

Rethinking Fundamental Interactions

Francesco Sannino

Asymptotic Safety in QFT

Francesco Sannino

The Standard Model ado

The Standard Model addendum

Fields:

Gauge fields + fermions + scalars

The Standard Model addendum

Fields:

Gauge fields + fermions + scalars

Interactions:

Gauge: $SU(3) \times SU(2) \times U(1)$ at EW scale

The Standard Model addendum

Fields:

Gauge fields + fermions + scalars

Interactions:

Gauge: $SU(3) \times SU(2) \times U(1)$ at EW scale

Yukawa: Fermion masses/Flavour

Culprit: Higgs

Scalar self-interaction

Gauge - Yukawa theories

$$L = -\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q + y(\bar{Q}_L H Q_R + \text{h.c.})$$

$$\text{Tr} [D H^\dagger D H] - \lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$$

Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + y(\bar{Q}_L H Q_R + \text{h.c.})$$

Gauge $\boxed{\text{Tr} [D H^\dagger D H]} - \lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge $\boxed{\text{Tr} [D H^\dagger D H]} - \lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge $\text{Tr} [D H^\dagger D H]$ - $\lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Scalar selfinteractions

Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge $\text{Tr} [D H^\dagger D H]$ - $\lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Scalar selfinteractions

4D: standard model, dark matter, ...

Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge $\text{Tr} [D H^\dagger D H]$ - $\lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Scalar selfinteractions

4D: standard model, dark matter, ...

3D: condensed matter, phase transitions

Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge $\text{Tr} [D H^\dagger D H]$ - $\lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Scalar selfinteractions

4D: standard model, dark matter, ...

3D: condensed matter, phase transitions

2D: graphene, ...

Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge $\text{Tr} [D H^\dagger D H]$ - $\lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Scalar selfinteractions

4D: standard model, dark matter, ...

3D: condensed matter, phase transitions

2D: graphene, ...

4plusD: extra dimensions, string theory, ...

Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge $\text{Tr} [D H^\dagger D H]$ - $\lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Scalar selfinteractions

4D: standard model, dark matter, ...

3D: condensed matter, phase transitions

2D: graphene, ...

4plusD: extra dimensions, string theory, ...

Universal description of physical phenomena

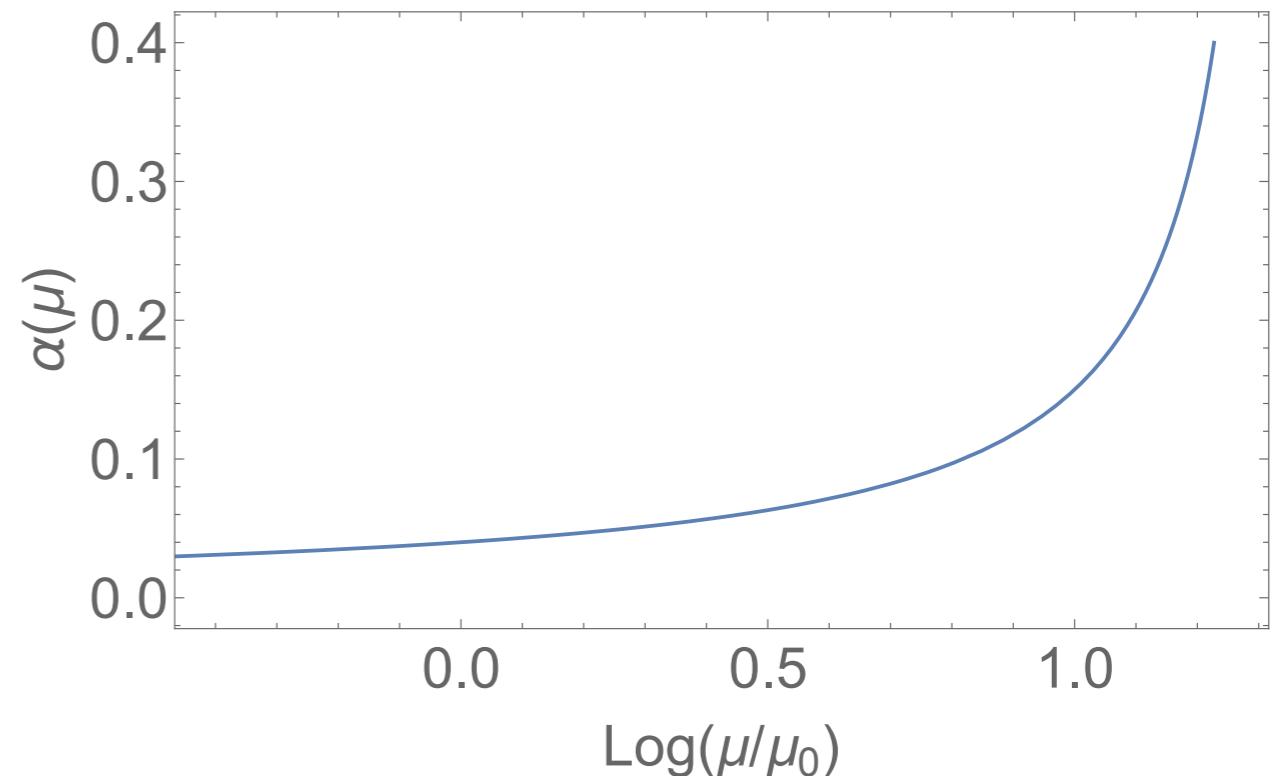
Two main issues

Two main issues

- ◆ EW scale stability
- ◆ UV triviality (Landau Pole)

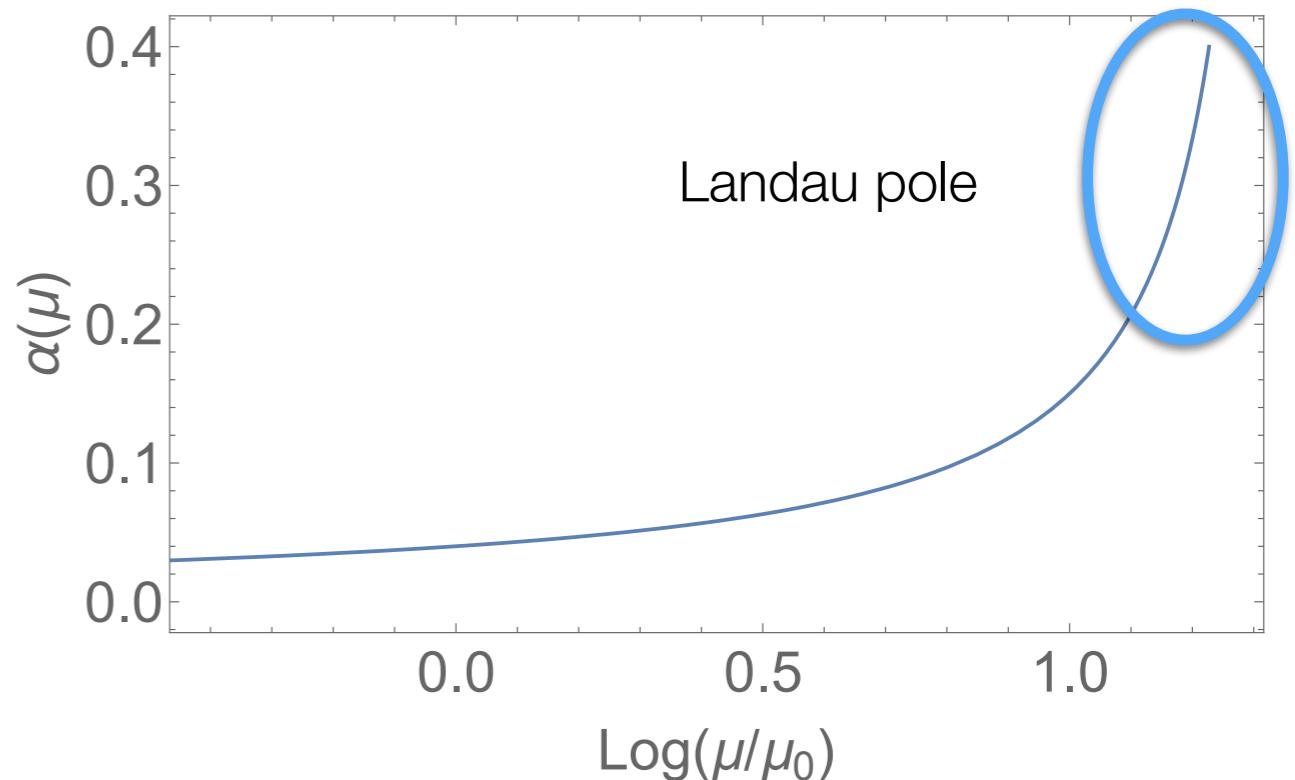
Two main issues

- ◆ EW scale stability
- ◆ UV triviality (Landau Pole)



Two main issues

- ◆ EW scale stability
- ◆ UV triviality (Landau Pole)



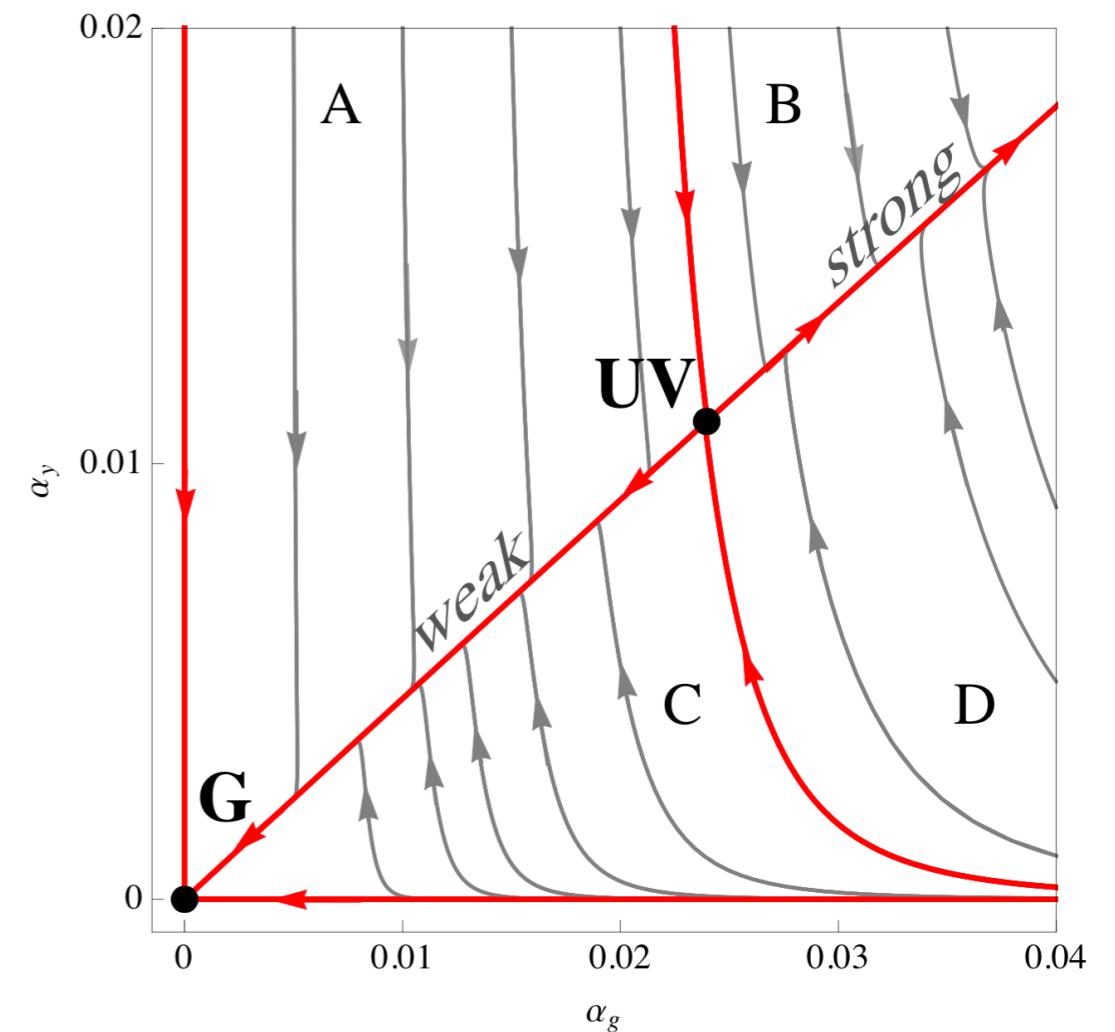
Fundamental theory

Fundamental theory

Wilson: A fundamental theory has an UV fixed point

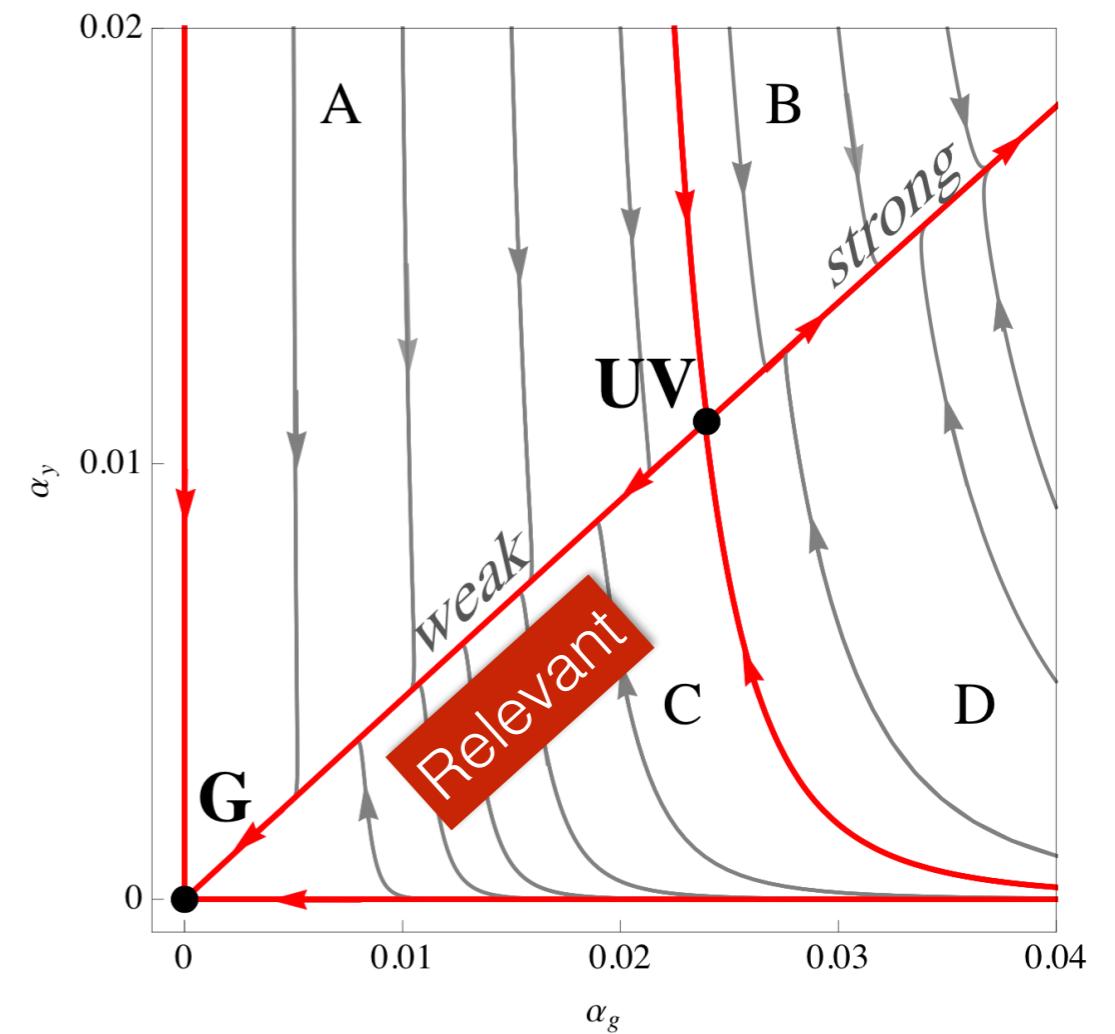
Fundamental theory

Wilson: A fundamental theory has an UV fixed point



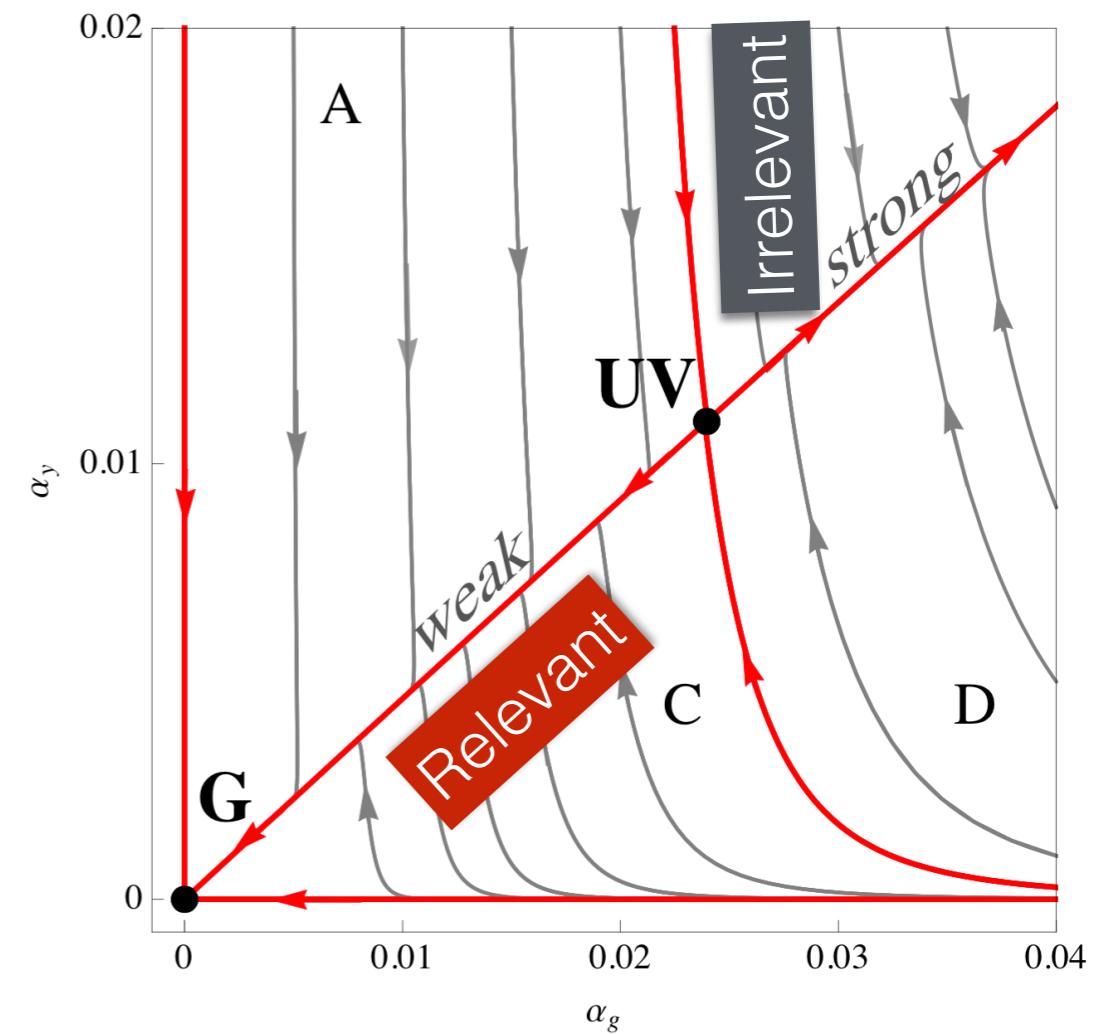
Fundamental theory

Wilson: A fundamental theory has an UV fixed point



Fundamental theory

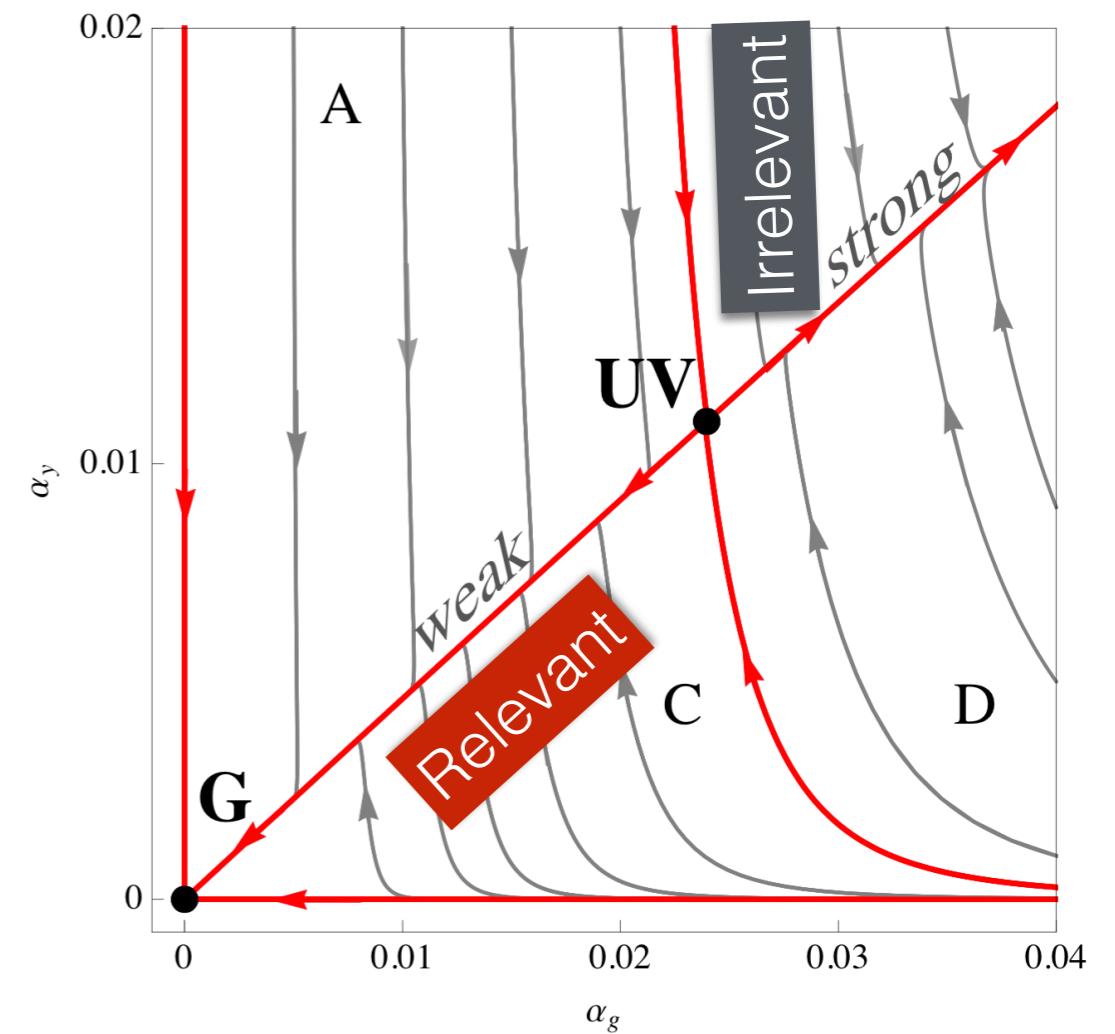
Wilson: A fundamental theory has an UV fixed point



Fundamental theory

Wilson: A fundamental theory has an UV fixed point

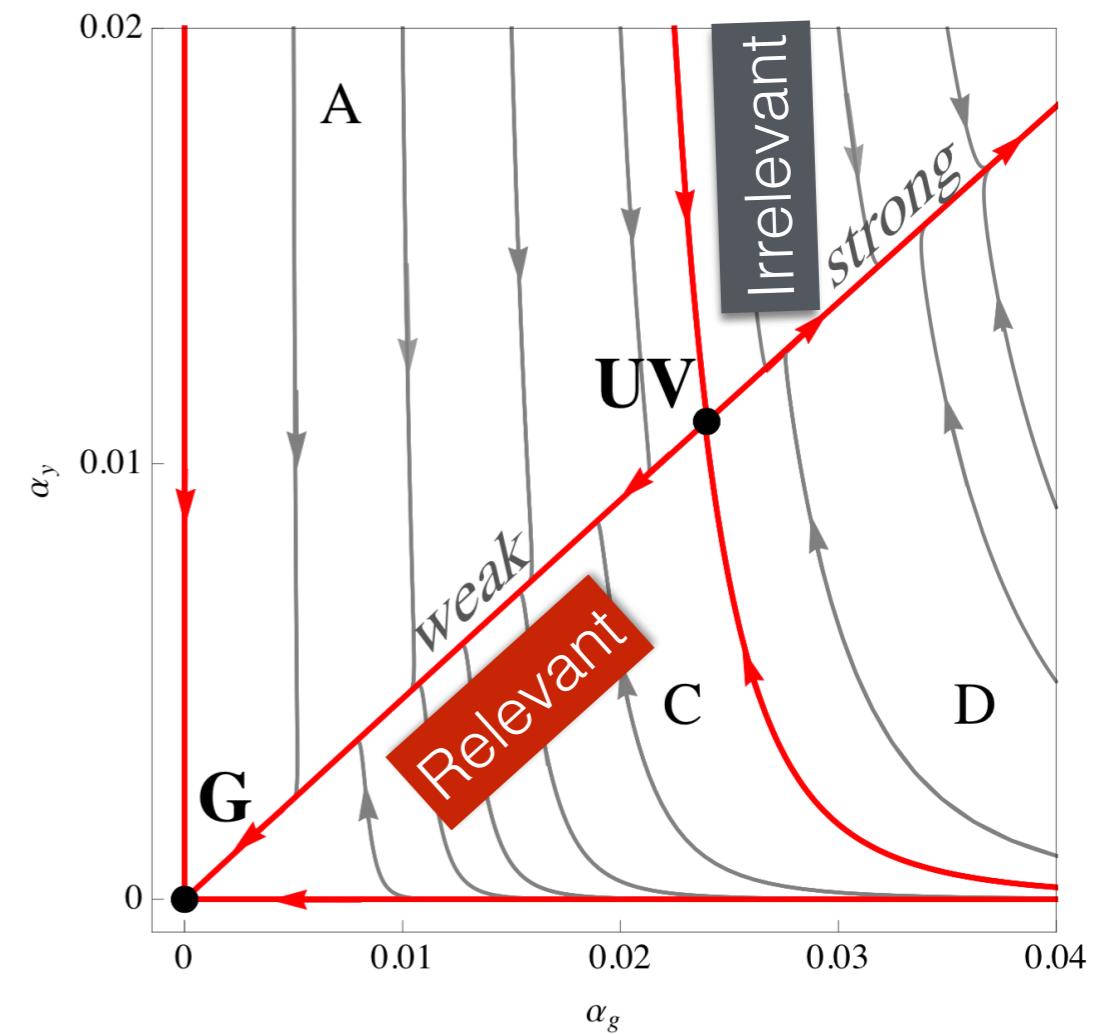
- ◆ Short distance conformality



Fundamental theory

Wilson: A fundamental theory has an UV fixed point

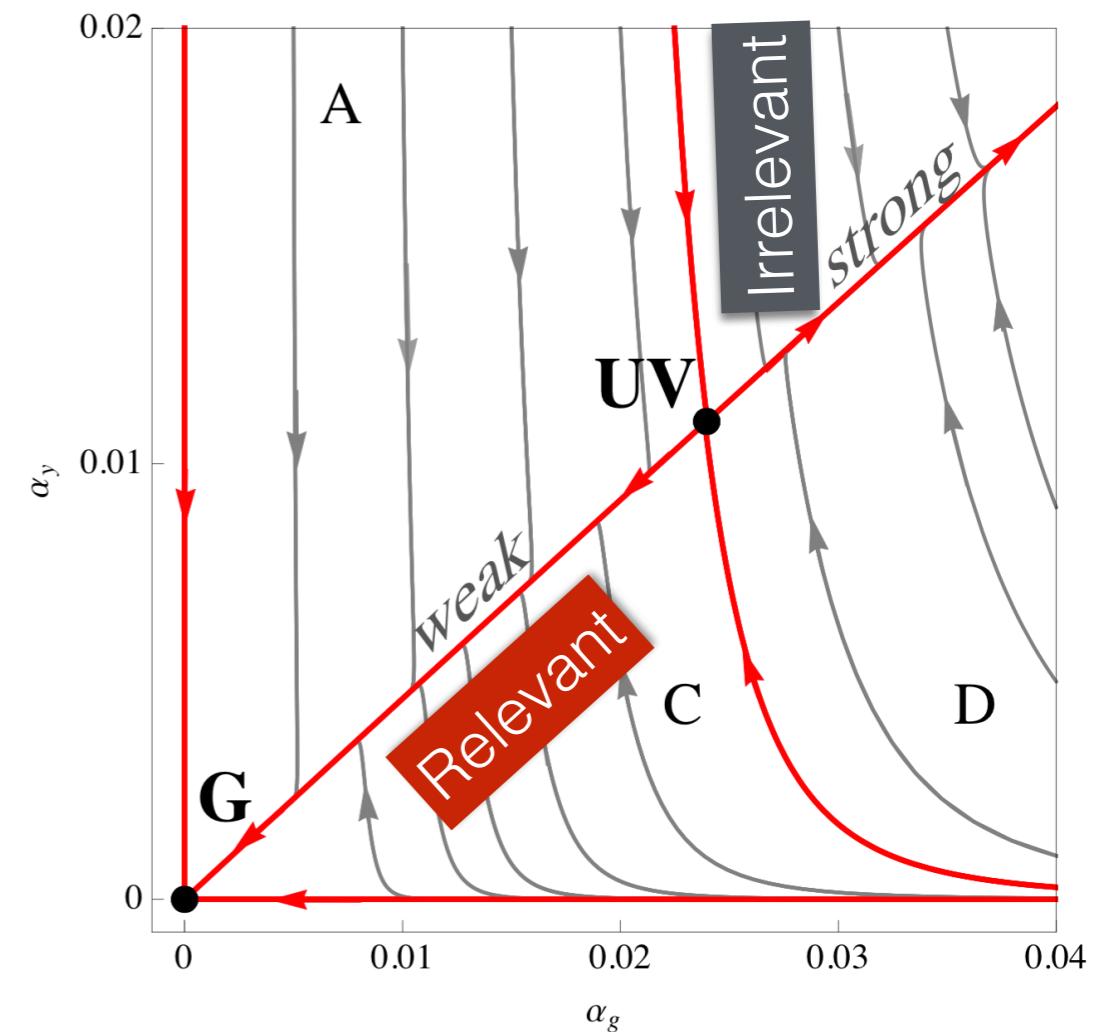
- ◆ Short distance conformality
- ◆ Continuum limit well defined



Fundamental theory

Wilson: A fundamental theory has an UV fixed point

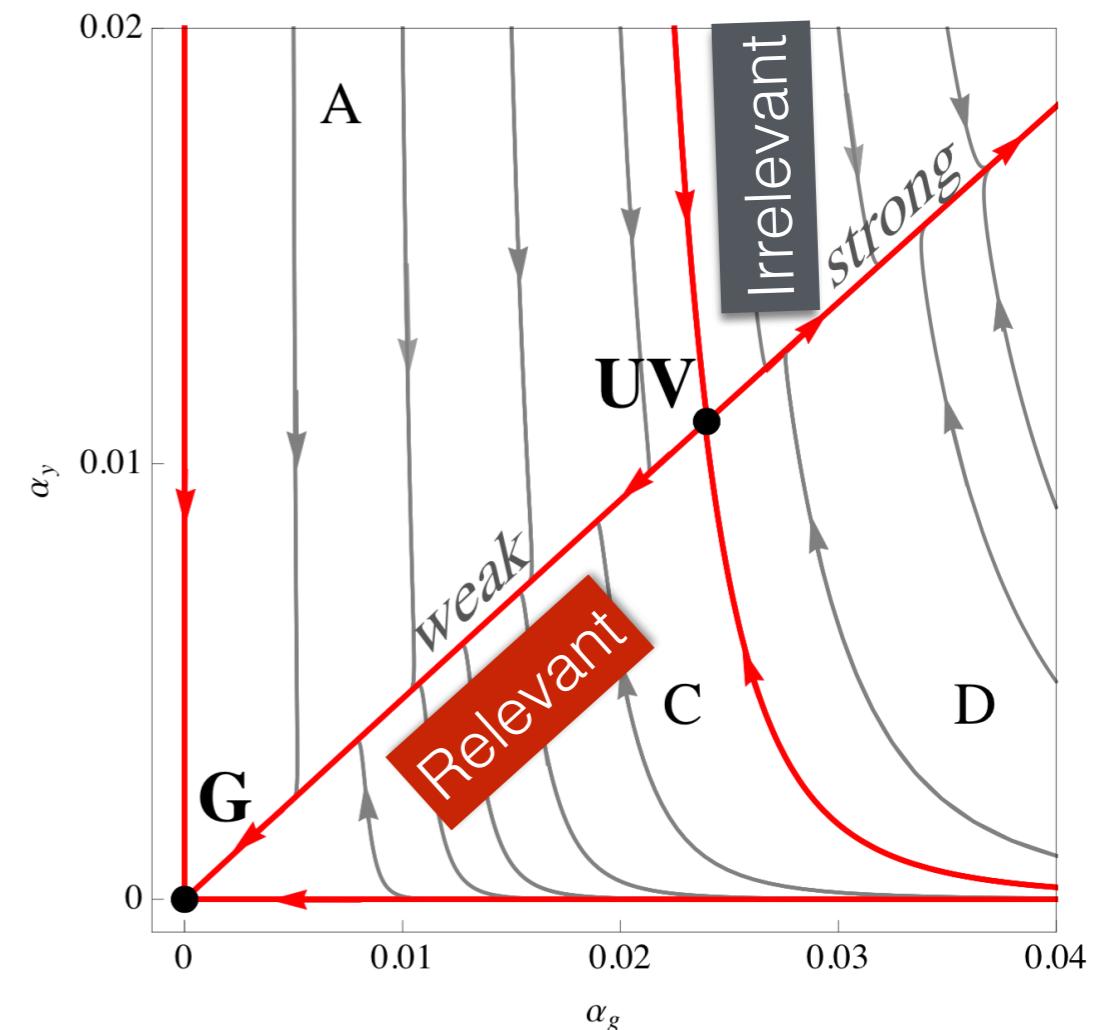
- ◆ Short distance conformality
- ◆ Continuum limit well defined
- ◆ Complete UV fixed point



Fundamental theory

Wilson: A fundamental theory has an UV fixed point

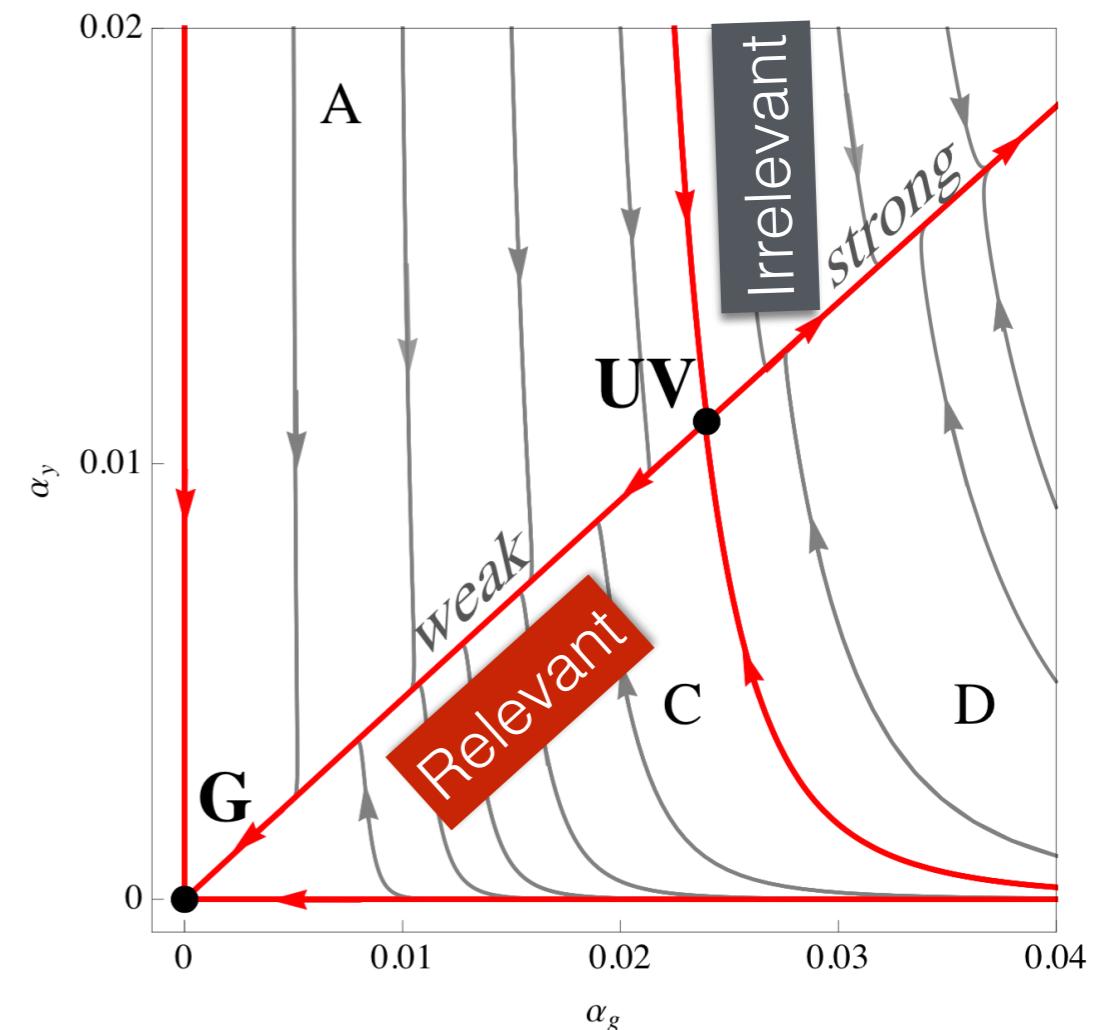
- ◆ Short distance conformality
- ◆ Continuum limit well defined
- ◆ Complete UV fixed point
- ◆ Smaller critical surface dim. = more IR predictiveness



Fundamental theory

Wilson: A fundamental theory has an UV fixed point

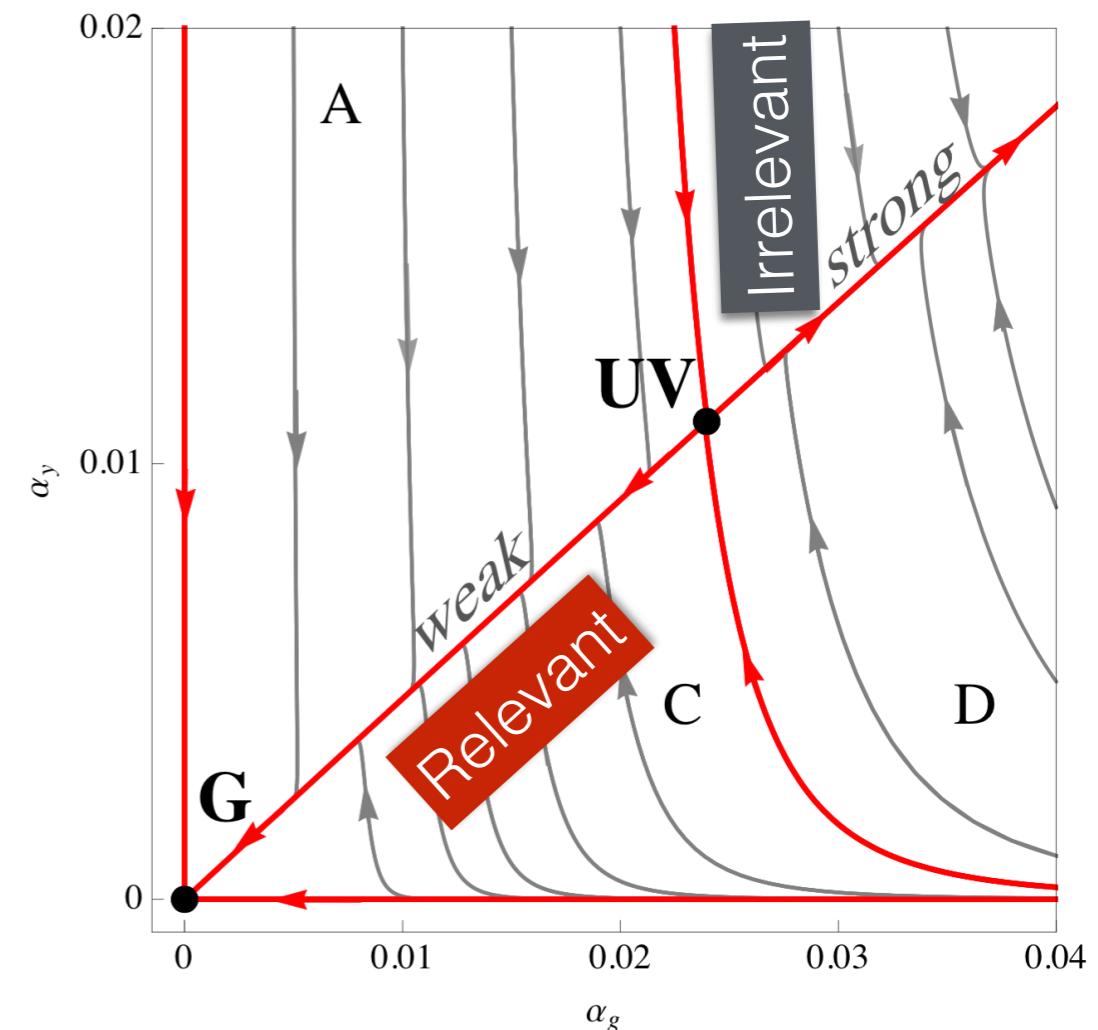
- ◆ Short distance conformality
- ◆ Continuum limit well defined
- ◆ Complete UV fixed point
- ◆ Smaller critical surface dim. = more IR predictiveness
- ◆ Mass operators relevant only for IR



Fundamental theory

Wilson: A fundamental theory has an UV fixed point

- ◆ Short distance conformality
- ◆ Continuum limit well defined
- ◆ Complete UV fixed point
- ◆ Smaller critical surface dim. = more IR predictiveness
- ◆ Mass operators relevant only for IR



The Standard Model is not a fundamental theory

Asymptotic Freedom

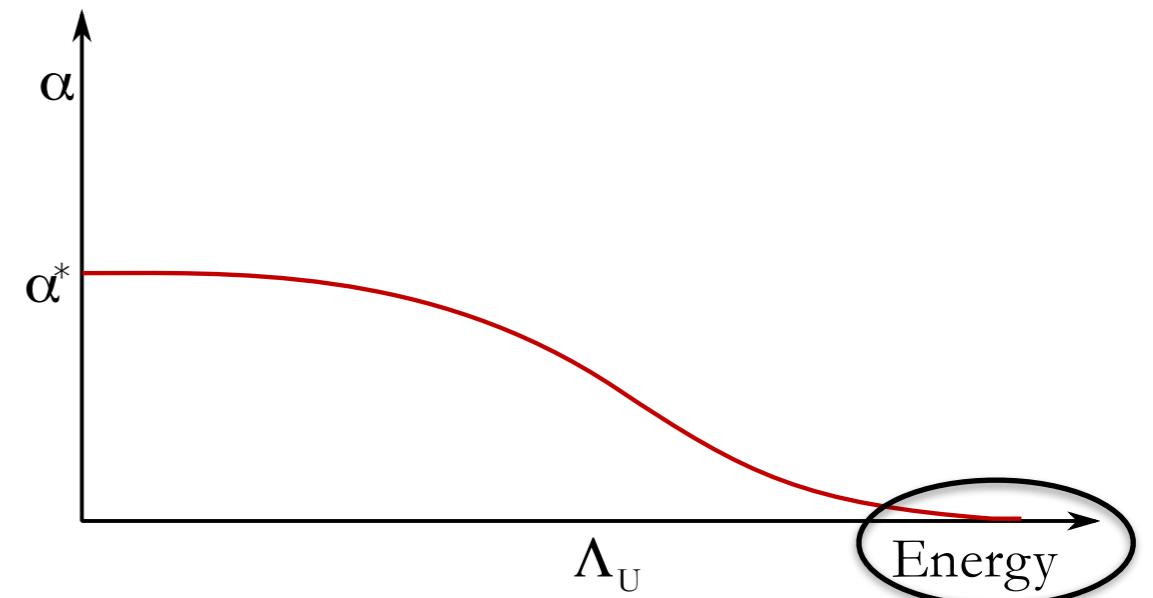
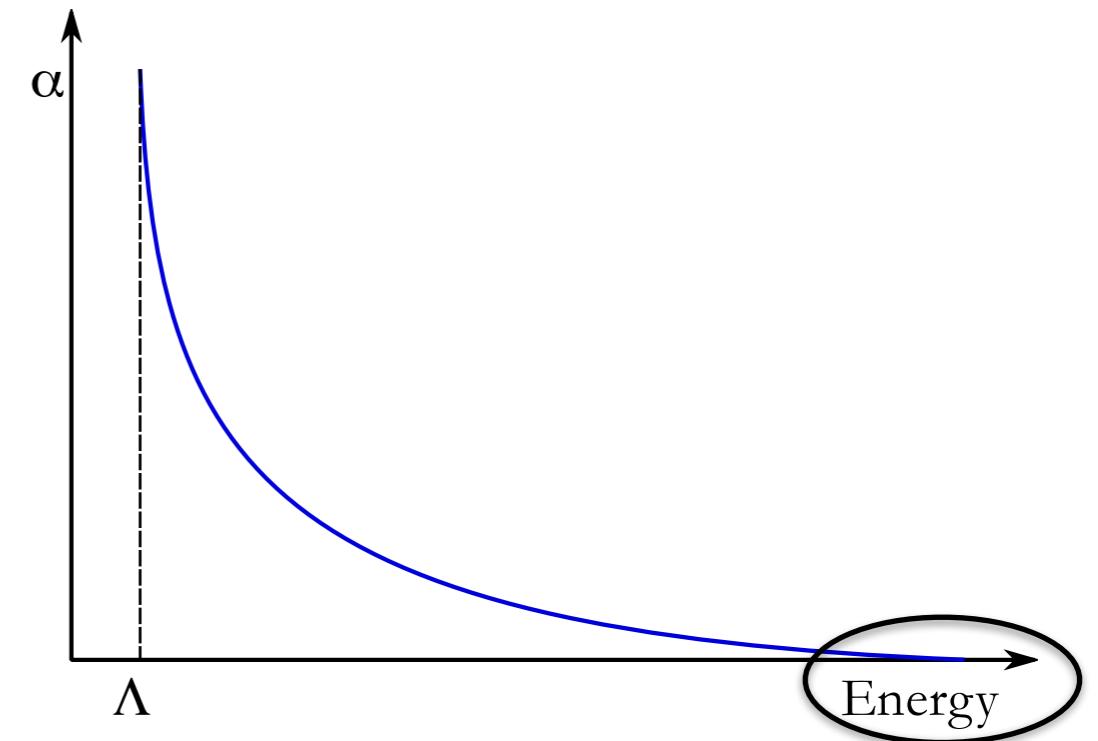
Asymptotic Freedom

Trivial UV fixed point

Asymptotic Freedom

Trivial UV fixed point

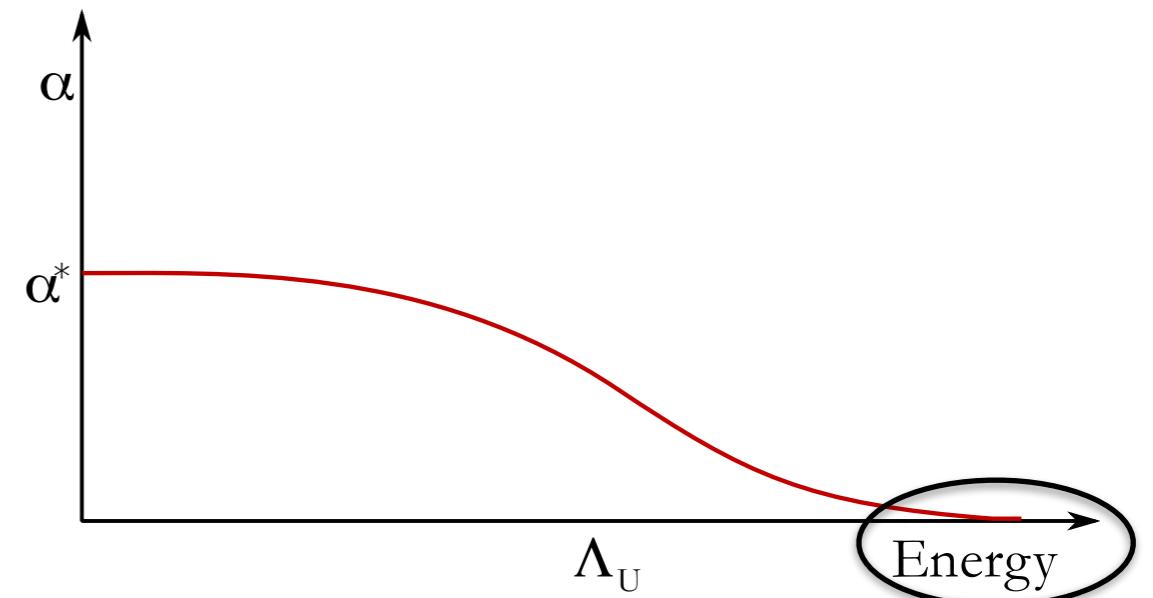
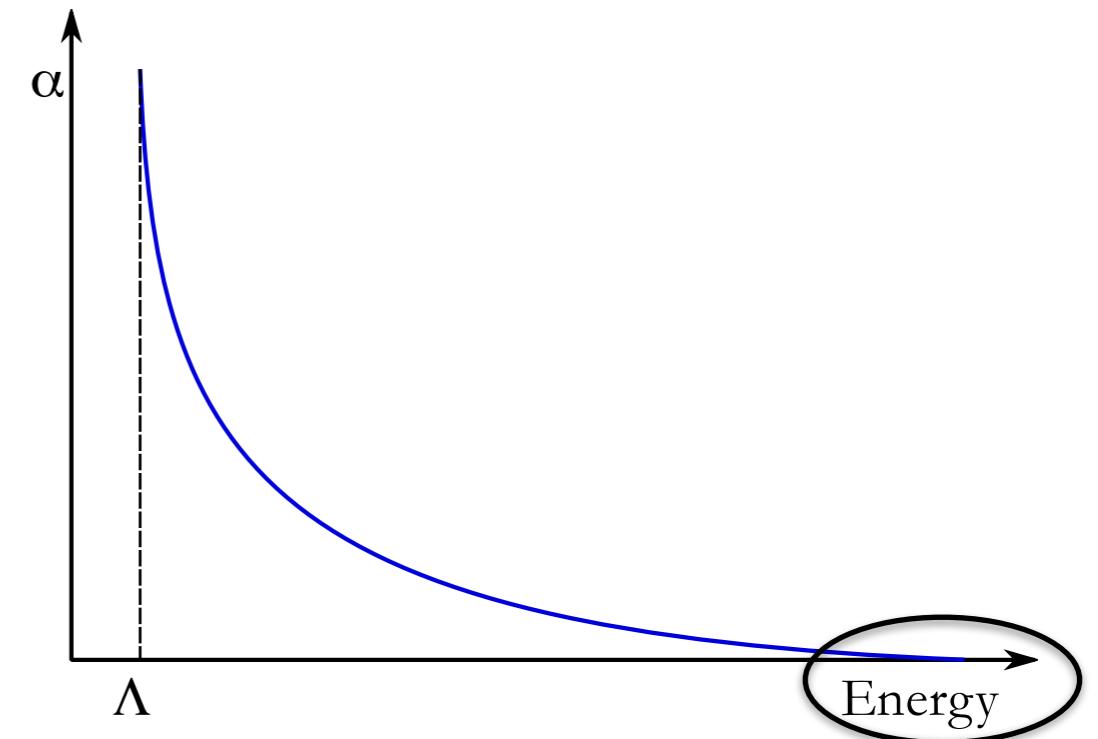
- Non-interacting in the UV



Asymptotic Freedom

Trivial UV fixed point

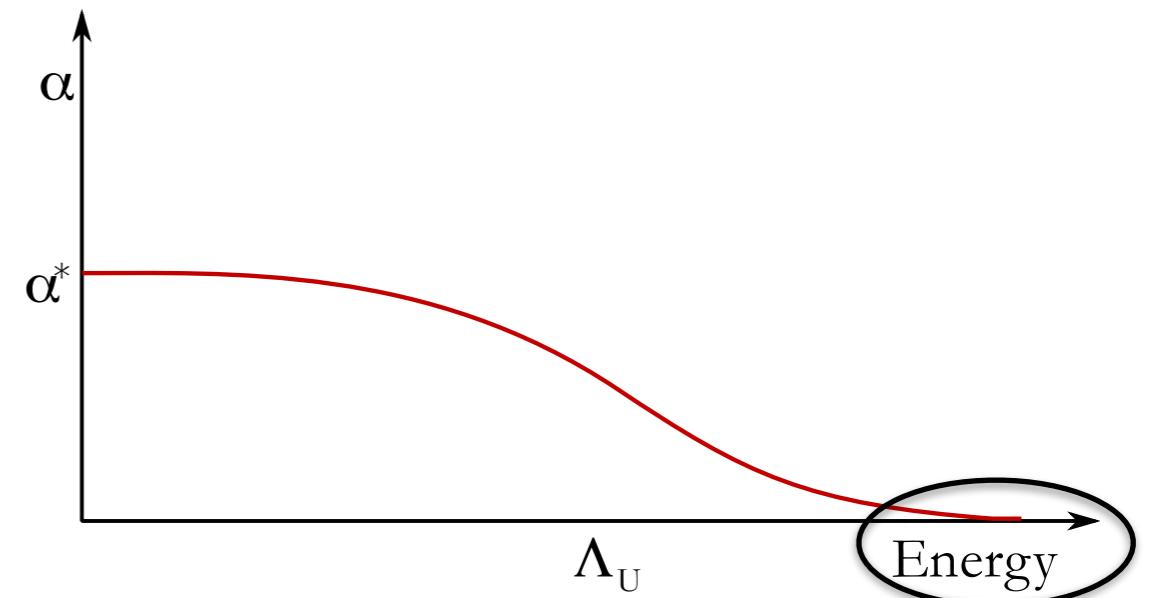
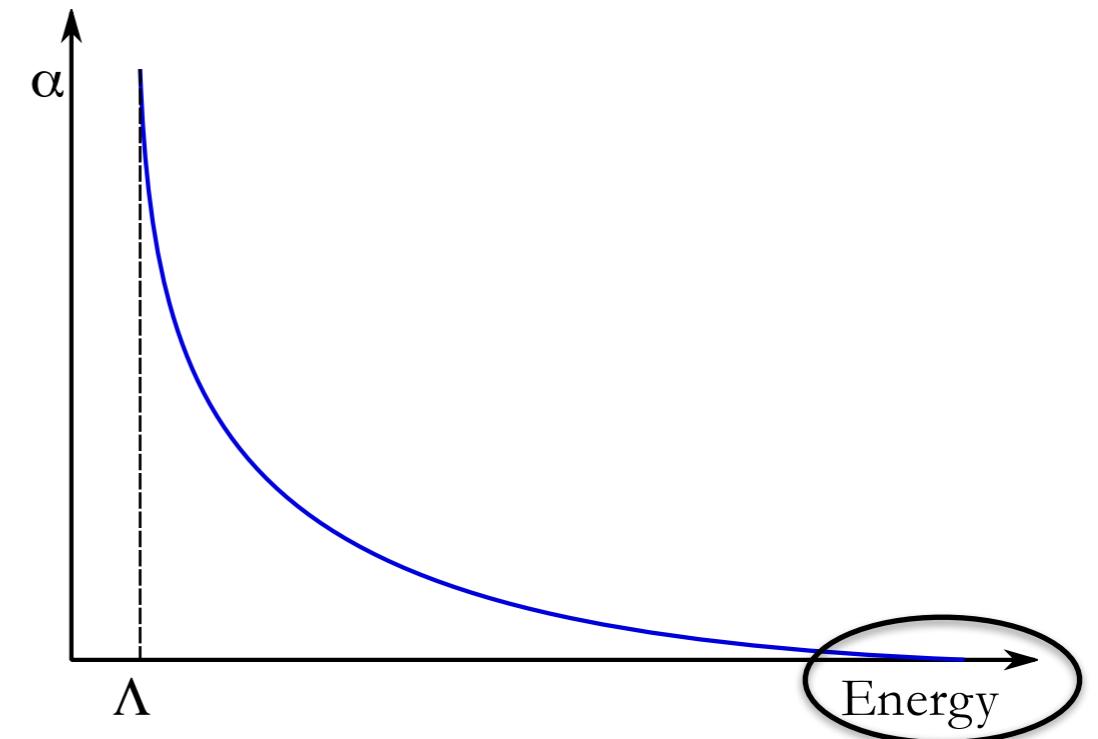
- ◆ Non-interacting in the UV
- ◆ UV logarithmic approach



Asymptotic Freedom

Trivial UV fixed point

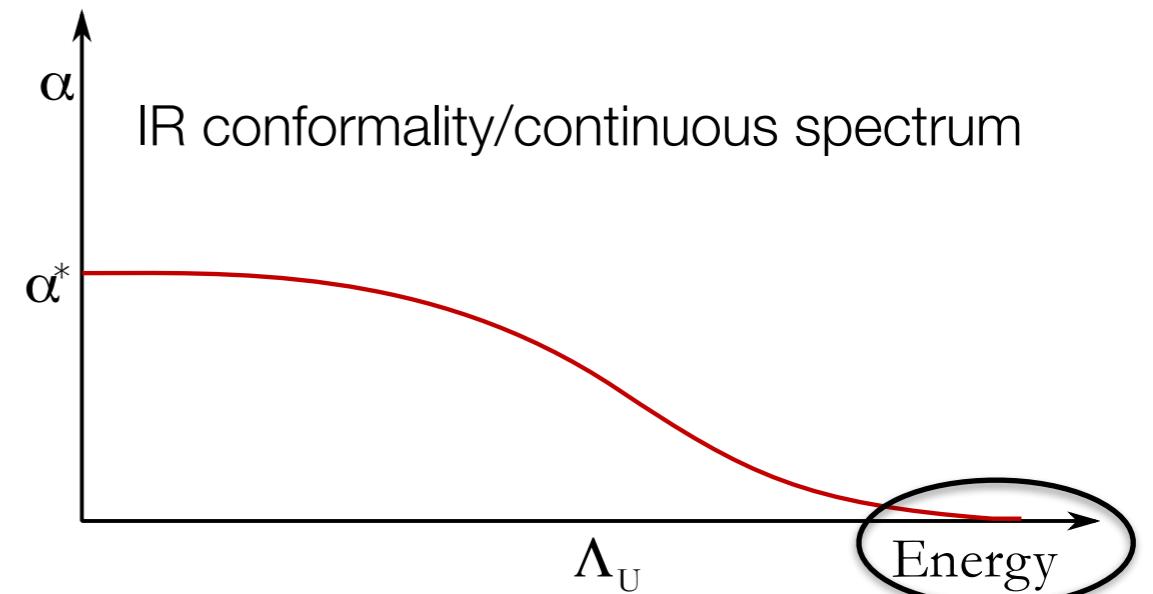
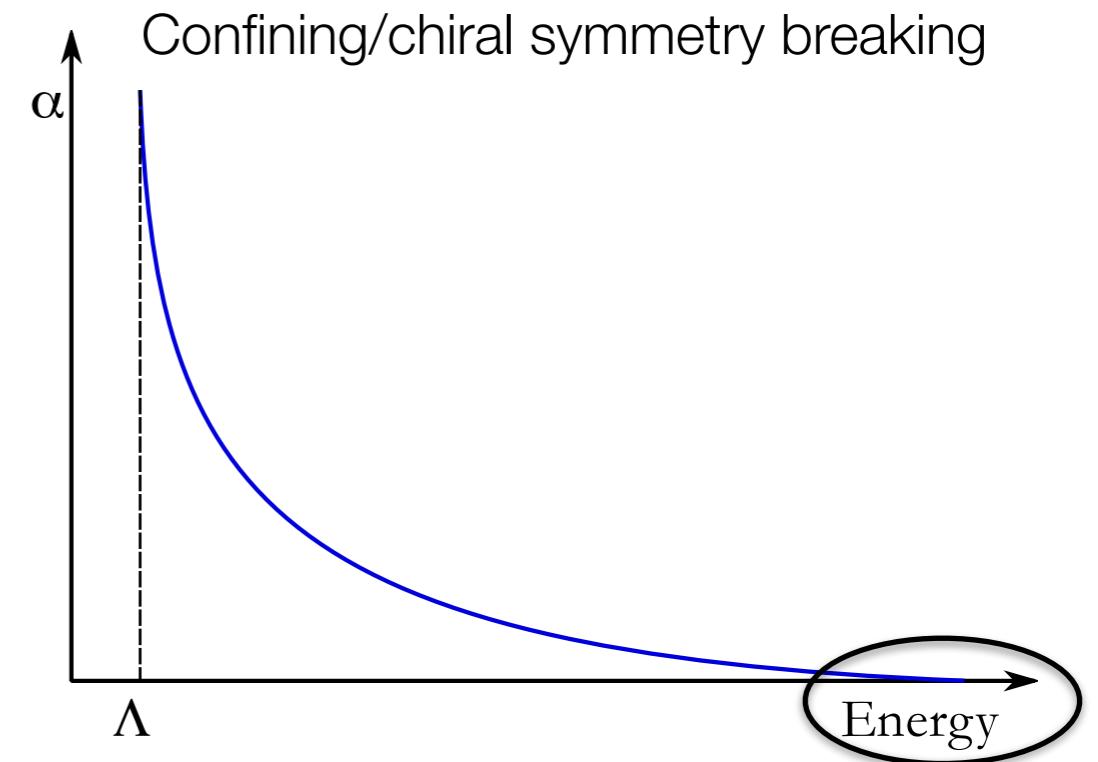
- ◆ Non-interacting in the UV
- ◆ UV logarithmic approach
- ◆ Perturbation theory in UV



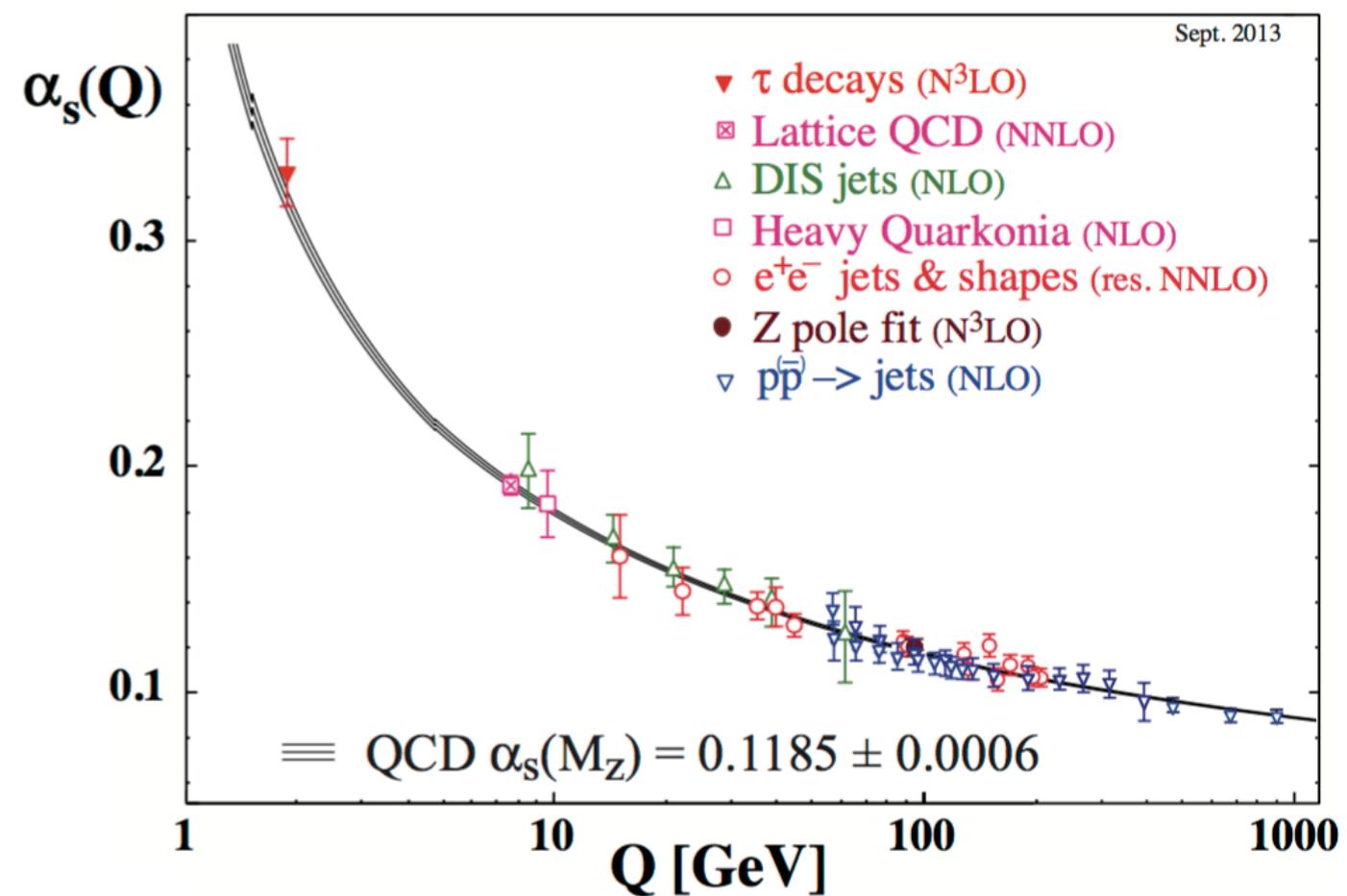
Asymptotic Freedom

Trivial UV fixed point

- ◆ Non-interacting in the UV
- ◆ UV logarithmic approach
- ◆ Perturbation theory in UV
- ◆ IR conformal or dyn. scale

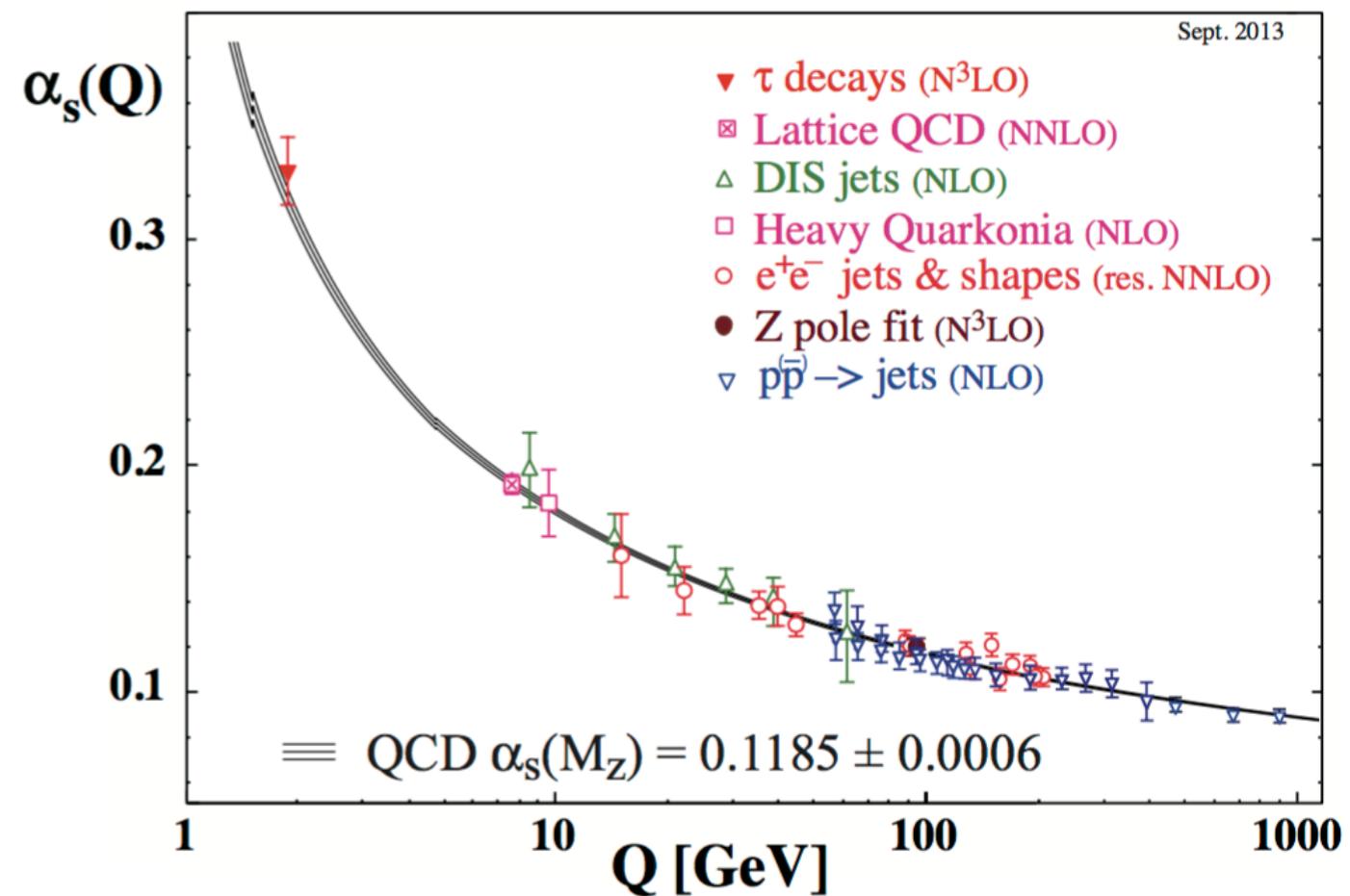


QCD



QCD

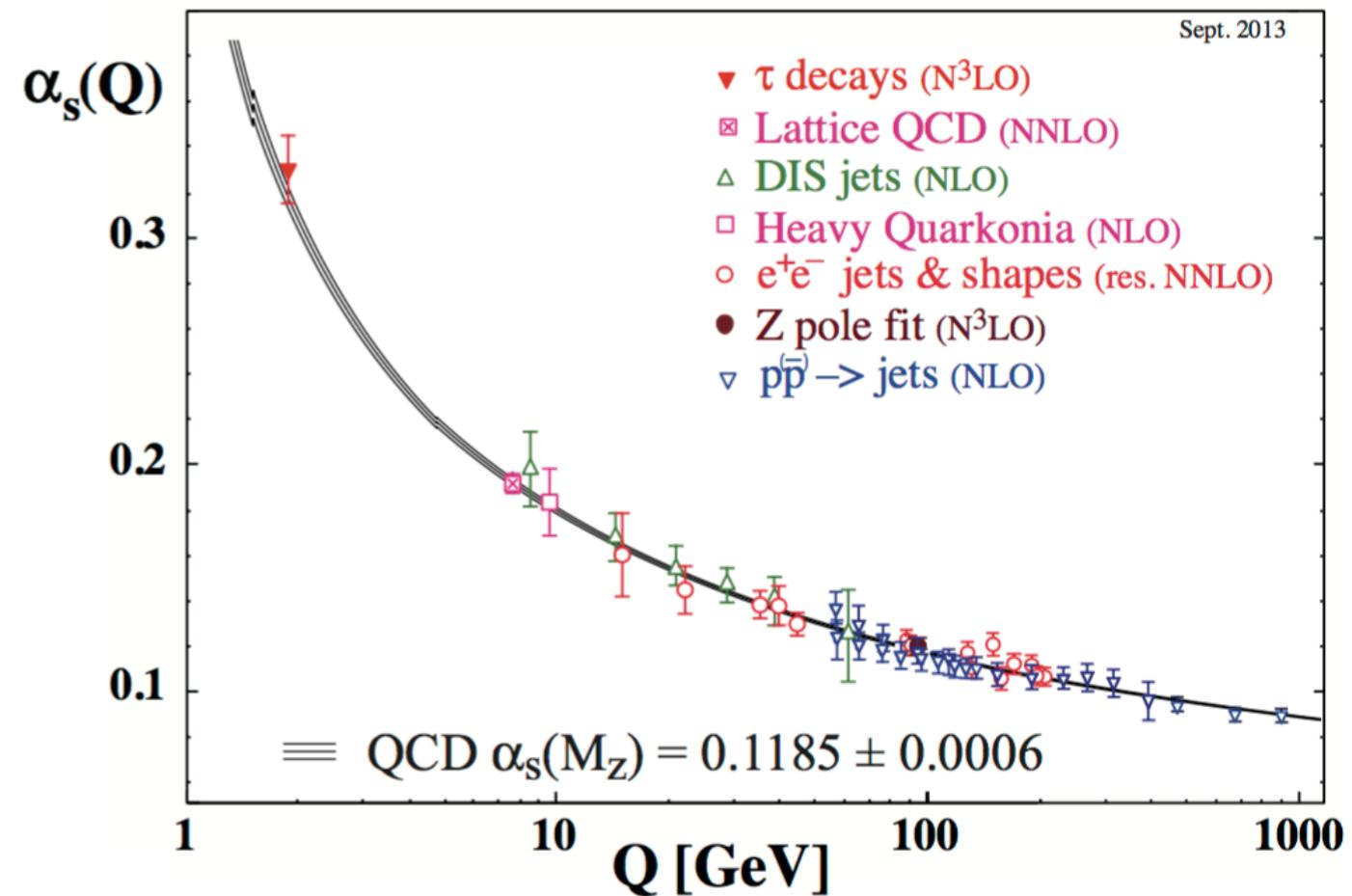
QCD is not IR conformal because



QCD

QCD is not IR conformal because

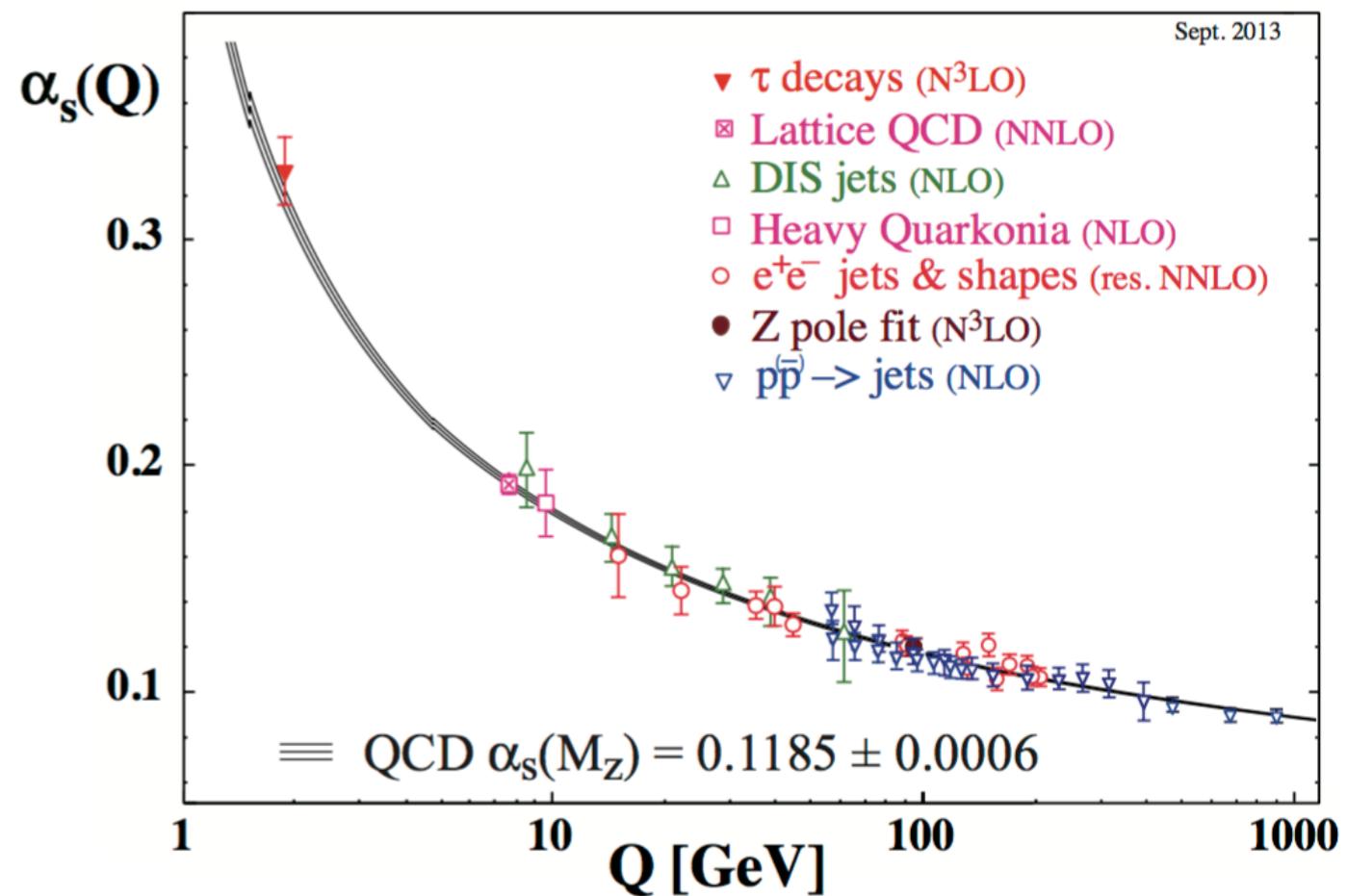
- ◆ Hadronic spectrum/dyn. mass



QCD

QCD is not IR conformal because

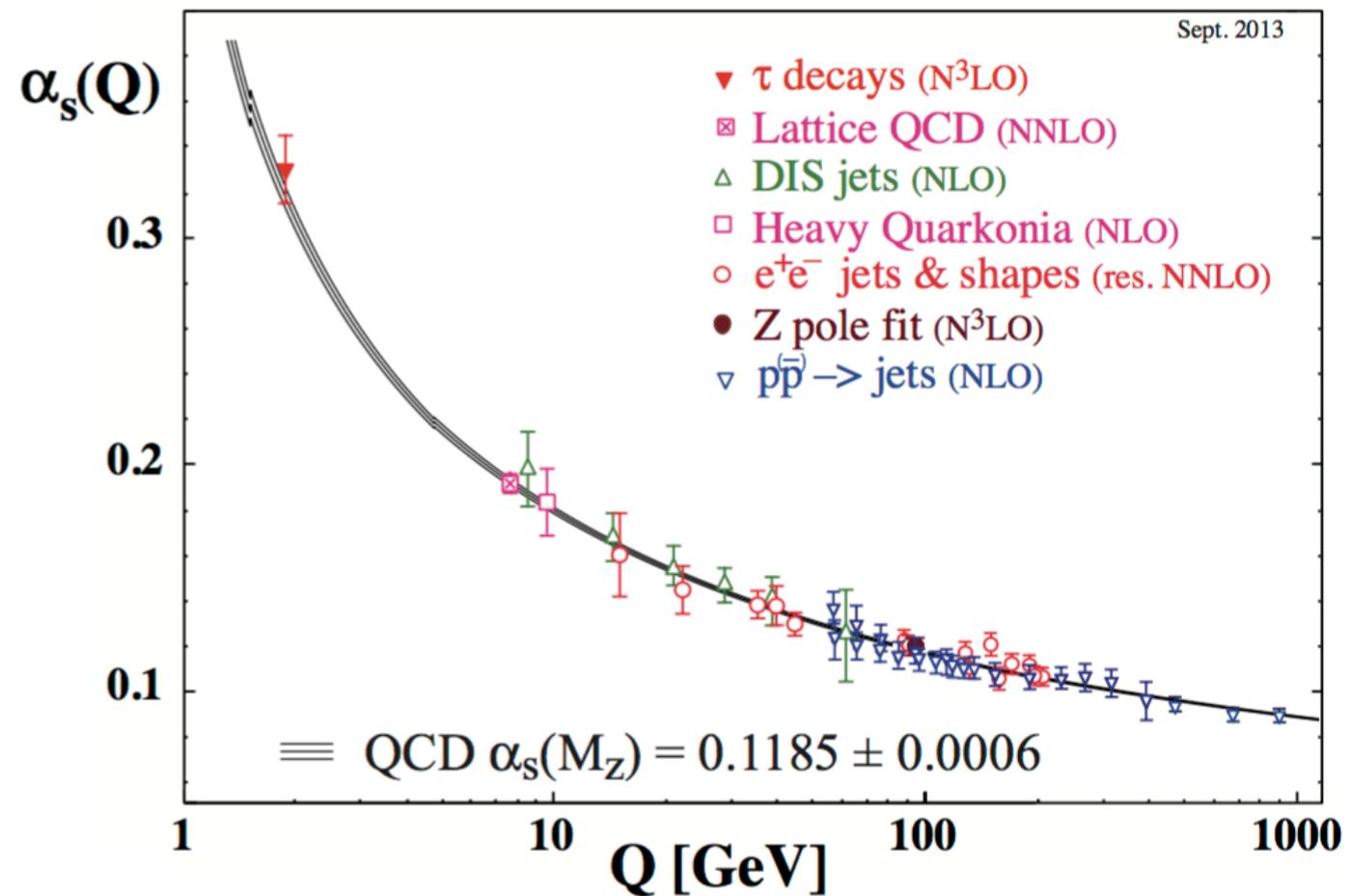
- ◆ Hadronic spectrum/dyn. mass
- ◆ Pions \leftrightarrow Spont. ChSB



QCD

QCD is not IR conformal because

- ◆ Hadronic spectrum/dyn. mass
- ◆ Pions \leftrightarrow Spont. ChSB

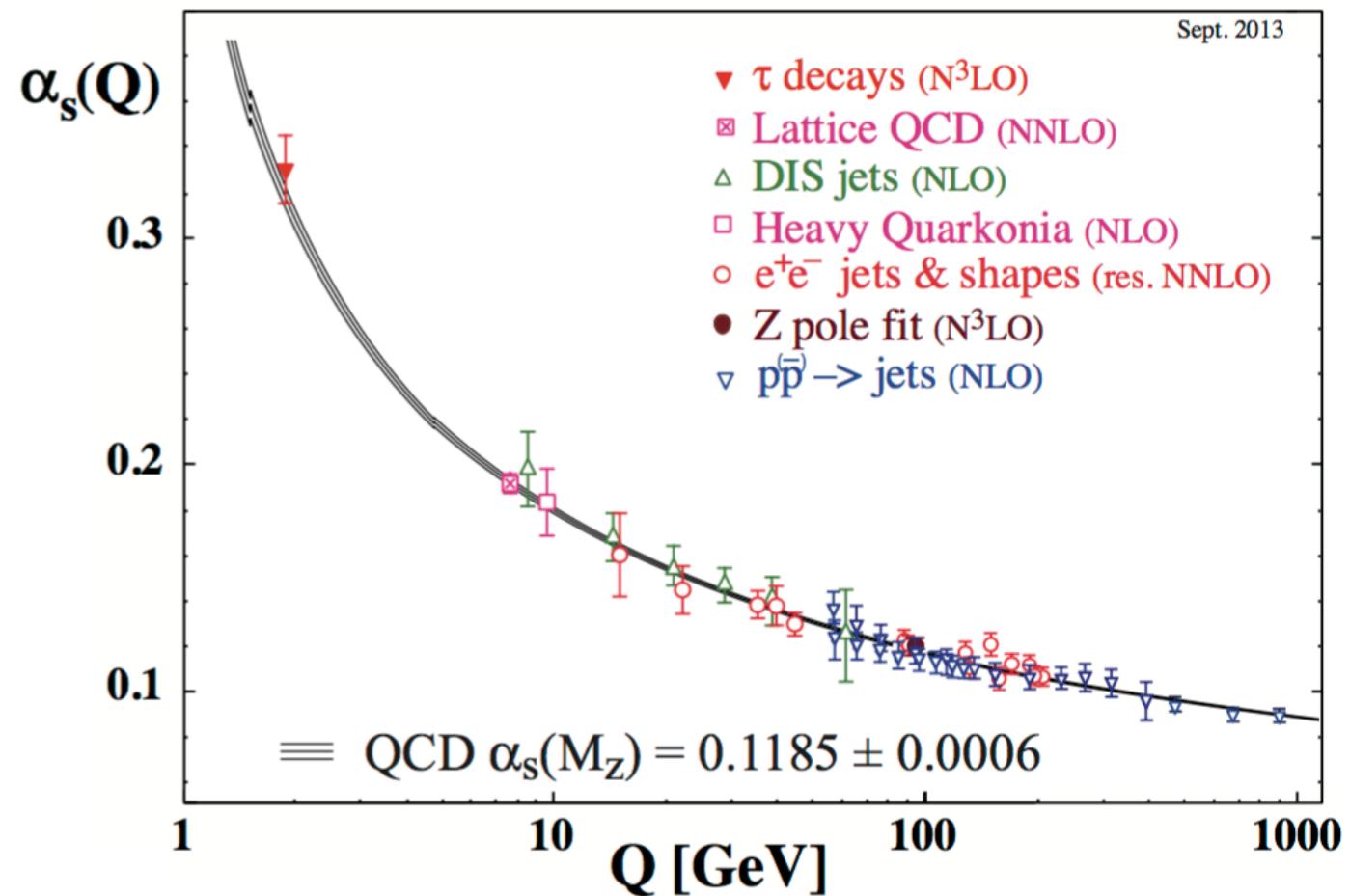


Asymptotic freedom verified $< TeV$

QCD

QCD is not IR conformal because

- ◆ Hadronic spectrum/dyn. mass
- ◆ Pions \leftrightarrow Spont. ChSB



Asymptotic freedom verified $< \text{TeV}$

If above TeV asymptotic freedom is lost, then what?

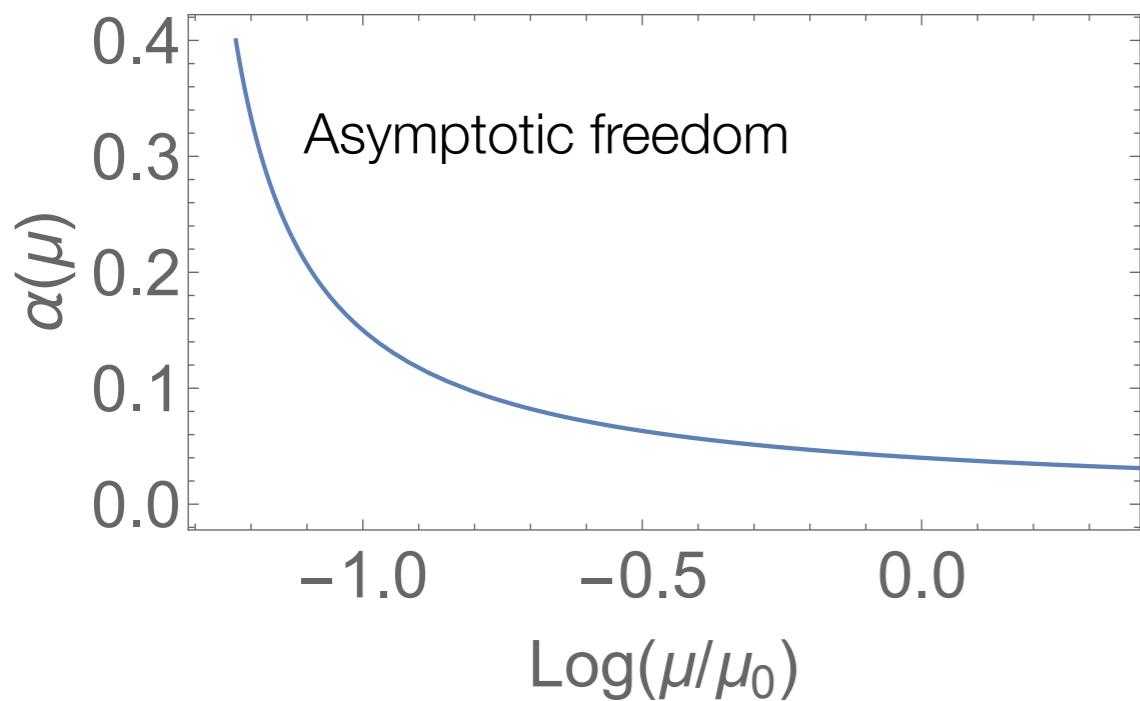
Beyond asymptotic freedom

Asymptotic Safety

Wilson: A fundamental theory has an UV fixed point

Trivial fixed point

- ◆ Non-interacting in the UV
- ◆ Logarithmic scale depend.



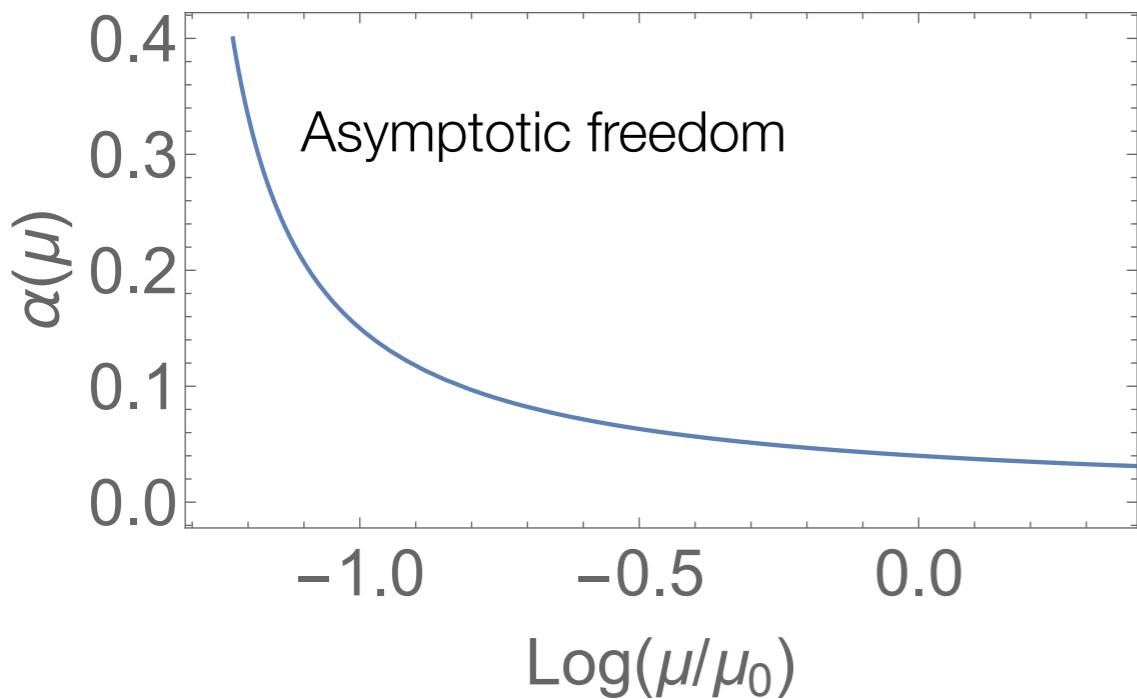
Asymptotic Safety

Wilson: A fundamental theory has an UV fixed point

Trivial fixed point

Interacting fixed point

- ◆ Non-interacting in the UV
- ◆ Logarithmic scale depend.

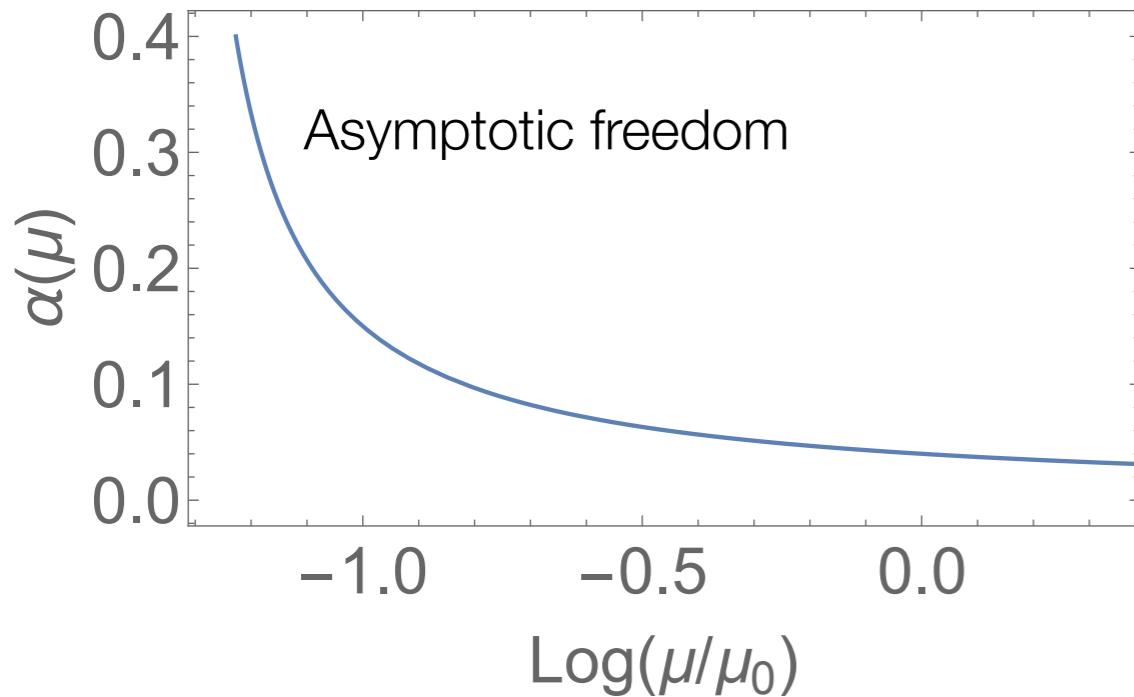


Asymptotic Safety

Wilson: A fundamental theory has an UV fixed point

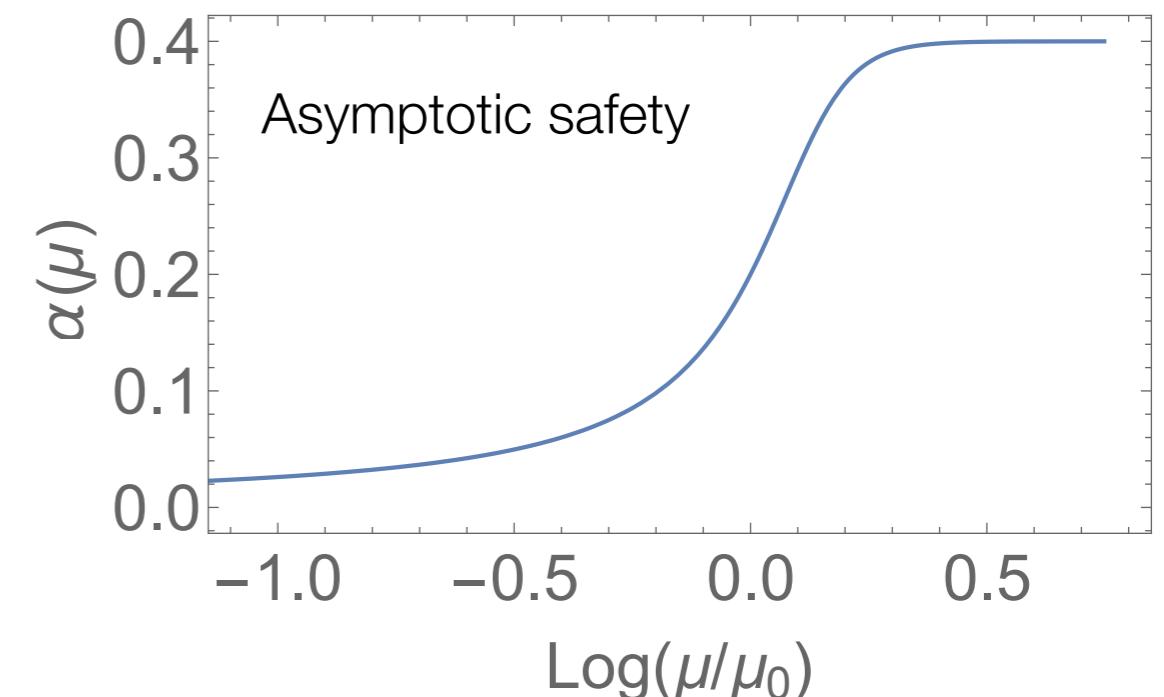
Trivial fixed point

- ◆ Non-interacting in the UV
- ◆ Logarithmic scale depend.

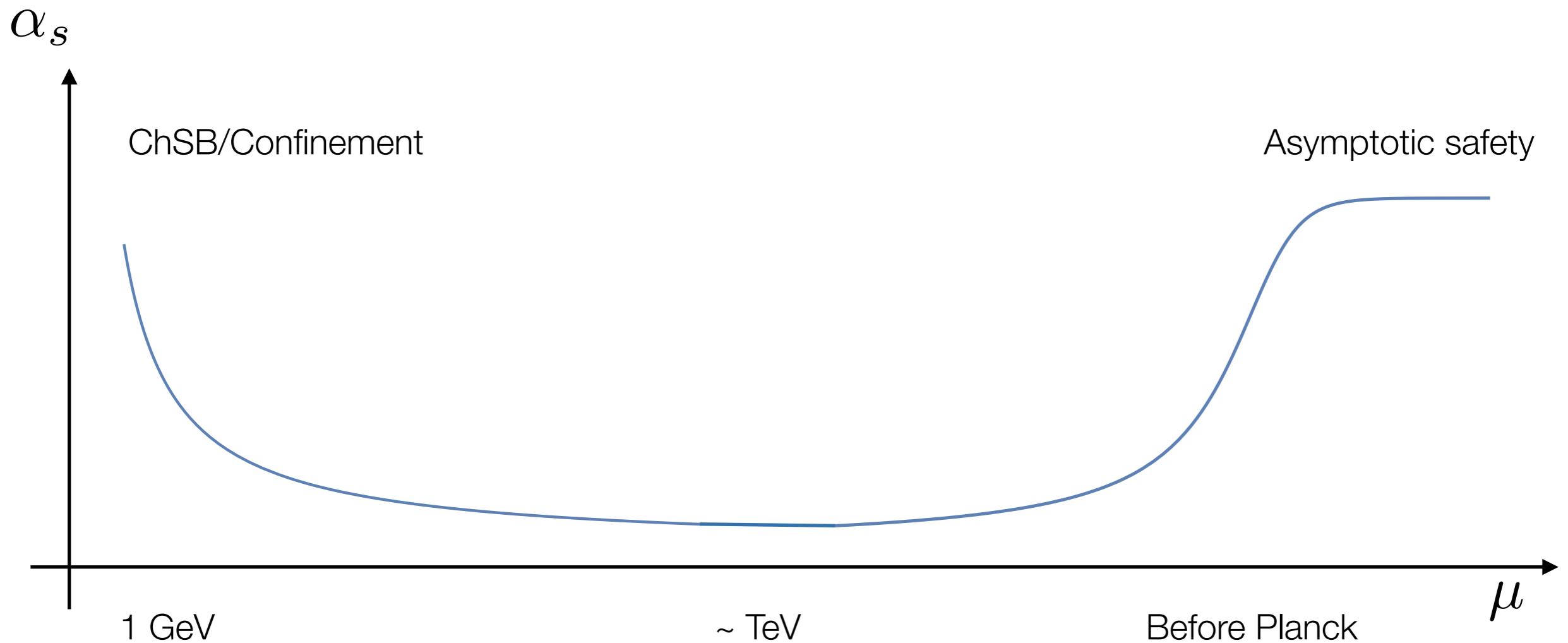


Interacting fixed point

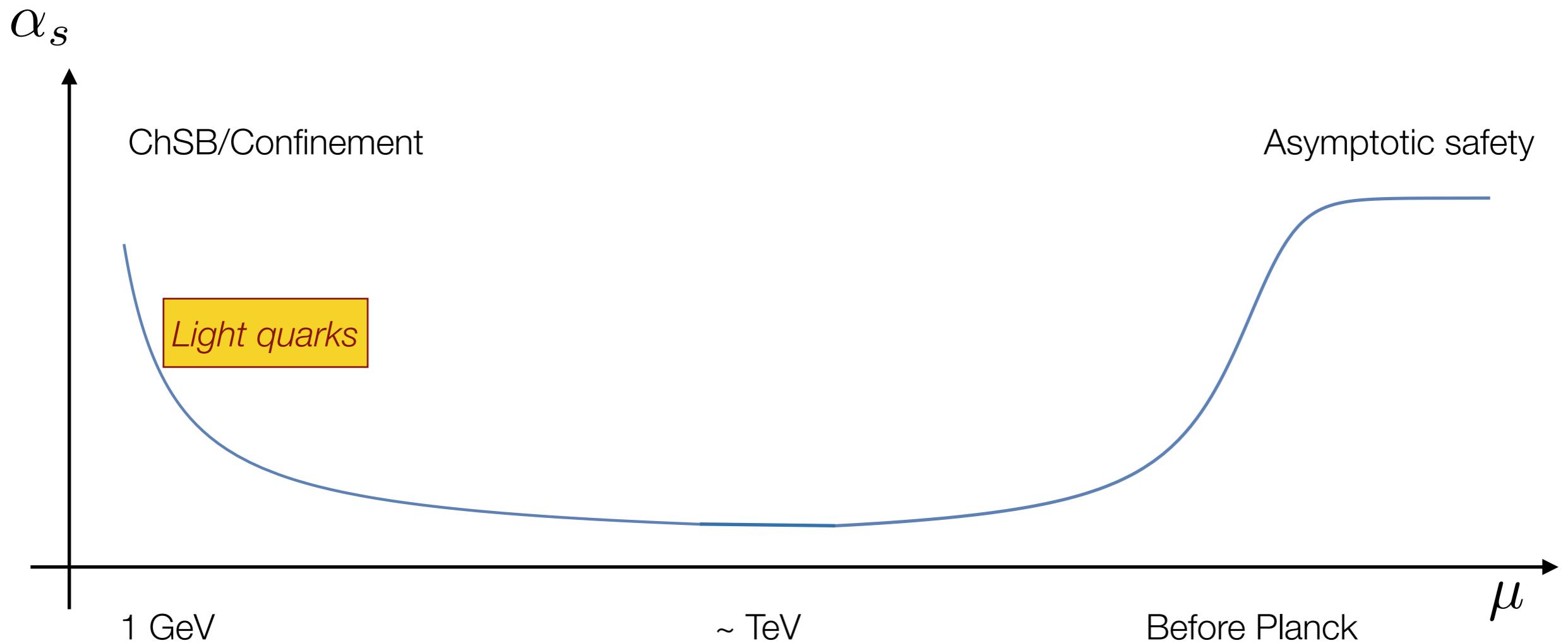
- ◆ Integrating in the UV
- ◆ Power law



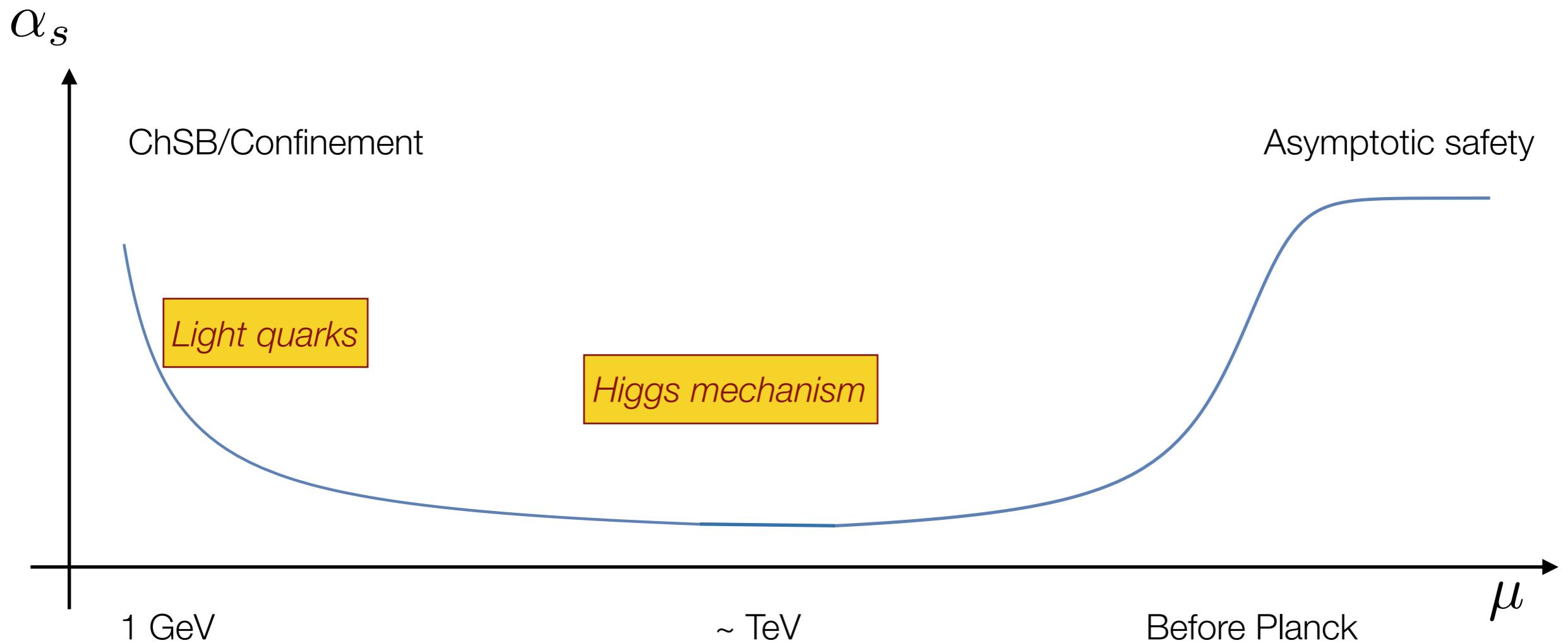
Safe QCD scenario



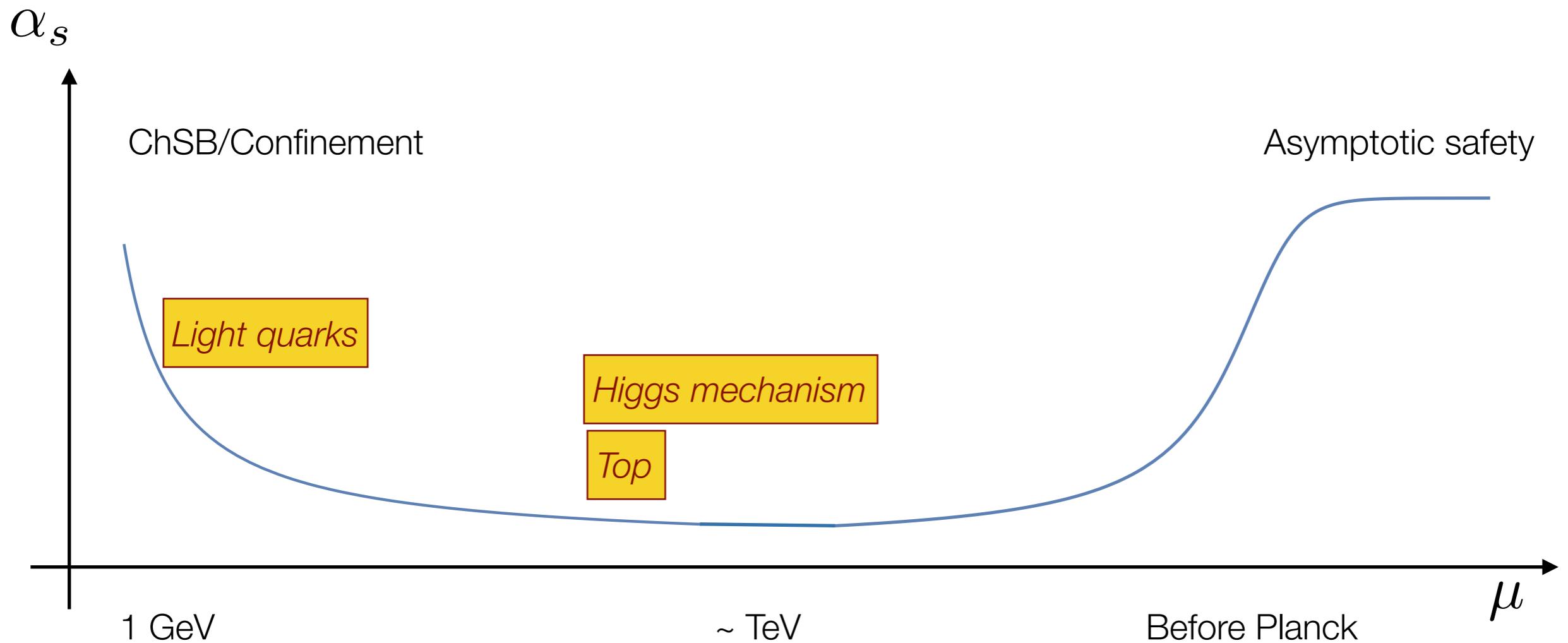
Safe QCD scenario



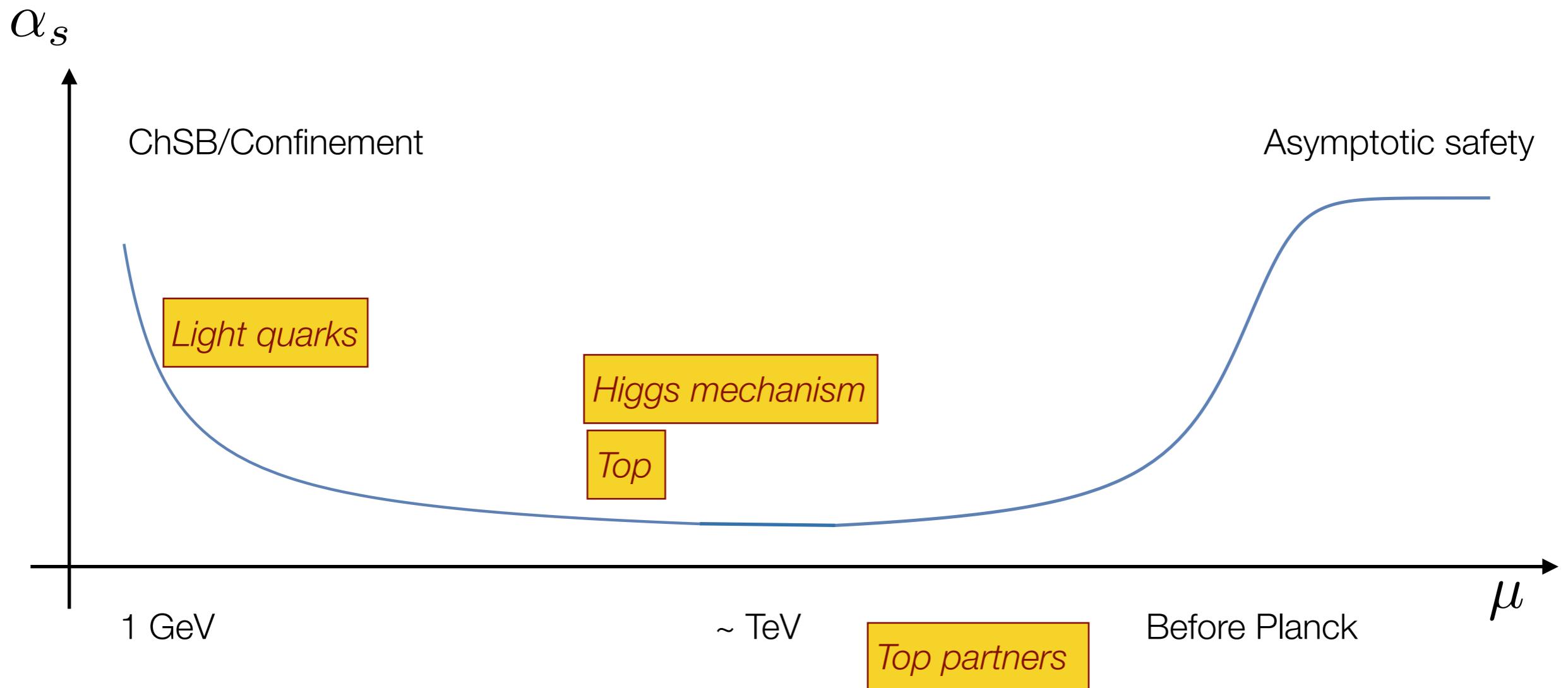
Safe QCD scenario



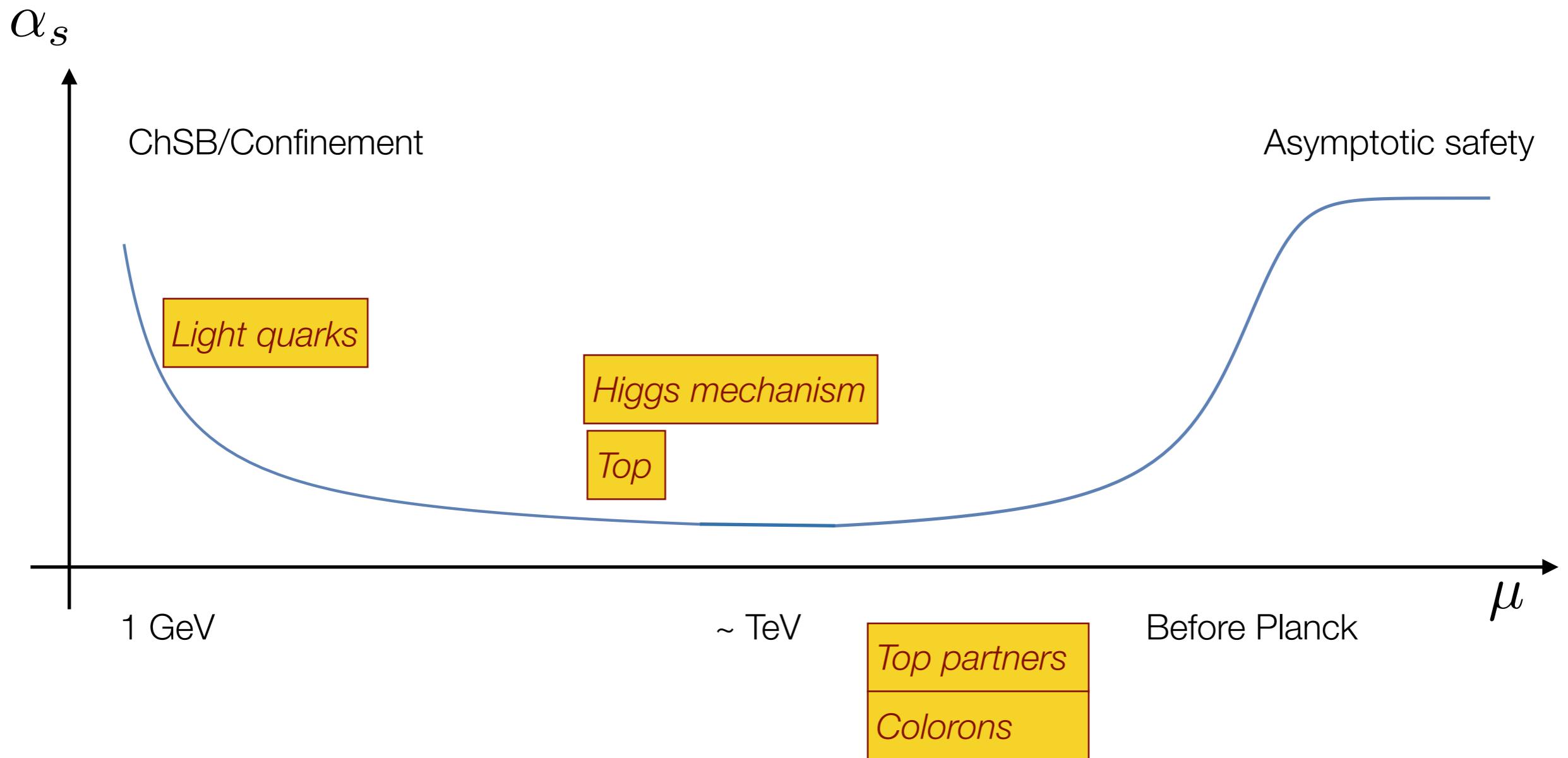
Safe QCD scenario



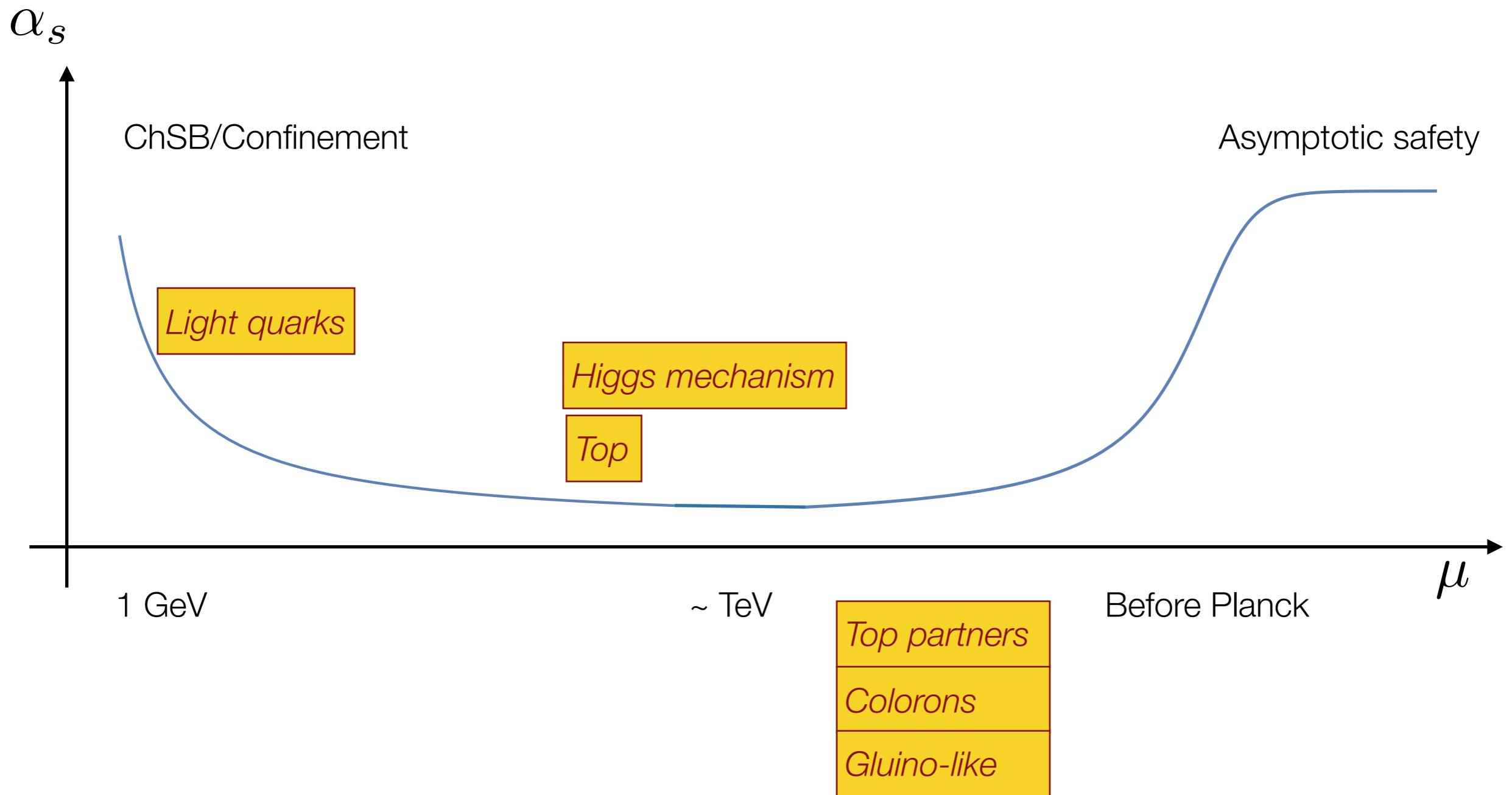
Safe QCD scenario



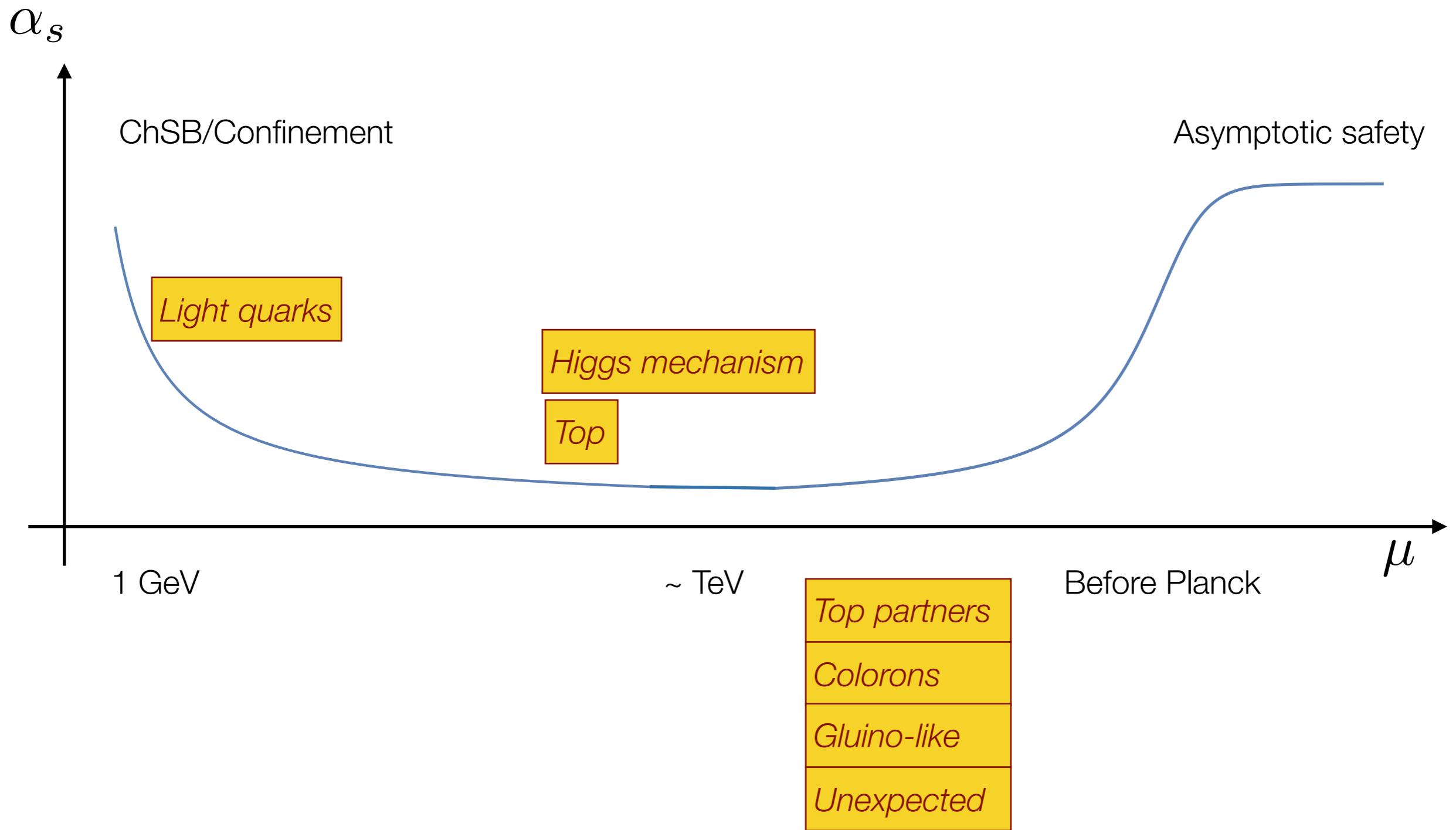
Safe QCD scenario



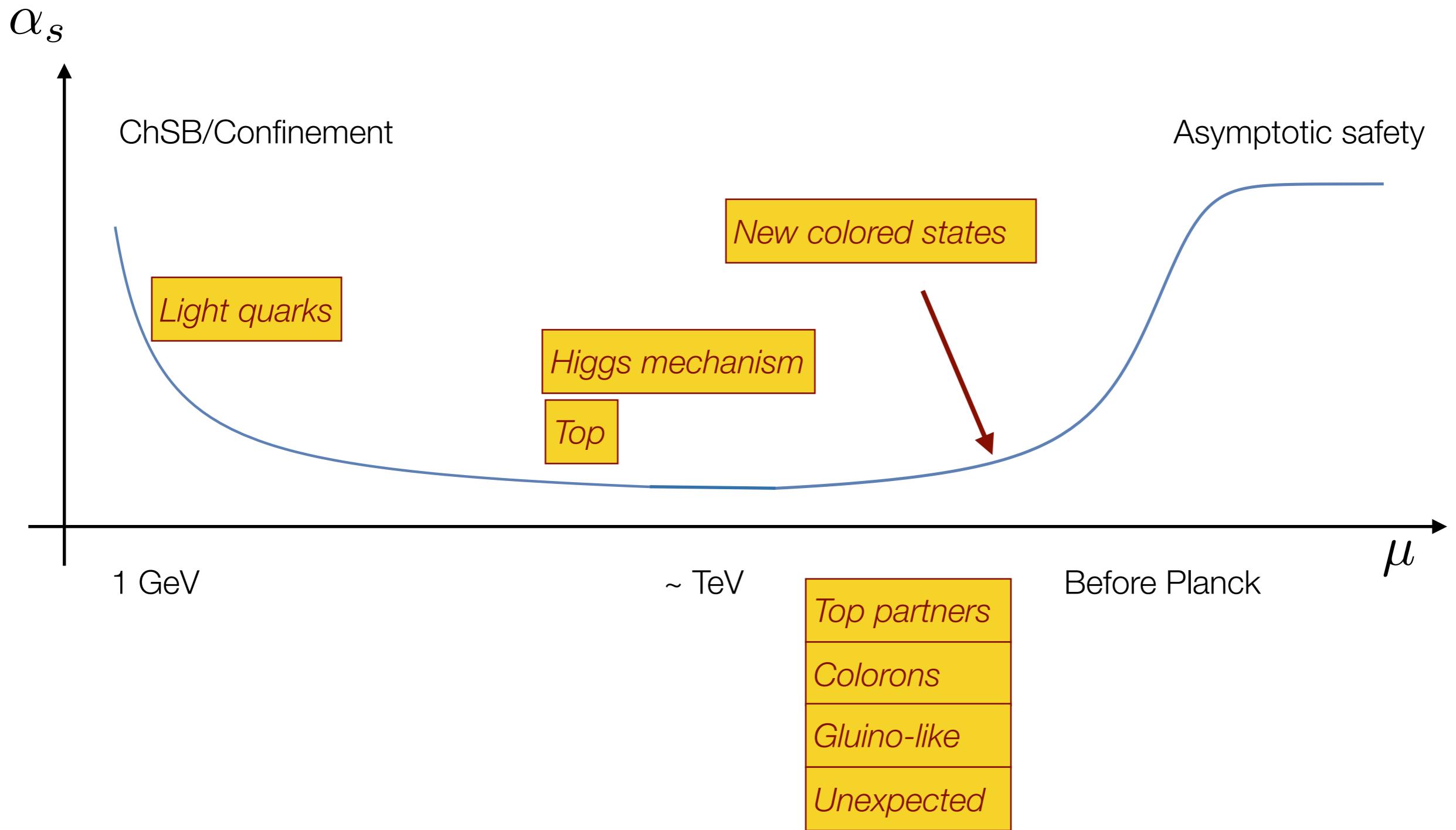
Safe QCD scenario



Safe QCD scenario



Safe QCD scenario



Does a theory like this exists?

Exact 4D Interacting UV Fixed Point

Litim and Sannino, 1406.2337, JHEP

Exact 4D Interacting UV Fixed Point

Litim and Sannino, 1406.2337, JHEP

$$L = -F^2 + i\bar{Q}\gamma \cdot DQ + y(\bar{Q}_L H Q_R + \text{h.c.}) +$$

$$\text{Tr} [\partial H^\dagger \partial H] - u \text{Tr} [(H^\dagger H)^2] - v \text{Tr} [(H^\dagger H)]^2$$

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$
G_μ	Adj	1	1	0
Q_L	\square	$\bar{\square}$	1	1
Q_R^c	$\bar{\square}$	1	\square	-1
H	1	\square	$\bar{\square}$	0

Veneziano Limit

- ◆ Normalised couplings

$$\alpha_g = \frac{g^2 N_C}{(4\pi)^2}, \quad \alpha_y = \frac{y^2 N_C}{(4\pi)^2}, \quad \alpha_h = \frac{u N_F}{(4\pi)^2}, \quad \alpha_v = \frac{v N_F^2}{(4\pi)^2}$$

$$\frac{v}{u} = \frac{\alpha_v}{\alpha_h N_F}$$

At large N $\frac{N_F}{N_C} \in \Re^+$

Non-Asymptotically Free

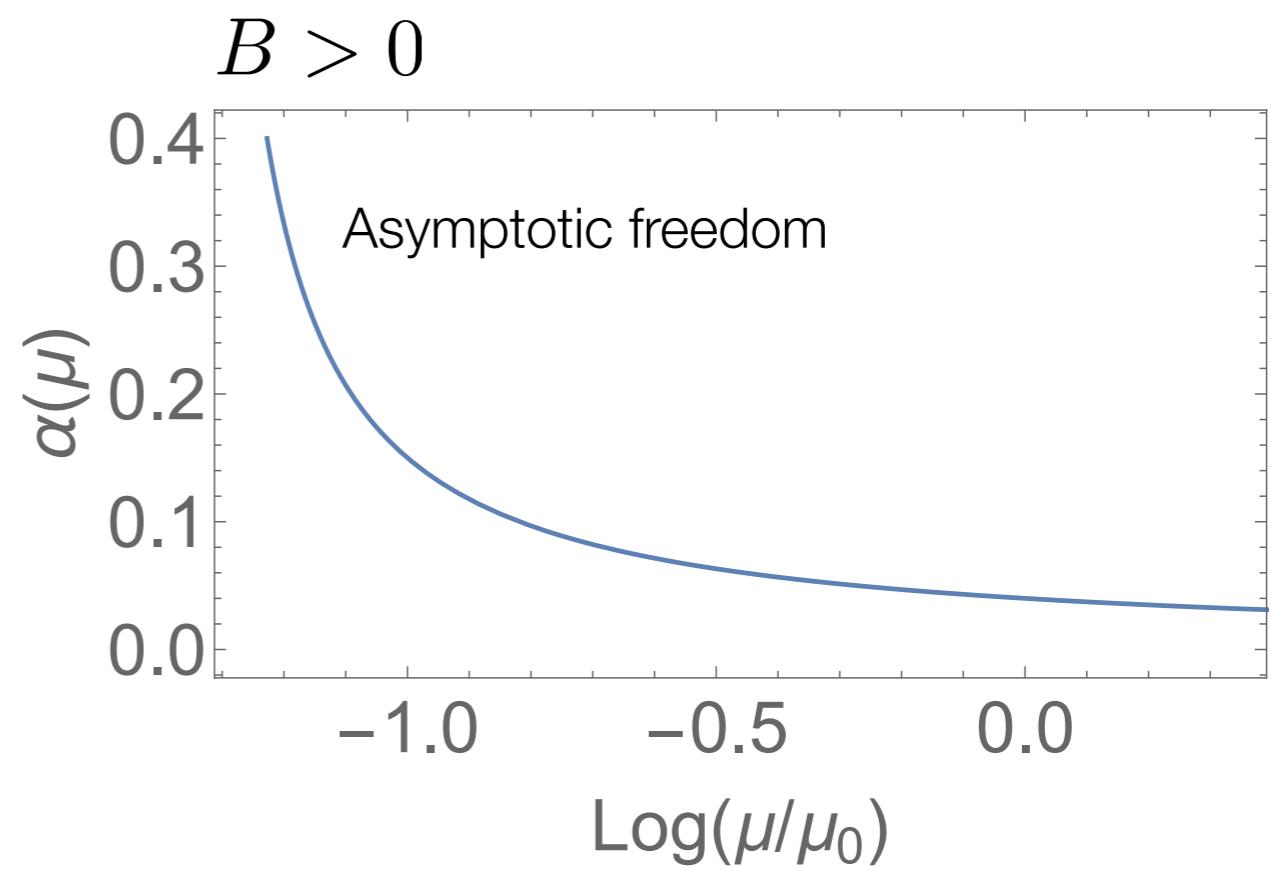
$$\beta_g = \partial_t \alpha_g = -B\alpha_g^2$$

$$t=\ln\frac{\mu}{\mu_0}$$

Non-Asymptotically Free

$$\beta_g = \partial_t \alpha_g = -B \alpha_g^2$$

$$t = \ln \frac{\mu}{\mu_0}$$

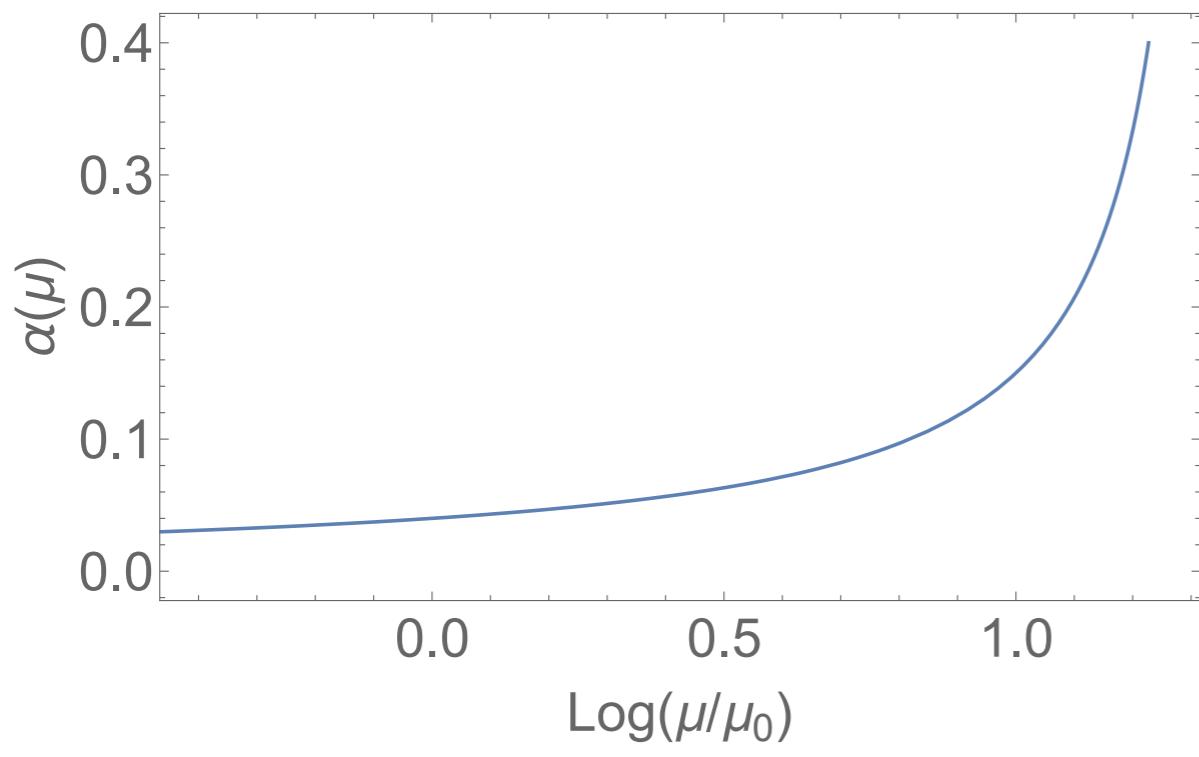


Non-Asymptotically Free

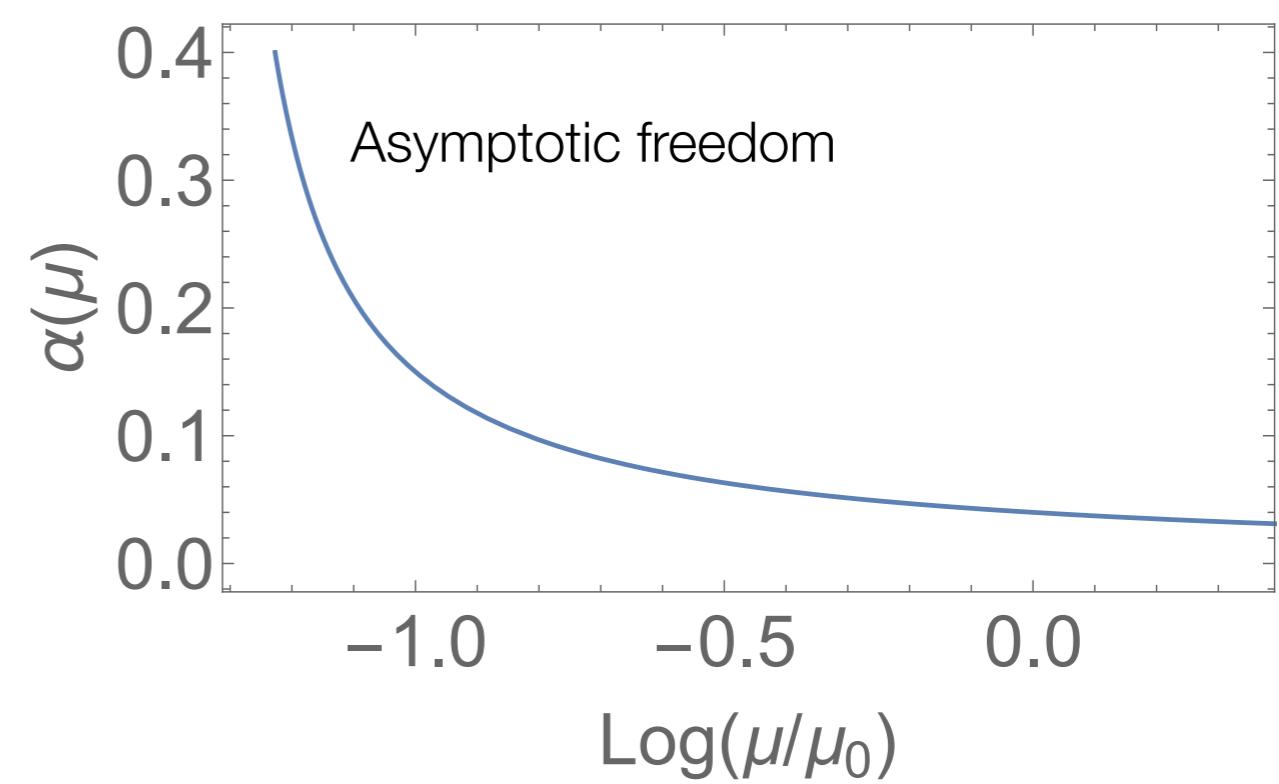
$$\beta_g = \partial_t \alpha_g = -B \alpha_g^2$$

$$t = \ln \frac{\mu}{\mu_0}$$

$$B < 0$$



$$B > 0$$

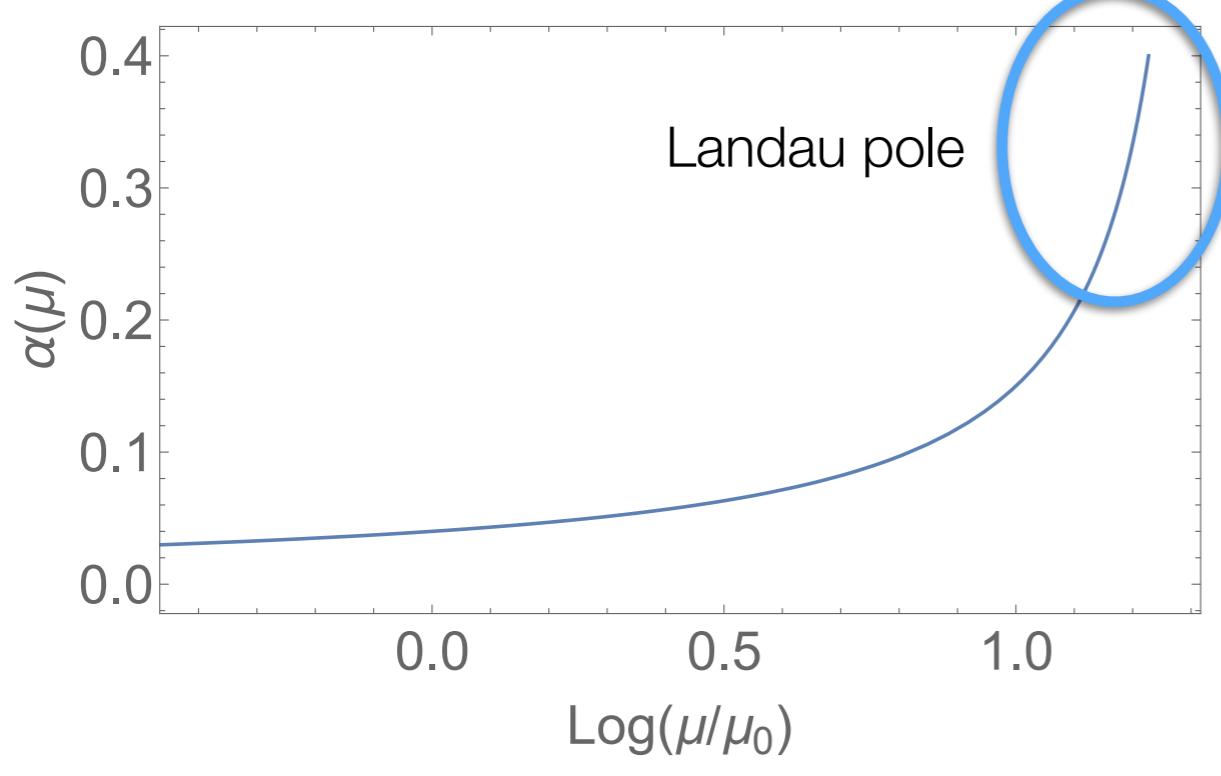


Non-Asymptotically Free

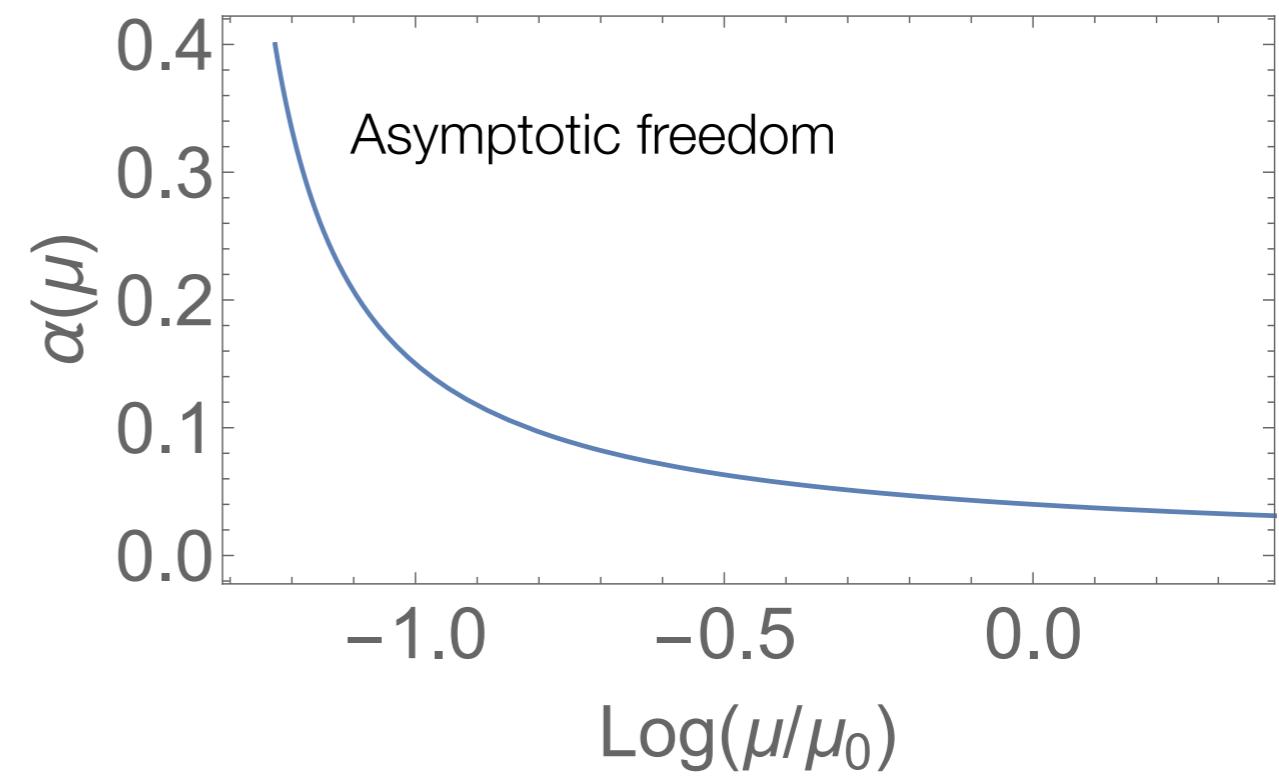
$$\beta_g = \partial_t \alpha_g = -B \alpha_g^2$$

$$t = \ln \frac{\mu}{\mu_0}$$

$$B < 0$$



$$B > 0$$



Small parameters

$$B = -\frac{4}{3}\epsilon$$

Small parameters

$$B = -\frac{4}{3}\epsilon$$

$$\epsilon = \frac{N_F}{N_C} - \frac{11}{2}$$

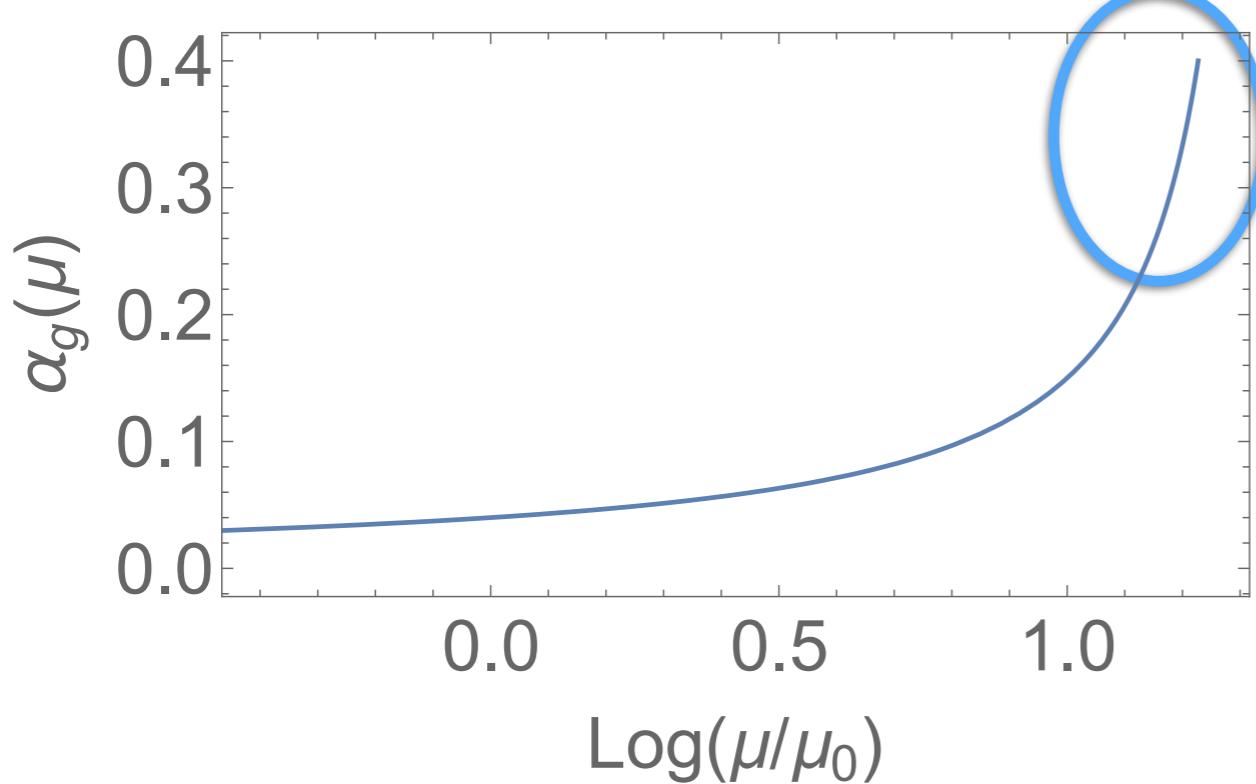
$$0 \leq \epsilon \ll 1$$

Small parameters

$$B = -\frac{4}{3}\epsilon$$

$$\epsilon = \frac{N_F}{N_C} - \frac{11}{2}$$

$$B < 0 \quad \epsilon > 0$$



$$0 \leq \epsilon \ll 1$$

Landau Pole ?

Can NL help?

$$\beta_g = -B\alpha_g^2 + C\alpha_g^3$$

$$B = -\frac{4}{3}\epsilon$$

Can NL help?

$$\beta_g = -B\alpha_g^2 + C\alpha_g^3$$

$$B = -\frac{4}{3}\epsilon$$

$$0 \leq \alpha_g^* \ll 1 \quad \text{iff} \quad C < 0$$

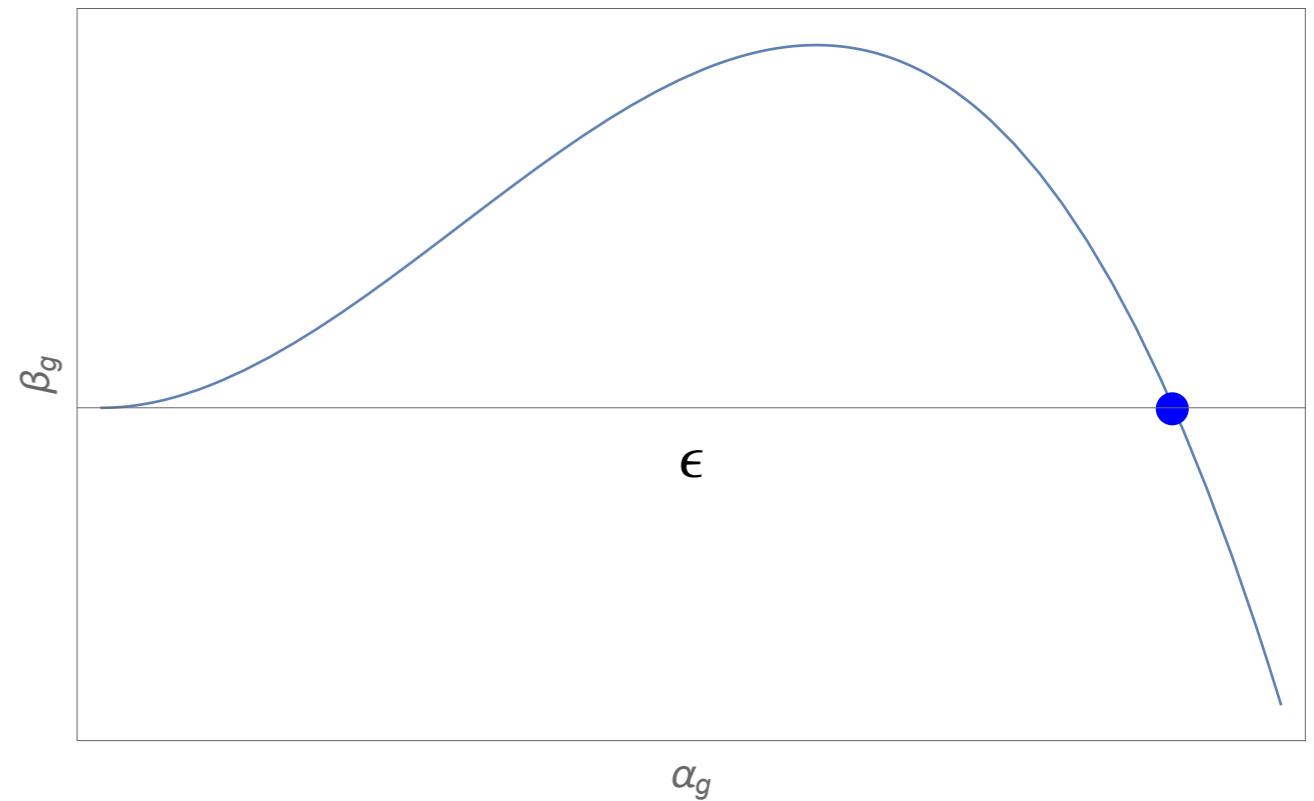
Can NL help?

$$\beta_g = -B\alpha_g^2 + C\alpha_g^3$$

$$B = -\frac{4}{3}\epsilon$$

$$0 \leq \alpha_g^* \ll 1 \quad \text{iff} \quad C < 0$$

$$\alpha_g^* = \frac{B}{C} \propto \epsilon$$



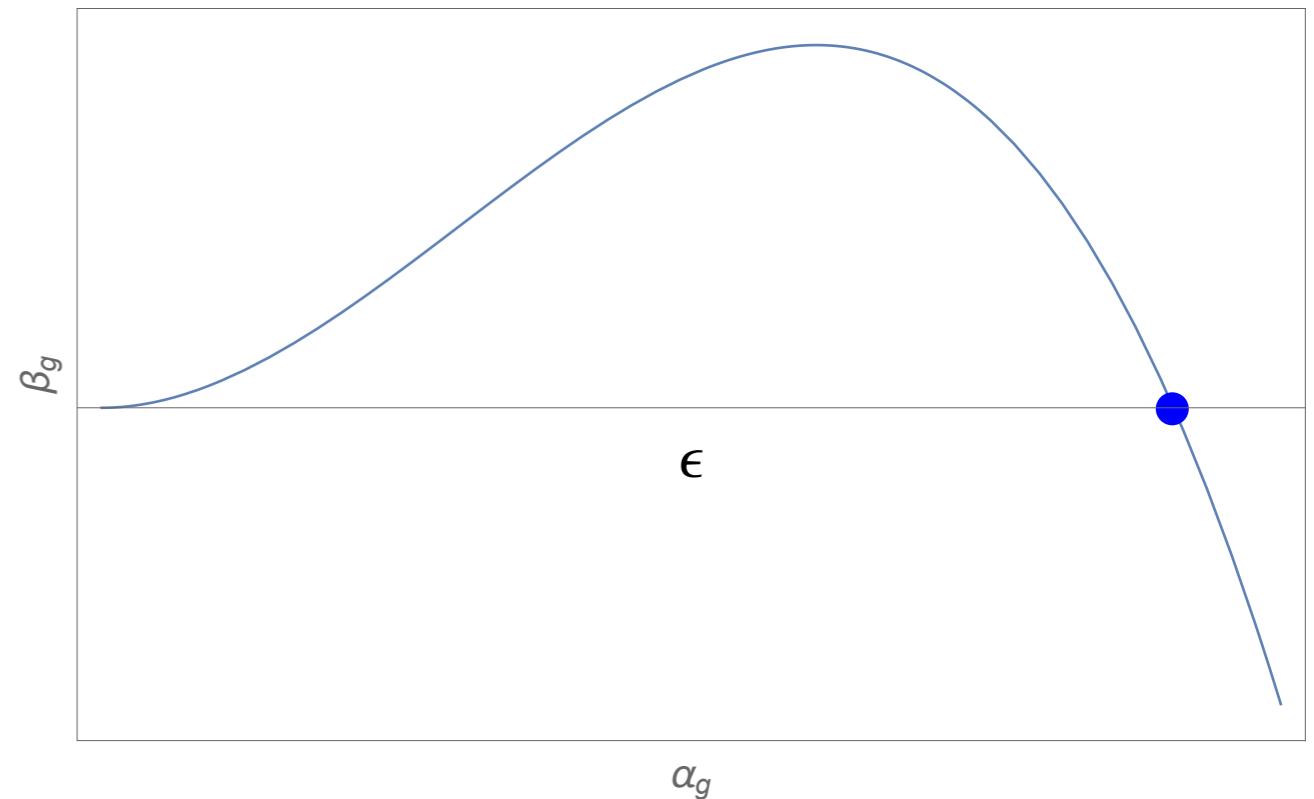
Can NL help?

$$\beta_g = -B\alpha_g^2 + C\alpha_g^3$$

$$B = -\frac{4}{3}\epsilon$$

$$0 \leq \alpha_g^* \ll 1 \quad \text{iff} \quad C < 0$$

$$\alpha_g^* = \frac{B}{C} \propto \epsilon$$



Impossible in Gauge Theories with Fermions alone

Caswell, PRL 1974

Add Yukawa

$$\beta_g = \alpha_g^2 \left[\frac{4}{3}\epsilon + \left(25 + \frac{26}{3}\epsilon \right) \alpha_g - 2 \left(\frac{11}{2} + \epsilon \right)^2 \alpha_y \right]$$

$$\beta_y = \alpha_y \left[\left(13 + 2\epsilon \right) \alpha_y - 6 \alpha_g \right]$$

Add Yukawa

$$\beta_g = \alpha_g^2 \left[\frac{4}{3}\epsilon + \left(25 + \frac{26}{3}\epsilon \right) \alpha_g - 2 \left(\frac{11}{2} + \epsilon \right)^2 \alpha_y \right]$$

$$\beta_y = \alpha_y \left[(13 + 2\epsilon) \alpha_y - 6 \alpha_g \right]$$

NLO - Fixed Points

- ◆ Gaussian fixed point

$$(\alpha_g^*, \alpha_y^*) = (0, 0)$$

NLO - Fixed Points

- ◆ Gaussian fixed point

$$(\alpha_g^*, \alpha_y^*) = (0, 0)$$

- ◆ Interacting fixed point

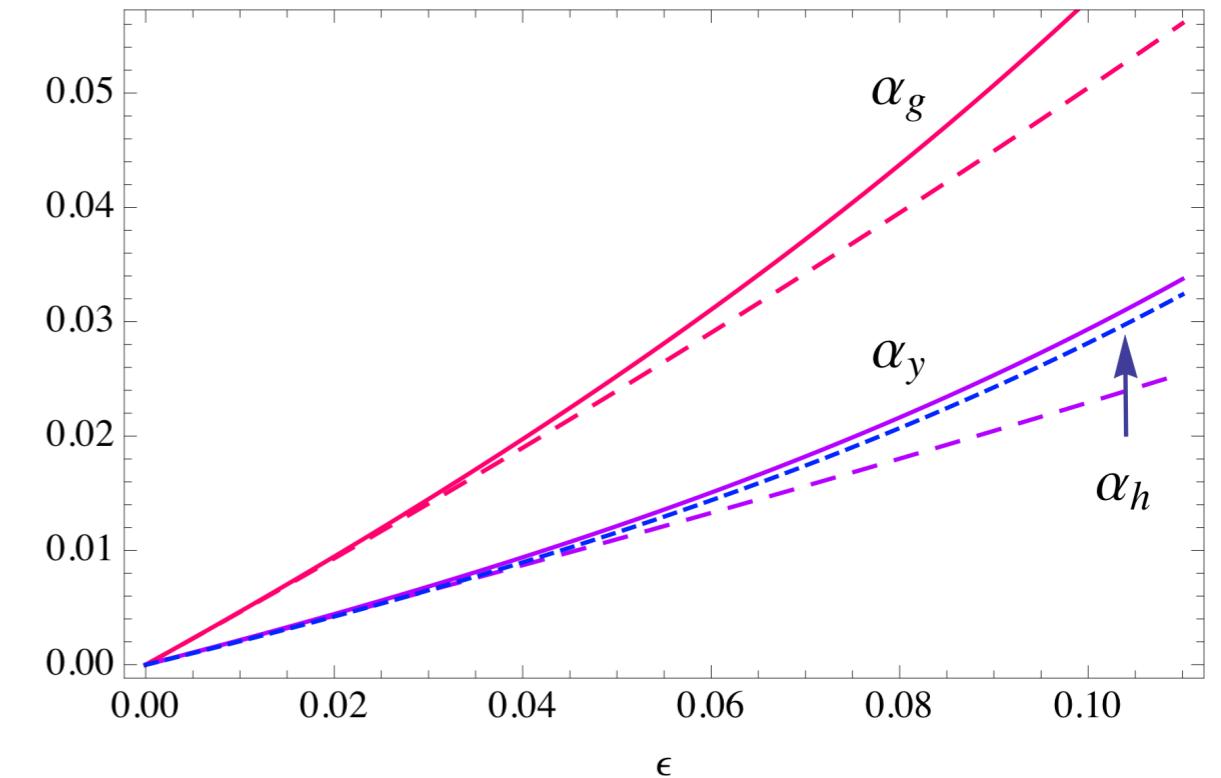
$$\alpha_g^* = \frac{26\epsilon + 4\epsilon^2}{57 - 46\epsilon - 8\epsilon^2} = \frac{26}{57}\epsilon + \frac{1424}{3249}\epsilon^2 + \frac{77360}{185193}\epsilon^3 + \mathcal{O}(\epsilon^4)$$

$$\alpha_y^* = \frac{12\epsilon}{57 - 46\epsilon - 8\epsilon^2} = \frac{4}{19}\epsilon + \frac{184}{1083}\epsilon^2 + \frac{10288}{61731}\epsilon^3 + \mathcal{O}(\epsilon^4).$$

NLO - Fixed Points

- ◆ Gaussian fixed point
 $(\alpha_g^*, \alpha_y^*) = (0, 0)$
- ◆ Interacting fixed point

$$\begin{aligned}\alpha_g^* &= \frac{26\epsilon + 4\epsilon^2}{57 - 46\epsilon - 8\epsilon^2} = \frac{26}{57}\epsilon + \frac{1424}{3249}\epsilon^2 + \frac{77360}{185193}\epsilon^3 + \mathcal{O}(\epsilon^4) \\ \alpha_y^* &= \frac{12\epsilon}{57 - 46\epsilon - 8\epsilon^2} = \frac{4}{19}\epsilon + \frac{184}{1083}\epsilon^2 + \frac{10288}{61731}\epsilon^3 + \mathcal{O}(\epsilon^4).\end{aligned}$$



Scaling exponents: UV completion

Scaling exponents: UV completion

$$\delta\alpha = (\alpha - \alpha_*) \propto \left(\frac{\mu}{\Lambda_c}\right)^\vartheta$$

Scaling exponents: UV completion

$$\delta\alpha = (\alpha - \alpha_*) \propto \left(\frac{\mu}{\Lambda_c} \right)^\vartheta$$

$$\begin{aligned}\vartheta_1 &= -\frac{104}{171}\epsilon^2 + \frac{2296}{3249}\epsilon^3 + \frac{1387768}{1666737}\epsilon^4 + \mathcal{O}(\epsilon^4) \\ \vartheta_2 &= \frac{52}{19}\epsilon + \frac{9140}{1083}\epsilon^2 + \frac{2518432}{185193}\epsilon^3 + \mathcal{O}(\epsilon^4).\end{aligned}$$

Scaling exponents: UV completion

$$\delta\alpha = (\alpha - \alpha_*) \propto \left(\frac{\mu}{\Lambda_c} \right)^\vartheta$$

$$\begin{aligned}\vartheta_1 &= -\frac{104}{171}\epsilon^2 + \frac{2296}{3249}\epsilon^3 + \frac{1387768}{1666737}\epsilon^4 + \mathcal{O}(\epsilon^4) \\ \vartheta_2 &= \frac{52}{19}\epsilon + \frac{9140}{1083}\epsilon^2 + \frac{2518432}{185193}\epsilon^3 + \mathcal{O}(\epsilon^4).\end{aligned}$$

$\vartheta_1 < 0$ Relevant direction

Scaling exponents: UV completion

$$\delta\alpha = (\alpha - \alpha_*) \propto \left(\frac{\mu}{\Lambda_c}\right)^\vartheta$$

$$\vartheta_1 = -\frac{104}{171}\epsilon^2 + \frac{2296}{3249}\epsilon^3 + \frac{1387768}{1666737}\epsilon^4 + \mathcal{O}(\epsilon^4)$$

$$\vartheta_2 = \frac{52}{19}\epsilon + \frac{9140}{1083}\epsilon^2 + \frac{2518432}{185193}\epsilon^3 + \mathcal{O}(\epsilon^4).$$

$\vartheta_1 < 0$ Relevant direction

$\vartheta_2 > 0$ Irrelevant direction

Scaling exponents: UV completion

$$\delta\alpha = (\alpha - \alpha_*) \propto \left(\frac{\mu}{\Lambda_c}\right)^\vartheta$$

$$\vartheta_1 = -\frac{104}{171}\epsilon^2 + \frac{2296}{3249}\epsilon^3 + \frac{1387768}{1666737}\epsilon^4 + \mathcal{O}(\epsilon^4)$$

$$\vartheta_2 = \frac{52}{19}\epsilon + \frac{9140}{1083}\epsilon^2 + \frac{2518432}{185193}\epsilon^3 + \mathcal{O}(\epsilon^4).$$

$\vartheta_1 < 0$ Relevant direction

$\vartheta_2 > 0$ Irrelevant direction

A true UV fixed point to this order

Scaling exponents: UV completion

$$\delta\alpha = (\alpha - \alpha_*) \propto \left(\frac{\mu}{\Lambda_c}\right)^\vartheta$$

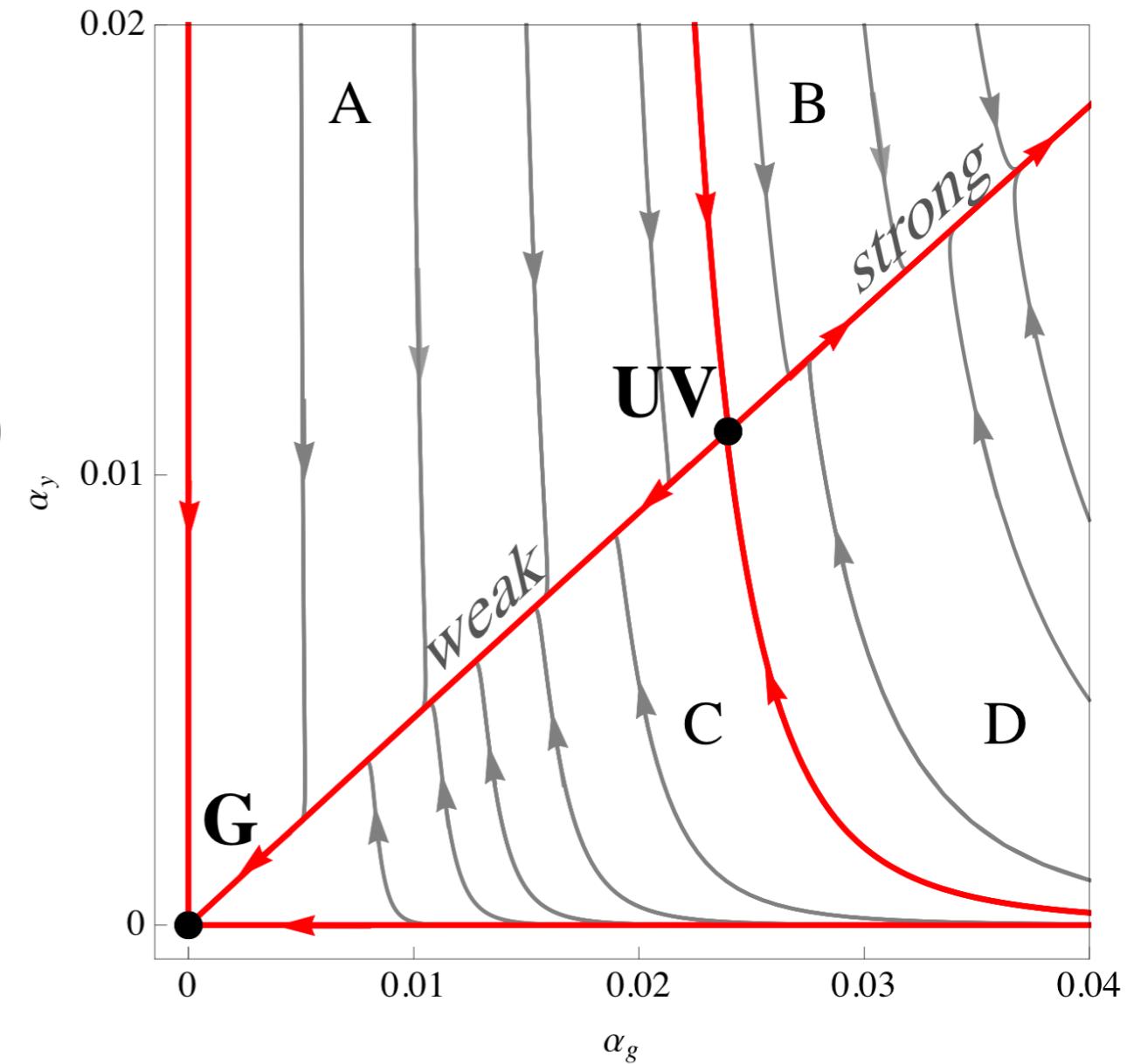
$$\vartheta_1 = -\frac{104}{171}\epsilon^2 + \frac{2296}{3249}\epsilon^3 + \frac{1387768}{1666737}\epsilon^4 + \mathcal{O}(\epsilon^4)$$

$$\vartheta_2 = \frac{52}{19}\epsilon + \frac{9140}{1083}\epsilon^2 + \frac{2518432}{185193}\epsilon^3 + \mathcal{O}(\epsilon^4).$$

$\vartheta_1 < 0$ Relevant direction

$\vartheta_2 > 0$ Irrelevant direction

A true UV fixed point to this order



Scaling exponents: UV completion

$$\delta\alpha = (\alpha - \alpha_*) \propto \left(\frac{\mu}{\Lambda_c}\right)^\vartheta$$

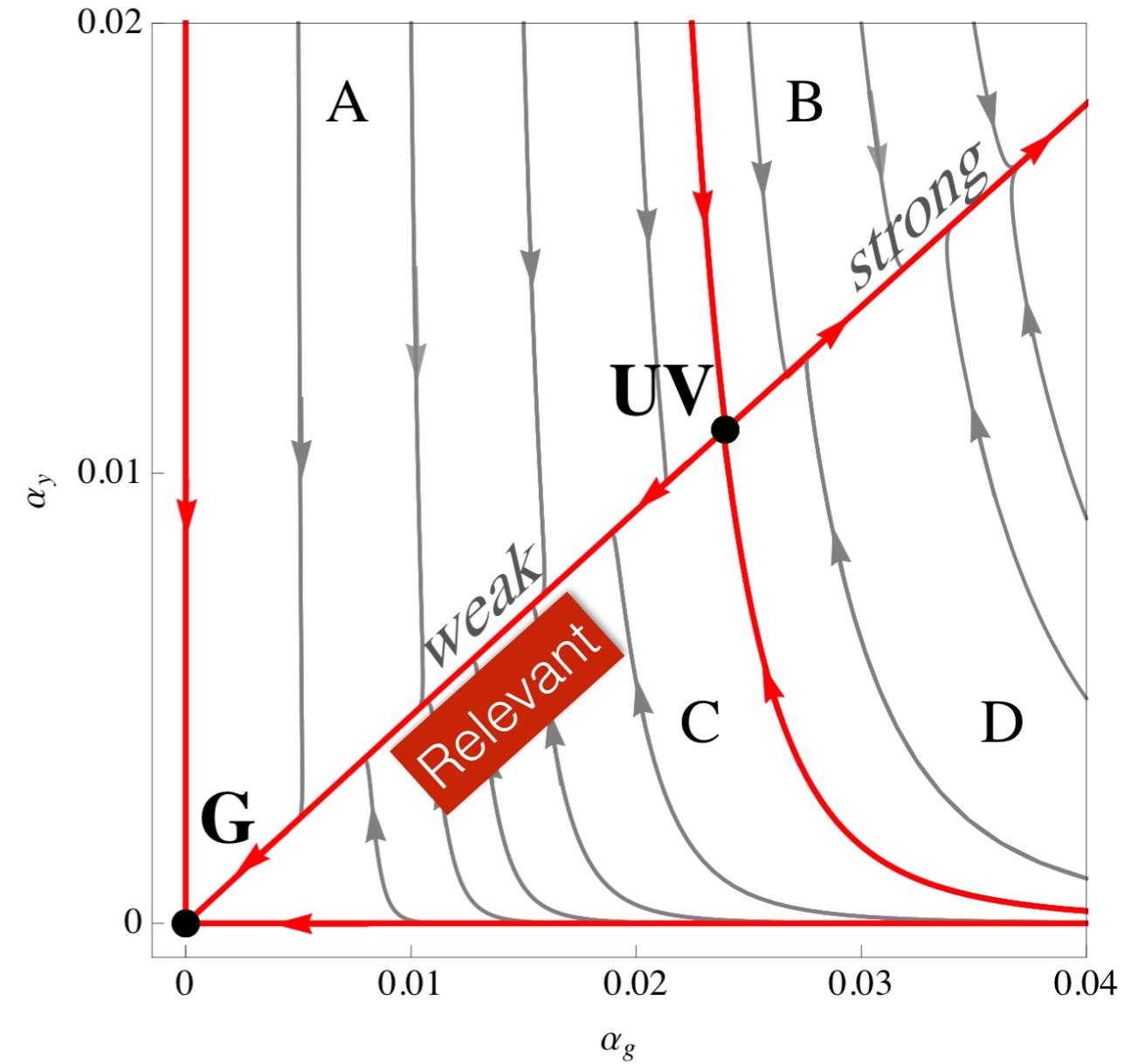
$$\vartheta_1 = -\frac{104}{171}\epsilon^2 + \frac{2296}{3249}\epsilon^3 + \frac{1387768}{1666737}\epsilon^4 + \mathcal{O}(\epsilon^4)$$

$$\vartheta_2 = \frac{52}{19}\epsilon + \frac{9140}{1083}\epsilon^2 + \frac{2518432}{185193}\epsilon^3 + \mathcal{O}(\epsilon^4).$$

$\vartheta_1 < 0$ Relevant direction

$\vartheta_2 > 0$ Irrelevant direction

A true UV fixed point to this order



Scaling exponents: UV completion

$$\delta\alpha = (\alpha - \alpha_*) \propto \left(\frac{\mu}{\Lambda_c}\right)^\vartheta$$

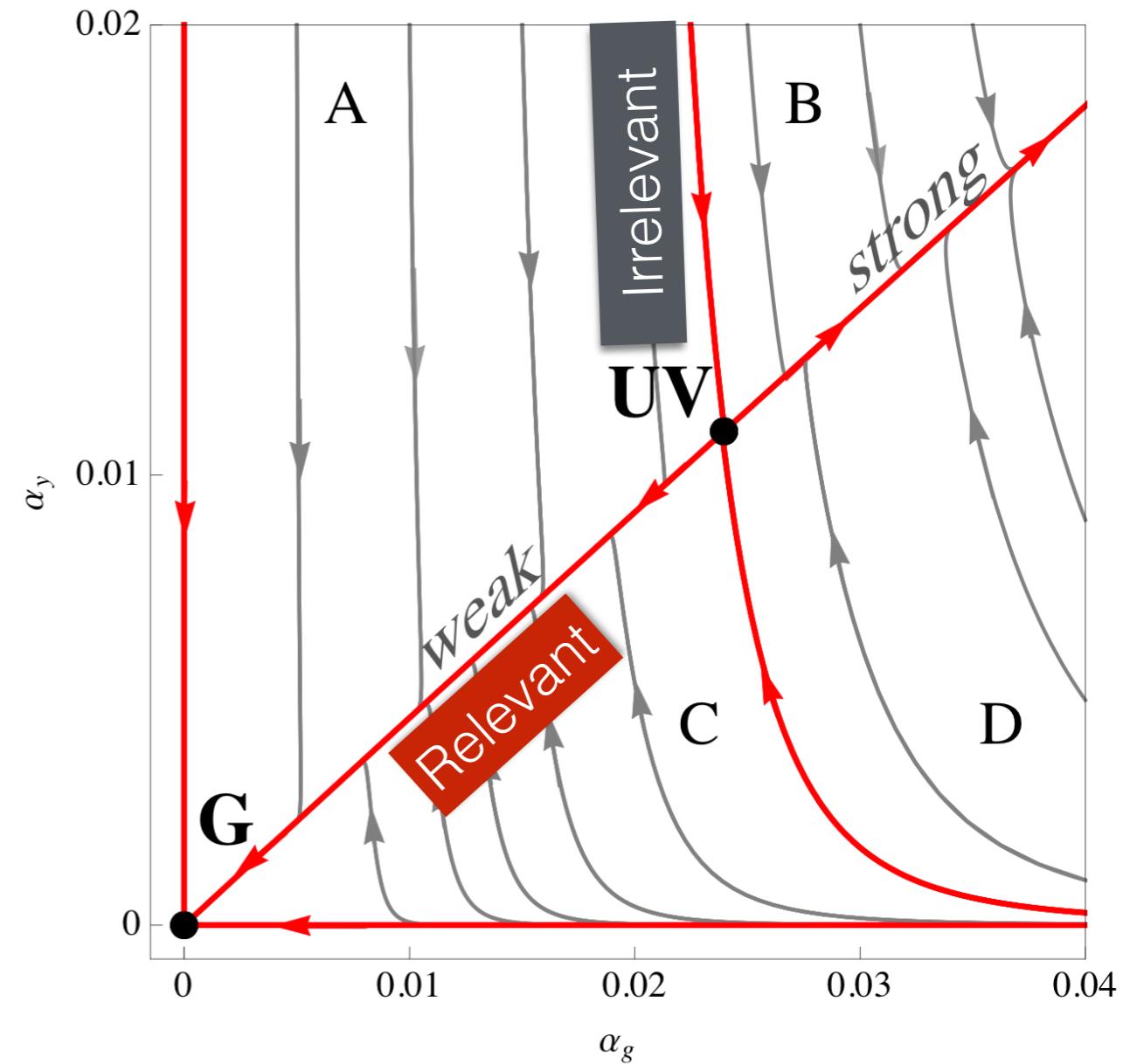
$$\vartheta_1 = -\frac{104}{171}\epsilon^2 + \frac{2296}{3249}\epsilon^3 + \frac{1387768}{1666737}\epsilon^4 + \mathcal{O}(\epsilon^4)$$

$$\vartheta_2 = \frac{52}{19}\epsilon + \frac{9140}{1083}\epsilon^2 + \frac{2518432}{185193}\epsilon^3 + \mathcal{O}(\epsilon^4).$$

$\vartheta_1 < 0$ Relevant direction

$\vartheta_2 > 0$ Irrelevant direction

A true UV fixed point to this order



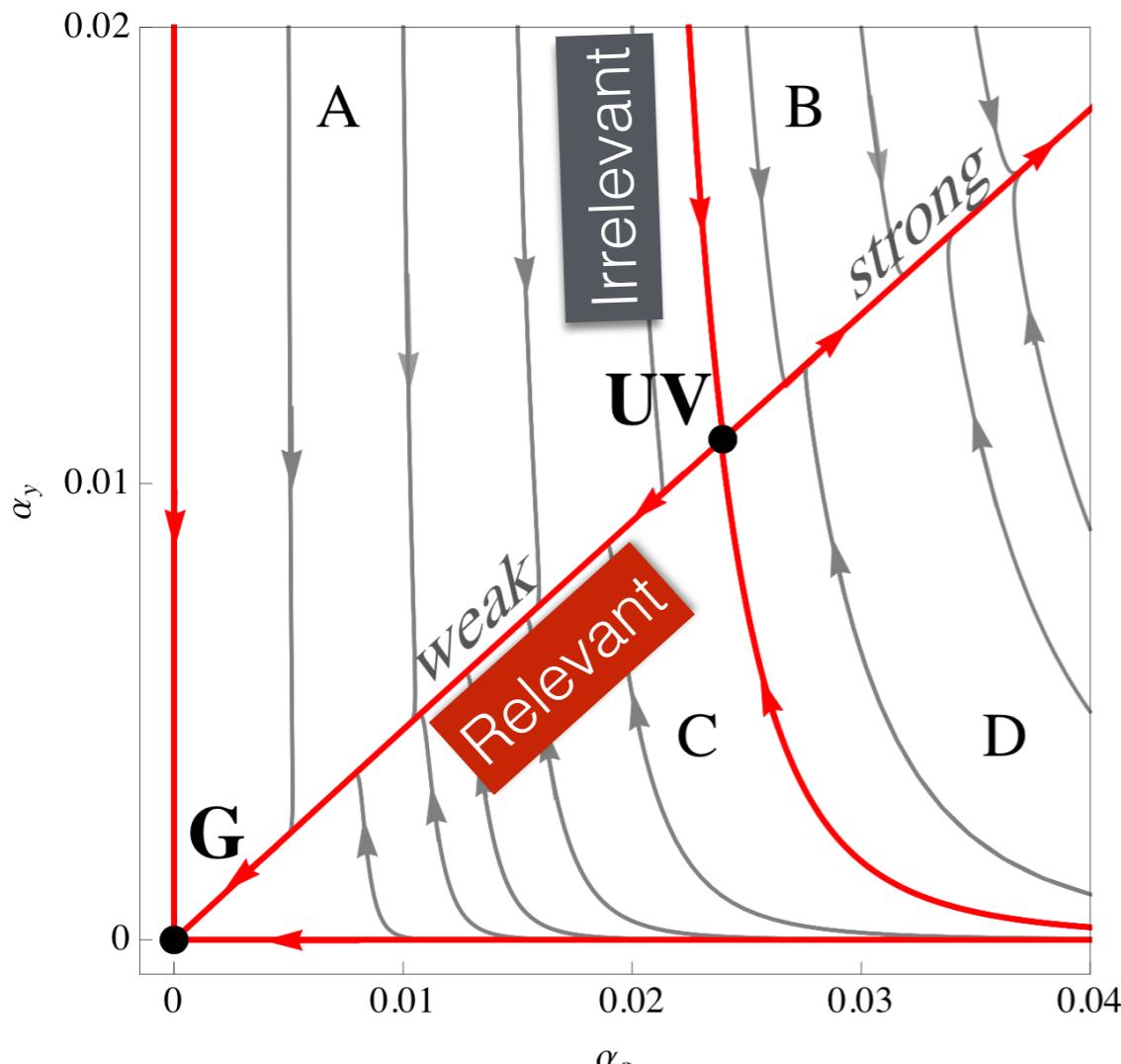
Phase Diagram

$$\vartheta_1 = -0.608 \epsilon^2 + \mathcal{O}(\epsilon^3)$$

$$\vartheta_2 = 2.737 \epsilon + \mathcal{O}(\epsilon^2)$$

$$\vartheta_3 = 4.039 \epsilon + \mathcal{O}(\epsilon^2)$$

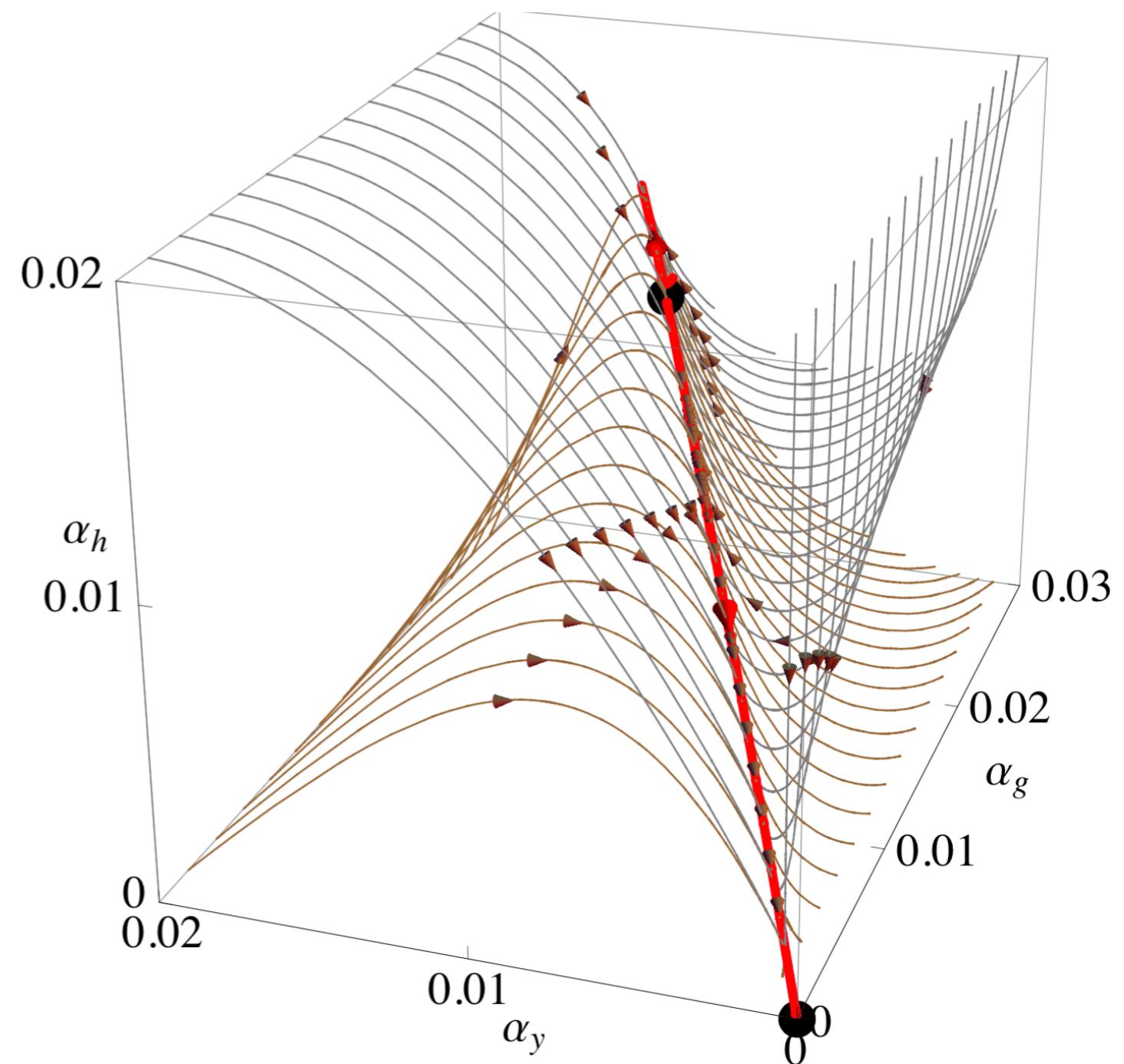
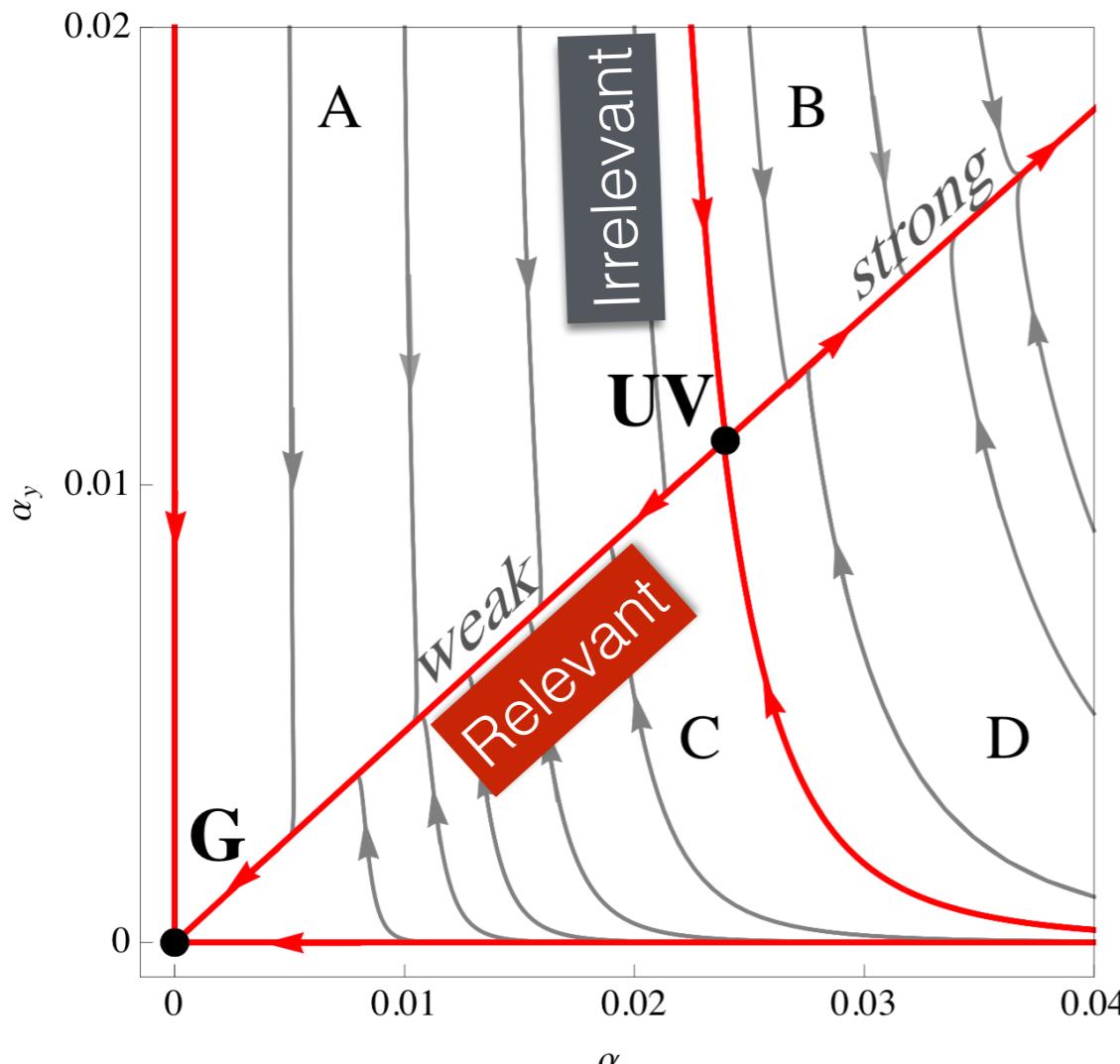
$$\vartheta_4 = 2.941 \epsilon + \mathcal{O}(\epsilon^2).$$



$$\vartheta_1 < 0 < \vartheta_2 < \vartheta_4 < \vartheta_3$$

Phase Diagram

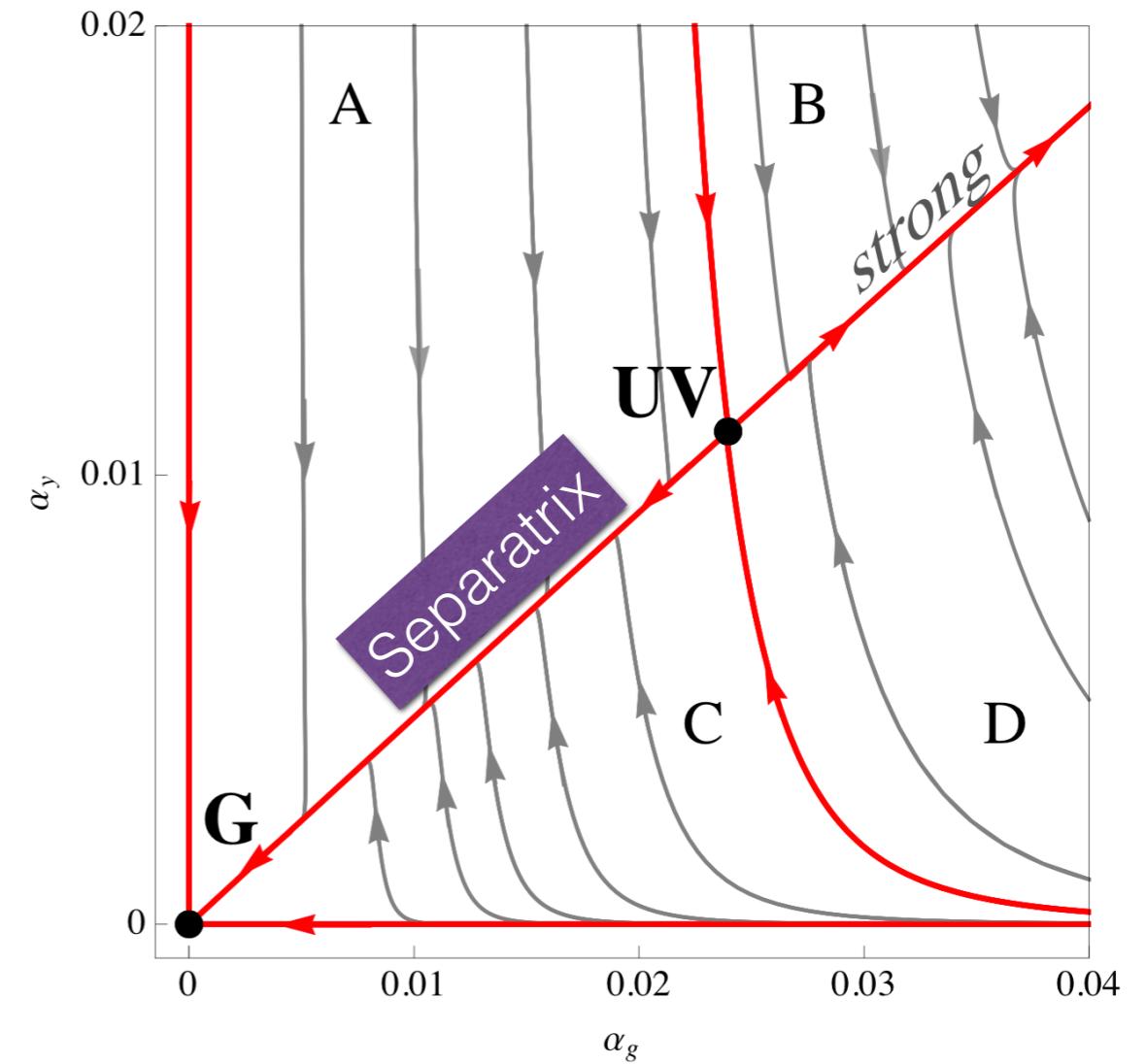
$$\begin{aligned}\vartheta_1 &= -0.608 \epsilon^2 + \mathcal{O}(\epsilon^3) \\ \vartheta_2 &= 2.737 \epsilon + \mathcal{O}(\epsilon^2) \\ \vartheta_3 &= 4.039 \epsilon + \mathcal{O}(\epsilon^2) \\ \vartheta_4 &= 2.941 \epsilon + \mathcal{O}(\epsilon^2).\end{aligned}$$



$$\vartheta_1 < 0 < \vartheta_2 < \vartheta_4 < \vartheta_3$$

Separatrix = Line of Physics

Globally defined line connecting two FPs

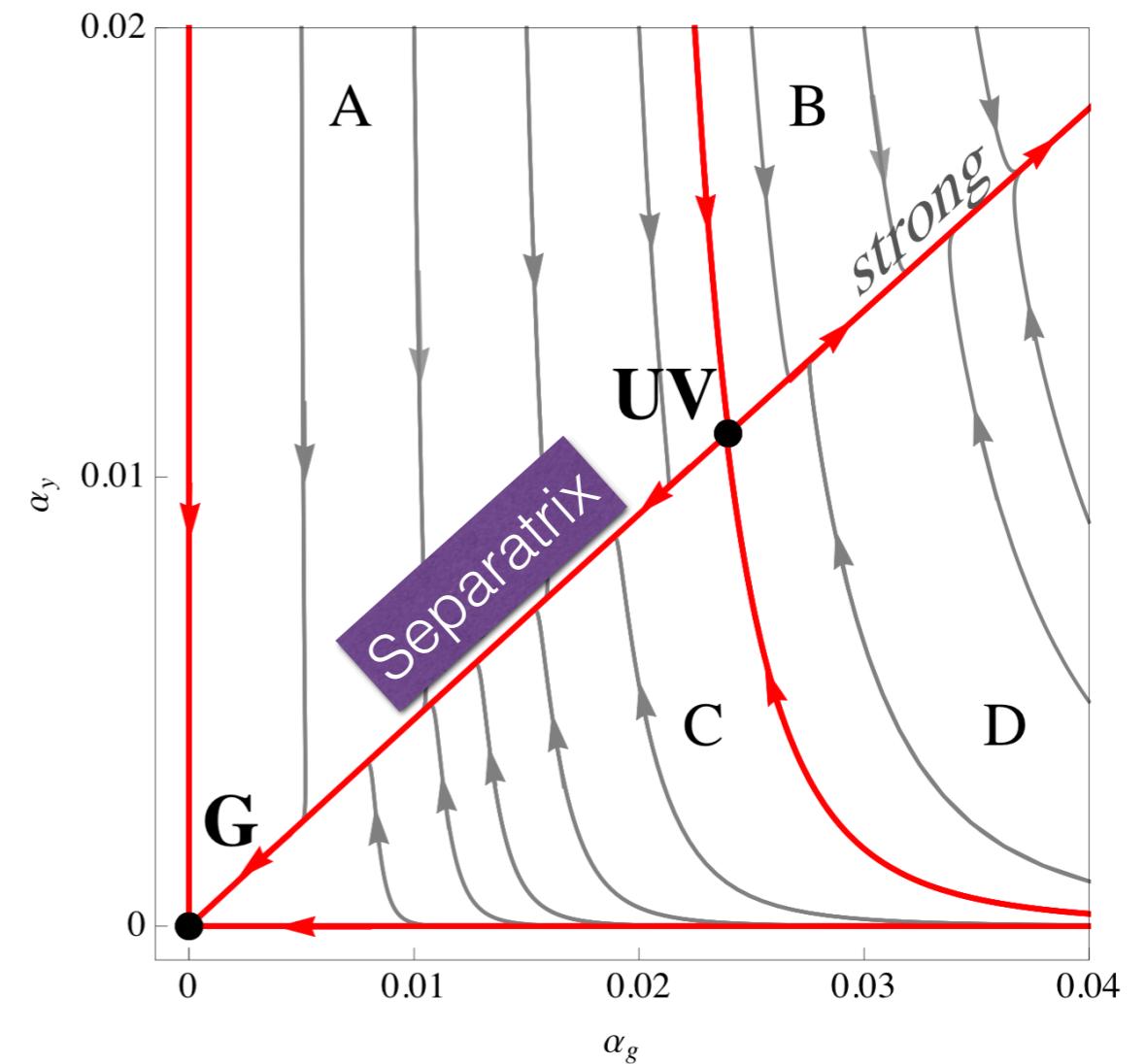


Separatrix = Line of Physics

Globally defined line connecting two FPs

$$\beta_g^{\text{sep}}(\alpha_g) \equiv \beta_g(\alpha_g, \alpha_y = F_y(\alpha_g))$$

$$\beta_y^{\text{sep}}(\alpha_g) \equiv \beta_y(\alpha_g, \alpha_y = F_y(\alpha_g))$$

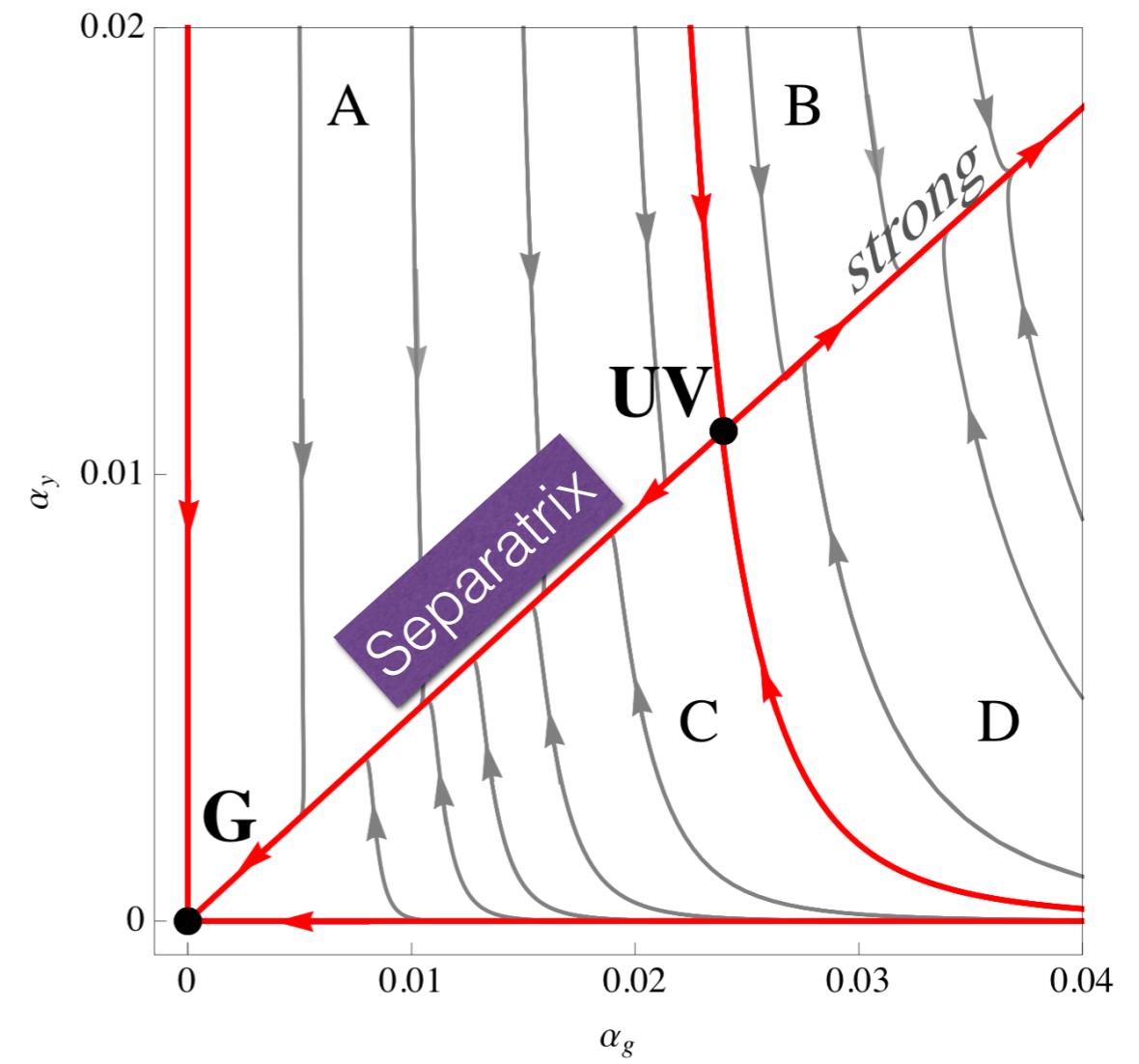
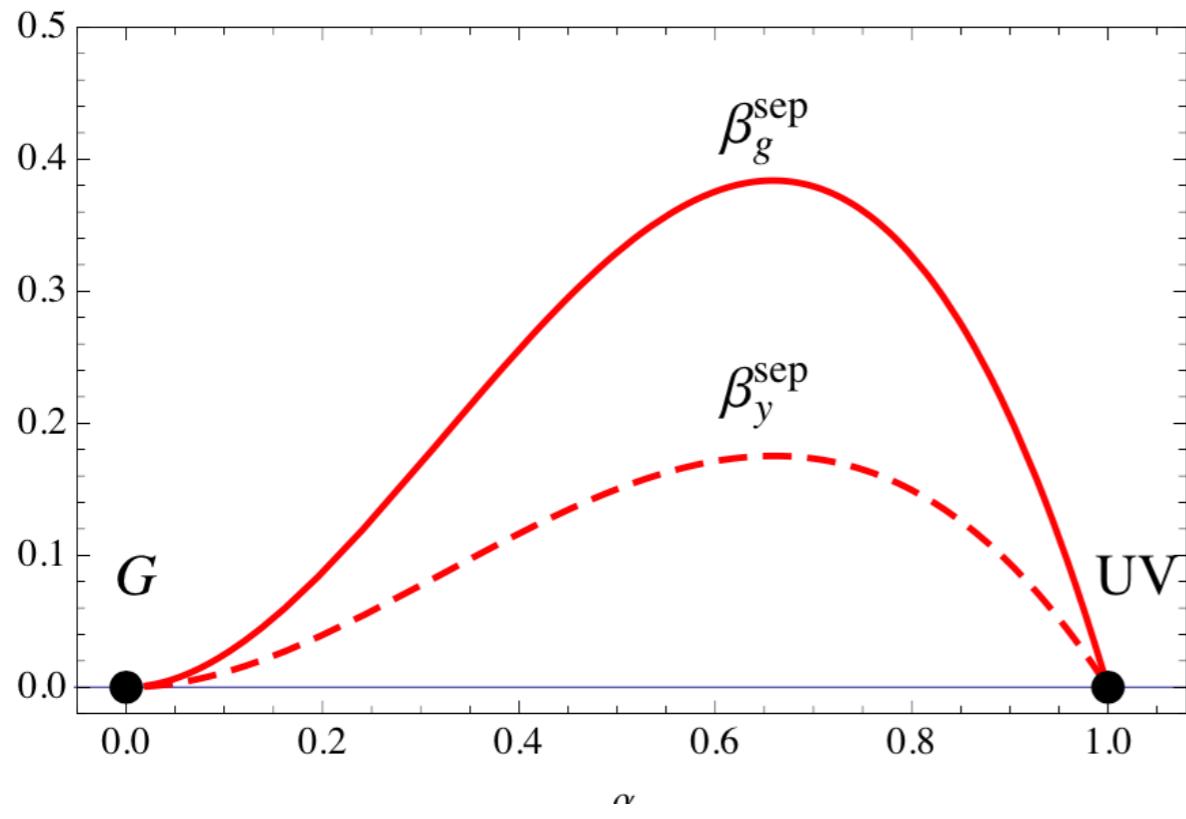


Separatrix = Line of Physics

Globally defined line connecting two FPs

$$\beta_g^{\text{sep}}(\alpha_g) \equiv \beta_g(\alpha_g, \alpha_y = F_y(\alpha_g))$$

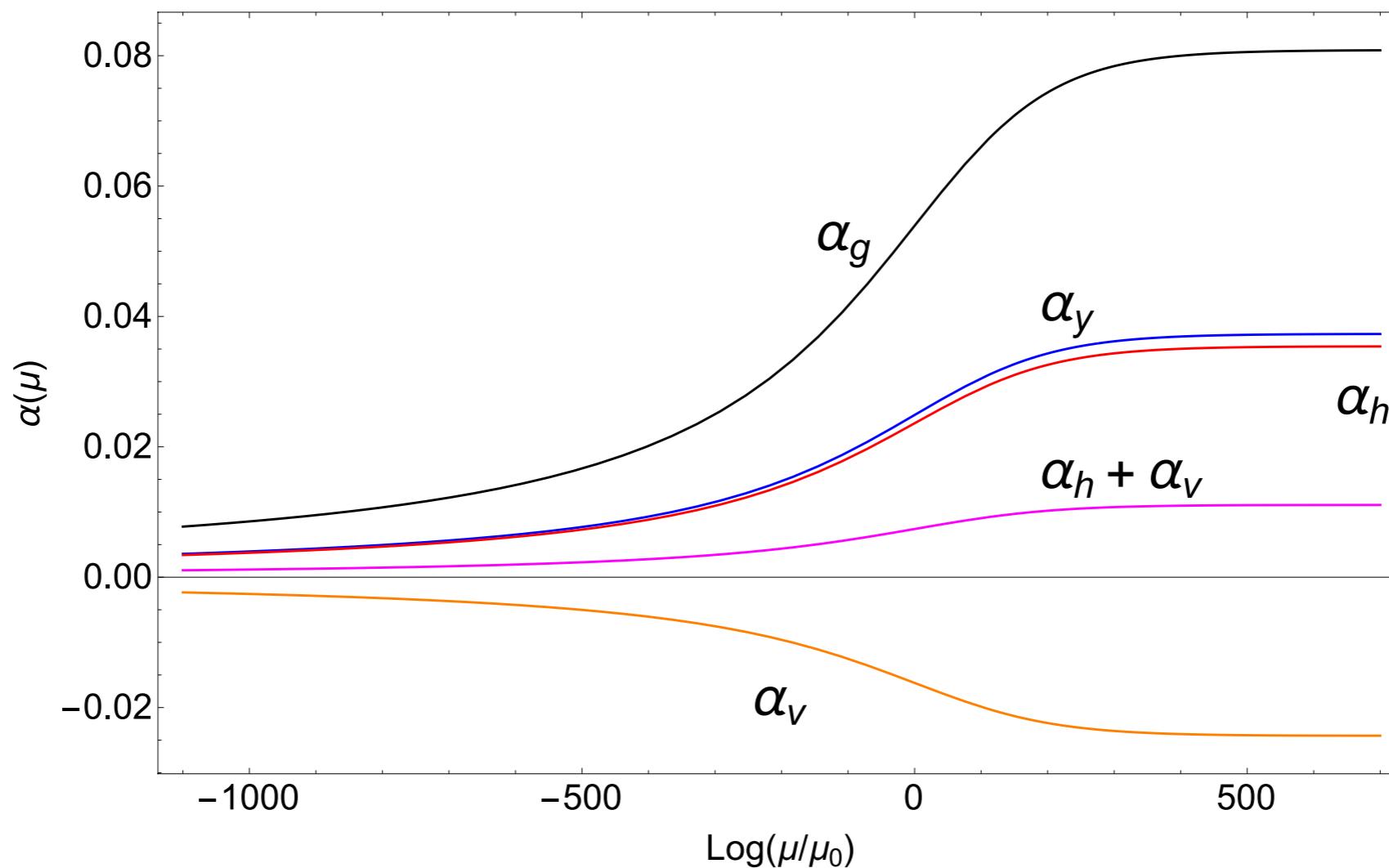
$$\beta_y^{\text{sep}}(\alpha_g) \equiv \beta_y(\alpha_g, \alpha_y = F_y(\alpha_g))$$



Complete asymptotic safety

Litim and Sannino, 1406.2337, JHEP

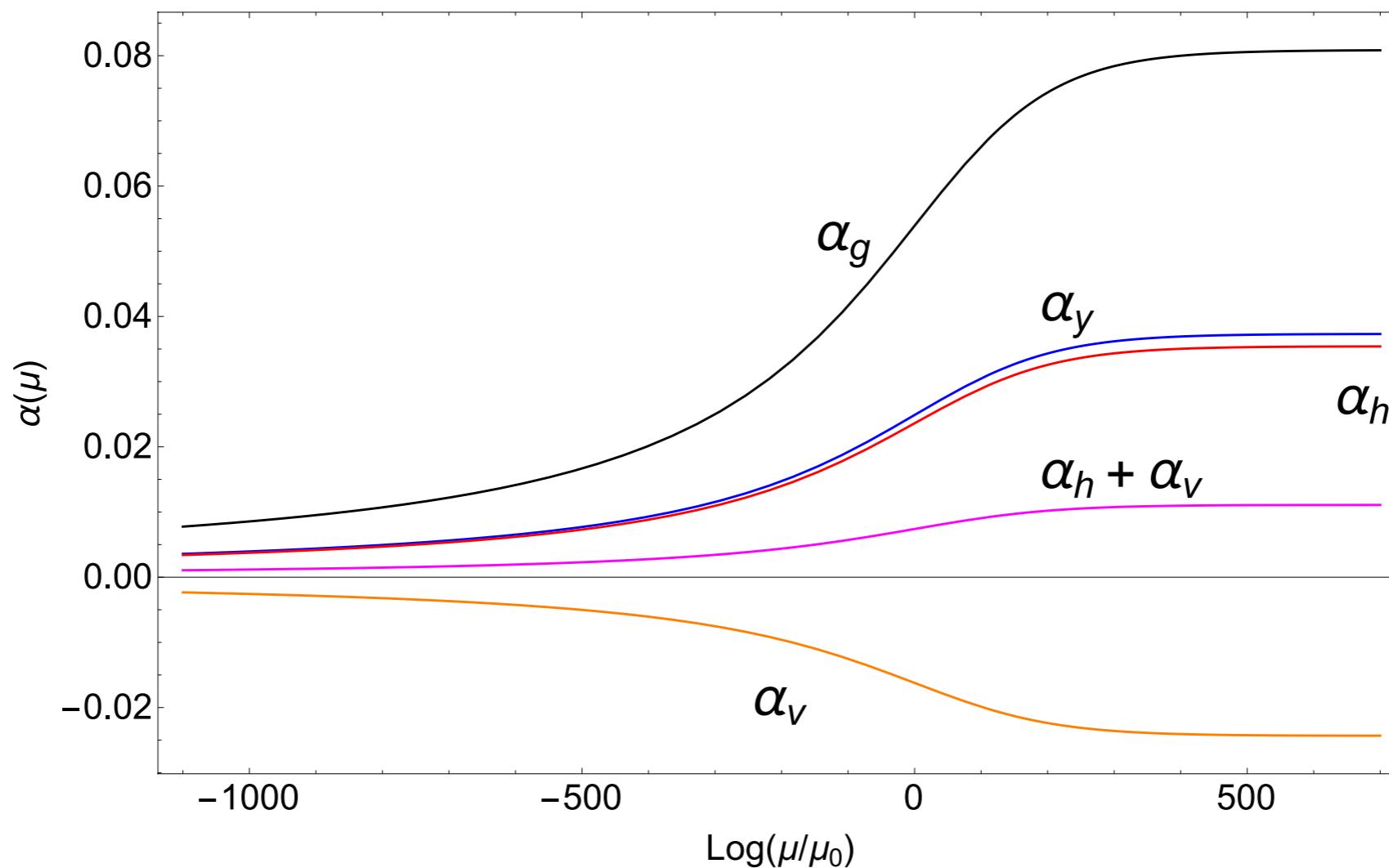
Gauge + fermion + scalars theories can be fund. at any energy scale



Complete asymptotic safety

Litim and Sannino, 1406.2337, JHEP

Gauge + fermion + scalars theories can be fund. at any energy scale



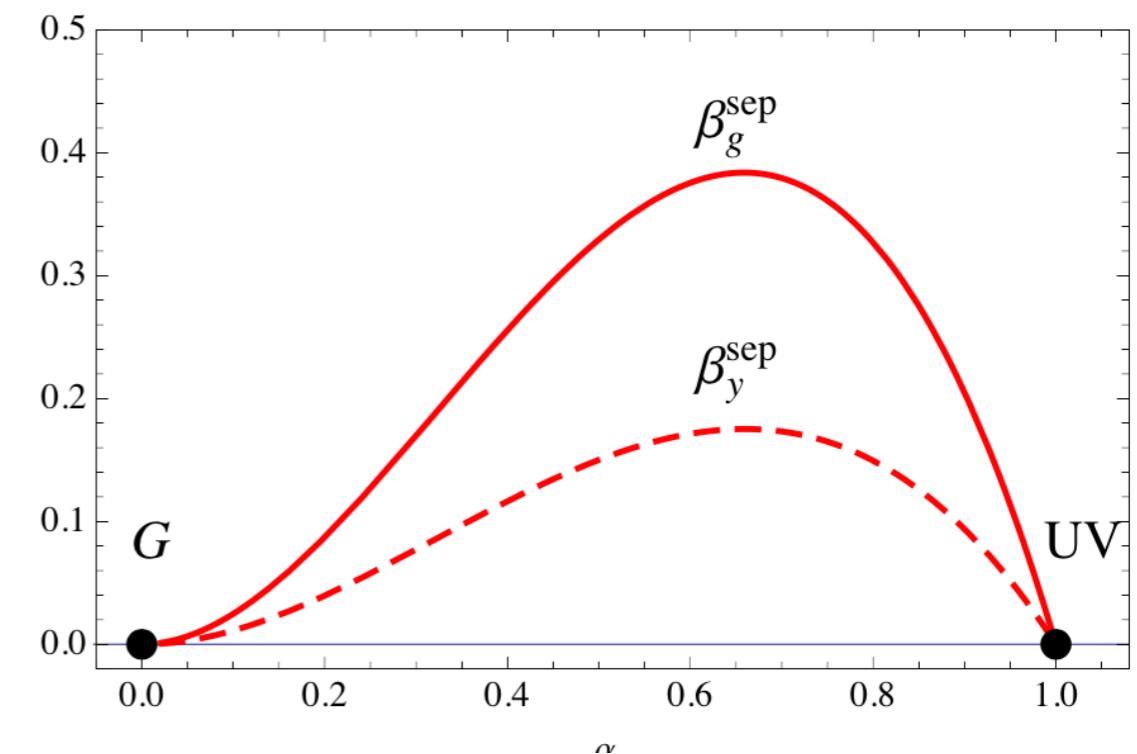
Scalars are needed to make the theory fundamental

Safe variation for the a-theorem function

Antipin, Gillioz, Mølgaard, Sannino 13

Safe variation for the a-theorem function

To leading order

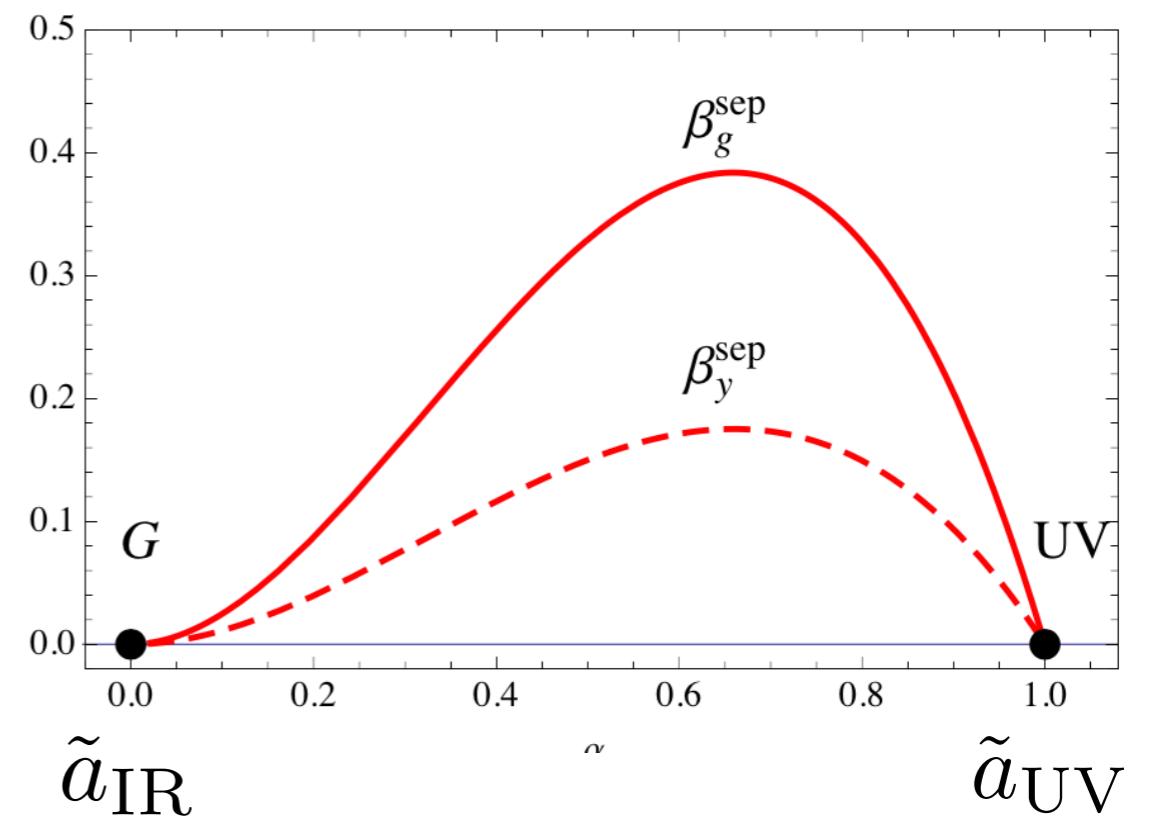


Safe variation for the a-theorem function

To leading order

$$\frac{\Delta \tilde{a}}{\chi_{gg}} = \frac{104}{171} \epsilon^2$$

$$\chi_{gg} = \frac{N_C^2 - 1}{128\pi^2}$$

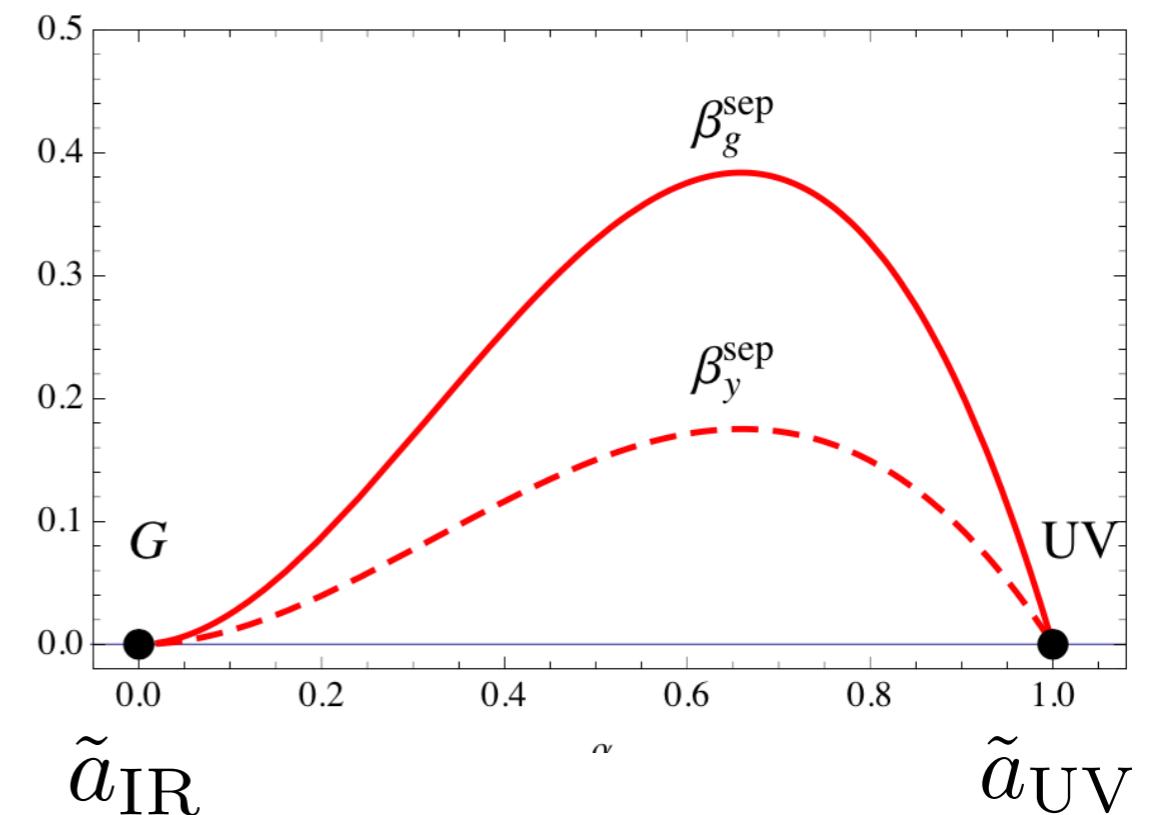


Safe variation for the a-theorem function

To leading order

$$\frac{\Delta \tilde{a}}{\chi_{gg}} = \frac{104}{171} \epsilon^2$$

$$\chi_{gg} = \frac{N_C^2 - 1}{128\pi^2}$$



Positive and growing with epsilon

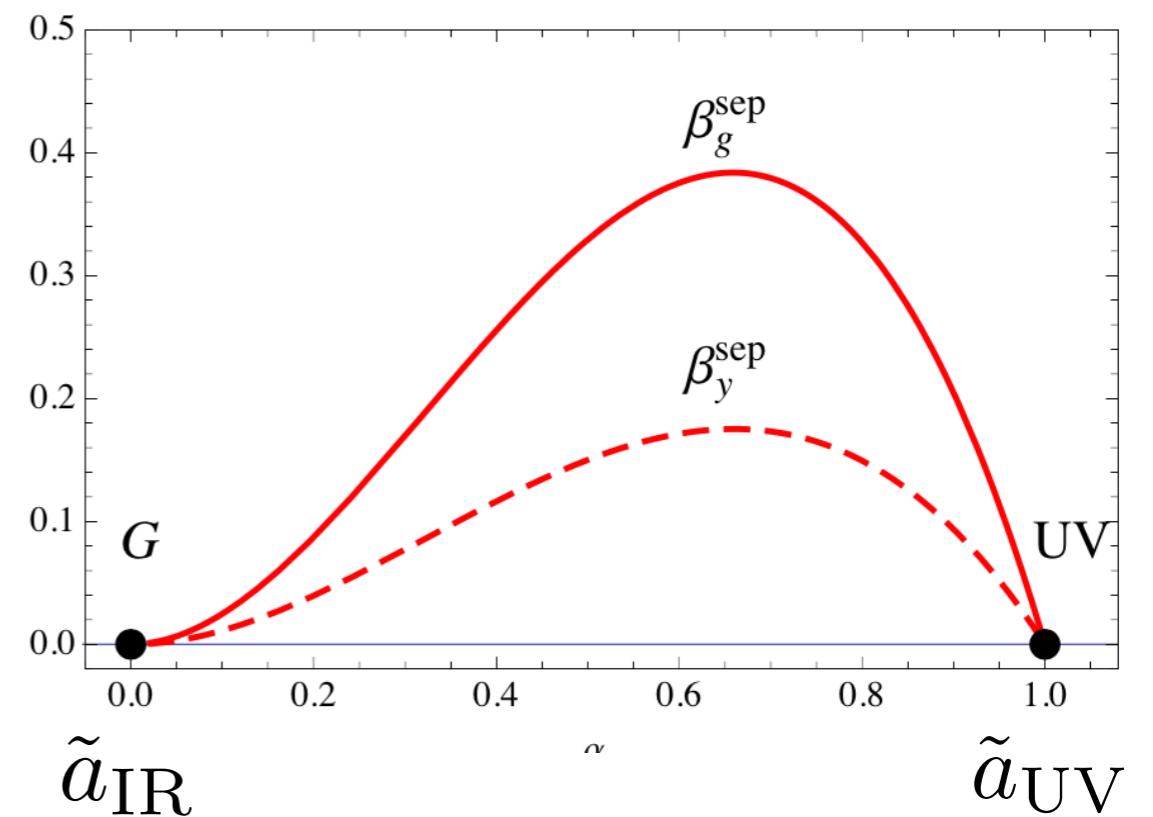
Antipin, Gillioz, Mølgaard, Sannino 13

Safe variation for the a-theorem function

To leading order

$$\frac{\Delta \tilde{a}}{\chi_{gg}} = \frac{104}{171} \epsilon^2$$

$$\chi_{gg} = \frac{N_C^2 - 1}{128\pi^2}$$



Positive and growing with epsilon

Antipin, Gillioz, Mølgaard, Sannino 13

Bootstrap and composite operators

Antipin, Mølgaard, Sannino 14

Supersymmetry is unsafe

Intriligator and Sannino, 1508.07413, JHEP

Martin and Wells, hep-ph/0011382, PRD

SQCD with H

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$	$U(1)_R$
W_α	Adj	1	1	0	1
Q	\square	$\bar{\square}$	1	1	$1 - \frac{N_c}{N_f}$
\tilde{Q}	$\bar{\square}$	1	\square	-1	$1 - \frac{N_c}{N_f}$
H	1	\square	$\bar{\square}$	0	$2\frac{N_c}{N_f}$

SQCD with \mathbb{H}

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$	$U(1)_R$	AF is lost
W_α	Adj	1	1	0	1	
Q	\square	$\bar{\square}$	1	1	$1 - \frac{N_c}{N_f}$	$N_f > 3N_c$
\tilde{Q}	$\bar{\square}$	1	\square	-1	$1 - \frac{N_c}{N_f}$	
H	1	\square	$\bar{\square}$	0	$2\frac{N_c}{N_f}$	

SQCD with H

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$	$U(1)_R$	AF is lost
W_α	Adj	1	1	0	1	
Q	\square	$\bar{\square}$	1	1	$1 - \frac{N_c}{N_f}$	$N_f > 3N_c$
\tilde{Q}	$\bar{\square}$	1	\square	-1	$1 - \frac{N_c}{N_f}$	
H	1	\square	$\bar{\square}$	0	$2\frac{N_c}{N_f}$	$W = y \operatorname{Tr} Q H \tilde{Q}$

SQCD with H

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$	$U(1)_R$
W_α	Adj	1	1	0	1
Q	□	□	1	1	$1 - \frac{N_c}{N_f}$
\tilde{Q}	□	1	□	-1	$1 - \frac{N_c}{N_f}$
H	1	□	□	0	$2\frac{N_c}{N_f}$

AF is lost

$$N_f > 3N_c$$

$$W = y \operatorname{Tr} Q H \tilde{Q}$$

$$\beta(\alpha_g) \approx -2\alpha_g^2 \left[3 - \frac{N_f}{N_c} + \left(6 - 4\frac{N_f}{N_c} \right) \alpha_g + 2\frac{N_f^2}{N_c^2} \alpha_y + \mathcal{O}(\alpha^2) \right]$$

$$\beta(\alpha_y) \approx 2\alpha_y \left[\left(2\frac{N_f}{N_c} + 1 \right) \alpha_y - 2\alpha_g + \mathcal{O}(\alpha^2) \right]$$

SQCD with H

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$	$U(1)_R$
W_α	Adj	1	1	0	1
Q	□	□	1	1	$1 - \frac{N_c}{N_f}$
\tilde{Q}	□	1	□	-1	$1 - \frac{N_c}{N_f}$
H	1	□	□	0	$2\frac{N_c}{N_f}$

AF is lost

$$N_f > 3N_c$$

$$W = y \operatorname{Tr} Q H \tilde{Q}$$

$$\beta(\alpha_g) \approx -2\alpha_g^2 \left[3 - \frac{N_f}{N_c} + \left(6 - 4\frac{N_f}{N_c} \right) \alpha_g + 2\frac{N_f^2}{N_c^2} \alpha_y + \mathcal{O}(\alpha^2) \right]$$

$$\beta(\alpha_y) \approx 2\alpha_y \left[\left(2\frac{N_f}{N_c} + 1 \right) \alpha_y - 2\alpha_g + \mathcal{O}(\alpha^2) \right]$$

No perturbative UV fixed point

$$\beta(\alpha_g) \approx 2\alpha_g^2 \left[\epsilon + \frac{6}{7}\alpha_g \right]$$

SQCD with H

SQCD with H

Assume a nonperturbative fixed point, however

SQCD with H

Assume a nonperturbative fixed point, however

$$D(H) = \frac{3}{2}R(H) = 3\frac{N_c}{N_f} < 1 \quad \text{for} \quad N_f > 3N_c$$

SQCD with H

Assume a nonperturbative fixed point, however

$$D(H) = \frac{3}{2}R(H) = 3\frac{N_c}{N_f} < 1 \quad \text{for} \quad N_f > 3N_c$$

Violates the unitarity bound

$$D(\mathcal{O}) \geq 1$$

SQCD with H

Assume a nonperturbative fixed point, however

$$D(H) = \frac{3}{2} R(H) = 3 \frac{N_c}{N_f} < 1 \quad \text{for} \quad N_f > 3N_c$$

Violates the unitarity bound

$$D(\mathcal{O}) \geq 1$$

Potential loophole: H is free and decouples at the fixed point

SQCD with H

Assume a nonperturbative fixed point, however

$$D(H) = \frac{3}{2} R(H) = 3 \frac{N_c}{N_f} < 1 \quad \text{for} \quad N_f > 3N_c$$

Violates the unitarity bound

$$D(\mathcal{O}) \geq 1$$

Potential loophole: H is free and decouples at the fixed point

Check if SQCD without H has an UV fixed point

SQCD

SQCD

Unitarity bound is not sufficient

SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q} \quad D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q}$$

$$D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

$$\mathcal{B} = Q^{N_c}$$

$$D_{SCFT}(\mathcal{B}) = D_{SCFT}(\tilde{\mathcal{B}}) = \frac{3}{2}R_{SCFT}(\mathcal{B}) = \frac{3}{2}N_c\frac{N_f - N_c}{N_f}$$

SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q}$$

$$D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

$$\mathcal{B} = Q^{N_c}$$

$$D_{SCFT}(\mathcal{B}) = D_{SCFT}(\tilde{\mathcal{B}}) = \frac{3}{2}R_{SCFT}(\mathcal{B}) = \frac{3}{2}N_c\frac{N_f - N_c}{N_f}$$

Can be ruled out via a-theorem $a(R) = 3\text{Tr}U(1)_R^3 - \text{Tr}U(1)_R$

SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q}$$

$$D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

$$\mathcal{B} = Q^{N_c}$$

$$D_{SCFT}(\mathcal{B}) = D_{SCFT}(\tilde{\mathcal{B}}) = \frac{3}{2}R_{SCFT}(\mathcal{B}) = \frac{3}{2}N_c\frac{N_f - N_c}{N_f}$$

Can be ruled out via a-theorem $a(R) = 3\text{Tr}U(1)_R^3 - \text{Tr}U(1)_R$

$$a_{\text{UV-safe}} - a_{\text{IR-safe}} < 0$$

SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q}$$

$$D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

$$\mathcal{B} = Q^{N_c}$$

$$D_{SCFT}(\mathcal{B}) = D_{SCFT}(\tilde{\mathcal{B}}) = \frac{3}{2}R_{SCFT}(\mathcal{B}) = \frac{3}{2}N_c\frac{N_f - N_c}{N_f}$$

Can be ruled out via a-theorem $a(R) = 3\text{Tr}U(1)_R^3 - \text{Tr}U(1)_R$

$$a_{\text{UV-safe}} - a_{\text{IR-safe}} < 0$$

SQCD with(out) H cannot be asymptotically safe

SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q}$$

$$D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

$$\mathcal{B} = Q^{N_c}$$

$$D_{SCFT}(\mathcal{B}) = D_{SCFT}(\tilde{\mathcal{B}}) = \frac{3}{2}R_{SCFT}(\mathcal{B}) = \frac{3}{2}N_c\frac{N_f - N_c}{N_f}$$

Can be ruled out via a-theorem $a(R) = 3\text{Tr}U(1)_R^3 - \text{Tr}U(1)_R$

$$a_{\text{UV-safe}} - a_{\text{IR-safe}} < 0$$

SQCD with(out) H cannot be asymptotically safe

*Generalisation to different susy theories using a-maximisation**

What did we achieve?

What did we achieve?

Gauge + fermion + scalars theories can be fund. at any energy scale

What did we achieve?

Gauge + fermion + scalars theories can be fund. at any energy scale

Precise results: independent on scheme choice

What did we achieve?

Gauge + fermion + scalars theories can be fund. at any energy scale

Precise results: independent on scheme choice

Higgs mass squared operator is UV irrelevant

What did we achieve?

Gauge + fermion + scalars theories can be fund. at any energy scale

Precise results: independent on scheme choice

Higgs mass squared operator is UV irrelevant

Existence of UV nontrivial Gauge-Yukawa theories

What did we achieve?

Gauge + fermion + scalars theories can be fund. at any energy scale

Precise results: independent on scheme choice

Higgs mass squared operator is UV irrelevant

Existence of UV nontrivial Gauge-Yukawa theories

Discovered UV complete Non-Abelian QED-like theories

What did we achieve?

Gauge + fermion + scalars theories can be fund. at any energy scale

Precise results: independent on scheme choice

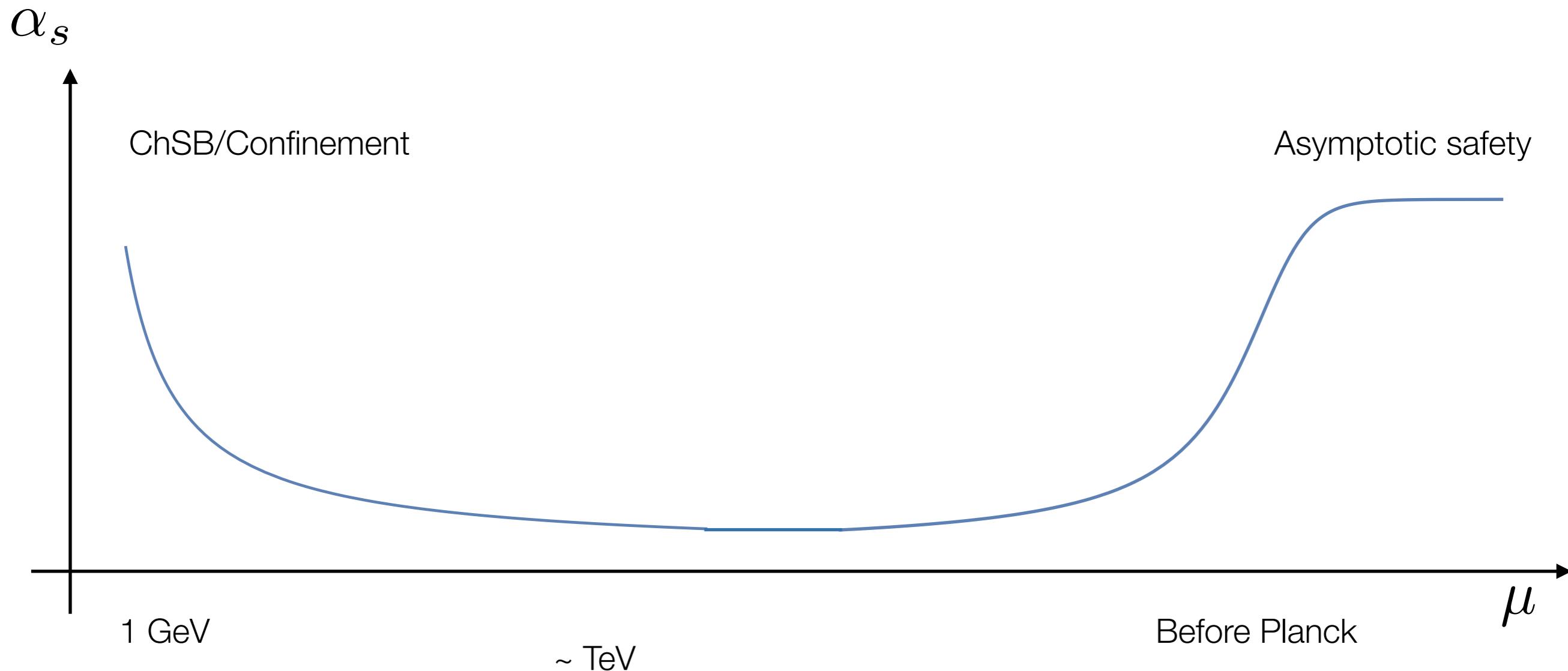
Higgs mass squared operator is UV irrelevant

Existence of UV nontrivial Gauge-Yukawa theories

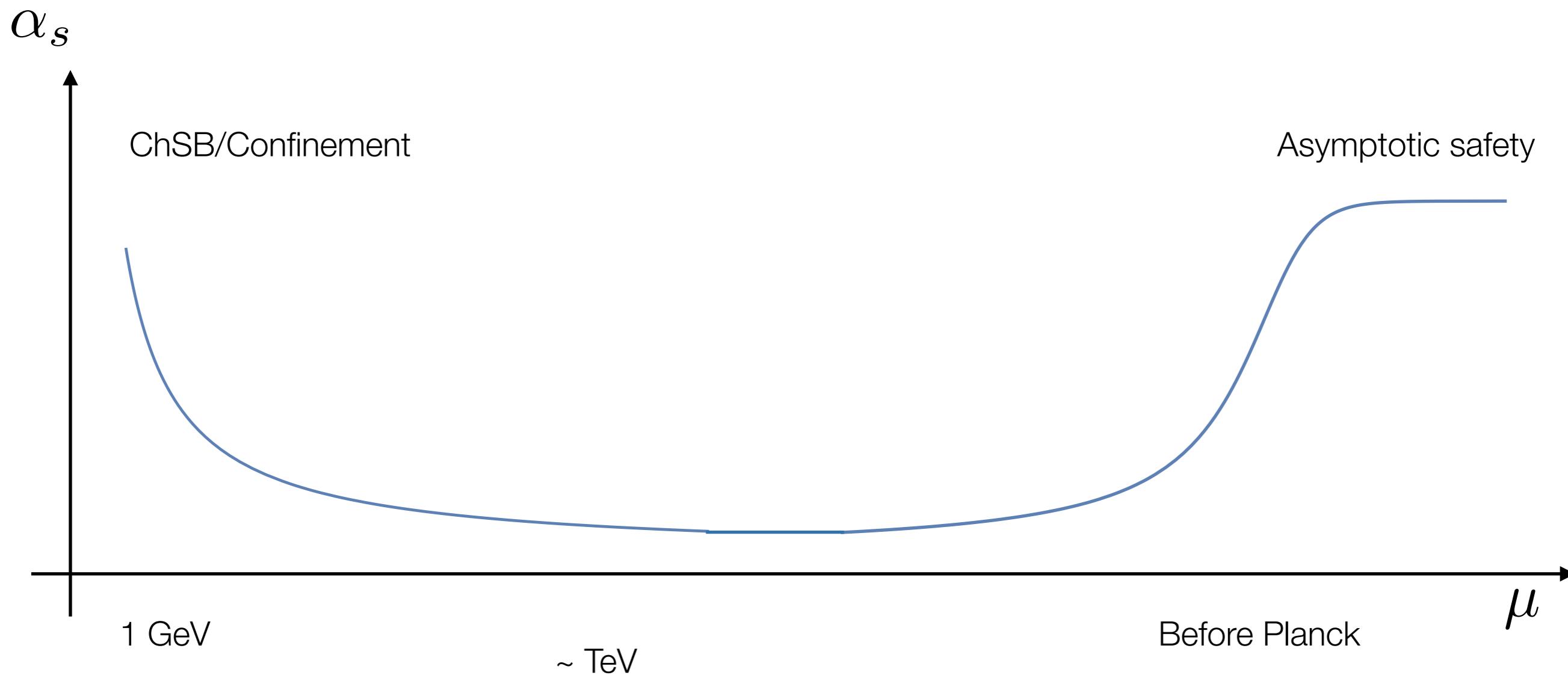
Discovered UV complete Non-Abelian QED-like theories

Scalars are needed to render the UV theory dynamically finite

Safe QCD scenario is testable

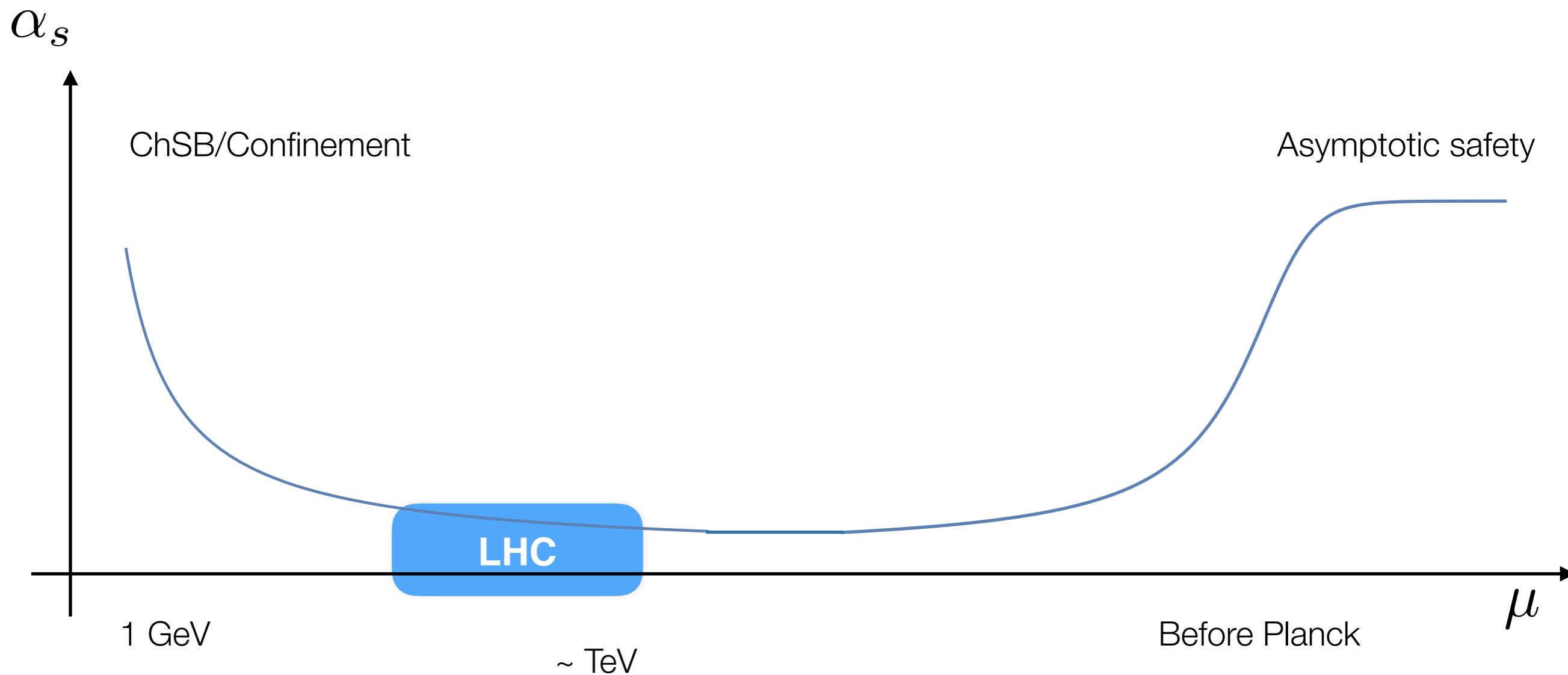


Safe QCD scenario is testable



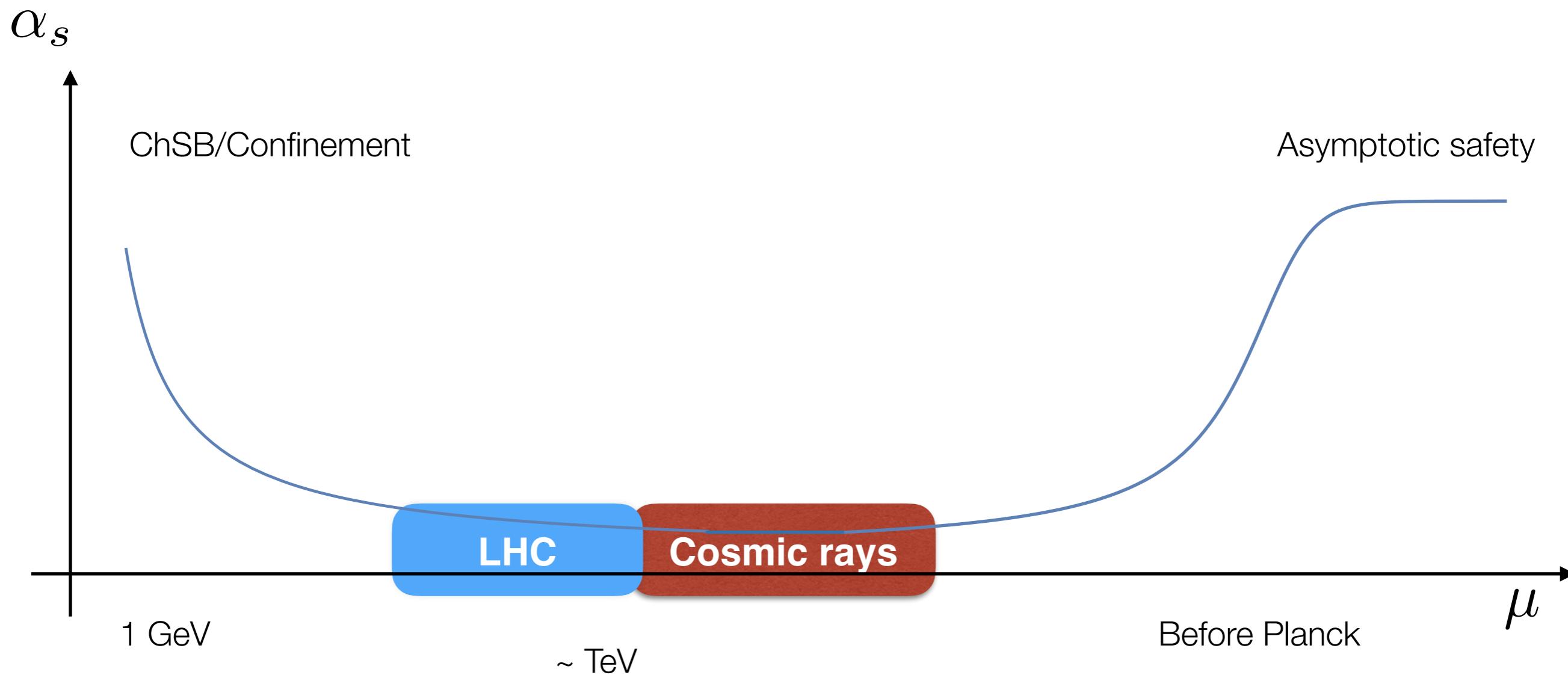
Asymptotic freedom is not a must for UV complete theories

Safe QCD scenario is testable



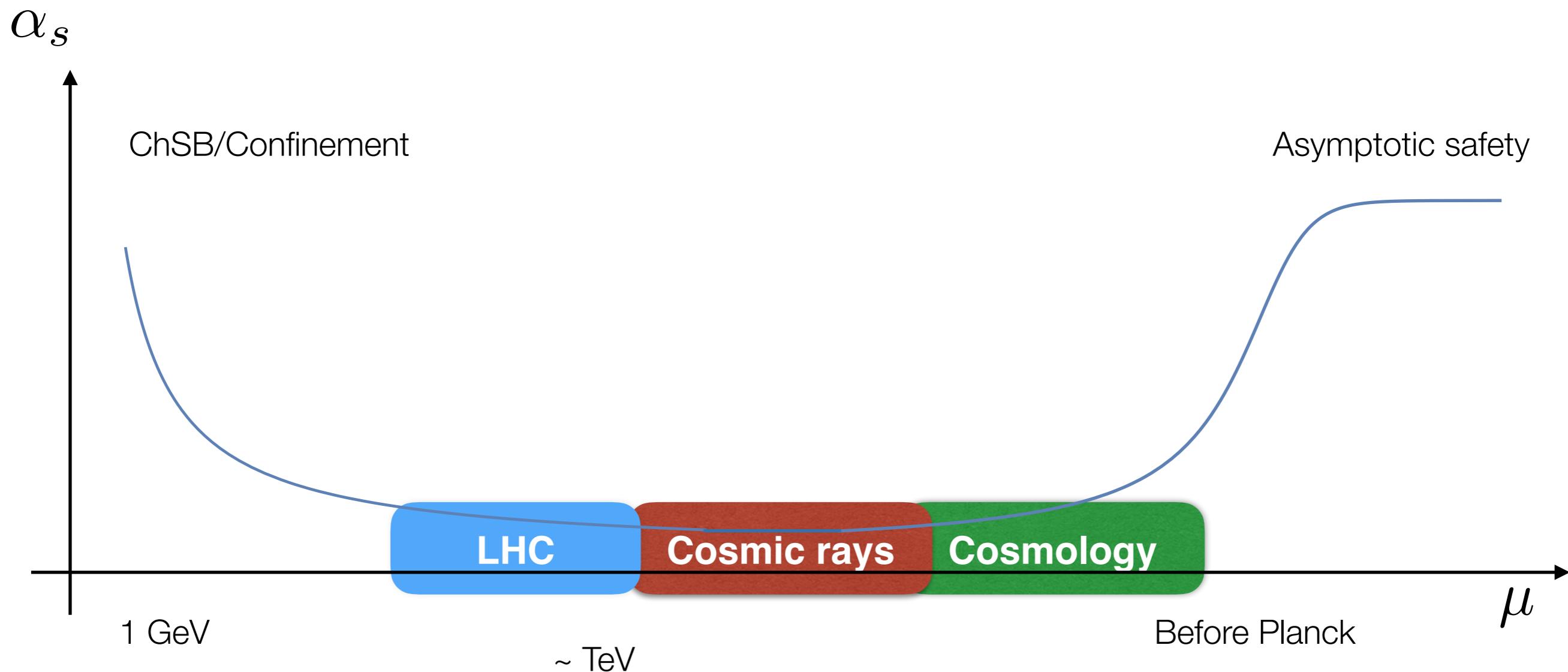
Asymptotic freedom is not a must for UV complete theories

Safe QCD scenario is testable



Asymptotic freedom is not a must for UV complete theories

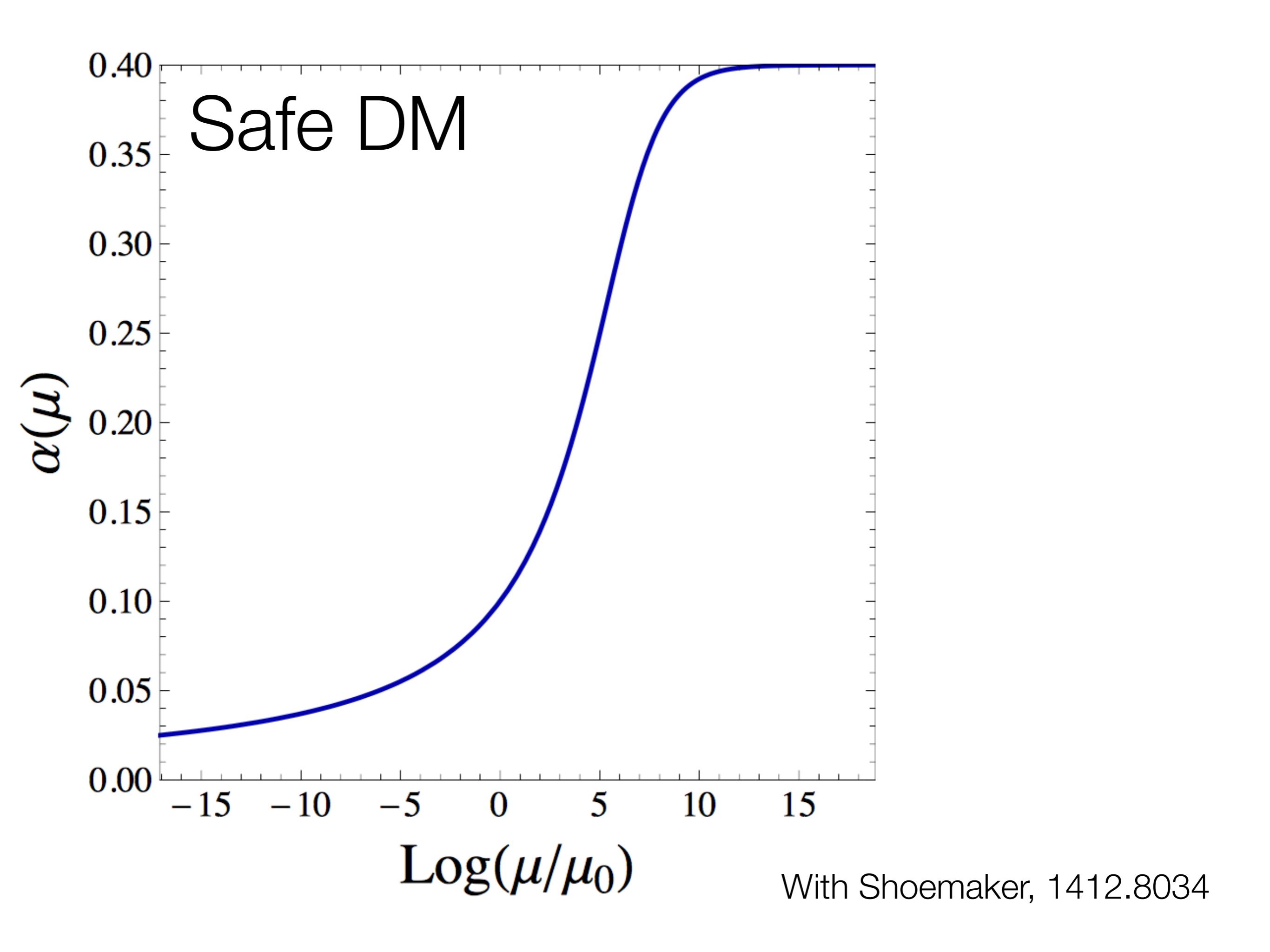
Safe QCD scenario is testable

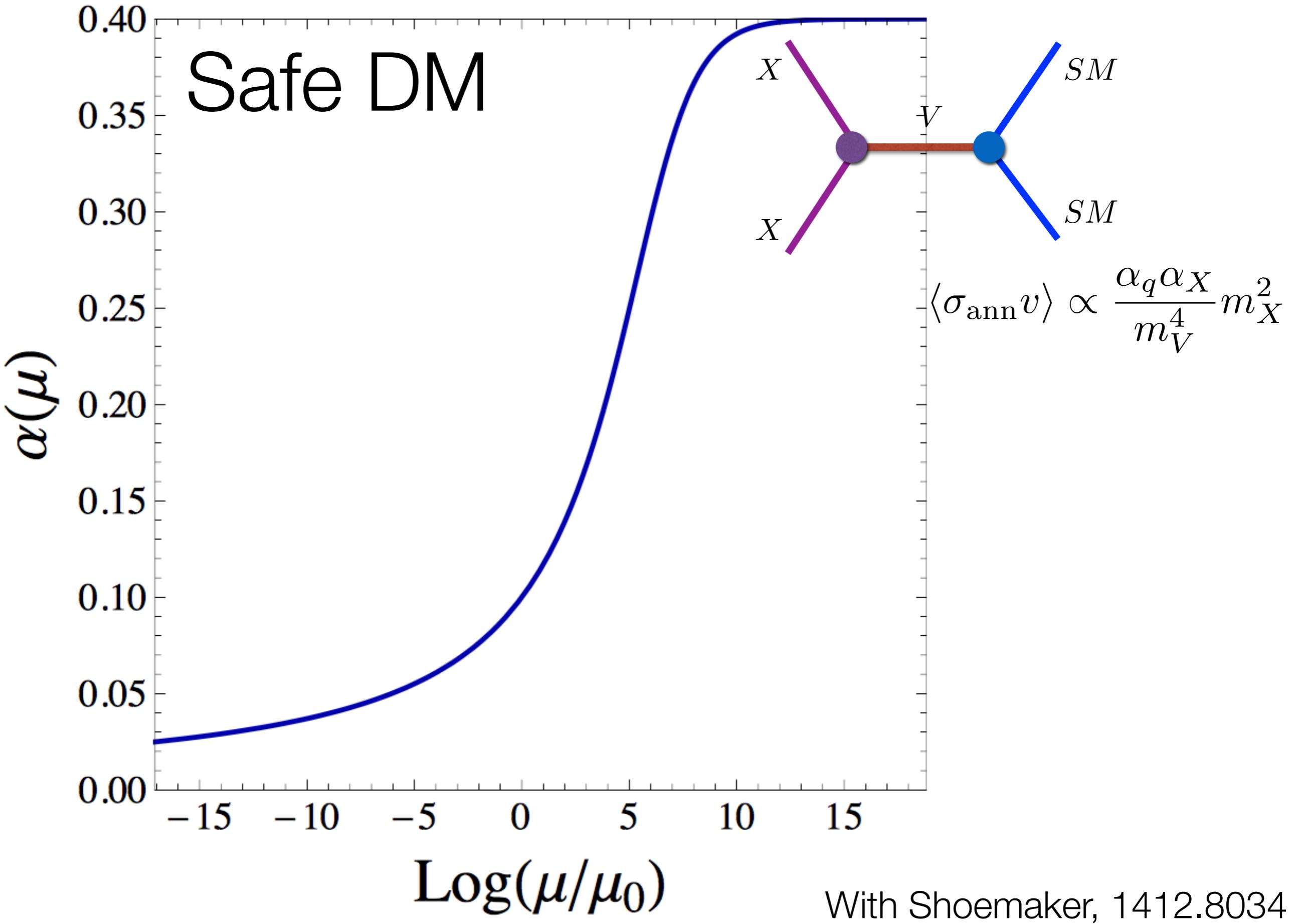


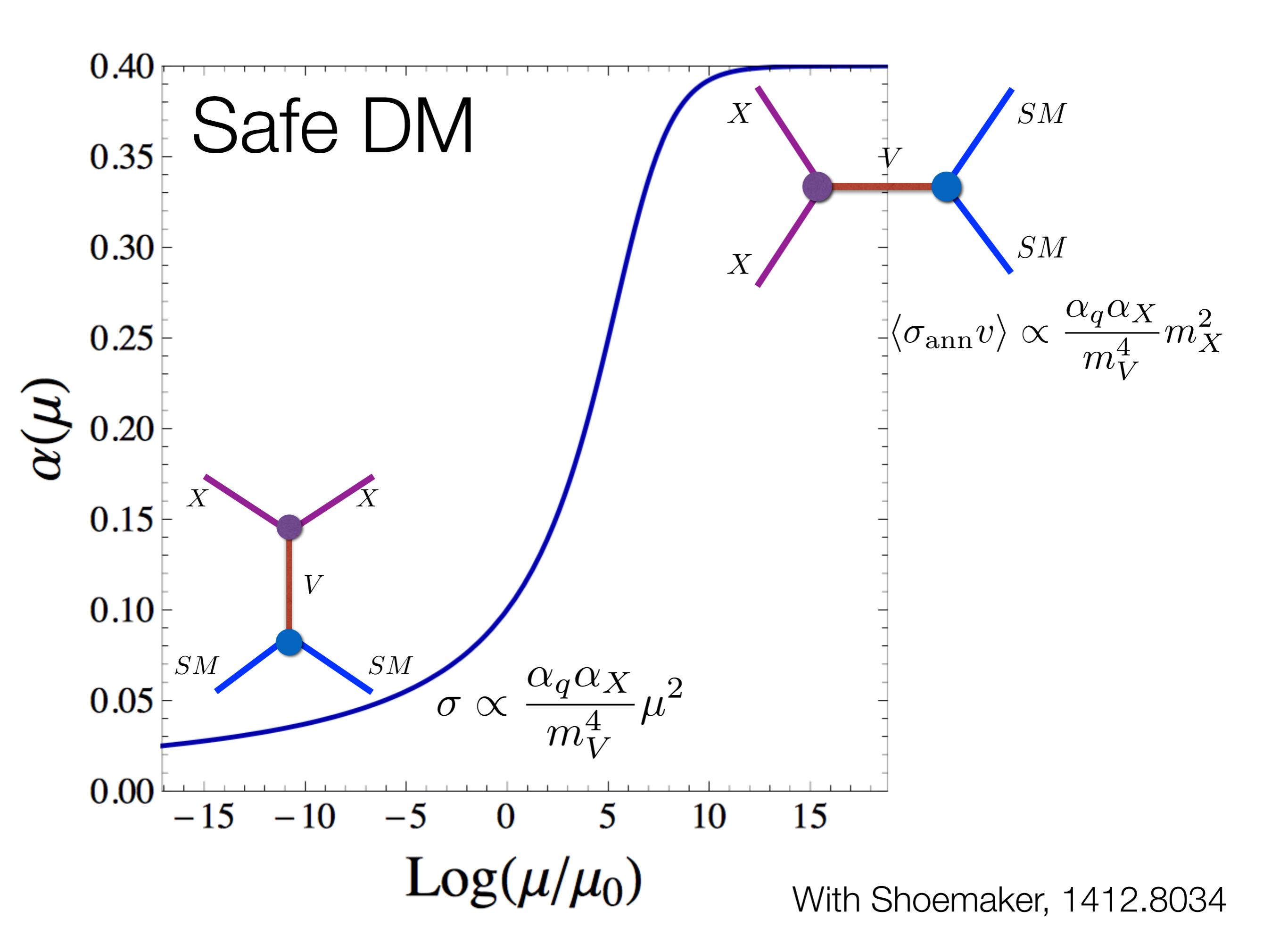
Asymptotic freedom is not a must for UV complete theories

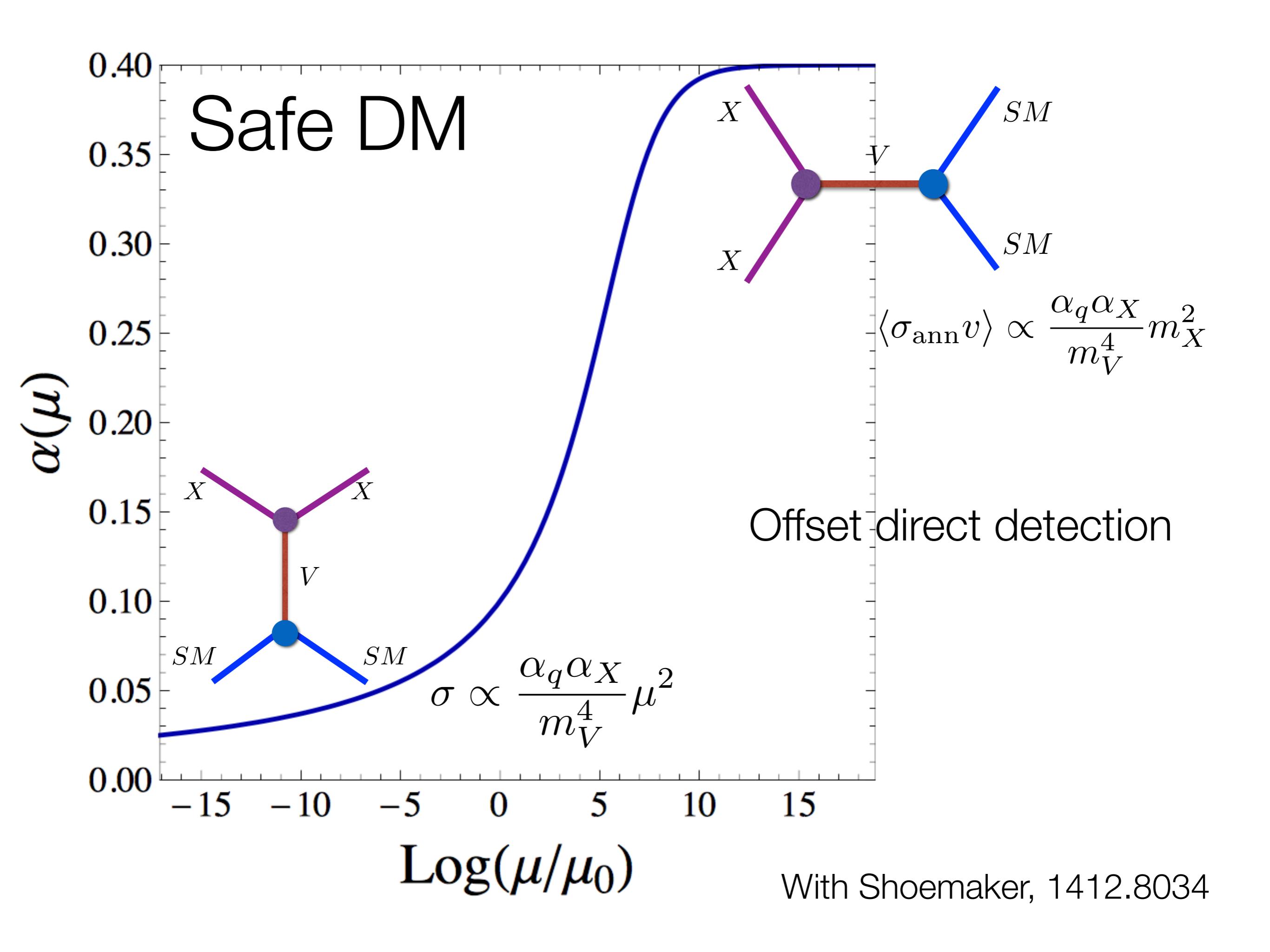
Safe DM

With Shoemaker, 1412.8034









Outlook

Outlook

- ◆ Novel models of DM/Inflation 1412.8034 & 1503.00702
- ◆ Asymptotically safe extensions of the SM (e.g. QCD)
- ◆ New ways to unify flavour, scalar and gauge interactions
- ◆ Super asymptotic safety is not guaranteed*
- ◆ Beyond P.T. (Lattice, dualities, holography, ...)
- ◆ Similarities and differences w.r.t. to N=4 (Wilson loops, MHV)
- ◆ Asymptotic safe quantum gravity **?

* Intriligator and Sannino, 1508.07411

** Weinberg

Future

Future

- ◆ Fundamental theories

Future

- ◆ Fundamental theories
- ◆ Beyond perturbation theory

Future

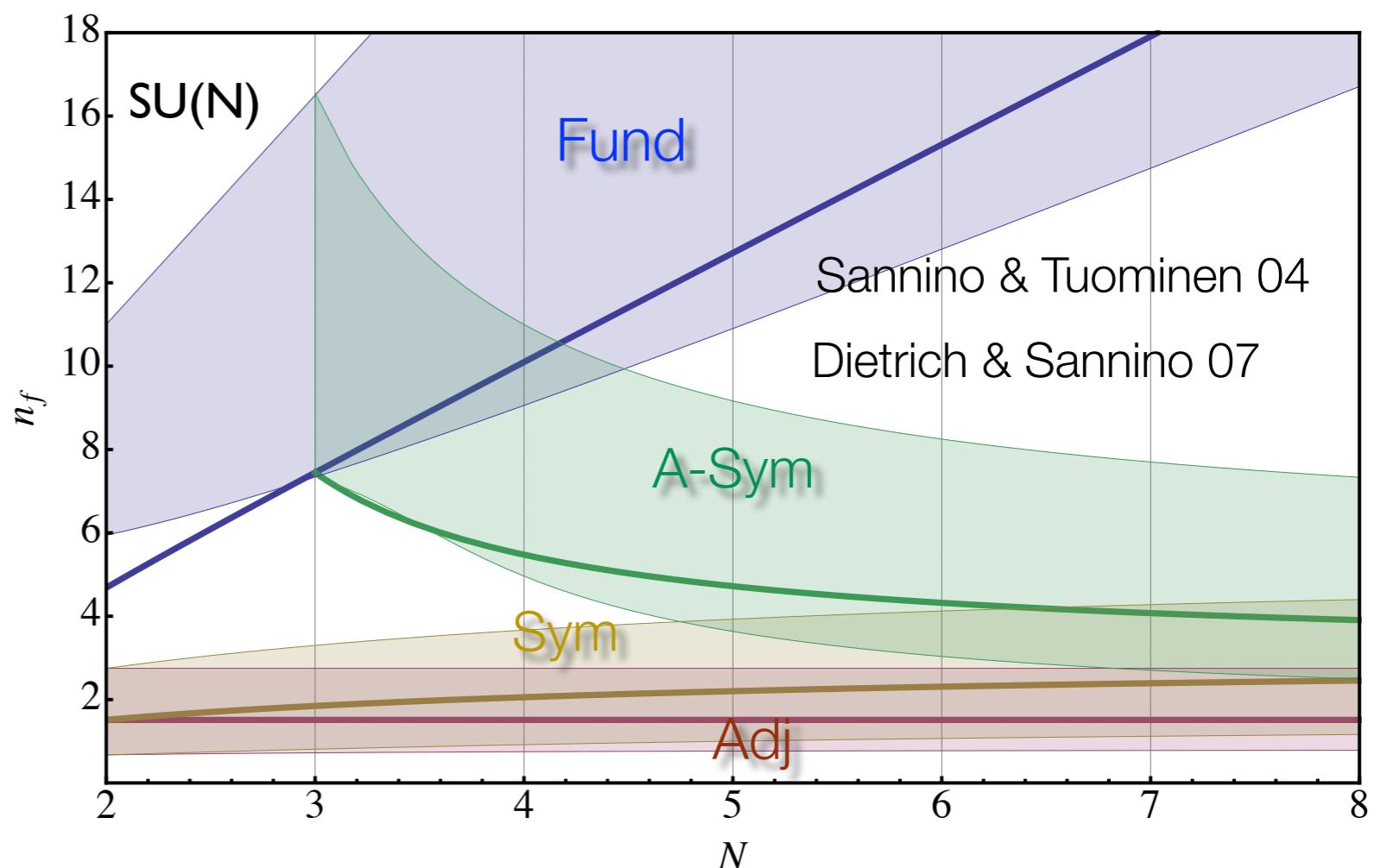
- ◆ Fundamental theories
- ◆ Beyond perturbation theory
- ◆ (Out of) equilibrium thermodynamics

Future

- ◆ Fundamental theories
- ◆ Beyond perturbation theory
- ◆ (Out of) equilibrium thermodynamics
- ◆ Phase diagram of fundamental interactions

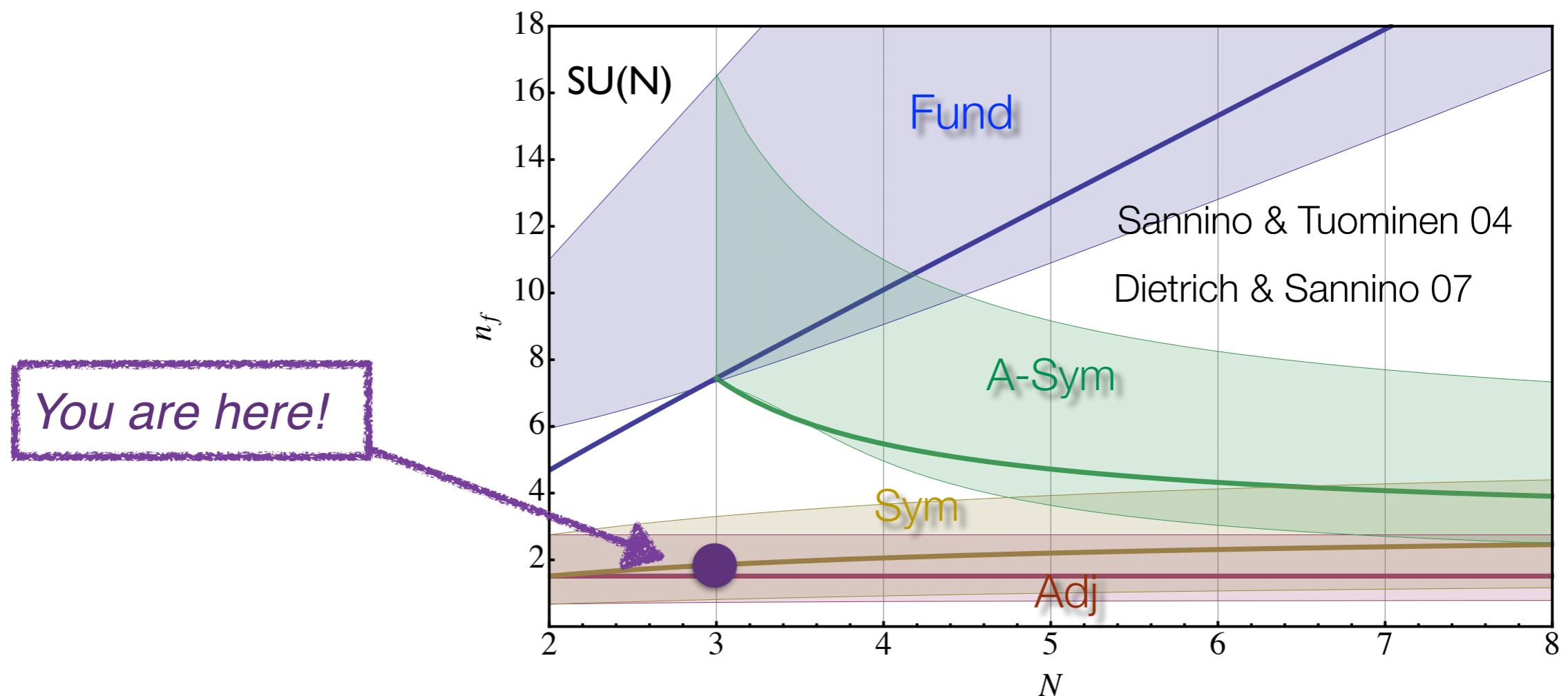
Future

- ◆ Fundamental theories
- ◆ Beyond perturbation theory
- ◆ (Out of) equilibrium thermodynamics
- ◆ Phase diagram of fundamental interactions

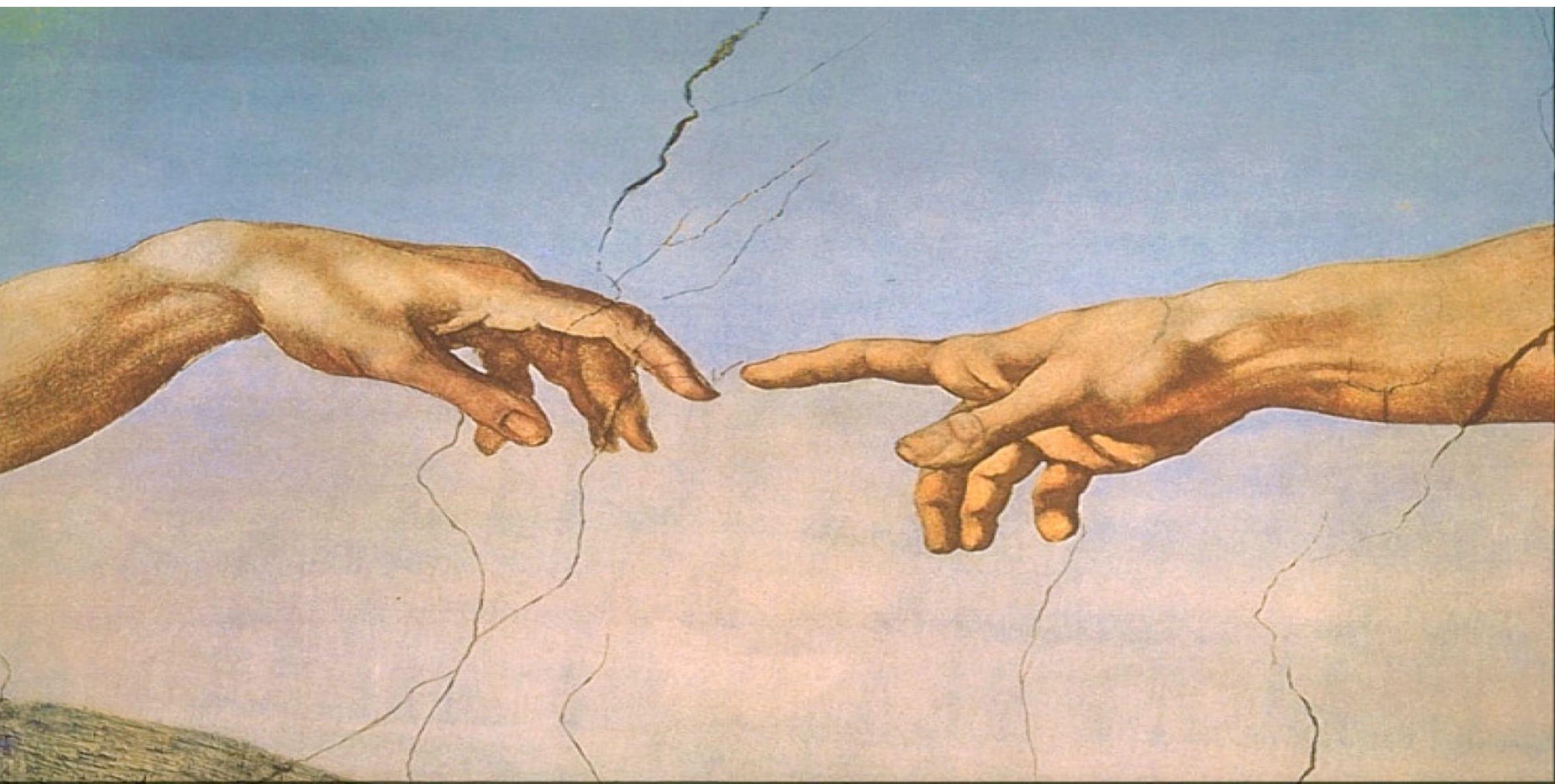


Future

- ◆ Fundamental theories
- ◆ Beyond perturbation theory
- ◆ (Out of) equilibrium thermodynamics
- ◆ Phase diagram of fundamental interactions



We pursue fundamental theories of Nature



'THE CREATION' - MICHELANGELO

Backup slides

a-theorem

$$\mathcal{L} = \mathcal{L}_{CFT} + g_i \mathcal{O}^i$$

a-theorem

$$\mathcal{L} = \mathcal{L}_{CFT} + g_i \mathcal{O}^i \longrightarrow \text{Quantum correct., marginal oper.}$$

a-theorem

$$\mathcal{L} = \mathcal{L}_{CFT} + g_i \mathcal{O}^i \longrightarrow$$

Quantum correct., marginal oper.

$$g_i = g_i(x)$$

Tool: Curved backgrounds

a-theorem

$$\mathcal{L} = \mathcal{L}_{CFT} + g_i \mathcal{O}^i \longrightarrow$$

Quantum correct., marginal oper.

$$g_i = g_i(x)$$

Tool: Curved backgrounds

$$\gamma_{\mu\nu} \rightarrow e^{2\sigma(x)} \gamma_{\mu\nu}$$

Conformal transformation

$$g_i(\mu) \rightarrow g_i(e^{-\sigma(x)} \mu)$$

a-theorem

$$\mathcal{L} = \mathcal{L}_{CFT} + g_i \mathcal{O}^i \longrightarrow$$

Quantum correct., marginal oper.

$$g_i = g_i(x)$$

Tool: Curved backgrounds

$$\gamma_{\mu\nu} \rightarrow e^{2\sigma(x)} \gamma_{\mu\nu}$$

Conformal transformation

$$g_i(\mu) \rightarrow g_i(e^{-\sigma(x)} \mu)$$

$$W = \log \left[\int \mathcal{D}\Phi e^{i \int d^4x \mathcal{L}} \right]$$

Variation of the generating functional

Weyl (anomaly) relations

$$\Delta_\sigma W \equiv \int d^4x \, \sigma(x) \left(2\gamma_{\mu\nu} \frac{\delta W}{\delta \gamma_{\mu\nu}} - \beta_i \frac{\delta W}{\delta g_i} \right) = \sigma \left(aE(\gamma) + \chi^{ij} \partial_\mu g_i \partial_\nu g_j G^{\mu\nu} \right) + \partial_\mu \sigma w^i \partial_\nu g_i G^{\mu\nu} + \dots$$

Weyl (anomaly) relations

$$\Delta_\sigma W \equiv \int d^4x \sigma(x) \left(2\gamma_{\mu\nu} \frac{\delta W}{\delta \gamma_{\mu\nu}} - \beta_i \frac{\delta W}{\delta g_i} \right) = \sigma \left(aE(\gamma) + \chi^{ij} \partial_\mu g_i \partial_\nu g_j G^{\mu\nu} \right) + \partial_\mu \sigma w^i \partial_\nu g_i G^{\mu\nu} + \dots$$

$$E(\gamma) = R^{\mu\nu\rho\sigma} R_{\mu\nu\rho\sigma} - 4R^{\mu\nu} R_{\mu\nu} + R^2 \quad \text{Euler density}$$

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2}\gamma^{\mu\nu}R \quad \text{Einstein tensor}$$

Weyl (anomaly) relations

$$\Delta_\sigma W \equiv \int d^4x \sigma(x) \left(2\gamma_{\mu\nu} \frac{\delta W}{\delta \gamma_{\mu\nu}} - \beta_i \frac{\delta W}{\delta g_i} \right) = \sigma \left(aE(\gamma) + \chi^{ij} \partial_\mu g_i \partial_\nu g_j G^{\mu\nu} \right) + \partial_\mu \sigma w^i \partial_\nu g_i G^{\mu\nu} + \dots$$

$$E(\gamma) = R^{\mu\nu\rho\sigma} R_{\mu\nu\rho\sigma} - 4R^{\mu\nu} R_{\mu\nu} + R^2 \quad \text{Euler density}$$

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2}\gamma^{\mu\nu}R \quad \text{Einstein tensor}$$

$$\beta_i \quad \text{Beta functions}$$

$$a, \quad \chi^{ij}, \quad \omega^i \quad \text{Functions of couplings}$$

Weyl (anomaly) relations

$$\Delta_\sigma W \equiv \int d^4x \sigma(x) \left(2\gamma_{\mu\nu} \frac{\delta W}{\delta \gamma_{\mu\nu}} - \beta_i \frac{\delta W}{\delta g_i} \right) = \sigma \left(aE(\gamma) + \chi^{ij} \partial_\mu g_i \partial_\nu g_j G^{\mu\nu} \right) + \partial_\mu \sigma w^i \partial_\nu g_i G^{\mu\nu} + \dots$$

$$E(\gamma) = R^{\mu\nu\rho\sigma} R_{\mu\nu\rho\sigma} - 4R^{\mu\nu} R_{\mu\nu} + R^2$$

Euler density

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2}\gamma^{\mu\nu}R$$

Einstein tensor

$$\beta_i$$

Beta functions

$$a, \quad \chi^{ij}, \quad \omega^i$$

Functions of couplings

Weyl relations from abelian nature of Weyl anomaly

$$\Delta_\sigma \Delta_\tau W = \Delta_\tau \Delta_\sigma W$$

Perturbative a-theorem

$$\tilde{a} \equiv a - w^i \beta_i \quad \frac{\partial \tilde{a}}{\partial g_i} = \left(-\chi^{ij} + \frac{\partial w^i}{\partial g_j} - \frac{\partial w^j}{\partial g_i} \right) \beta_j$$

Perturbative a-theorem

$$\tilde{a} \equiv a - w^i \beta_i \quad \frac{\partial \tilde{a}}{\partial g_i} = \left(-\chi^{ij} + \frac{\partial w^i}{\partial g_j} - \frac{\partial w^j}{\partial g_i} \right) \beta_j$$

$$\frac{d}{d\mu} \tilde{a} = -\chi^{ij} \beta_i \beta_j$$

Perturbative a-theorem

$$\tilde{a} \equiv a - w^i \beta_i \quad \frac{\partial \tilde{a}}{\partial g_i} = \left(-\chi^{ij} + \frac{\partial w^i}{\partial g_j} - \frac{\partial w^j}{\partial g_i} \right) \beta_j$$

$$\frac{d}{d\mu} \tilde{a} = -\chi^{ij} \beta_i \beta_j$$

a-tilde is RG monotonically decreasing if chi is positive definite

Perturbative a-theorem

$$\tilde{a} \equiv a - w^i \beta_i \quad \frac{\partial \tilde{a}}{\partial g_i} = \left(-\chi^{ij} + \frac{\partial w^i}{\partial g_j} - \frac{\partial w^j}{\partial g_i} \right) \beta_j$$

$$\frac{d}{d\mu} \tilde{a} = -\chi^{ij} \beta_i \beta_j$$

a-tilde is RG monotonically decreasing if chi is positive definite

Cardy 88, conjecture

Perturbative a-theorem

$$\tilde{a} \equiv a - w^i \beta_i \quad \frac{\partial \tilde{a}}{\partial g_i} = \left(-\chi^{ij} + \frac{\partial w^i}{\partial g_j} - \frac{\partial w^j}{\partial g_i} \right) \beta_j$$

$$\frac{d}{d\mu} \tilde{a} = -\chi^{ij} \beta_i \beta_j$$

a-tilde is RG monotonically decreasing if chi is positive definite

Cardy 88, conjecture

True in lowest order PT

Osborn 89 & 91, Jack & Osborn 90

Perturbative a-theorem

$$\tilde{a} \equiv a - w^i \beta_i \quad \frac{\partial \tilde{a}}{\partial g_i} = \left(-\chi^{ij} + \frac{\partial w^i}{\partial g_j} - \frac{\partial w^j}{\partial g_i} \right) \beta_j$$

$$\frac{d}{d\mu} \tilde{a} = -\chi^{ij} \beta_i \beta_j$$

a-tilde is RG monotonically decreasing if chi is positive definite

Cardy 88, conjecture

True in lowest order PT

Osborn 89 & 91, Jack & Osborn 90

Analyticity: a-tilde bigger in UV

Komargodski & Schwimmer 11, Komargodski 12

Gauge - Yukawa theories/Gradient Flow

Antipin, Gillioz, Mølgaard, Sannino 13

Jack and Poole 15

Gauge - Yukawa theories/Gradient Flow

omega is an exact form

Osborn 89 & 91, Jack & Osborn 90

Antipin, Gillioz, Mølgaard, Sannino 13

Jack and Poole 15

Gauge - Yukawa theories/Gradient Flow

omega is an exact form

Osborn 89 & 91, Jack & Osborn 90

$$\frac{\partial \tilde{a}}{\partial g_i} = \left(-\chi^{ij} + \frac{\partial w^i}{\partial g_j} - \frac{\partial w^j}{\partial g_i} \right) \beta_j \quad \Rightarrow \quad \frac{\partial \tilde{a}}{\partial g_i} = -\beta^i, \quad \beta^i \equiv \chi^{ij} \beta_j$$

Antipin, Gillioz, Mølgaard, Sannino 13

Jack and Poole 15

Gauge - Yukawa theories/Gradient Flow

omega is an exact form

Osborn 89 & 91, Jack & Osborn 90

$$\frac{\partial \tilde{a}}{\partial g_i} = \left(-\chi^{ij} + \frac{\partial w^i}{\partial g_j} - \frac{\partial w^j}{\partial g_i} \right) \beta_j \quad \Rightarrow \quad \frac{\partial \tilde{a}}{\partial g_i} = -\beta^i, \quad \beta^i \equiv \chi^{ij} \beta_j$$

Gradient flow fundamental relation

$$\frac{\partial \beta^j}{\partial g_i} = \frac{\partial \beta^i}{\partial g_j},$$

Antipin, Gillioz, Mølgaard, Sannino 13

Jack and Poole 15

Gauge - Yukawa theories/Gradient Flow

omega is an exact form

Osborn 89 & 91, Jack & Osborn 90

$$\frac{\partial \tilde{a}}{\partial g_i} = \left(-\chi^{ij} + \frac{\partial w^i}{\partial g_j} - \frac{\partial w^j}{\partial g_i} \right) \beta_j \quad \Rightarrow \quad \frac{\partial \tilde{a}}{\partial g_i} = -\beta^i, \quad \beta^i \equiv \chi^{ij} \beta_j$$

Gradient flow fundamental relation

$$\frac{\partial \beta^j}{\partial g_i} = \frac{\partial \beta^i}{\partial g_j},$$

Relations among the modified β of different couplings

Precise prescription for expanding beta functions in perturb. theory

Antipin, Gillioz, Mølgaard, Sannino 13

Jack and Poole 15

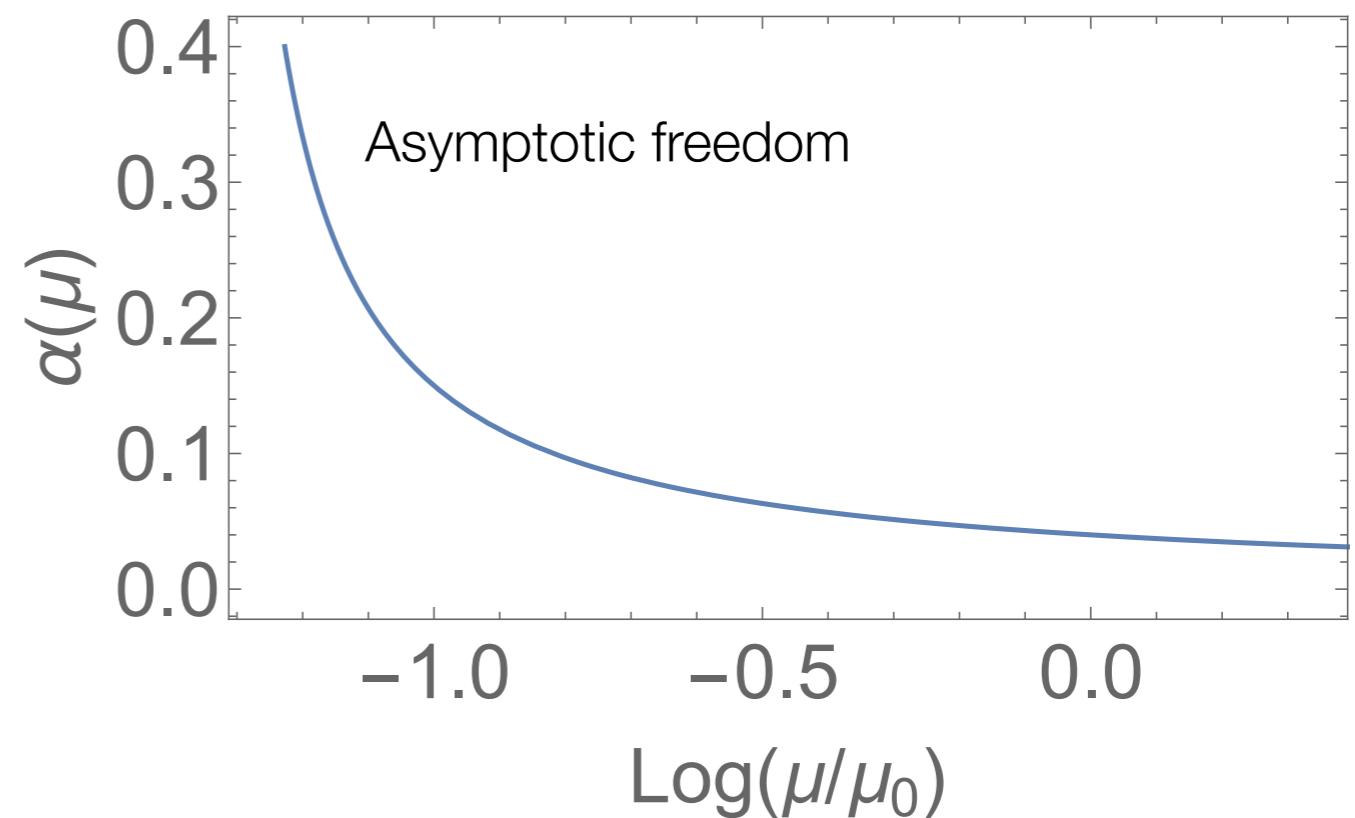
The Compositeness Solution

The Compositeness Solution

- ◆ EW scale = Composite scale
- ◆ UV non-interacting

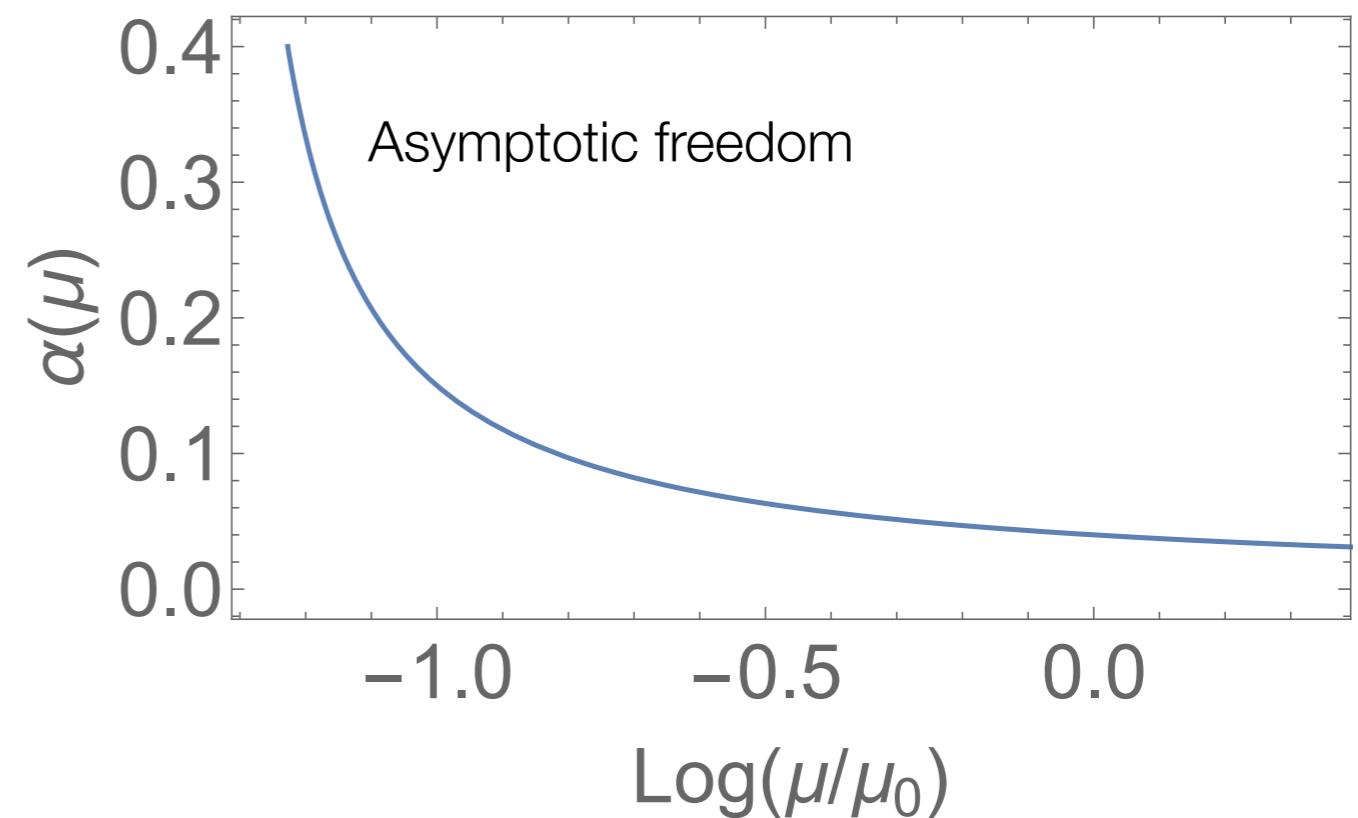
The Compositeness Solution

- ◆ EW scale = Composite scale
- ◆ UV non-interacting



The Compositeness Solution

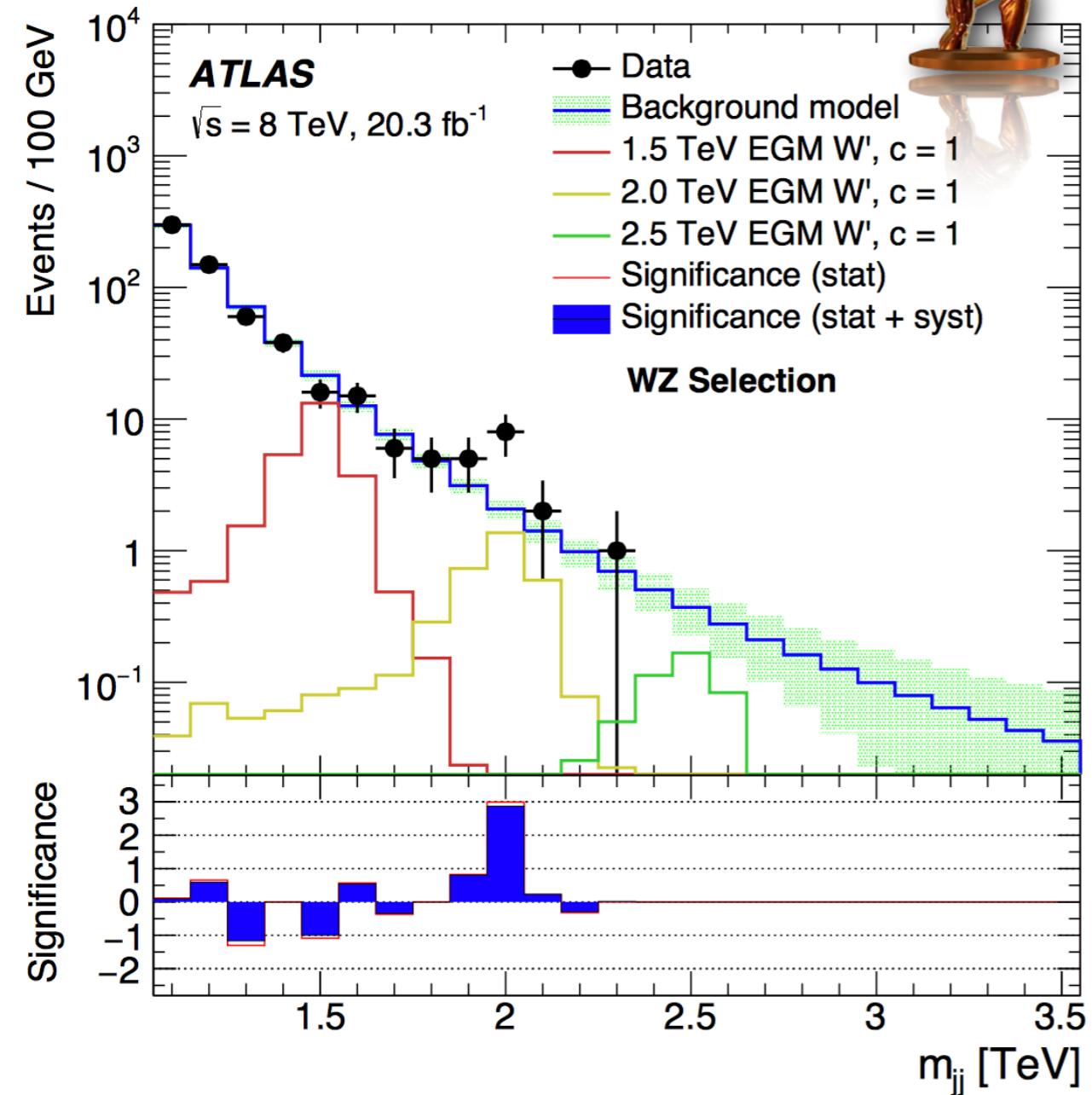
- ◆ EW scale = Composite scale
- ◆ UV non-interacting
- ◆ Not ruled out



Arbey et al. 1502.04718



Diboson excesses

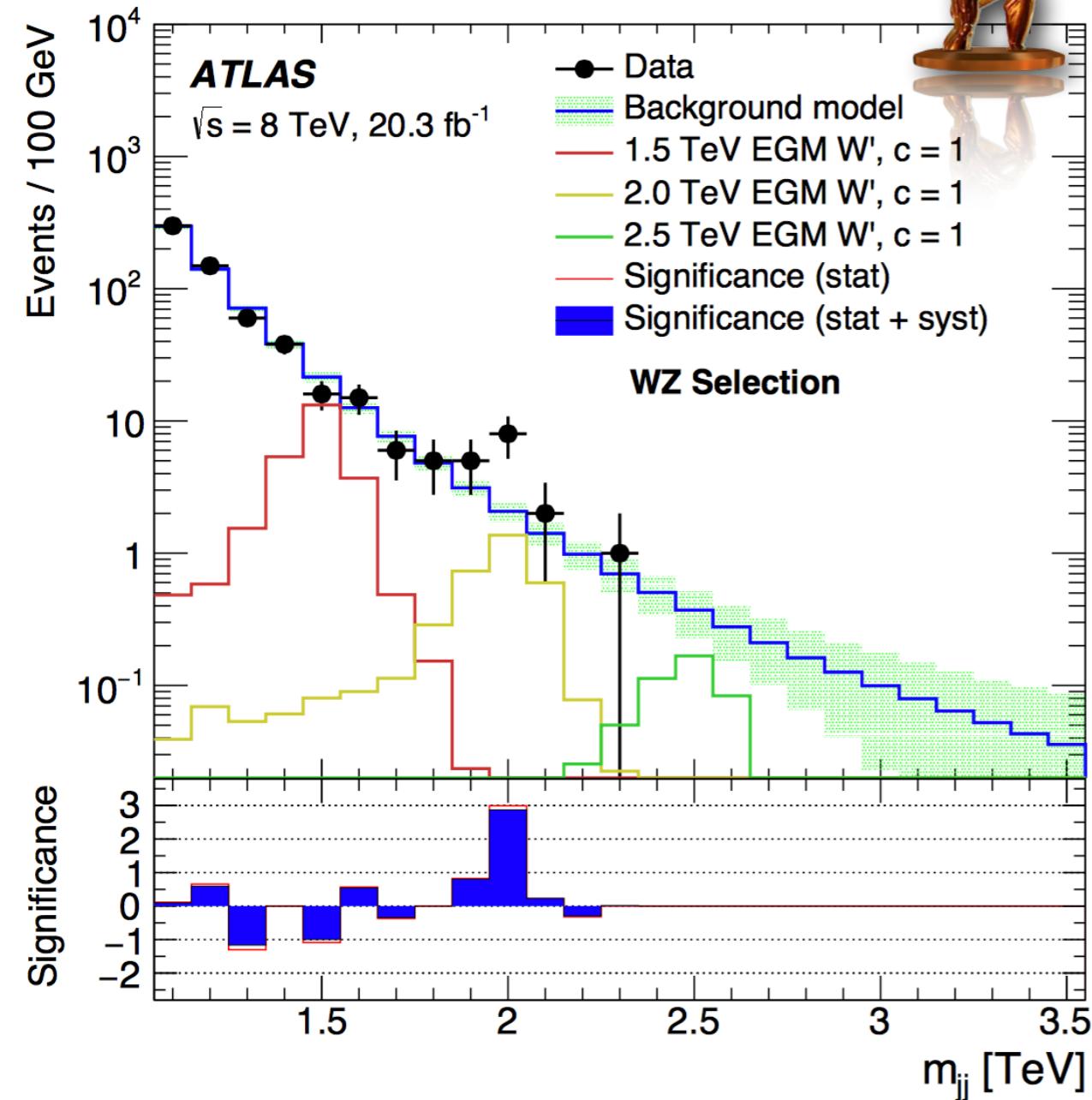


- ◆ WZ, WW and ZZ excesses

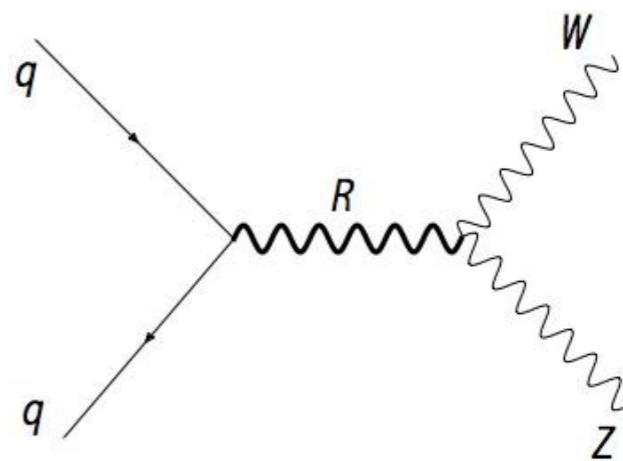
3.4, 2.6, 2.9 σ



Diboson excesses



Frandsen, Franzosi, Sannino hep-ph/1506.04392



General effective Lagrangian analysis

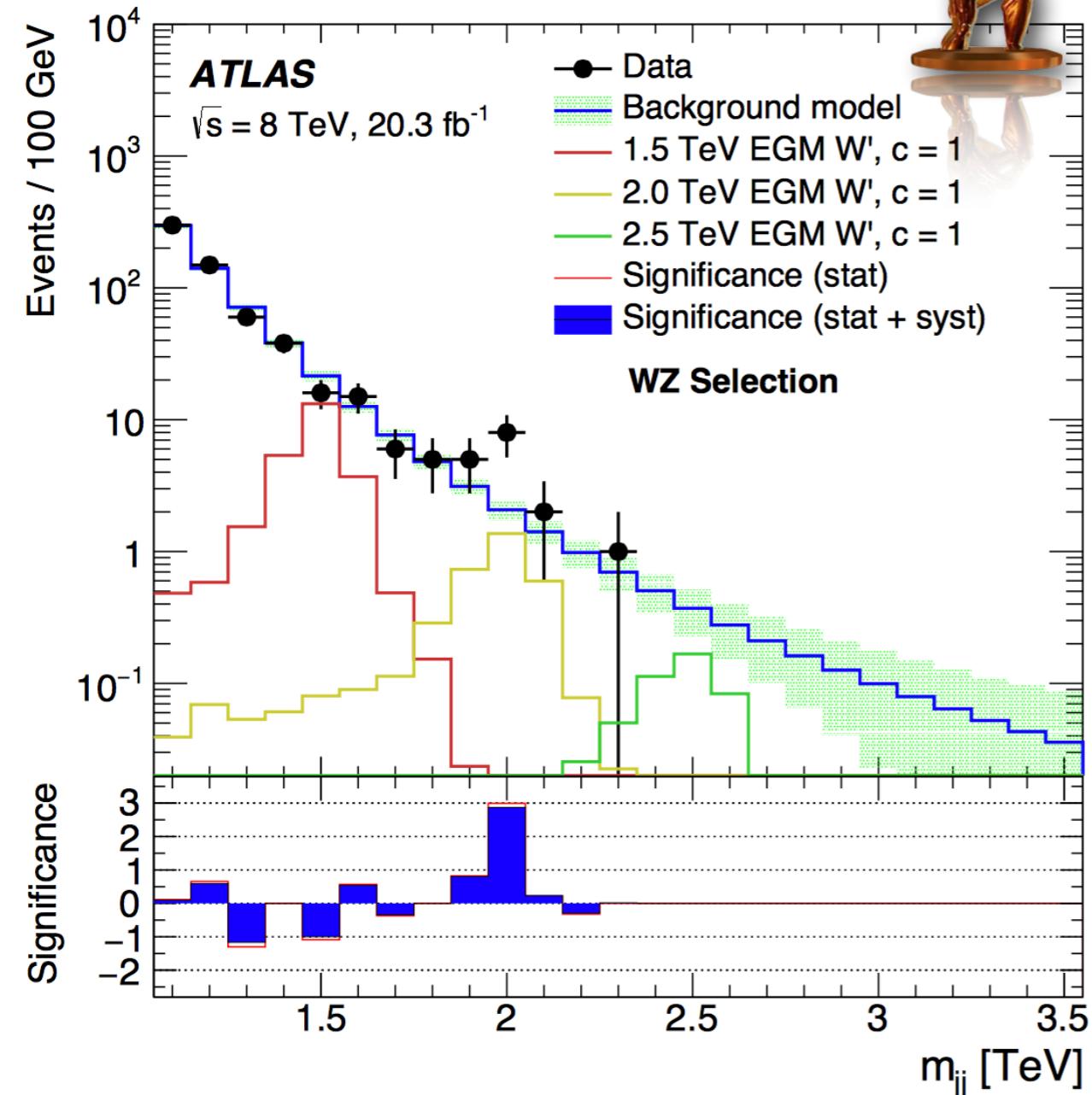
$$\sigma(pp \rightarrow R \rightarrow WZ) \sim (10 - 30) \text{ fb}$$

- ◆ WZ, WW and ZZ excesses

3.4, 2.6, 2.9 σ



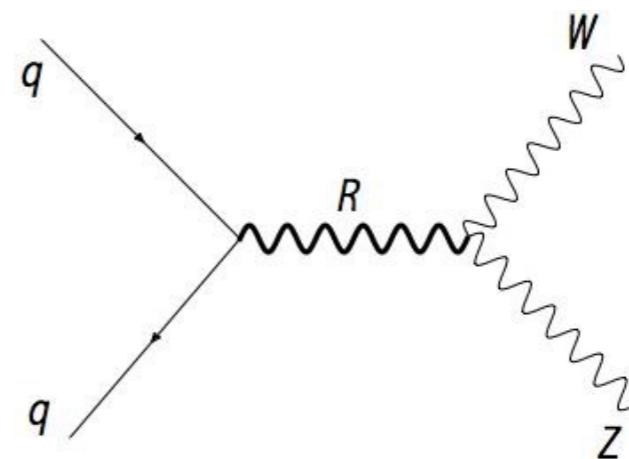
Diboson excesses



- ◆ WZ, WW and ZZ excesses

3.4, 2.6, 2.9 σ

Frandsen, Franzosi, Sannino hep-ph/1506.04392



General effective Lagrangian analysis

$$\sigma(pp \rightarrow R \rightarrow WZ) \sim (10 - 30) \text{ fb}$$

- ✓ Spin-one mass range
- ✓ Vector decay constant
- ✓ Dileptons constraints
- ✓ Dijets constraints

Lattice predictions

Lattice predictions

Minimal (Walking) Technicolor

$SU(3)$ symmetric $f_\pi \simeq 246$ GeV

$m_{R_V} \simeq 1.75 \pm 0.1$ TeV

$m_{R_A} \simeq 2.3 \pm 0.1$ TeV

Sannino, Tuominen 0405209

Dietrich, Sannino, Tuominen 0405209

Fodor, Holland, Kuti, Nogradi, Schroeder, Wong, 1209.0391

Lattice predictions

Minimal (Walking) Technicolor

$SU(3)$ symmetric $f_\pi \simeq 246$ GeV

$m_{R_V} \simeq 1.75 \pm 0.1$ TeV

$m_{R_A} \simeq 2.3 \pm 0.1$ TeV

Sannino, Tuominen 0405209

Dietrich, Sannino, Tuominen 0405209

Fodor, Holland, Kuti, Nogradi, Schroeder, Wong, 1209.0391

(Ultra) Minimal Technicolor & Composite (Goldstone) Higgs

$SU(2)$ fundamental

$m_{R_V} \simeq (2.5 \pm 0.5) \frac{\text{TeV}}{\sin \theta}$

Ryttov, Sannino, 0809.0713

Cacciapaglia, Sannino, 1402.0233

$m_{R_A} \simeq (3.3 \pm 0.7) \frac{\text{TeV}}{\sin \theta}$

Lewis, Pica, Sannino, 1109.3513

Knobs



Knobs

Gauge Group: SU, SO, SP, Exceptional



Knobs



Gauge Group: SU, SO, SP, Exceptional
Matter Representation

Knobs



Gauge Group: SU, SO, SP, Exceptional

Matter Representation

of Flavors per Representation

Knobs



Gauge Group: SU, SO, SP, Exceptional

Matter Representation

of Flavors per Representation

4 Fermi interactions

Knobs

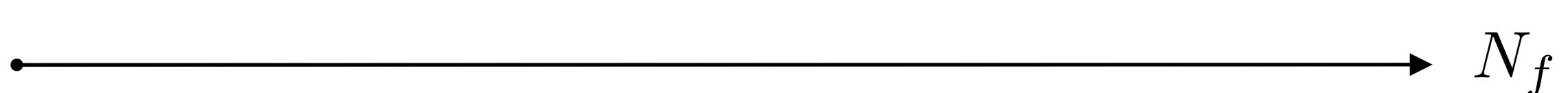


Gauge Group: SU, SO, SP, Exceptional

Matter Representation

of Flavors per Representation

4 Fermi interactions



Knobs

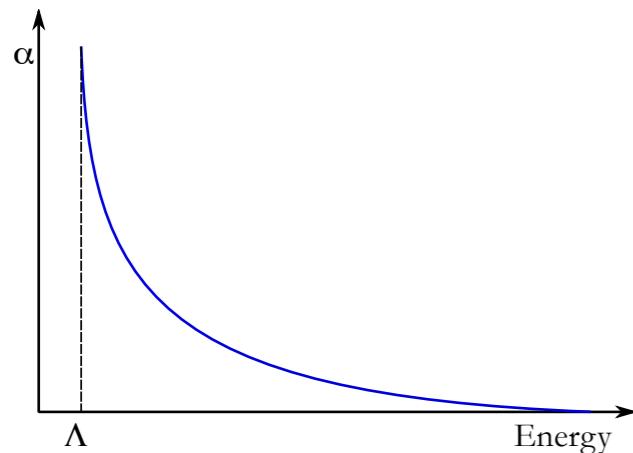
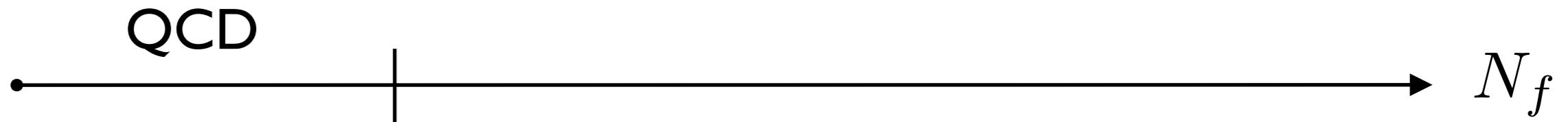


Gauge Group: SU, SO, SP, Exceptional

Matter Representation

of Flavors per Representation

4 Fermi interactions



Knobs

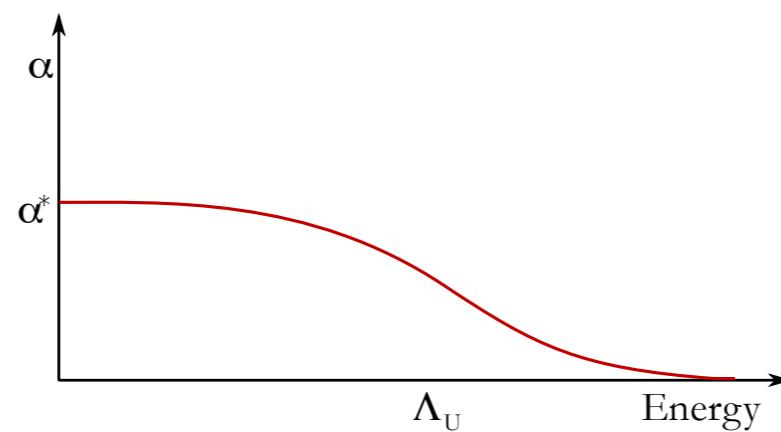
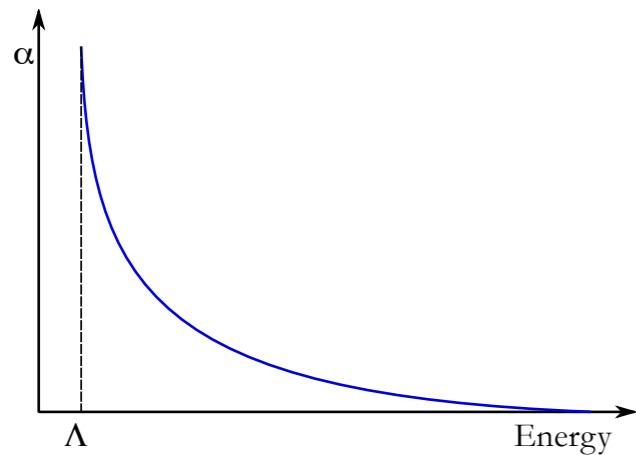


Gauge Group: SU, SO, SP, Exceptional

Matter Representation

of Flavors per Representation

4 Fermi interactions



Knobs

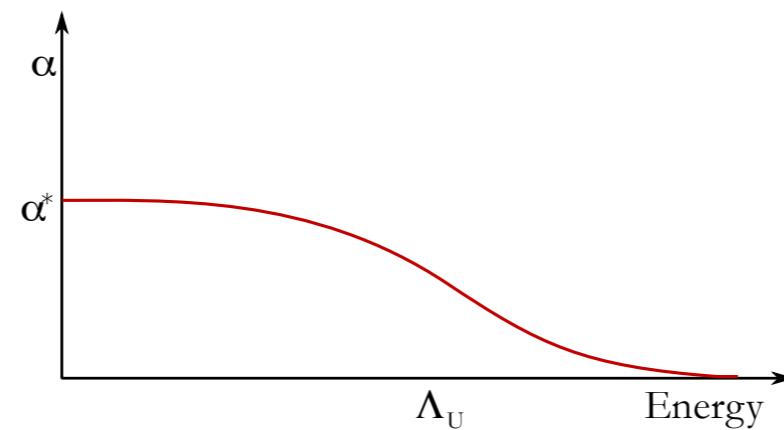
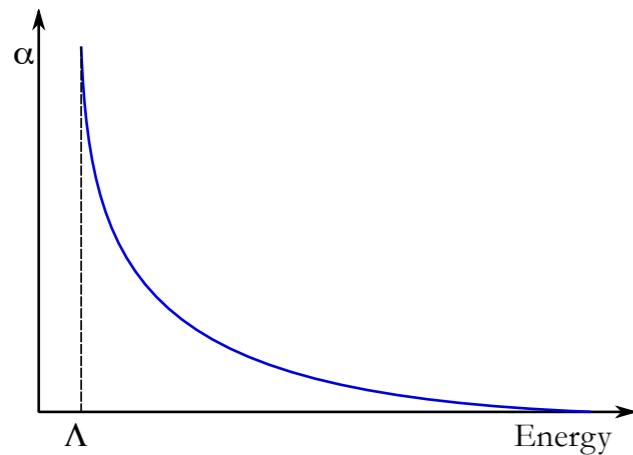


Gauge Group: SU, SO, SP, Exceptional

Matter Representation

of Flavors per Representation

4 Fermi interactions



Knobs

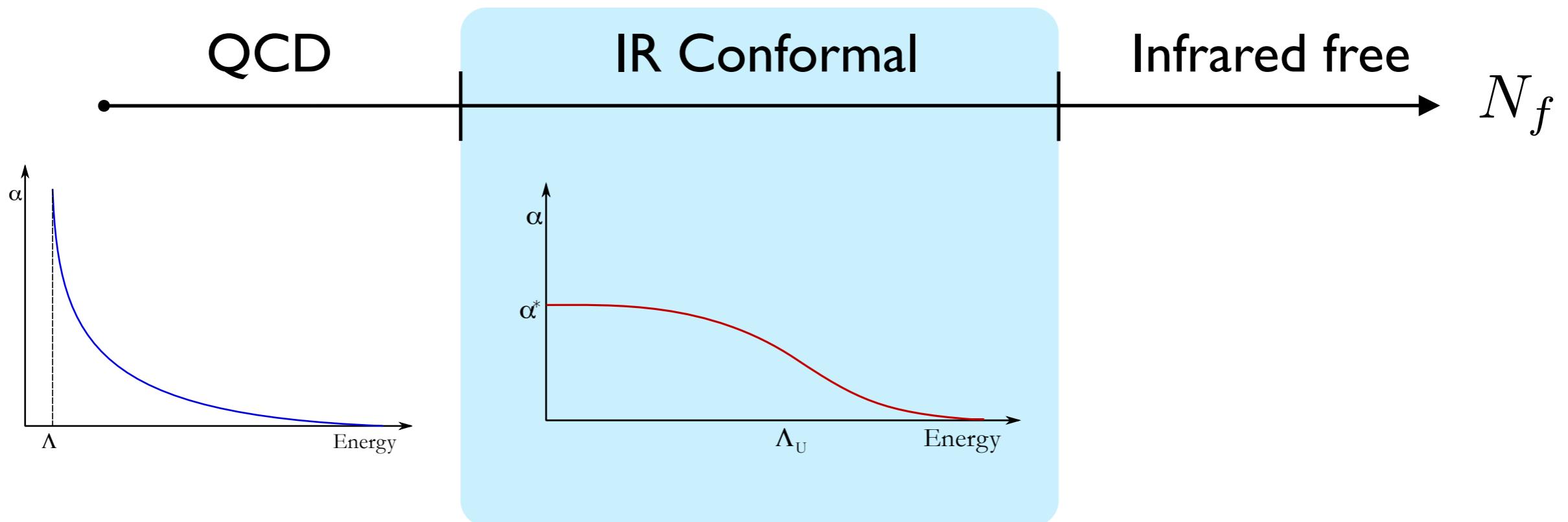


Gauge Group: SU, SO, SP, Exceptional

Matter Representation

of Flavors per Representation

4 Fermi interactions



Knobs

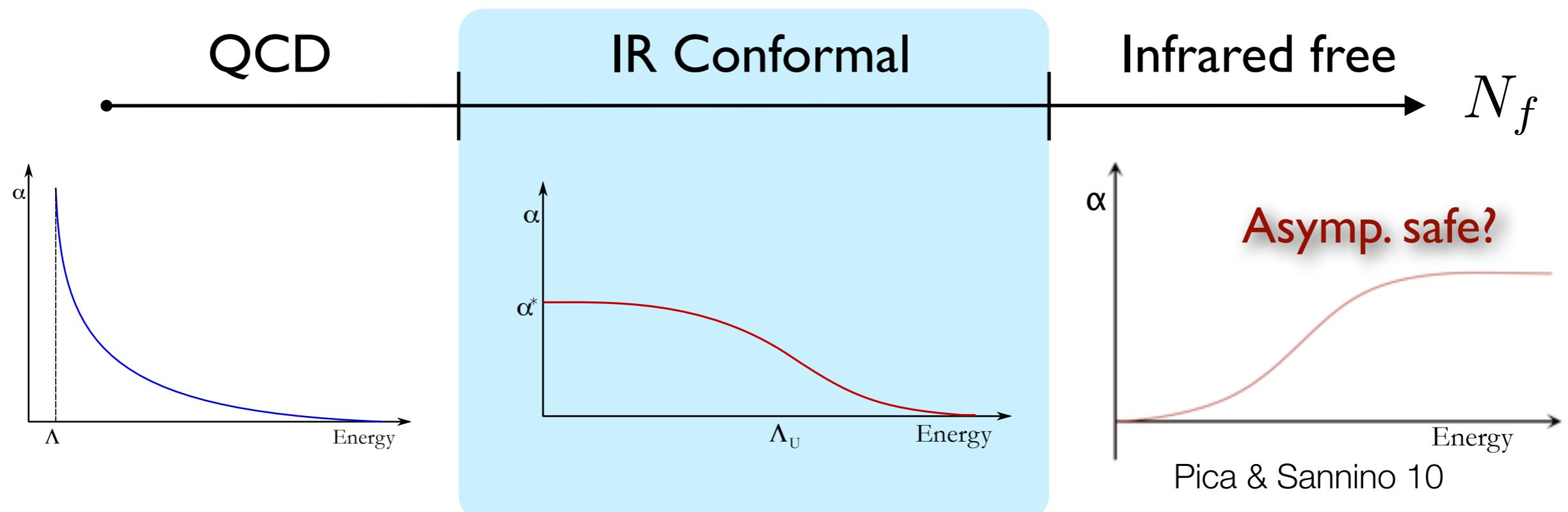


Gauge Group: SU, SO, SP, Exceptional

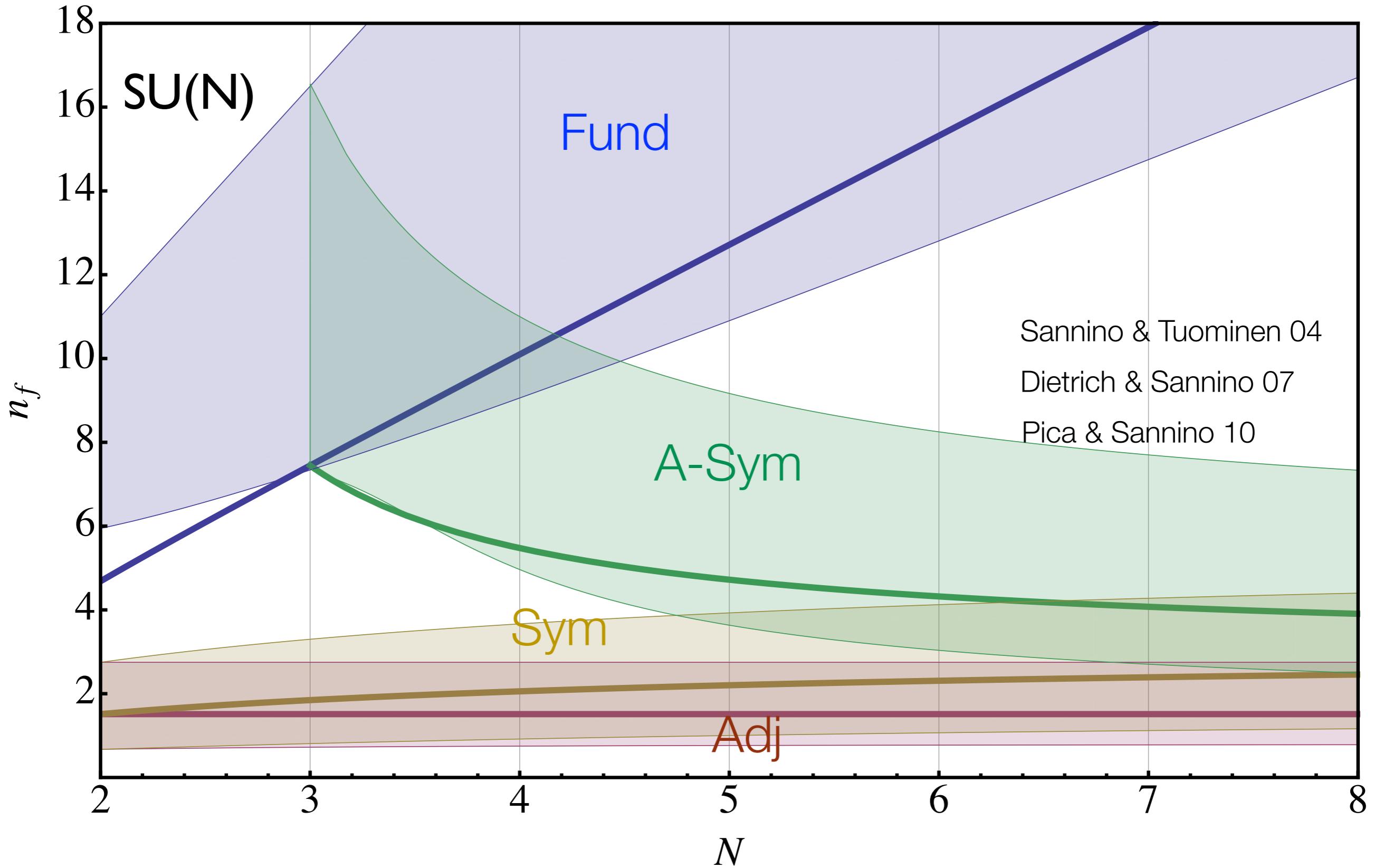
Matter Representation

of Flavors per Representation

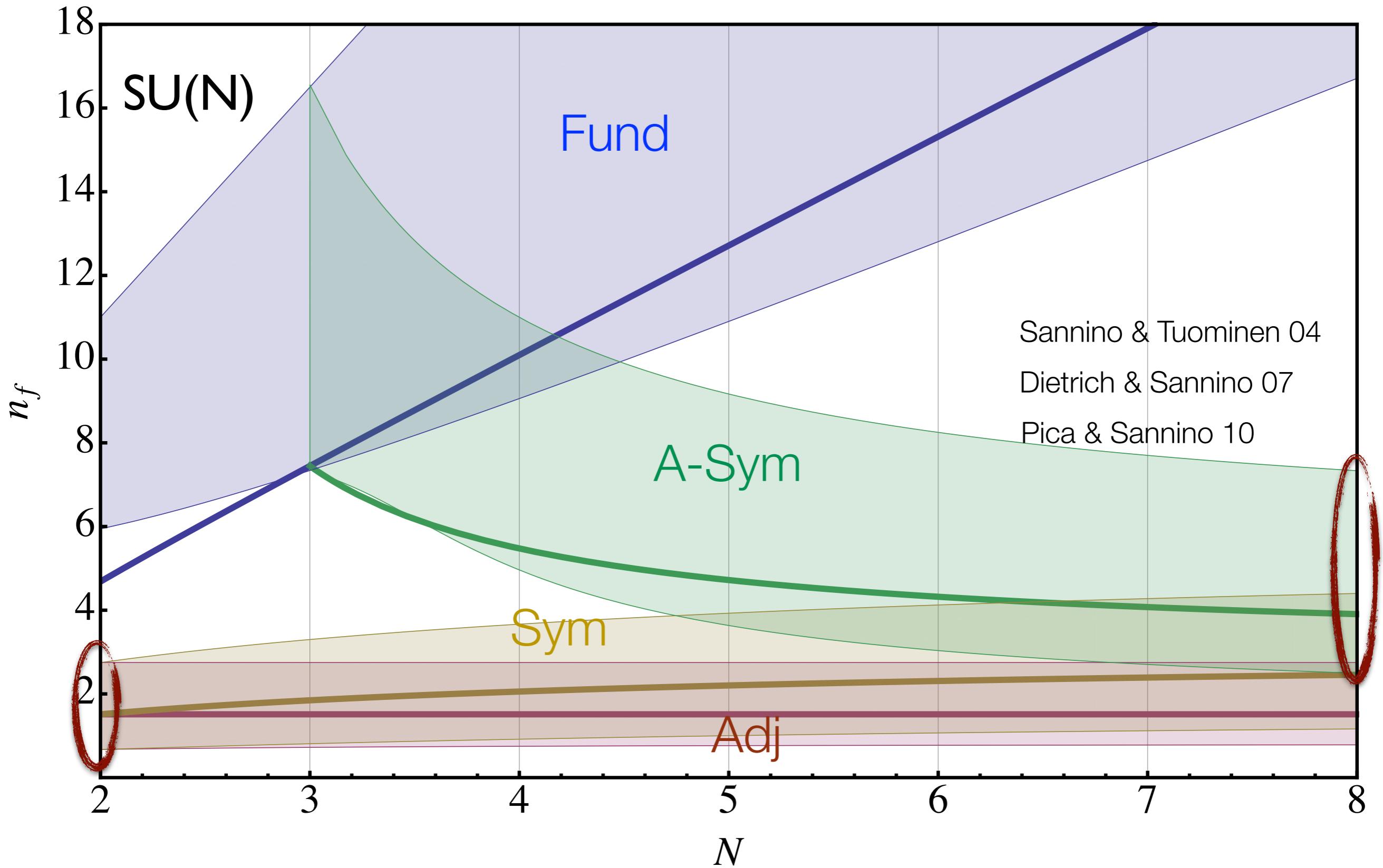
4 Fermi interactions



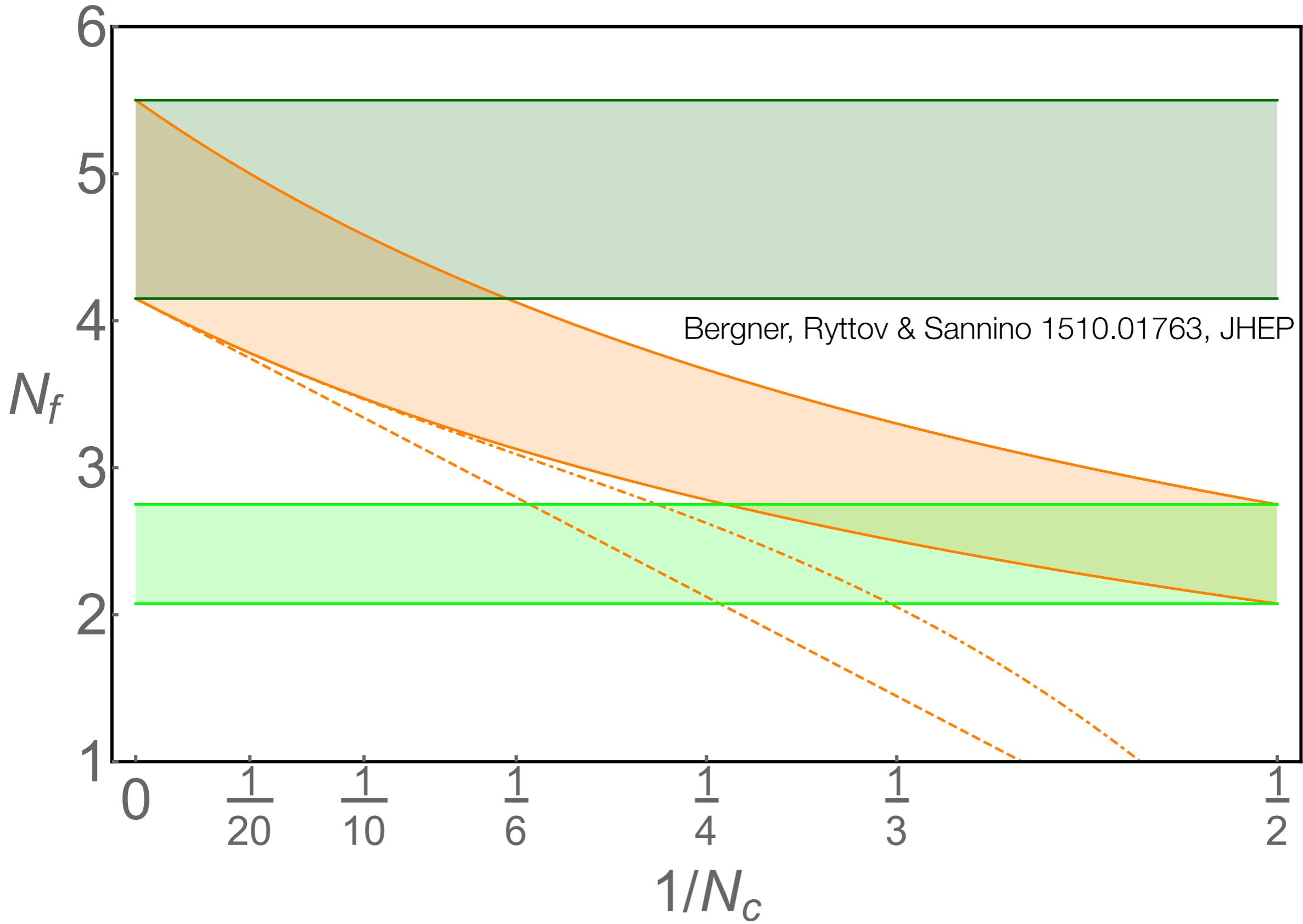
SU(N) Phase Diagram



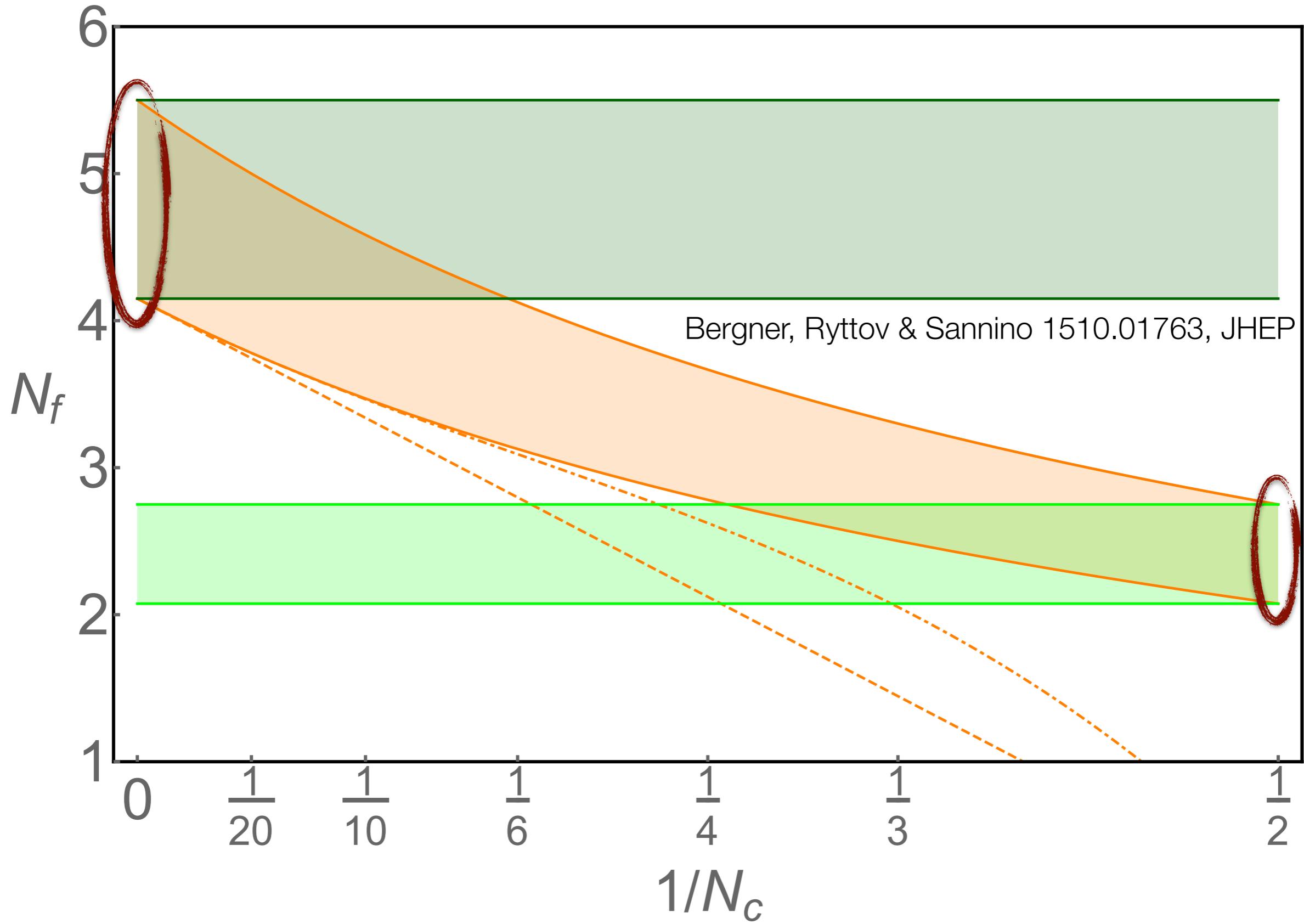
SU(N) Phase Diagram



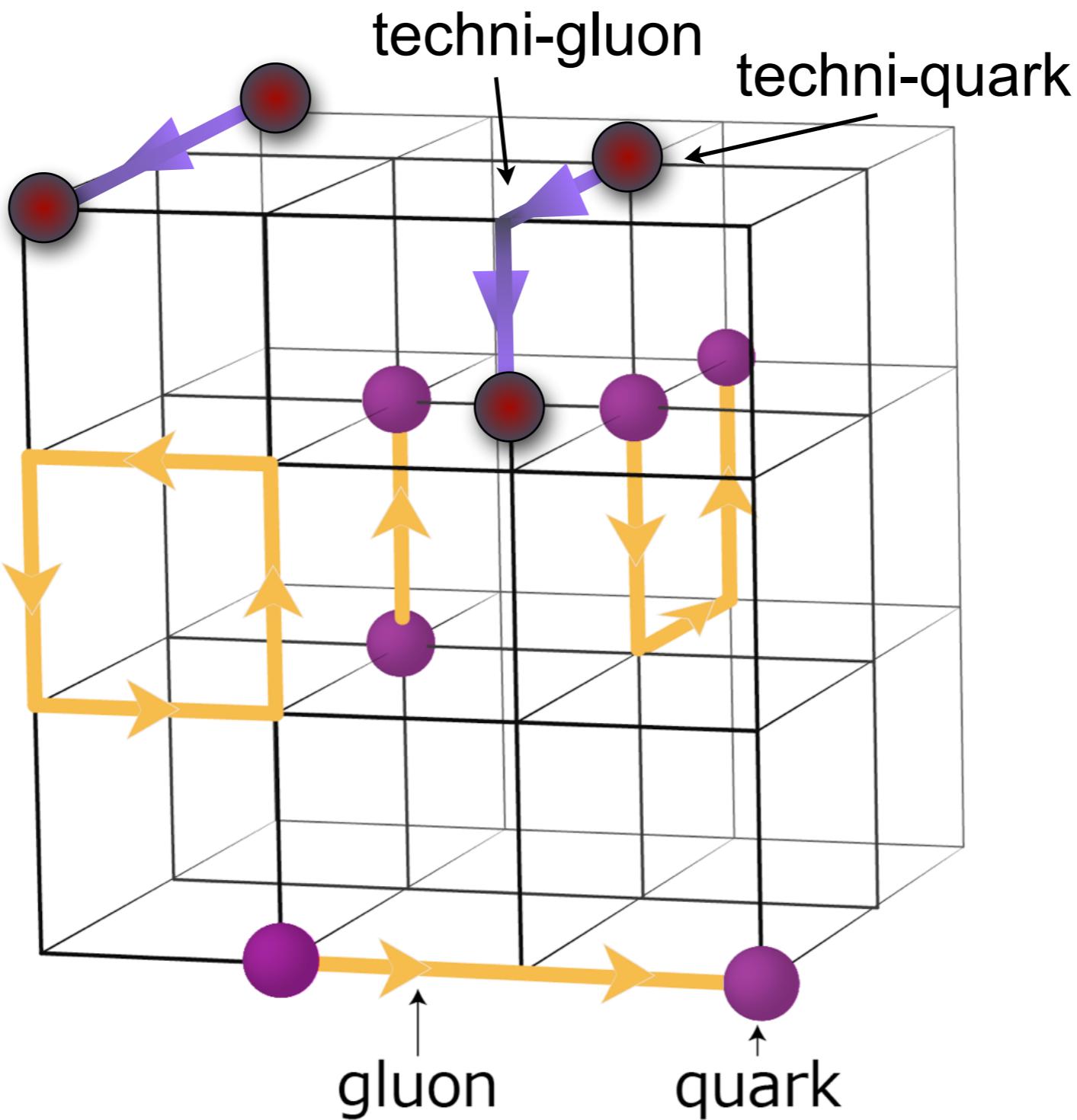
From 2 to infinite N & exact results



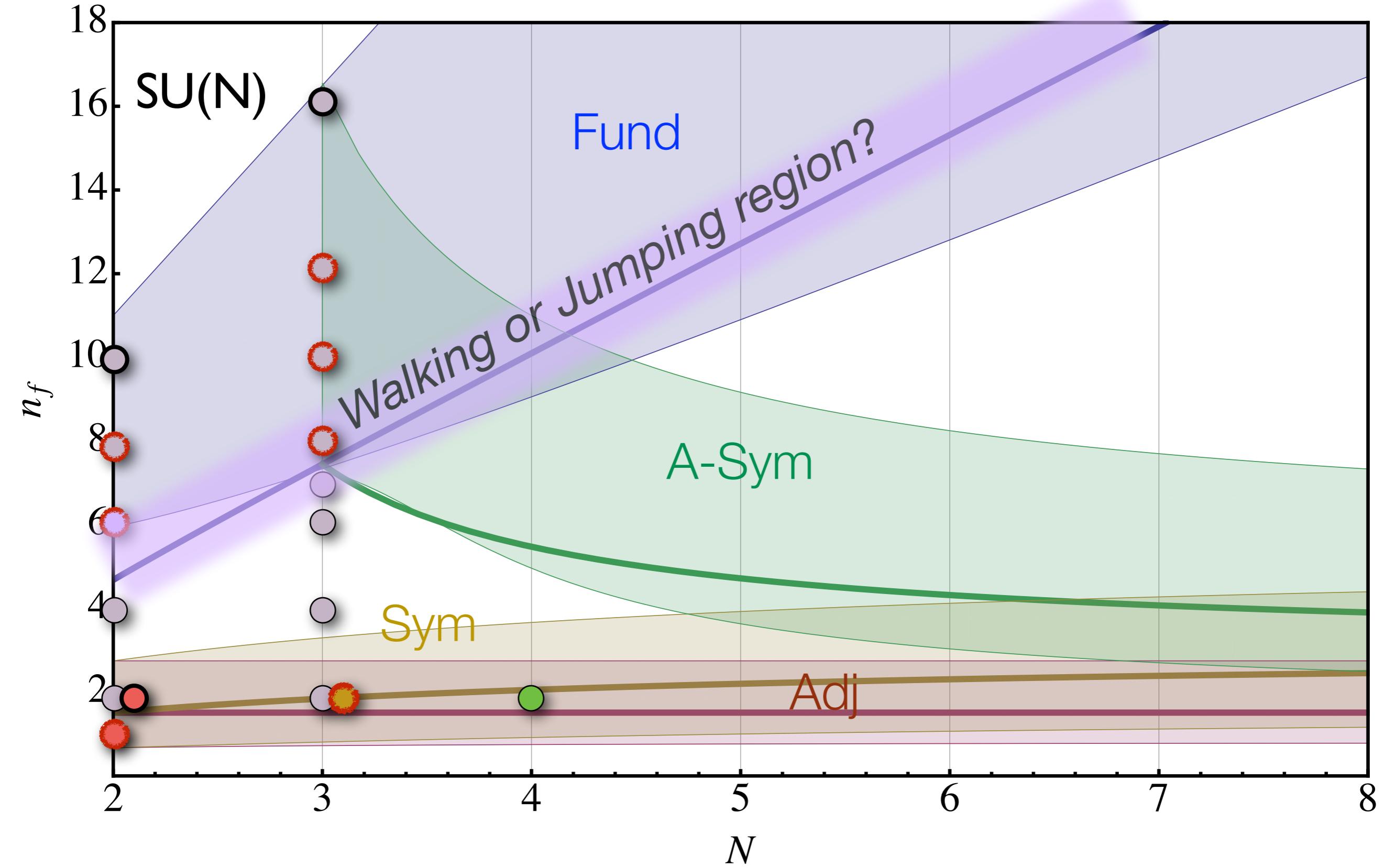
From 2 to infinite N & exact results



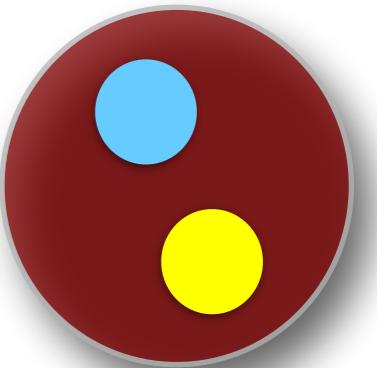
Composite Dynamics on the Lattice



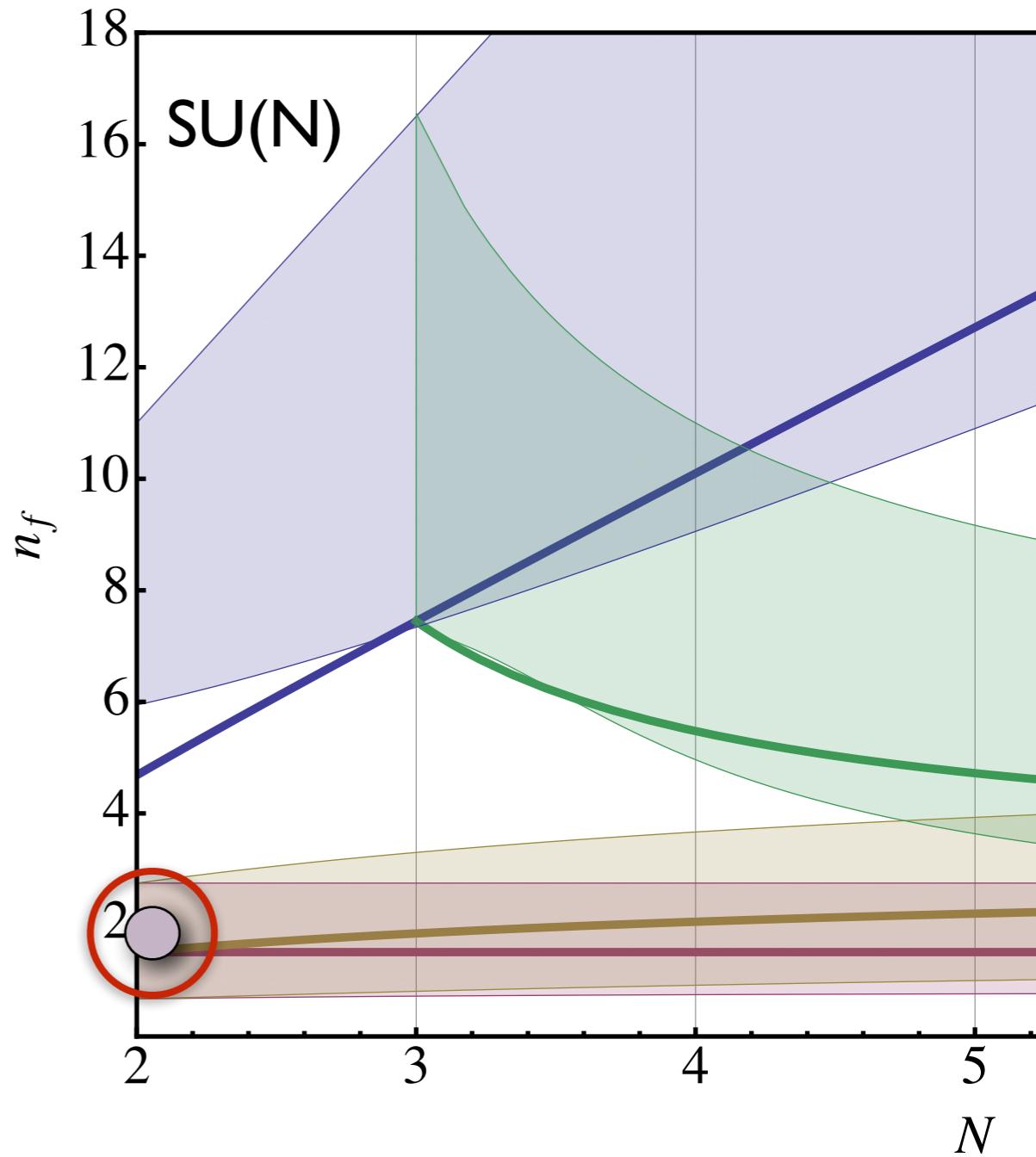
Lattice SU(N) Phase Diagram



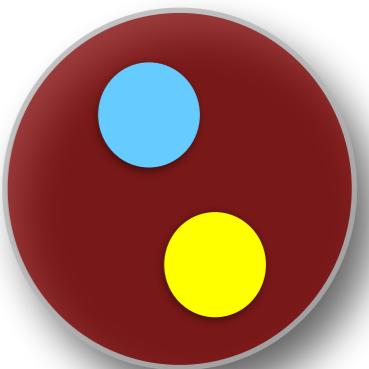
$SU(2) = Sp(2)$ with 2 Flavors



Minimal Fund. Gauge Theory



$SU(2) = Sp(2)$ with 2 Flavors



Minimal Fund. Gauge Theory

- ◆ Unified TC & Comp. Goldstone Higgs
- ◆ TC Meson DM
- ◆ Stealth DM
- ◆ Dark Nuclei
- ◆ SIMPlest Miracle
- ◆ ...

