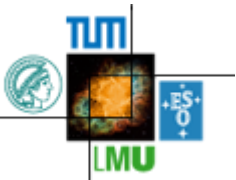


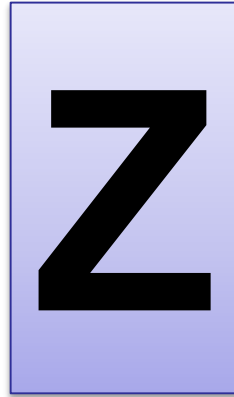
Z-Mediated New Physics, Vector-Like Quark Models and Beyond

Andrzej J. Buras
(Technical University Munich TUM-IAS)

Planck 2017, Warsaw, May 2017

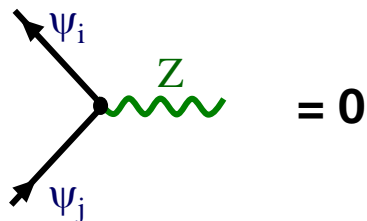


Overture



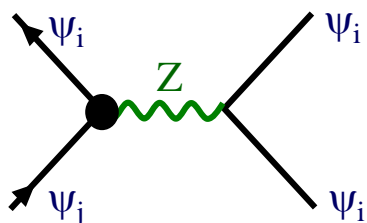
**After no signs of new particles
at the LHC Z-boson is among SM
particles particularly suited to be
a messenger of New Physics at
and beyond the LHC scales**

Z-Boson at Work (SM)



GIM

$\psi_i = \text{quarks, leptons}$



Z-Penguin enters leptonic, semi-leptonic, non-leptonic decays of mesons

gauge
 $C(x_t, \xi)$
Inami-Lim (1981)

Gauge-independent functions:

Buchalla, AJB, Harlander (1990)

$$X(x_t) = C(x_t) - 4B(x_t)$$

$$Y(x_t) = C(x_t) - B(x_t)$$

$$Z(x_t) = C(x_t) + \frac{1}{4} D(x_t)$$

Box

Govern most of rare $K, B_{s,d}$ decays

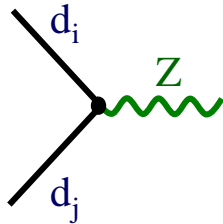
$$B_{s,d} \rightarrow \mu^+ \mu^-, K \rightarrow \pi \nu \bar{\nu},$$

$$B \rightarrow K(K^*) l^+ l^-, \varepsilon'/\varepsilon \text{ etc.}$$

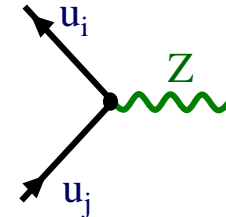
Photon penguin

Z-Boson at Work (BSM)

FC couplings generated by some NP



$$[\Delta_{L,R}^d(\mathbf{Z})]_{ij}$$



$$[\Delta_{L,R}^u(\mathbf{Z})]_{ij}$$

Example:

**Vector-like quarks mixing
with SM quarks during
electroweak symmetry breaking**

(Nir; Aquila et al; Ishikawa, Ligeti, Wise)

**Some recent phenomenological applications of
General Z-scenario (simplified Z-models)**

AJB, De Fazio, Girrbach (1211.1896)

AJB, Buttazzo, Knecht (1507.08672)

AJB (1601.00005)

Gauge-Invariant Standard Model Effective Theory

Buchmüller, Wyler, 1990

Grzadkowski, Iskrzynski, Misiak, Rosiek, 2010

Warsaw basis

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda^2} \sum_K \mathbf{c}_K^{(6)} \mathbf{O}_K^{(6)} \leftarrow \text{Dimension 6 operators}$$

59 Operators when flavour indices omitted

2499 Operators with flavour indices Jenkins, Manohar, Trott

**2499 x 2499 Anomalous dim matrix governs
Renormalization Group evolution**

from μ_{EW} to $\Lambda = \text{scale of new physics}$

But: **Useful results for analyses of specific models**

New Physics Model

Λ_{NP}



Integrate out Heavy Particles



**Renormalization Group
Running: QCD, Yukawa, etc.**



**Decay Amplitudes
at $\mu \approx 0(M_W)$**



QCD + QED Renormalization Group



**Low energy decay amplitudes
Lattice Calculations**

**Known
at LO,
NLO,
NNLO**

**Recent
Progress**

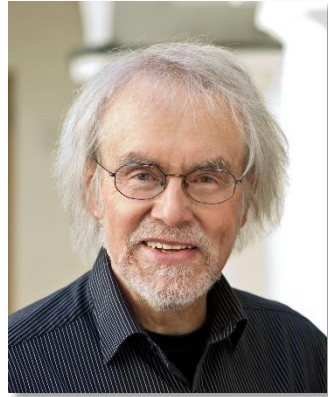
Basic Questions for Next 24 min

- 1.** What is the pattern of New Physics mediated by Z-boson?
- 2.** How is it described by SMEFT?
- 3.** What can models with Vector-like Quarks offer in this context?
- 4.** Z' at work

BBCJ Collaboration



Christoph Bobeth



AJB



Alejandro Celis



Martin Jung

1703.04753 (General Z Models) (40 pages)

1609.04783 (Vector-like Quark Models) (74 pages)

AJB 1601.00005 (Z' models) (67 pages)

Main Actors

SMEFT

$$\mathbf{O}_{\text{Hd}} = \left(\mathbf{H}^+ i \vec{\mathbf{D}}_{\mu} \mathbf{H} \right) \left[\bar{\mathbf{d}}_{\text{R}}^i \gamma^{\mu} \mathbf{d}_{\text{R}}^j \right]$$

(RH Scenario)

$$\mathbf{O}_{\text{Hq}}^{(1)} = \left(\mathbf{H}^+ i \vec{\mathbf{D}}_{\mu} \mathbf{H} \right) \left[\bar{\mathbf{q}}_{\text{L}}^i \gamma^{\mu} \mathbf{q}_{\text{L}}^j \right]$$

$$\mathbf{O}_{\text{Hq}}^{(3)} = \left(\mathbf{H}^+ i \vec{\mathbf{D}}_{\mu}^a \mathbf{H} \right) \left[\bar{\mathbf{q}}_{\text{L}}^i \sigma^a \gamma^{\mu} \mathbf{q}_{\text{L}}^j \right]$$

(LH Scenario)

$$\left[\mathbf{C}_{\text{Hd}} \right]_{ij}, \quad \left[\mathbf{C}_{\text{Hq}}^{(1)} \right]_{ij}, \quad \left[\mathbf{C}_{\text{Hq}}^{(3)} \right]_{ij}$$

Complex Couplings

$$\mathcal{O} \left(\frac{1}{\Lambda^2} \right)$$

Generated Z-Couplings :

$$\left[\Delta_{\text{R}}^{\text{d}}(\mathbf{z}) \right]_{ij} = -\frac{g_{\text{z}}}{2} v^2 \left[\mathbf{C}_{\text{Hd}} \right]_{ij}$$

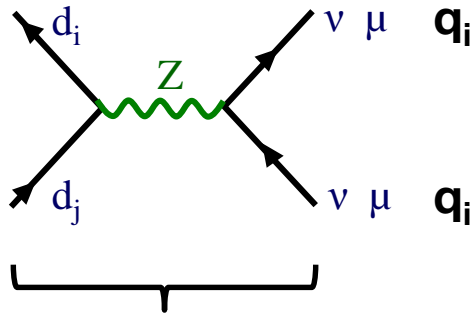
$$\left[\Delta_{\text{L}}^{\text{d,u}}(\mathbf{z}) \right]_{ij} = -\frac{g_{\text{z}}}{2} v^2 \left[\mathbf{C}_{\text{Hq}}^{(1)} \pm \mathbf{C}_{\text{Hq}}^{(3)} \right]_{ij}$$

$$\Lambda = \Lambda_{\text{NP}}$$

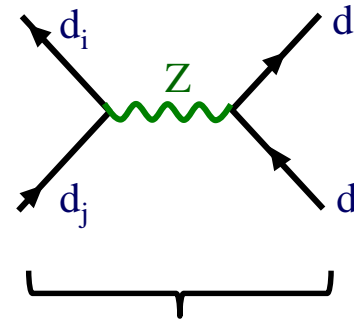
$$g_{\text{z}} = \sqrt{g_1^2 + g_2^2}$$

Simplified Z Scenarios

$\Delta F=1$



$\Delta F=2$



{ **Dim 6 Contributions**
 v^2/Λ^2 }

{ **Dim 8 Contributions**
 v^4/Λ^4 }

(only relevant for $\Lambda < 3\text{TeV}$)

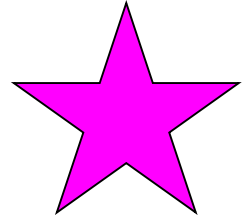
What are the Dim 6 Contributions to $\Delta F=2$ Processes in Z-Scenarios?

Use SMEFT to find out.

4 Lessons from SMEFT

(BBCJ)

Lesson 1 (RH Scenario)



$[C_{Hd}]_{ij} \neq 0$ generates through RG evolution $\Lambda \rightarrow \mu_{EW}$ due to Yukawa couplings

$\Delta F=2$ LR operators

which are further enhanced through QCD RG evolution and large hadronic matrix elements.

$$\lambda_t^{ij} = \mathbf{V}_{ti}^* \mathbf{V}_{tj}$$

$$\mathbf{x}_t = \frac{m_t^2}{M_W^2}$$

$$O_{LR,1}^{ij} = [\bar{d}_i \gamma_\mu P_L d_j] [\bar{d}_i \gamma^\mu P_R d_j]$$
$$C_{LR,1}^{ij}(\mu_{EW}) \propto v^2 \lambda_t^{ij} [C_{Hd}]_{ij} \mathbf{x}_t \ln \frac{\Lambda}{\mu_{EW}}$$

$i, j = d, s, b$

scale dependence



Very strong constraint on rare decays!

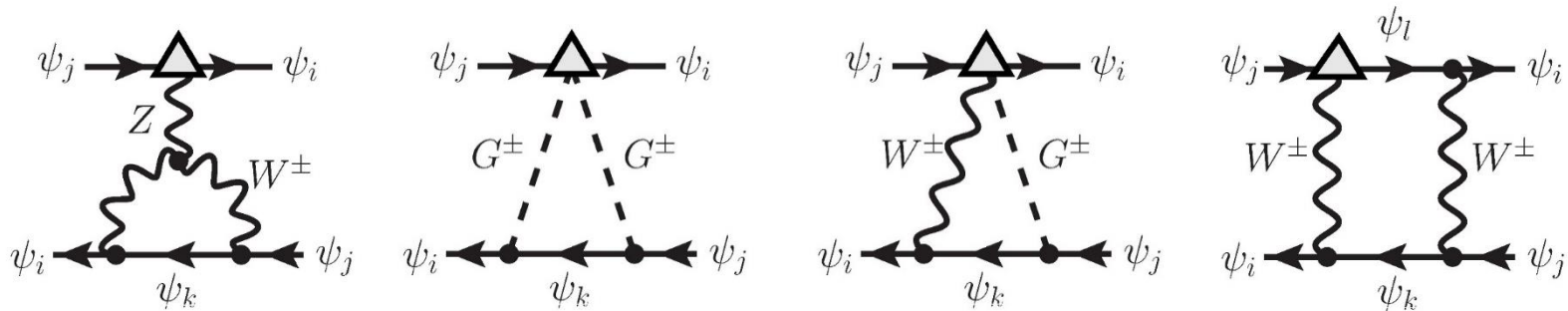
(in particular in K-system)

Lesson 2

(BBCJ)



μ_{EW} dependence cancelled by NLO Corrections



Corresponds to Z-penguin with $d^i \gamma^\mu P_L d^i$ replaced by $\Delta_R^{ij}(\mathbf{Z})$

Cancel gauge dependence of the first diagram

New gauge-independent Function $H_1(x_t)$

Analog to $X(x_t)$, $Y(x_t)$, $Z(x_t)$ in SM (confirmed by Endo et al. V2)

(Endo, Kitahara, Mishima, Yamamoto (1612.08839))

Lesson 3

(BBCJ)



In LH scenario no new $\Delta F=2$ operators generated

But **two couplings** required to describe all effects.

$$\left[\Delta_L^d(\mathbf{Z}) \right]_{ij} \sim \mathbf{C}_{Hq}^{(+)}$$

$$\mathbf{C}_{Hq}^{(\pm)} = \mathbf{C}_{Hq}^{(1)} \pm \mathbf{C}_{Hq}^{(3)}$$

$$\left[\Delta_L^u(\mathbf{Z}) \right]_{ij} \sim \mathbf{C}_{Hq}^{(-)}$$

$$i, j = d, s, b$$

$$\left[\Delta \mathbf{C}_{VLL}(\mu_{EW}) \right]^{ij} \propto \lambda_t^{ij} \mathbf{v}^2 \mathbf{x}_t \left[\mathbf{C}_{Hq}^{(-)} \ln \frac{\Lambda}{\mu_{EW}} + \mathbf{F}(\mathbf{C}_{Hq}^{(\pm)}, \mathbf{x}_t, \mu_{EW}) \right]$$

Cancellation of μ_{EW} dependence.

2 gauge independent Functions $H_1(x_t), H_2(x_t)$

Connection between up and down through $SU(2)_L$ invariance

$$\mathbf{O}_{VLL}^{ij} = \left[\bar{\mathbf{d}}_i \gamma_\mu \mathbf{P}_L \mathbf{d}_j \right]^2$$

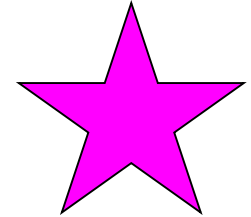
(SM operator)

Lesson 4

(BBCJ)

in down sector

In rare decays ($\Delta F=1$) only $C_{Hq}^{(+)}$ enters
but in $\Delta F=2$ $C_{Hq}^{(+)}$ and $C_{Hq}^{(-)}$



$$\left[\Delta_L^d(\mathbf{Z}) \right]_{ij}$$

insufficient

No model independent correlation
between $\Delta F=1$ and $\Delta F=2$ transitions

In contrast to
simplified models



Weak constraints on rare decays from $\Delta F=2$



Still correlations between various $\Delta F=1$ transitions
present



In specific models also

$$(\Delta F=2) \leftrightarrow (\Delta F=1)$$

B Physics Anomalies

Many papers:

Violation of lepton flavour universality

New flavour violating interactions:

**Z' , Leptoquarks, Vector-like quarks,
General 2HDM, $U(2)$, W' , H^+ , ...**

But no particular signs of new sources of CP-violation!

Yet also anomaly in CP-violation in K-physics (ϵ'/ϵ)

$\epsilon' =$ CP-violation in Decay ($K_L \rightarrow \pi\pi$)

$\epsilon =$ CP-violation in $K^0 - \bar{K}^0$ Mixing



ε'/ε Anomaly

$$\left(\varepsilon'/\varepsilon\right)_{\text{exp}} = (16.6 \pm 2.3) \cdot 10^{-4}$$

NA48 (CERN) (2001)
KTeV (Fermilab)

$$\left(\varepsilon'/\varepsilon\right)_{\text{SM}} = (1.4 \pm 6.9) \cdot 10^{-4}$$

(RBC-UKQCD) (Lattice)

Use RBC-QCD

$$\left(\varepsilon'/\varepsilon\right)_{\text{SM}} = (1.9 \pm 4.5) \cdot 10^{-4}$$

(AJB, Gorbahn, Jamin, Jäger)

$$\left(\varepsilon'/\varepsilon\right)_{\text{SM}} = (1.1 \pm 5.1) \cdot 10^{-4}$$

(Kitahara, Nierste, Tremper)

$$\left(\varepsilon'/\varepsilon\right)_{\text{SM}} \leq (6.0 \pm 2.4) \cdot 10^{-4}$$

(AJB, Gérard)
(Dual QCD Approach, I/N)
(1507.06326)

$$\left(\varepsilon'/\varepsilon\right) = \left(\varepsilon'/\varepsilon\right)_{\text{SM}} + \left(\varepsilon'/\varepsilon\right)_{\text{NP}}$$

$$\left(\varepsilon'/\varepsilon\right)_{\text{NP}} = \kappa_{\varepsilon'} \cdot 10^{-3}$$



ε'/ε Anomaly

Largest anomaly
in Flavour Physics

$$\left(\varepsilon'/\varepsilon\right)_{\text{exp}} = (16.6 \pm 2.3) \cdot 10^{-4}$$

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KTeV (Fermilab)

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$K^+ \rightarrow \pi^+ \nu \bar{\nu}$ and $K_L \rightarrow \pi^0 \nu \bar{\nu}$ in the SM

QCD Corrections:

NLO Buchalla, AJB; Misiak, Urban (93, 98)
 NNLO AJB, Gorbahn, Haisch, Nierste (2005)

NLO EW Corrections:

Large m_t : Buchalla, AJB (1997)
 Exact NLO (m_t): Brod, Gorbahn, Stamou (2010)
 " " (m_c): Brod, Gorbahn (2008)

LD Effects:

Isidori, Mescia, Smith (2005)
 Mescia, Smith (2007)

+ Isospin breaking corrections



TH uncertainties at the level of 2% in BR

Unique in Flavour Physics !!

But significant parametric uncertainties

due to

$|V_{ub}|, |V_{cb}|, \gamma$

Data

Waiting for NA62, KOPIO

SM:

$\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.4 \pm 1.0) \cdot 10^{-11}$ $\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu}) = (3.4 \pm 0.6) \cdot 10^{-11}$	$\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (17.3 \pm 11) \cdot 10^{-11}$ $\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu}) \leq 2.6 \cdot 10^{-8}$
---	---

$$\mathbf{K^+ \rightarrow \pi^+ \nu \bar{\nu} \text{ vs } K_L \rightarrow \mu^+ \mu^-} \quad (\mathbf{Z})$$

(BBCJ)

The fate of $\mathbf{K^+ \rightarrow \pi^+ \nu \bar{\nu}}$ in LH scenarios depends on the sign of the interference between SD and LD part of the dispersive part of $\mathbf{K_L \rightarrow \mu^+ \mu^-}$

$$\left\{ \begin{array}{l} \text{D'Ambrosio, Portoles (9610244)} \\ \text{Gérard, Smith, Trine (0508189)} \end{array} \right\} \Rightarrow \mathbf{Br(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \leq 2 Br(K^+ \rightarrow \pi^+ \nu \bar{\nu})_{SM}}$$

using

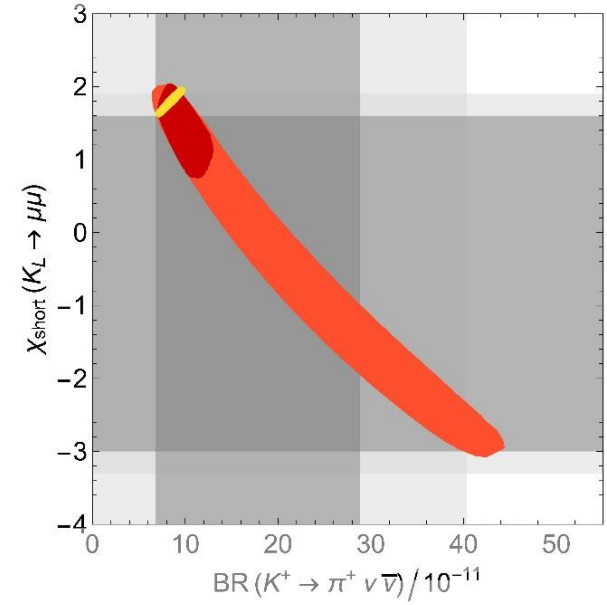
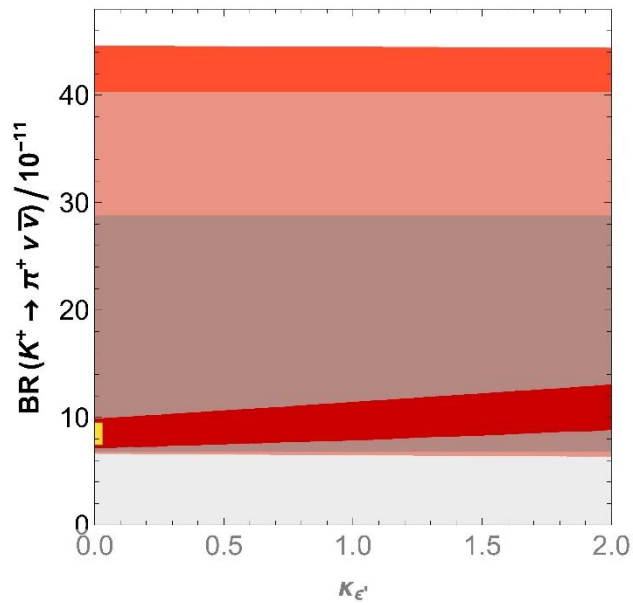
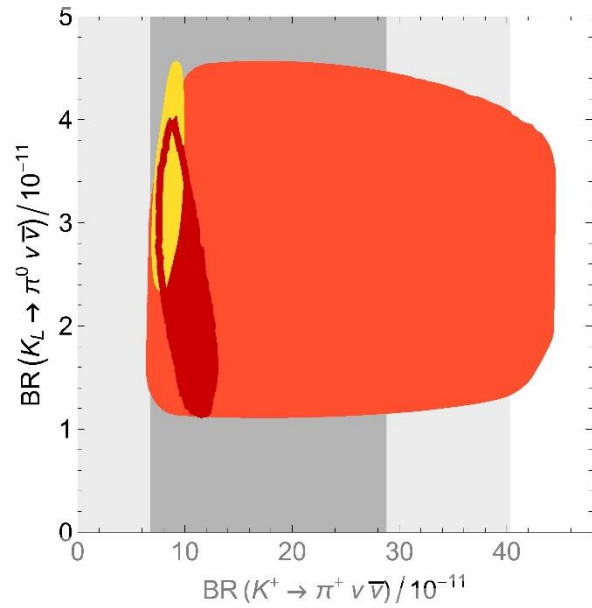
$$\left\{ \begin{array}{l} \text{Dumm, Pich (9801298)} \\ \text{Isidori, Unterdorfer (0311084)} \end{array} \right\} \Rightarrow \mathbf{Br(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \leq Br(K^+ \rightarrow \pi^+ \nu \bar{\nu})_{SM}}$$

$\mathbf{K_L \rightarrow \mu^+ \mu^-}$ less relevant for RH scenarios because ε_K stronger (LR operators)

$$\Rightarrow \mathbf{Br(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \leq 1.5 Br(K^+ \rightarrow \pi^+ \nu \bar{\nu})_{SM}}$$

General RH Scenario

(BBCJ)

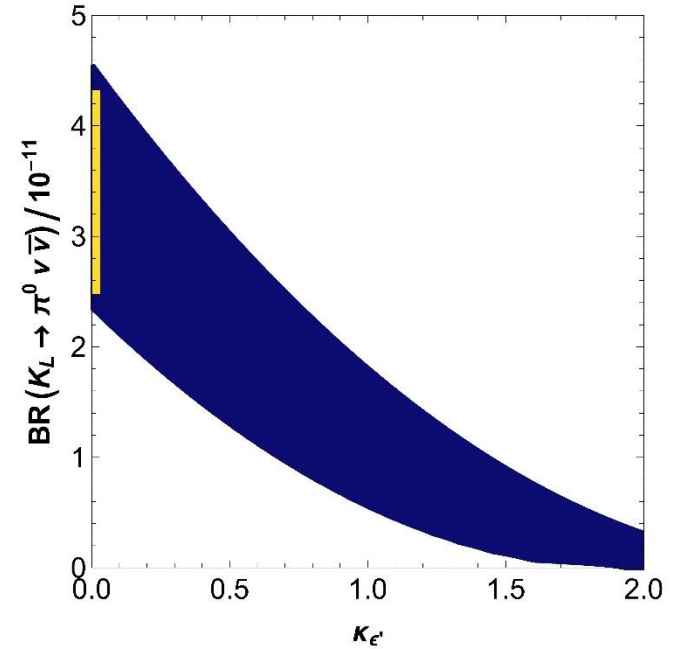
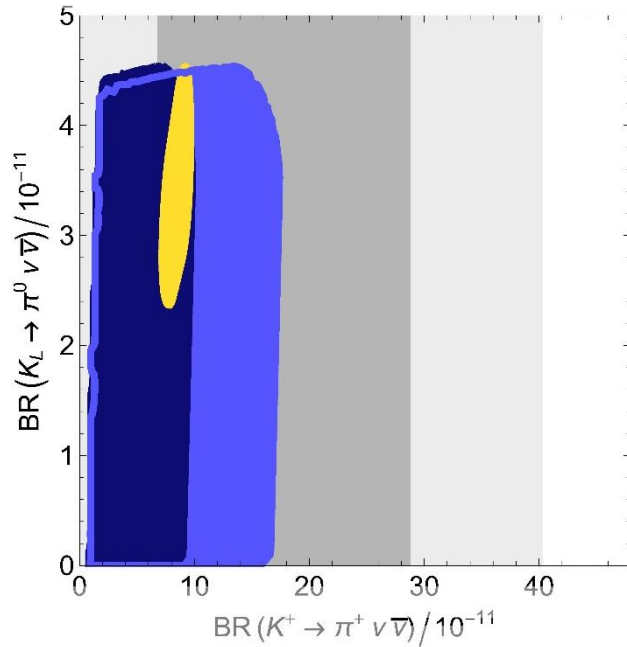


- No ϵ_K
- ϵ_K
- SM

ϵ'/ϵ anomaly solved

General LH Scenario

(BBCJ)



 strong $K_L \rightarrow \mu^+ \mu^-$

 SM

 weak $K_L \rightarrow \mu^+ \mu^-$

ϵ'/ϵ anomaly
solved

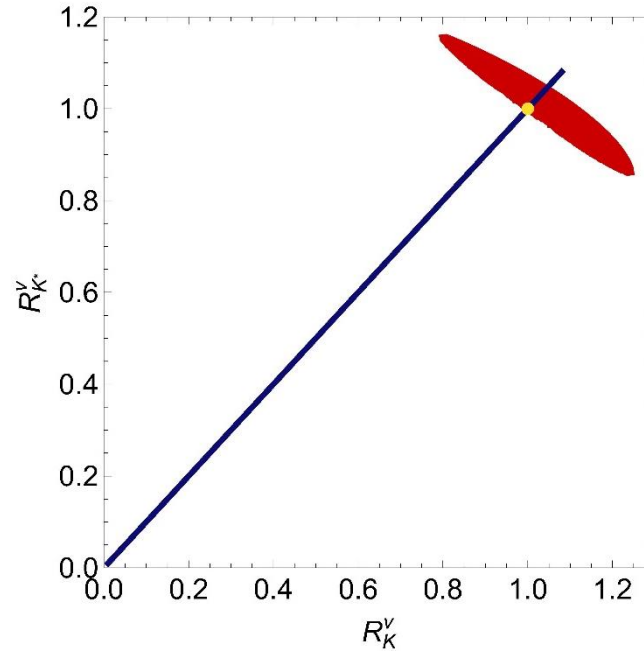
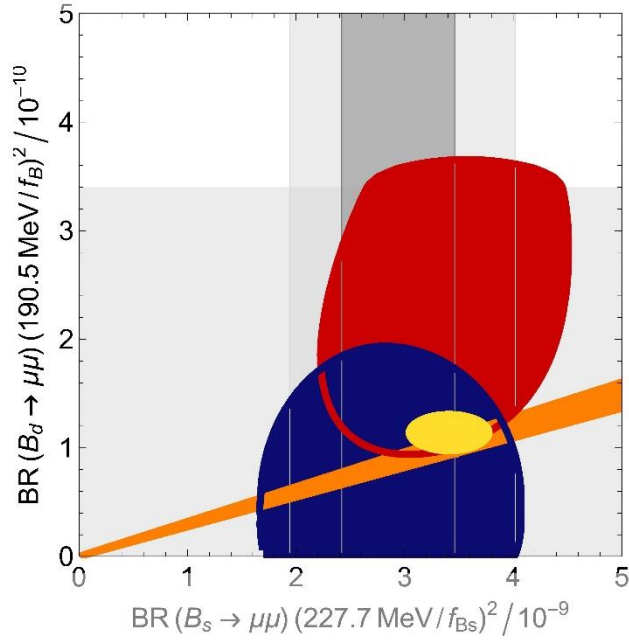
LH vs. RH



LH



RH




SM



MFV

$B \rightarrow K(K^*) \nu \bar{\nu}$



- 1.** ε'/ε -anomaly can easily be explained 
- 2.** $K_L \rightarrow \pi^0 \nu \bar{\nu}$ can only be suppressed unless both LH and RH couplings at work. **AJB (1601.0005)**
Endo et al (1612.08839)
- 3.** $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ can be enhanced but only by a factor of 2
- 4.** Interesting effects in $B_{s,d} \rightarrow \mu^+ \mu^-$ (in particular in RH scenario)
- 5.** Less interesting in $B \rightarrow K(K^*) \nu \bar{\nu}$
- 6.** No explanation of $P'_5, R_K, R_{K^*}, R(D), R(D^*)$ anomalies !

11 Vector-like Quark (VLQ) Models

Bobeth, AJB, Celis, Jung
1609.04783

- | | | | |
|-----|---|---|-------------------------------------|
| (5) | { | $\mathbf{G}_{\text{SM}} = \mathbf{SU}(3)_C \otimes \mathbf{SU}(2)_L \otimes \mathbf{U}(1)_Y$ | Ishikawa, Ligeti, Wise (1506.03484) |
| (2) | | $\mathbf{G}'_{\text{SM}}(\mathbf{S}) = \mathbf{G}_{\text{SM}} \otimes \mathbf{U}(1)_{L_\mu - L_\tau}$ | Altmannshofer et al. (1403.1269) |
| (4) | | $\mathbf{G}'_{\text{SM}}(\Phi) = \mathbf{G}_{\text{SM}} \otimes \mathbf{U}(1)_{L_\mu - L_\tau}$ | BBCJ |

B-Physics Anomalies

ε'/ε - Anomaly

\mathbf{G}_{SM}	(Z, box, RG)	No	★	3 LH, 2 RH
$\mathbf{G}'_{\text{SM}}(\mathbf{S})$	(Z', box)	★	No	1 LH, 1 RH
$\mathbf{G}'_{\text{SM}}(\Phi)$	(Z, Z', box)	★	★	3 LH, 1 RH

RG \equiv Yukawa effects

box \equiv boxes with VLQ + Scalars

Combination of $\Delta F=2$ and $\Delta F=1$ processes can determine New Physics scale Λ_{NP}

(BBCJ)



$B_s^0 - \bar{B}_s$ mixing

$B_s \rightarrow \mu^+ \mu^-$

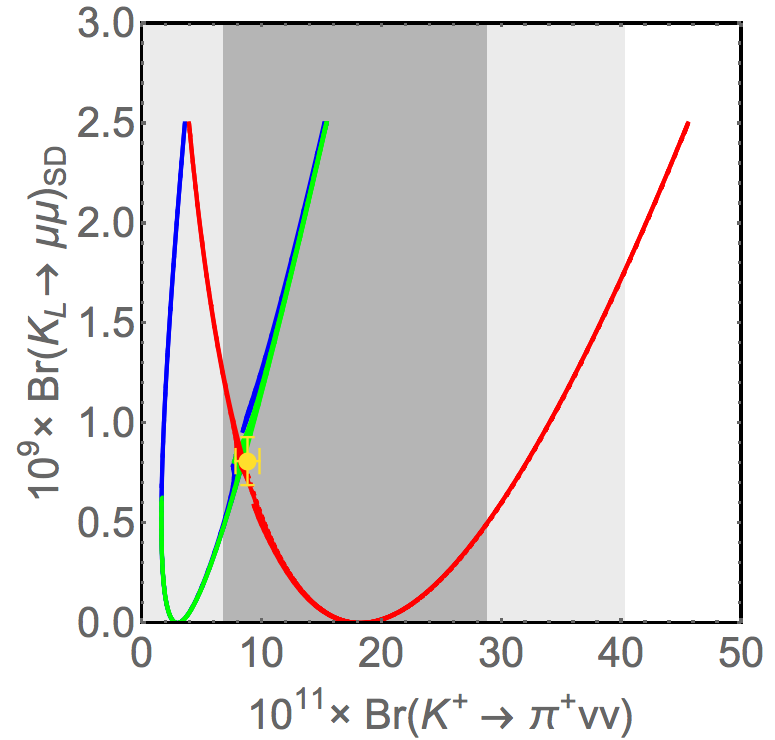
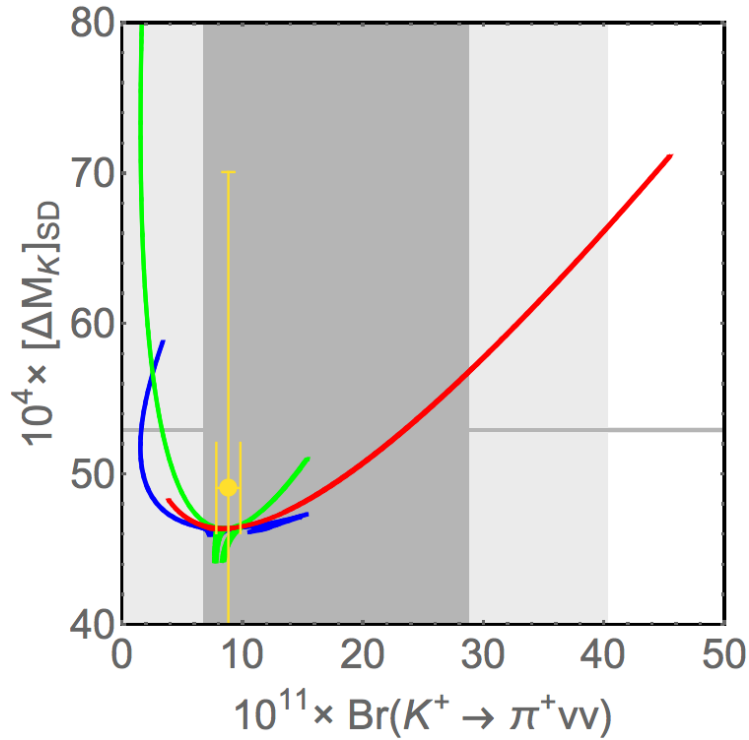
$$\frac{\sqrt{\Delta S_{NP}}}{\Delta Y_{NP}} = a \left[\frac{\Lambda_{NP}}{1\text{TeV}} \right]$$

See also
AJB, De Fazio, Girschbach,
Carlucci, 1211.1237
(331 Models)

a: independent of CKM and Yukawa couplings
but dependent on quantum numbers of VLQs

$G'_{SM}(\Phi)$

ΔM_K strikes back



from Christoph Bobeth

Messages on VLQ Models

- 1.** G_{SM} models: **Similar to Z-models but more specific in LH-models**
Only Z
- 2.** $G'_{SM}(S)$: **Interesting effects only in $\Delta F=2$, $B \rightarrow K(K^*)\mu^+\mu^-$ (solve anomalies)**
Only Z'
Fail with ε'/ε , $K \rightarrow \pi\nu\bar{\nu}$, $B_{s,d} \rightarrow \mu^+\mu^-$
- 3.** $G'_{SM}(\Phi)$: **Large effects in RH models in $K^+ \rightarrow \pi^+\nu\bar{\nu}$, $B_{s,d} \rightarrow \mu^+\mu^-$, ΔM_K smaller, but significant in LH. ε'/ε easily solved.**
Z, Z'
 $B \rightarrow K(K^*)\mu^+\mu^-$ anomalies only partly solved.
- 4.** **Large enhancement of $Br(B \rightarrow K(K^*)\nu\bar{\nu})$ would require several VLQ representations at work.**

General Z' at Work

Can solve anomalies in R_K, R_{K^*}, P_5'
(many papers)

Here :

$\varepsilon'/\varepsilon, K^+ \rightarrow \pi^+ \nu \bar{\nu}, K_L \rightarrow \pi^0 \nu \bar{\nu}, \Delta M_K$



Q_6, Q_6' – QCD Penguin operators

Q_8, Q_8' – Electroweak Penguin operators

ε'/ε within SM

$$\varepsilon'/\varepsilon \sim \left[\frac{\text{Re } A_2}{\text{Re } A_0} \text{Im } C_6 \langle Q_6 \rangle_0 - \text{Im } C_8 \langle Q_8 \rangle_2 + \text{smaller contributions} \right]$$

$$\left\{ \frac{\text{Re } A_2}{\text{Re } A_0} \approx \frac{1}{22} \quad \frac{\text{Im } C_6}{\text{Im } C_8} \approx 90 \quad \frac{\langle Q_8 \rangle_2}{\langle Q_6 \rangle_0} \approx 2 \right\} \Rightarrow \text{strong cancellations}$$

ε'/ε beyond SM (Q_6, Q_8, Q'_6, Q'_8)

1. Generally Q_8 wins over Q_6 because $\left(\frac{\text{Im } C_6}{\text{Im } C_8} \right)^{\text{NP}} \approx 0(1)$ but can provide $\Delta(\varepsilon'/\varepsilon) > 0$
2. Q_6 wins over Q_8 in the presence of a flavour symmetry forbidding Q_8

Strategy (Z')

AJB (1601.00005)

$$\left(\varepsilon'/\varepsilon\right)^{\text{NP}} = \kappa_{\varepsilon'} \cdot 10^{-3}$$

$$0.5 \leq \kappa_{\varepsilon'} \leq 1.5$$

(Im)

$$\varepsilon_{\kappa}^{\text{NP}} = \kappa_{\varepsilon} \cdot 10^{-3}$$

$$0.1 \leq \kappa_{\varepsilon} \leq 0.4$$

(Im, Re)

ε_{κ} more important than $K_L \rightarrow \mu^+\mu^-$ in Z' models

Re and Im Parts: Z' Couplings

$\Delta_L^{\text{sd}}(Z'), \Delta_R^{\text{sd}}(Z')$

$$\begin{array}{cccc}
 K^+ \rightarrow \pi^+ \nu \bar{\nu}, & K_L \rightarrow \pi^0 \nu \bar{\nu}, & K_L \rightarrow \mu^+ \mu^-, & \Delta M_K \\
 (\text{Re, Im}) & (\text{Im}) & (\text{Re}) & (\text{Im, Re})
 \end{array}$$

Basic Structure of NP Contributions

AJB (1601.00005)

$$\begin{array}{ll} (\varepsilon'/\varepsilon)^{\text{NP}} \rightarrow \text{Im} & \varepsilon_{\text{K}}^{\text{NP}} \rightarrow \text{Im} \cdot \text{Re} \\ (\kappa_{\varepsilon'} \geq 0.5) & (\kappa_{\varepsilon} \geq 0.1) \\ \Delta M_{\text{K}}^{\text{NP}} \sim \left[(\text{Re})^2 - (\text{Im})^2 \right] \end{array}$$

Dominance of $Q_6 (Q_6')$ \Rightarrow $\text{Im} \gg \text{Re} \Rightarrow \left\{ \Delta M_{\text{K}}^{\text{NP}} < 0 \right\}$
(large)

Dominance of $Q_8 (Q_8')$ \Rightarrow $\text{Re} \gg \text{Im} \Rightarrow \left\{ \Delta M_{\text{K}}^{\text{NP}} > 0 \right\}$
(small)



**Distinction between
these scenarios**

Main Message

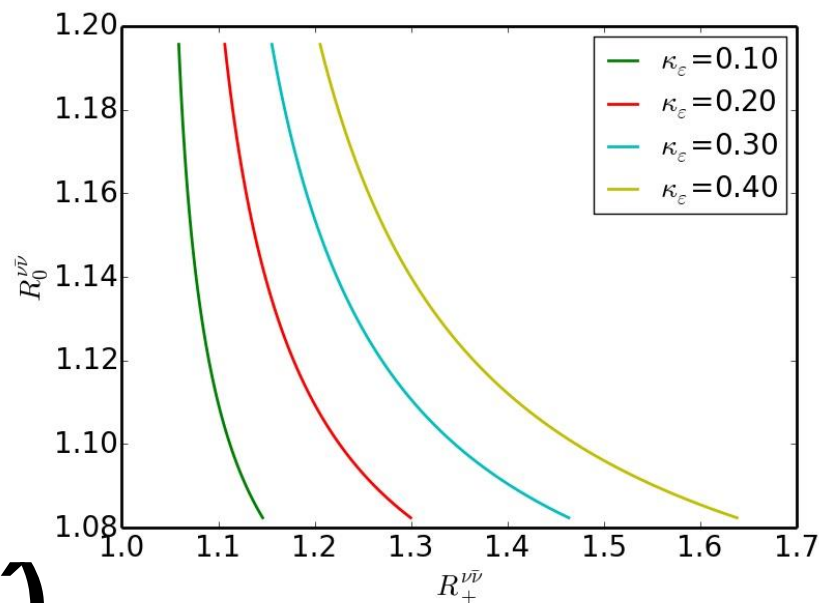
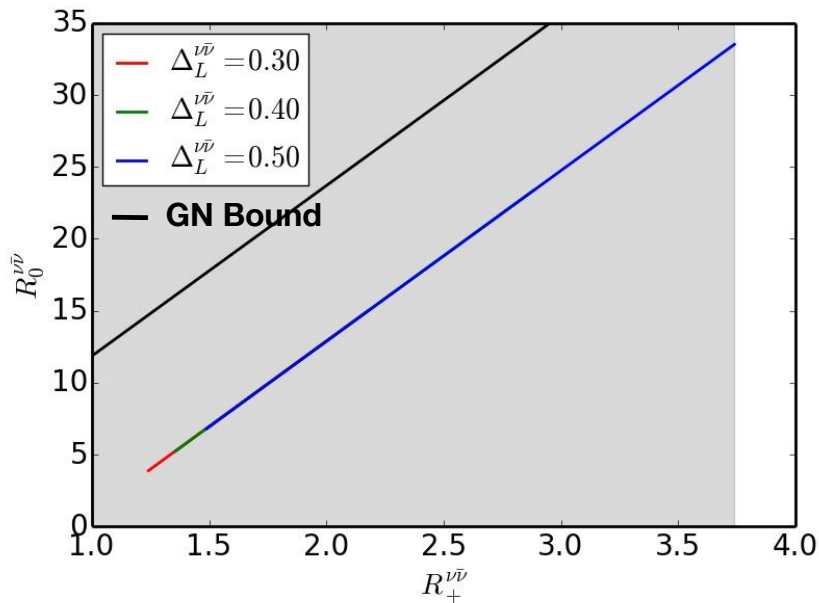


**Correlation between ε'/ε and $K \rightarrow \pi\nu\bar{\nu}$
in Z' scenarios depends on whether
QCP Penguin (Q_6) or EWP (Q_8)
dominates NP in ε'/ε**

QCDP (Q₆)

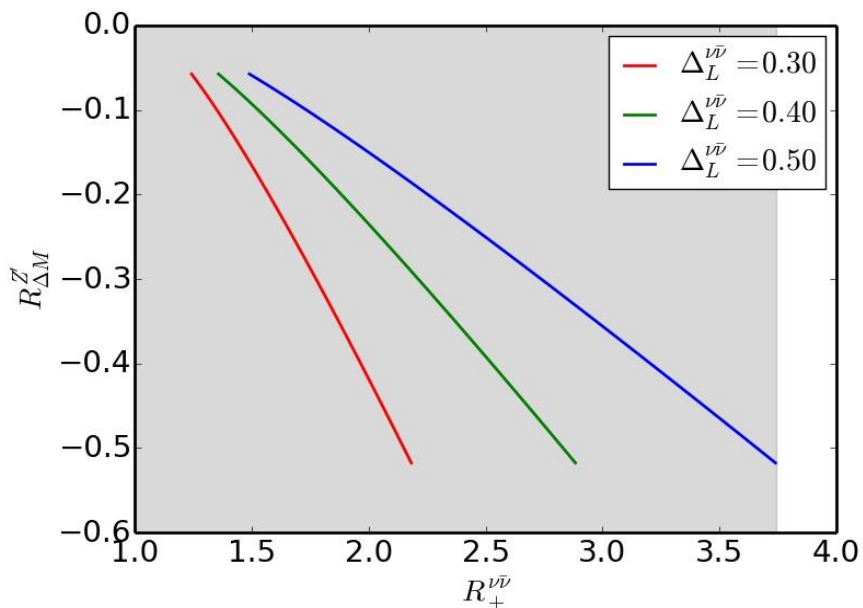
(0.5 < κ_{ε'} < 1.5)

EWP (Q₈)



(Z)

(R_{ΔM}^{Z'} > 0 but small)

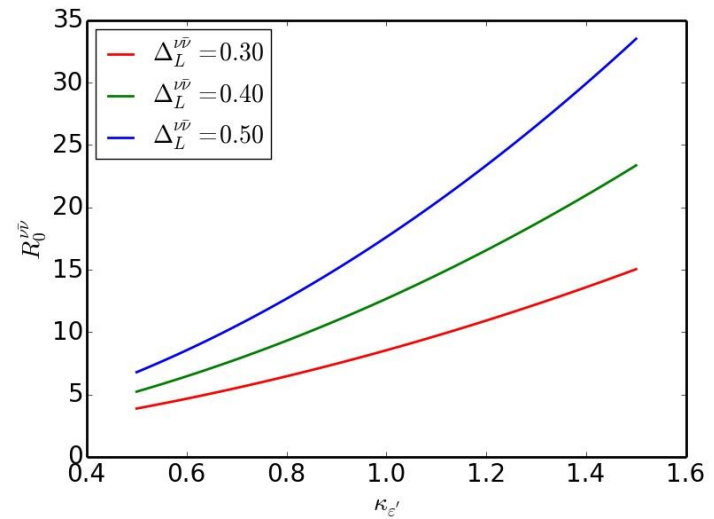
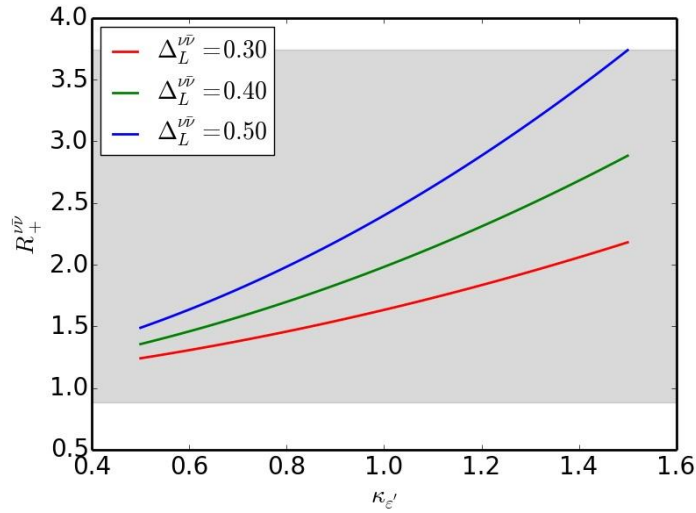


$$R_+^{v\bar{v}} = \frac{\text{Br}(K^+ \rightarrow \pi^+ v\bar{v})}{\text{Br}(K^+ \rightarrow \pi^+ v\bar{v})_{\text{SM}}}$$

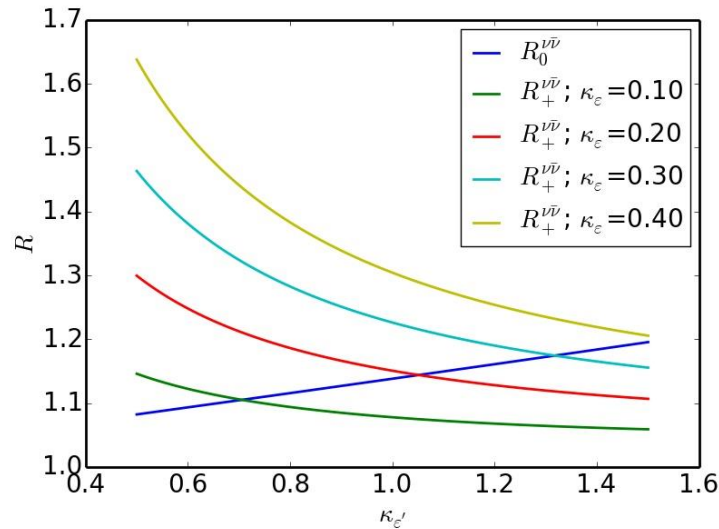
$$R_0^{v\bar{v}} = \frac{\text{Br}(K_L \rightarrow \pi^0 v\bar{v})}{\text{Br}(K_L \rightarrow \pi^0 v\bar{v})_{\text{SM}}}$$

$$R_{\Delta M}^{Z'} = \frac{(\Delta M_K)^{\text{NP}}}{(\Delta M_K)^{\text{exp}}}$$

QCD Penguin (Q_6)



Electroweak Penguin (Q_8)



(Z)

What about $\Delta I = 1/2$ Rule?

$$\frac{\text{Re } A_0}{\text{Re } A_2} \approx 22.4$$

Since 1955

Gell-Mann
Pais

1986, 2014

Large N including
1/N corrections

Quark Evolution $1 \text{ GeV} \leq \mu \leq M_W$
Meson Evolution $0 \leq \mu \leq 1 \text{ GeV}$

Correct value
of $\text{Re}A_2$

$$\left(\frac{\text{Re } A_0}{\text{Re } A_2} \right)_{1/N} \approx 16.0 \pm 1.5^*)$$

Dominance
of current-
current
operators

Correct value
of $\text{Re}A_2$

$$\left(\frac{\text{Re } A_0}{\text{Re } A_2} \right)_{\text{Lattice}} \approx 31 \pm 11$$

RBC-QCD
(2013, 2015)

*) G' with particular couplings ($M_{G'} \approx 3.5 \text{ TeV}$)
could be responsible for the missing piece

AJB
De Fazio
Girrbach-Noe
1404.3824

Main Messages

see



**Exciting Times are just
ahead of us !!!**

Coming Years

Flavour Precision Era

LHC Upgrade
E = 14 TeV
(CERN)

Precision
B_{d,s} – Meson
Decays
LHCb, CMS
KEK (Japan)

★
 $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ ($\sim 10^{-10}$) (CERN)
 $K_L \rightarrow \pi^0 \nu \bar{\nu}$ ($\sim 3 \cdot 10^{-11}$) J-PARC
(Japan)

Lepton Flavour
Violation

$$\mu \rightarrow e \gamma$$

$$\mu \rightarrow e e e$$

$$\tau \rightarrow \mu \gamma, \tau \rightarrow 3 \mu$$

Electric
Dipole
Moments

★
 $(g-2)_\mu$

Improved
Lattice
Gauge Theory
Calculations

★
 ϵ'/ϵ

★
 $\Delta I = 1/2$ Rule,
 ΔM_K

Neutrinos

2017-2025 : Expedition
Attouniverse → Zeptouniverse
 $10^{-18}\text{m} \rightarrow 10^{-21}\text{m}$

The Return of Kaon Flavour Physics

Looking for Anomalies in Kaon Flavour Physics



Photo: Gurli Buras

Anomalies in Kaon Flavour Physics

Finding Anomalies in Kaon Flavour Physics

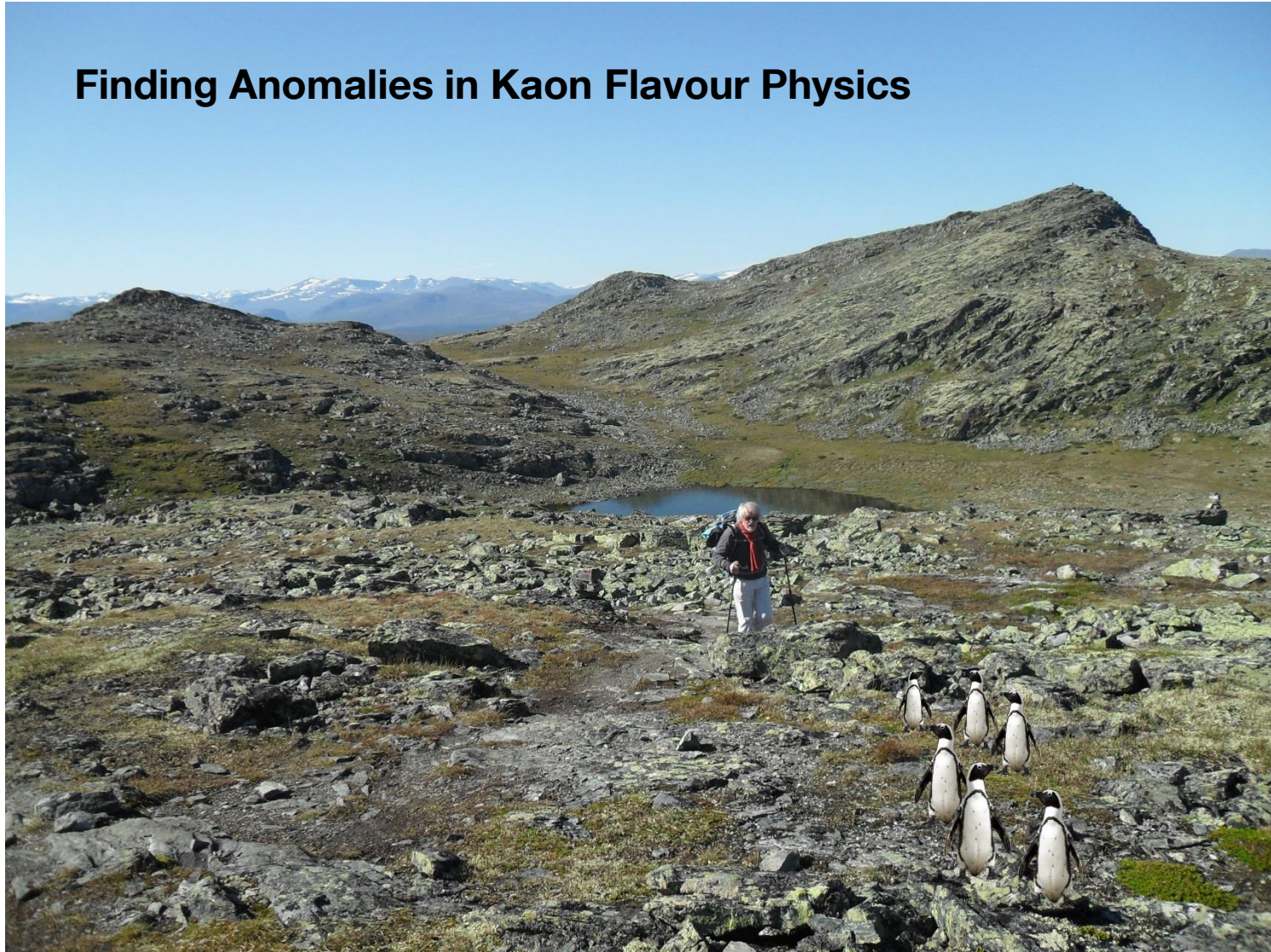


Photo: Gurli Buras

Backup

CKM Uncertainties

AJB, Buttazzo,
Girrbach-Noe,
Knegjens
1503.02693

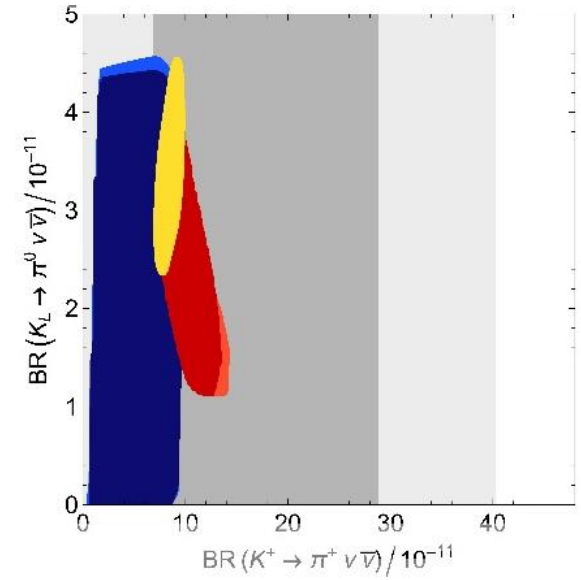
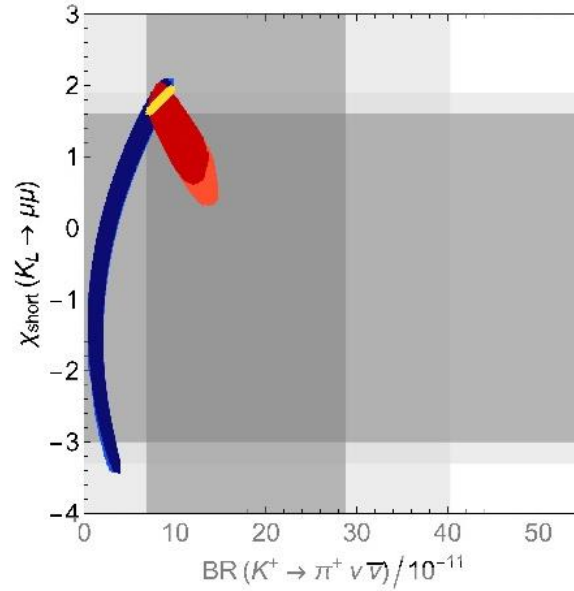
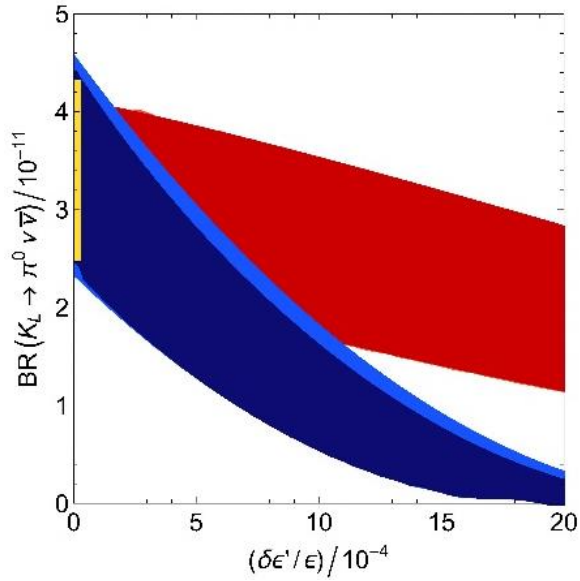
$$\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.39 \pm 0.30) \cdot 10^{-11} \left[\frac{|\mathbf{V}_{cb}|}{0.0407} \right]^{2.8} \left[\frac{\gamma}{73.2^\circ} \right]^{0.74}$$
$$\text{Br}(\text{K}_L \rightarrow \pi^0 \nu \bar{\nu}) = (3.36 \pm 0.05) \cdot 10^{-11} \left[\frac{|\mathbf{V}_{ub}|}{3.88 \cdot 10^{-3}} \right]^2 \left[\frac{|\mathbf{V}_{cb}|}{0.0407} \right]^2 \left[\frac{\sin \gamma}{\sin(73.2)} \right]^2$$

$$\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.39 \pm 0.58) \cdot 10^{-11} \left[\frac{\gamma}{73.2^\circ} \right]^{0.81} \left[\frac{\bar{\text{Br}}(\text{B}_s \rightarrow \mu^+ \mu^-)}{3.4 \cdot 10^{-9}} \right]^{1.42} \left[\frac{227.7}{F_{\text{B}_s}} \right]^{2.84}$$
$$\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.39 \pm 1.11) \cdot 10^{-11} \left[\frac{|\varepsilon_K|}{2.23 \cdot 10^{-3}} \right]^{1.07} \left[\frac{\gamma}{73.2^\circ} \right]^{-0.11} \left[\frac{|\mathbf{V}_{ub}|}{3.88 \cdot 10^{-3}} \right]^{-0.95}$$

$$\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.4 \pm 1.0) \cdot 10^{-11}$$
$$\text{Br}(\text{K}_L \rightarrow \pi^0 \nu \bar{\nu}) = (3.4 \pm 0.6) \cdot 10^{-11}$$

Vector-Like Quark Models

(BBCJ)



LH

RH

10 TeV

SM

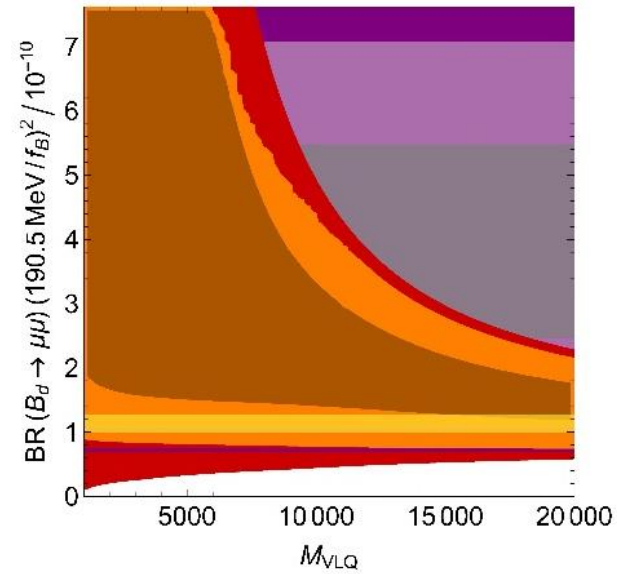
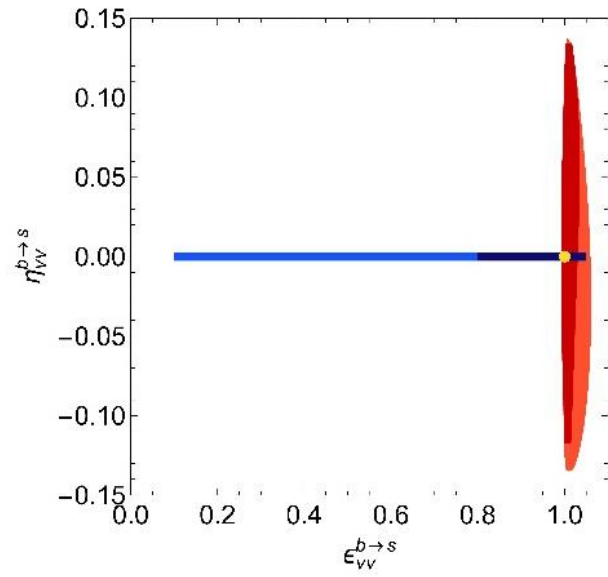
LH

RH

1 TeV

Vector-Like Quark Models

(BBCJ)



Basic Questions in Flavour Physics

**New Flavour
violating
CPV phases?**

**Flavour Conserving
CPV phases?**

**Non-MFV
Interactions?**

**Right-Handed
Charged
Currents?**

**Scalars H^0 , H^\pm
and related
FCNC's?**

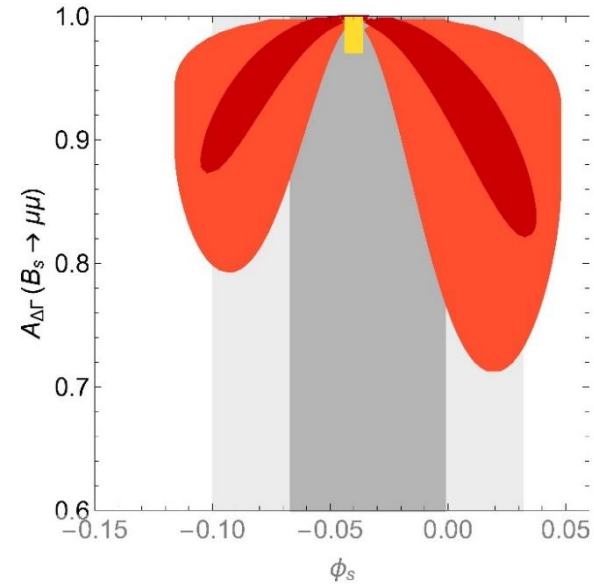
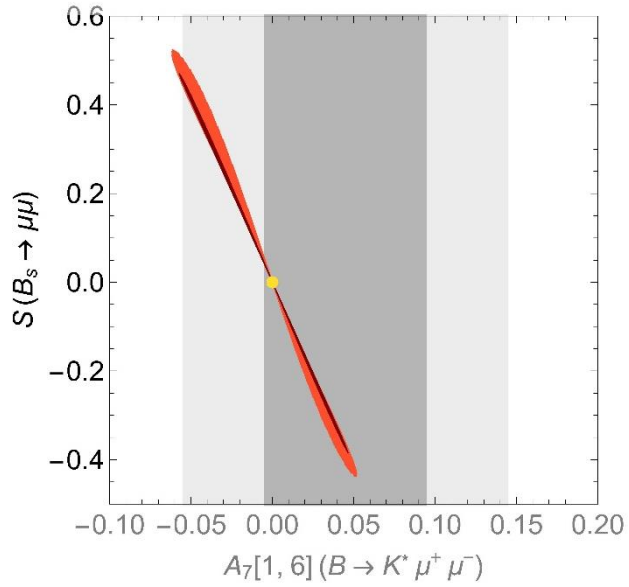
**New Fermions?
New Gauge
Bosons?**



**How to explain dynamically 22 free
Parameters in the Flavour Sector ?**

General RH Scenario

(BBCJ)

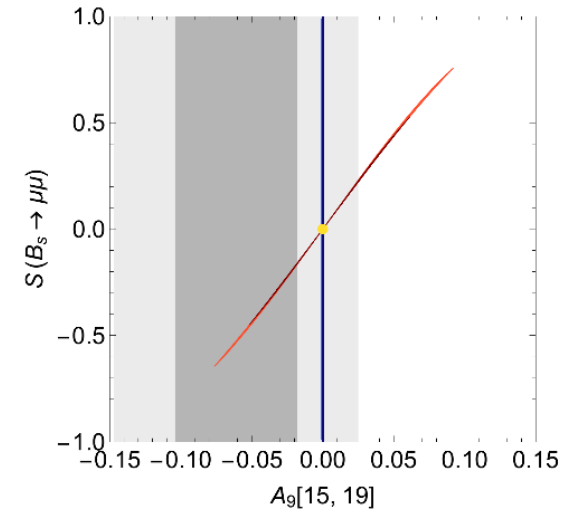
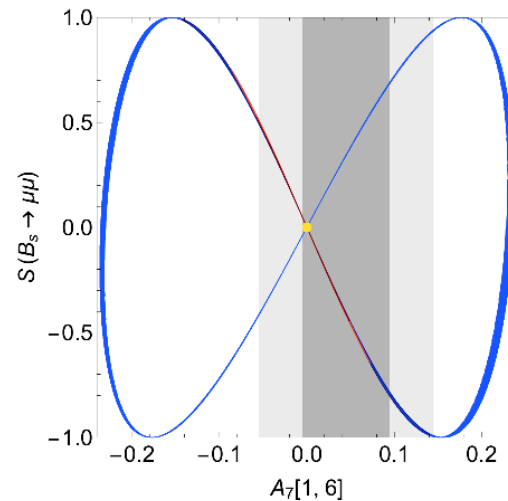
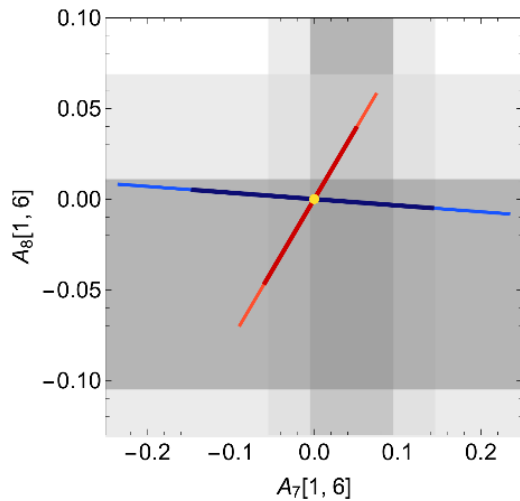


- No ε_K
- ε_K
- SM

Vector-Like Quark Models

(G_{SM})

Large CP-Violating Effects



 LH

 RH

10 TeV

 SM

 LH

 RH

1 TeV