MSSM Higgses

very preliminary

