

# Relaxing the Higgs with cosmology

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# The Higgs: Now what? What's Next?

*"The experiment worked better than expected and the analysis uncovered a very difficult to find signal"*

the words of a string theorist



**Great success...**

...but the experimentalists haven't found what the BSM theorists told them they will find in addition to the Higgs boson:  
no susy, no BH, no extra dimensions, nothing ...



Have the theorists been lying for so many years?  
Have the EXP's been too naive to believe the TH's?



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**HEP future**

exploration/discovery era or consolidation/measurement era?

# Naturalness & TeV scale new physics

Following the arguments of Wilson, 't Hooft (and others):

only small numbers associated to the breaking of a symmetry survive quantum corrections  
( others are not necessarily theoretically inconsistent  
but they require some conspiracy at different scales )

Field	Symmetry as $m \rightarrow 0$	Implication
Spin-1/2 $m\Psi\bar{\Psi}$	$\Psi \rightarrow e^{i\theta}\Psi$ $\bar{\Psi} \rightarrow e^{-i\theta}\bar{\Psi}$ (chiral symmetry)	$\delta m \propto m$ <b>Natural!</b>
Spin-1 $m^2 A_\mu A^\mu$	$A_\mu \rightarrow A_\mu + \partial_\mu \alpha$ (gauge invariance)	$\delta m \propto m$ <b>Natural!</b>

courtesy to N. Craig @ Blois '15

The Higgs mass in the SM doesn't break any (quantum\*) symmetry

\* it does break classical scale invariance, as the running of the gauge couplings does too!



# Naturalness principle @ work

Following the arguments of Wilson, 't Hooft (and others):

only small numbers associated to the breaking of a symmetry survive quantum corrections

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Beautiful examples of naturalness to understand the need of “new” physics

see for instance Giudice '13 (and refs. therein) for an account

- ▶ the need of the positron to screen the electron self-energy:  $\Lambda < m_e/\alpha_{\text{em}}$
- ▶ the rho meson to cutoff the EM contribution to the charged pion mass:  $\Lambda^2 < \delta m_\pi^2/\alpha_{\text{em}}$
- ▶ the kaon mass difference regulated by the charm quark:  $\Lambda^2 < \frac{\delta m_K}{m_K} \frac{6\pi^2}{G_F^2 f_K^2 \sin^2 \theta_C}$
- ▶ the light Higgs boson to screen the EW corrections to gauge bosons self-energies
- ▶ ...
- ▶ new physics at the weak scale to cancel the UV sensitivity of the Higgs mass?

Apparent fine-tunings have always pointed to new degrees of freedom

# The Darwinian solution to the Hierarchy

Other origin of small/large numbers according to Weyl and Dirac:  
hierarchies are induced/created by time evolution/the age of the Universe

Can this idea be formulated in a QFT language?

In which sense is it addressing the stability of small numbers at the quantum level?

Graham, Kaplan, Rajendran '15

- ▶ Higgs mass-squared promoted to a field
- ▶ The field evolves in time in the early universe and scans a vast range of Higgs mass
- ▶ The Higgs mass-squared relaxes to a small negative value
- ▶ The electroweak symmetry breaking stops the time-evolution of the dynamical system

## Self-organized criticality

when the Higgs mass becomes negative, it back-reacts and generates a potential barrier that stops the evolution of the scanning field



# Higgs-axion cosmological relaxation

Graham, Kaplan, Rajendran '15

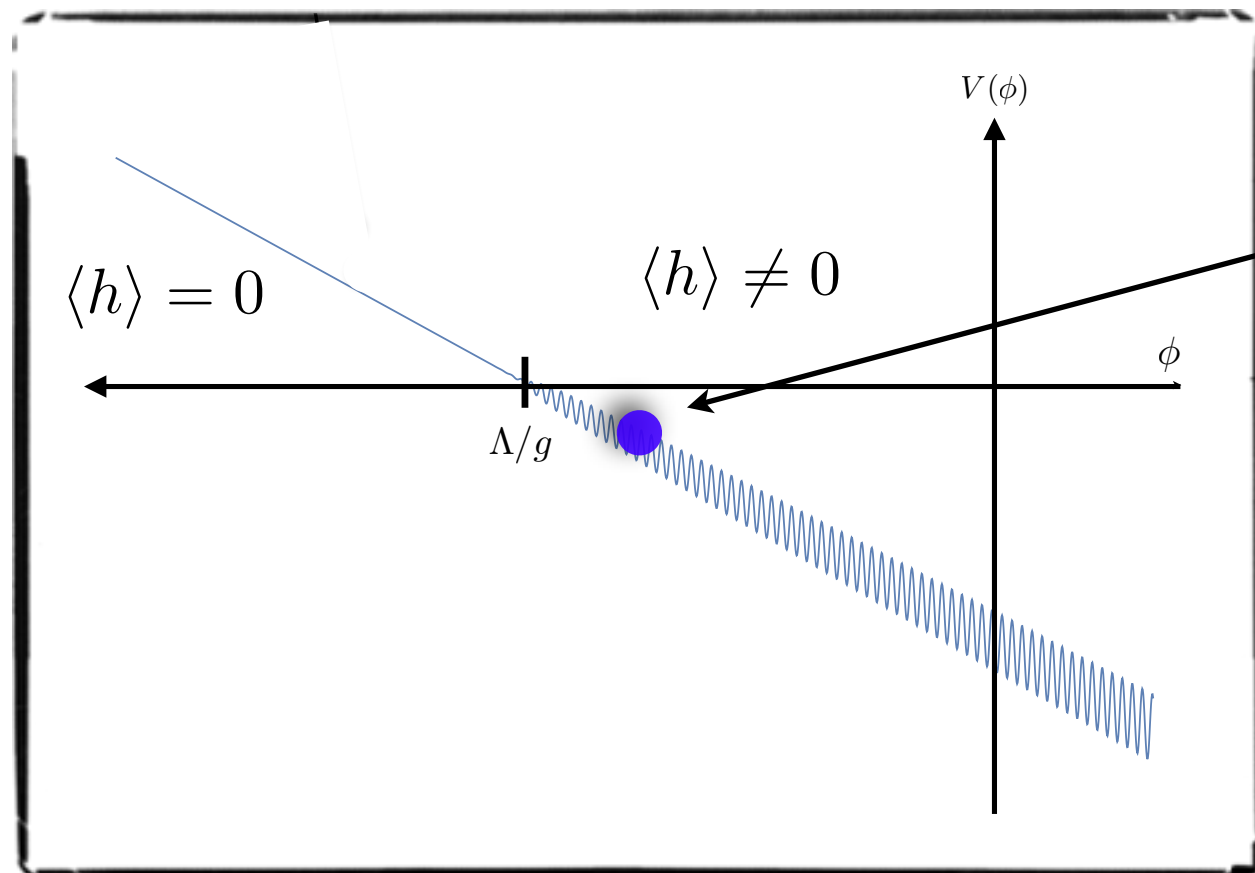
$\phi$  slowly rolling field (inflation provides friction) that scans the Higgs mass

$$\Lambda^2 \left( -1 + f \left( \frac{g\phi}{\Lambda} \right) \right) |H|^2 + \Lambda^4 V \left( \frac{g\phi}{\Lambda} \right) + \frac{1}{32\pi^2} \frac{\phi}{f} \tilde{G}^{\mu\nu} G_{\mu\nu}$$

Higgs mass  
depends on  $\phi$

potential needed to force  
 $\phi$  to roll-down in time  
(during inflation)

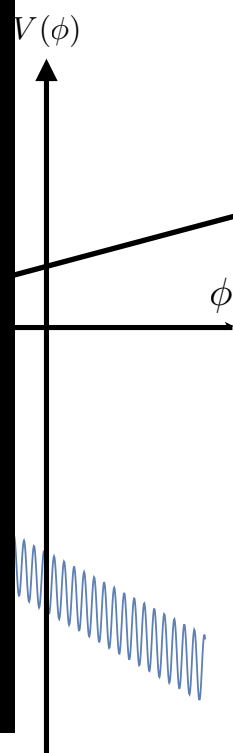
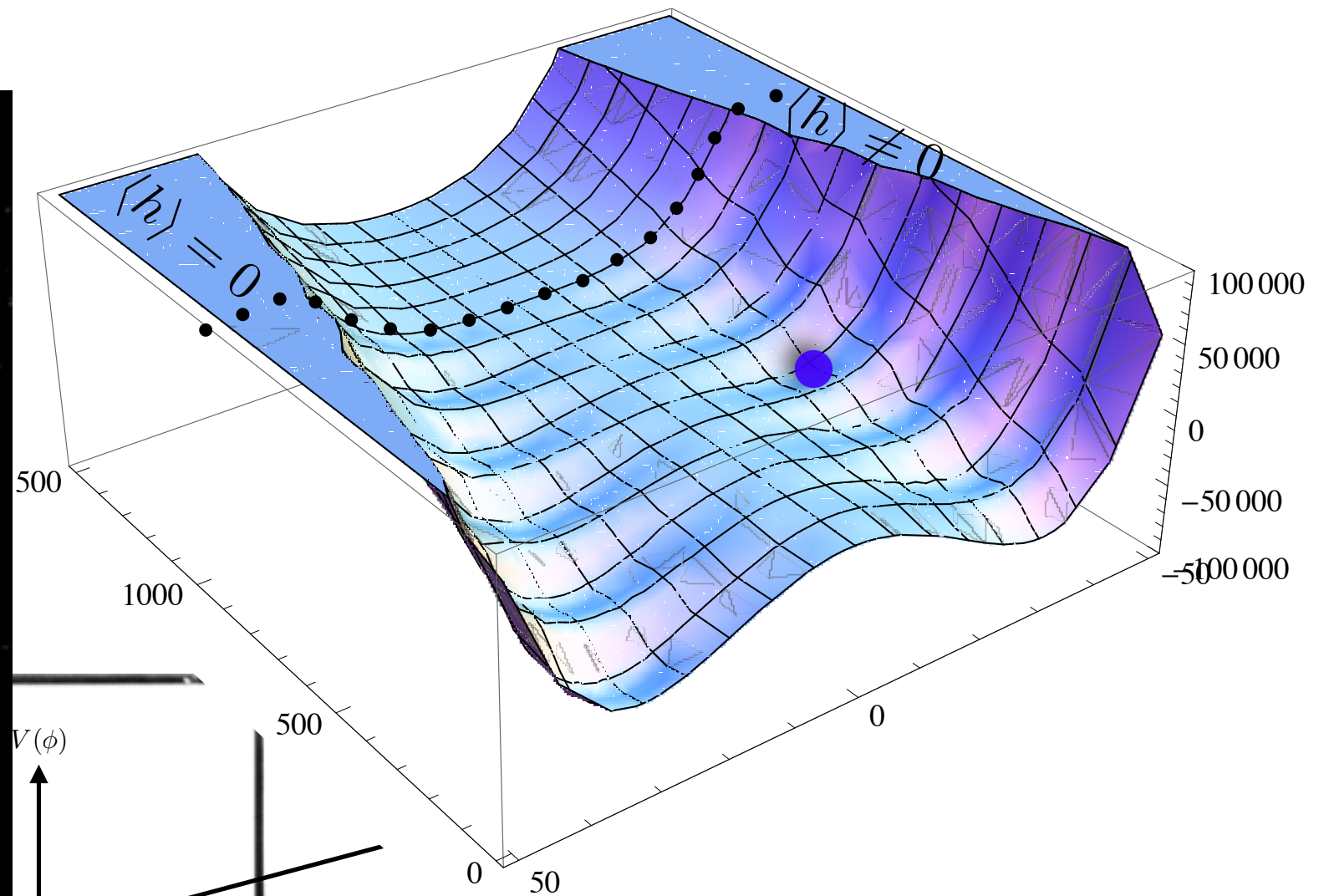
axion-like coupling  
that will seed the potential barrier  
stopping the rolling when the Higgs  
develops its vev  
 $\Lambda_{\text{QCD}}^3 h \cos \frac{\phi}{f}$



# Higgs-axion cosmological relaxation

Graham, Kaplan, Rajendran '15

$\phi$



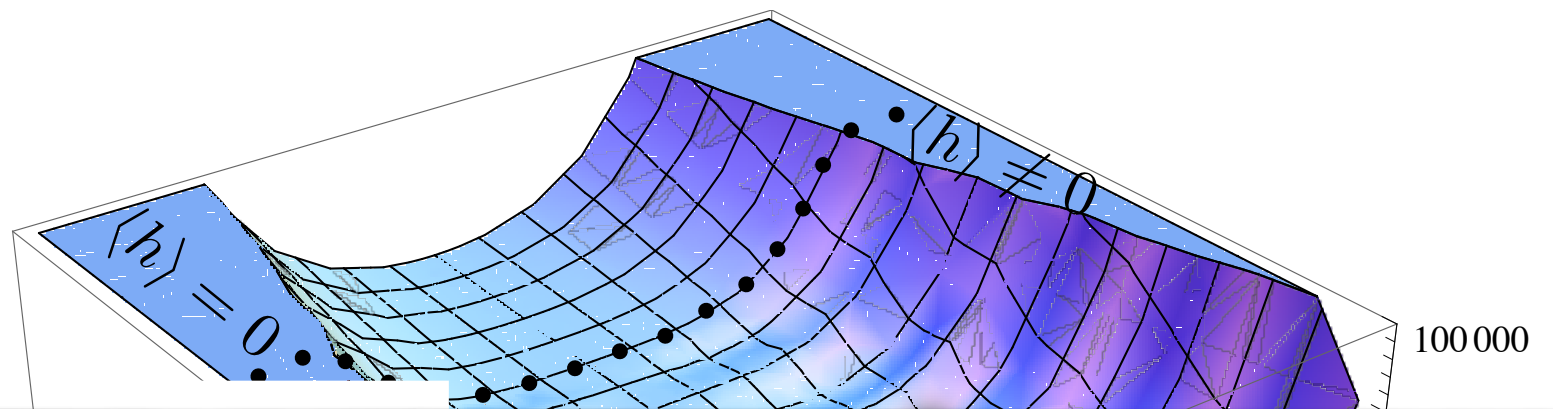
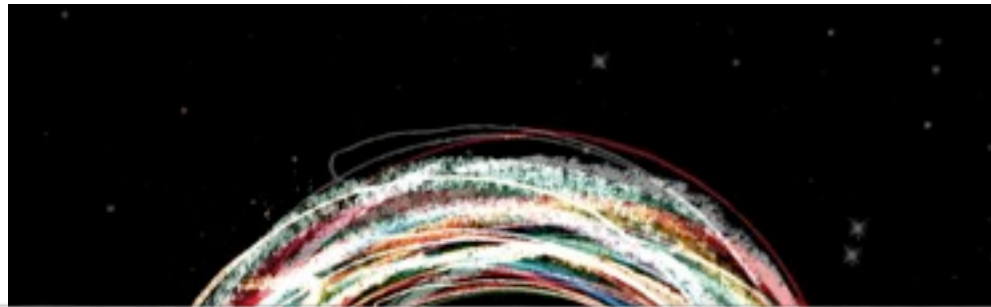
If  $\phi$  continues rolling, the Higgs vev increases, the potential barrier increases and ultimately prevents  $\phi$  from rolling down further



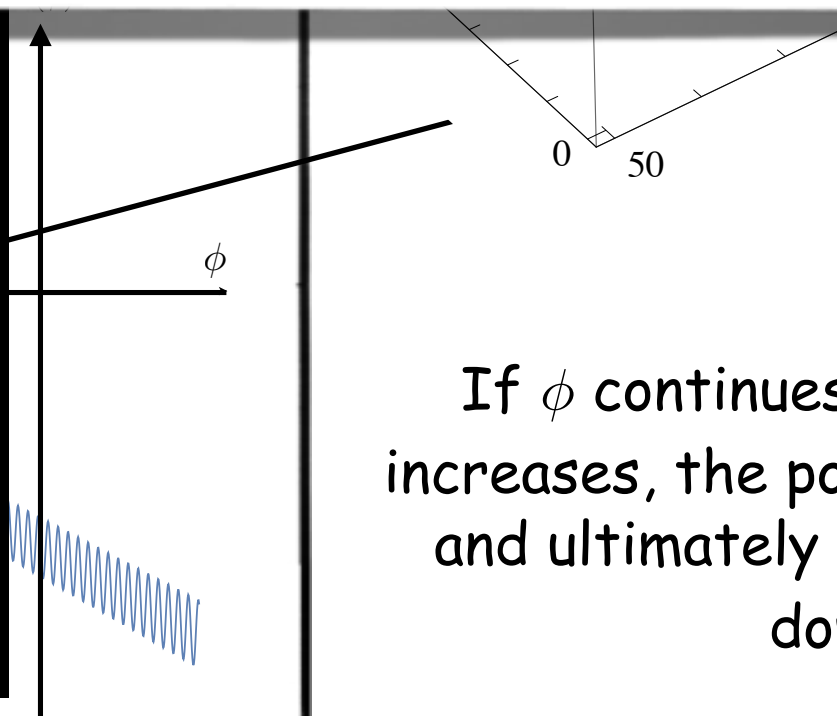
# Higgs-axion cosmological relaxation

Graham, Kaplan, Rajendran '15

$\phi$



Hierarchy problem solved  
by light weakly coupled new physics  
and not by TeV scale physics



If  $\phi$  continues rolling, the Higgs vev increases, the potential barrier increases and ultimately prevents  $\phi$  from rolling down further

# Higgs-axion cosmological relaxation

- Higgs vev stops cosmological rolling

$$\Lambda_{\text{QCD}}^3 \frac{v}{f} \sim \frac{\partial}{\partial \phi} (\Lambda^4 V(g\phi/\Lambda)) \simeq g\Lambda^3$$

**note:**  $v \ll \Lambda$  provided that  $g \ll 1$ . It doesn't explain why the coupling is small (that question can be postponed to higher energies) but it ensures that the solution is stable under quantum correction.

- Slow rolling:  $H_I > \frac{\Lambda^2}{M_P}$  ensures that the energy density stored in  $\phi$  does not affect inflation

- Classical rolling:  $H_I^3 < g\Lambda^3$

classical displacement  
over one Hubble time

$$\frac{1}{H_I} \frac{d\phi}{dt} = \frac{1}{H_I^2} \frac{dV}{d\phi} = \frac{g\Lambda^3}{H_I^2}$$

quantum fluctuation

>

$$H_I$$

$$\frac{\Lambda^6}{M_P^3} < g\Lambda^3 = \Lambda_{\text{QCD}}^3 \frac{v}{f}$$

i.e.

$$\Lambda < 10^7 \text{ GeV} \left( \frac{10^9 \text{ GeV}}{f} \right)^{1/6}$$

Important issue:  $\theta_{\text{QCD}} \sim 1 \gg 10^{-10}$ . Can be solved but  $\Lambda < 30 \text{ TeV}$



# Classifying relaxing Lagrangians...

$$V(\phi, h) = \Lambda^3 g \phi - \frac{1}{2} \Lambda^2 \left( 1 - \frac{g\phi}{\Lambda} \right) h^2 + \epsilon \Lambda_c^4 \left( \frac{h}{\Lambda_c} \right)^n \cos(\phi/f) + \dots$$

► n=1: need another source of EWSB

► QCD condensate  $\langle qq \rangle \sim \Lambda_{\text{QCD}}$

► new strongly-coupled sector à la Technicolor

⊢ new physics @ TeV, coincidence problem? ⊣

► n=2: no extra source of EWSB needed

► quantum stability? h-loops generate extra interactions that will stop  $\phi$  before the Higgs vev develops unless  $\Lambda_c \ll v$  (coincidence pb and new physics @ TeV again?)

$$\epsilon \Lambda_c^4 \cos(\phi/f) , \quad \epsilon \Lambda_c^3 g \phi \cos(\phi/f)$$

► our solution: make the envelop of the oscillatory potential a field

# Cosmological Higgs-Axion Interplay (CHAIN)

Espinosa, Grojean, Panico, Pomarol, Pujolas, Servant '15

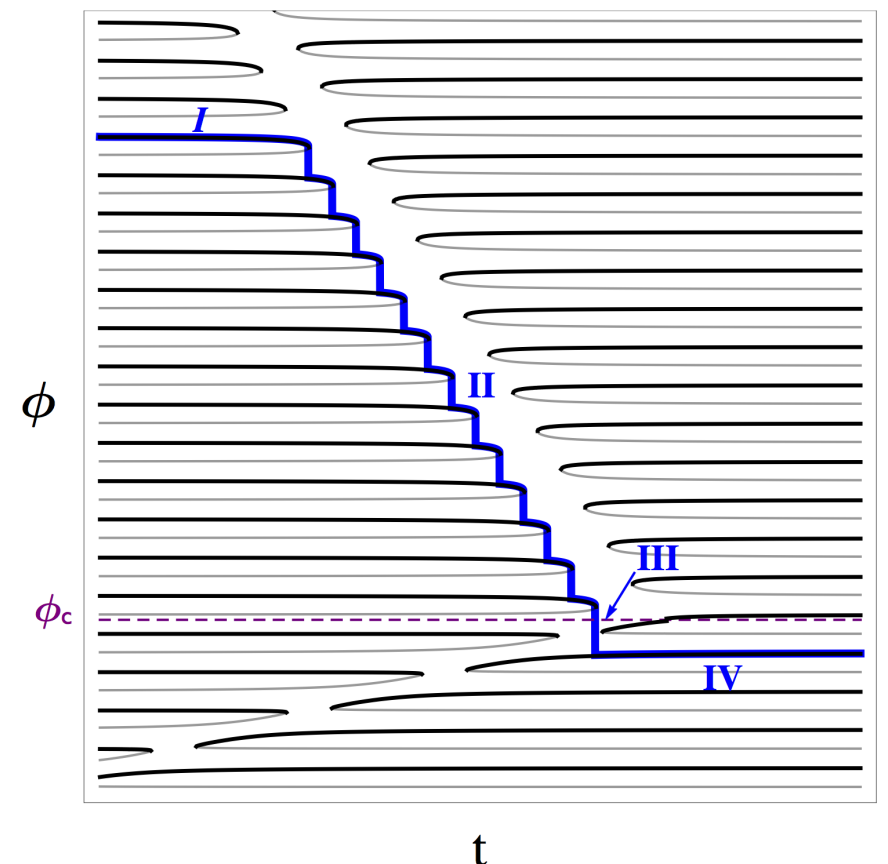
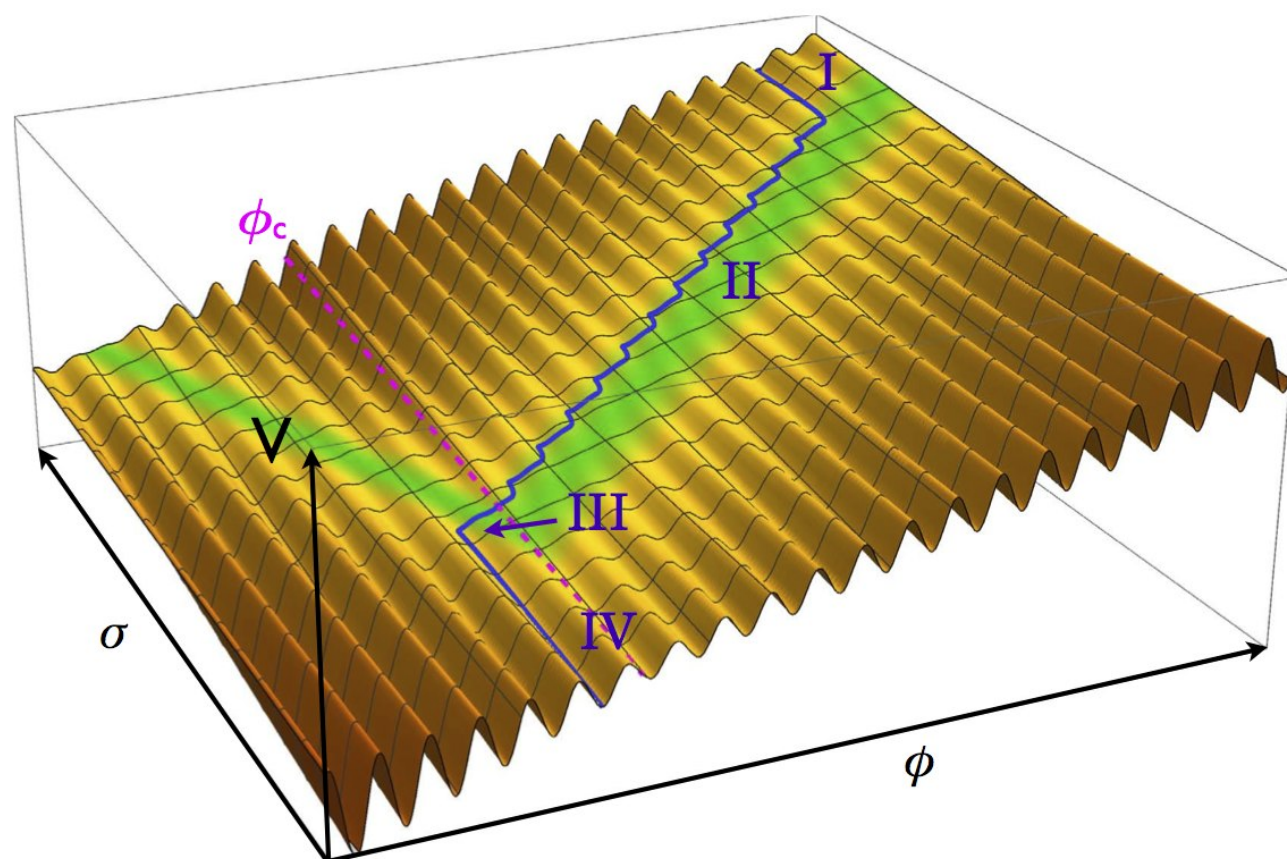
$$V(\phi, \sigma, H) = \Lambda^4 \left( \frac{g\phi}{\Lambda} + \frac{g_\sigma \sigma}{\Lambda} \right) - \Lambda^2 \left( \alpha - \frac{g\phi}{\Lambda} \right) |H|^2 + \frac{1}{2} \lambda |H|^4 + A(\phi, \sigma, H) \cos(\phi/f)$$

$$A(\phi, \sigma, H) \equiv \epsilon \Lambda^4 \left( \beta + c_\phi \frac{g\phi}{\Lambda} - c_\sigma \frac{g_\sigma \sigma}{\Lambda} + \frac{|H|^2}{\Lambda^2} \right)$$

quantum generated  
new terms from  
the  $|H|^2 \cos(\phi/f)$  term

the new interaction  
that saves our day

original relaxion-type  
term

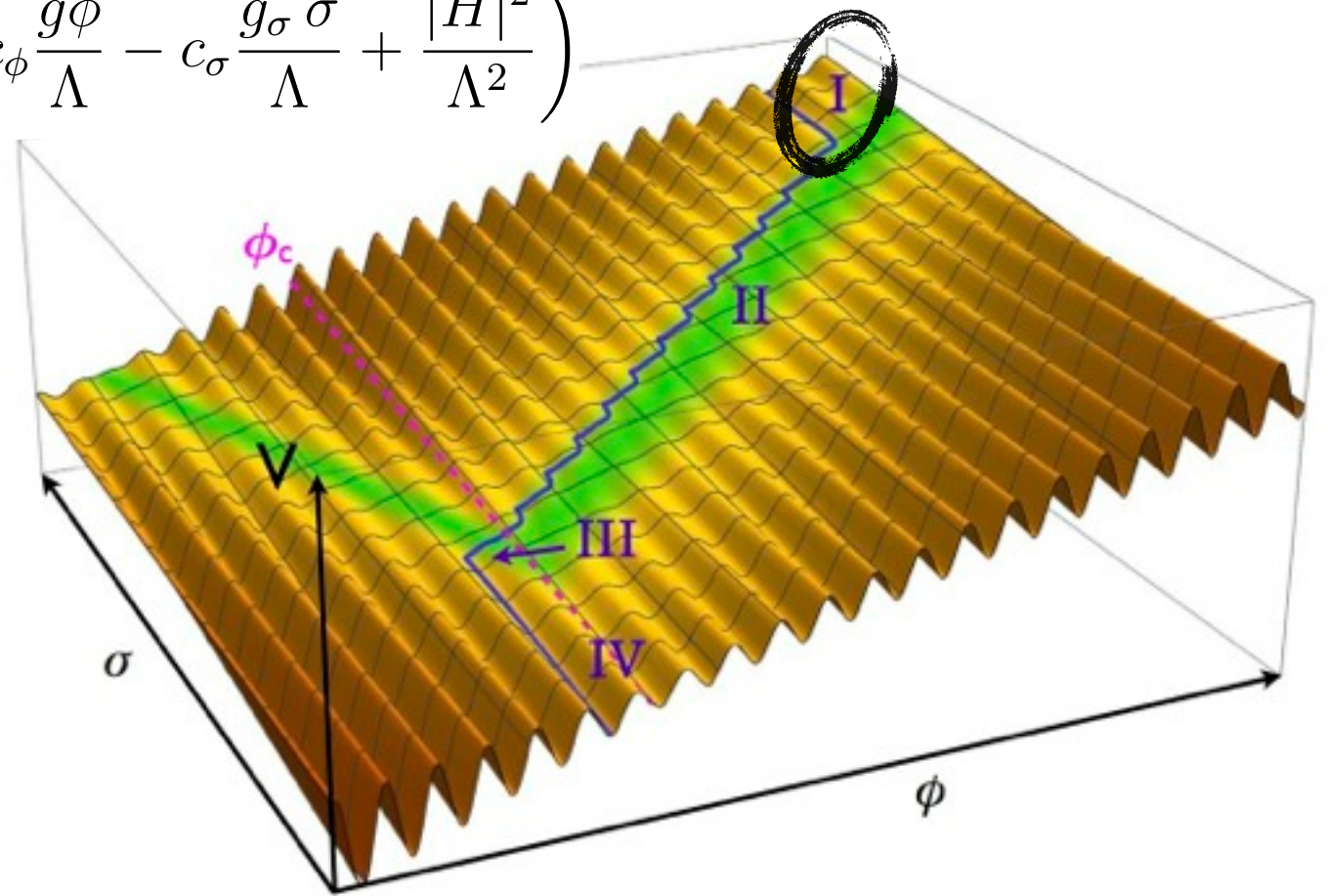


# Cosmological Higgs-Axion Interplay dynamics

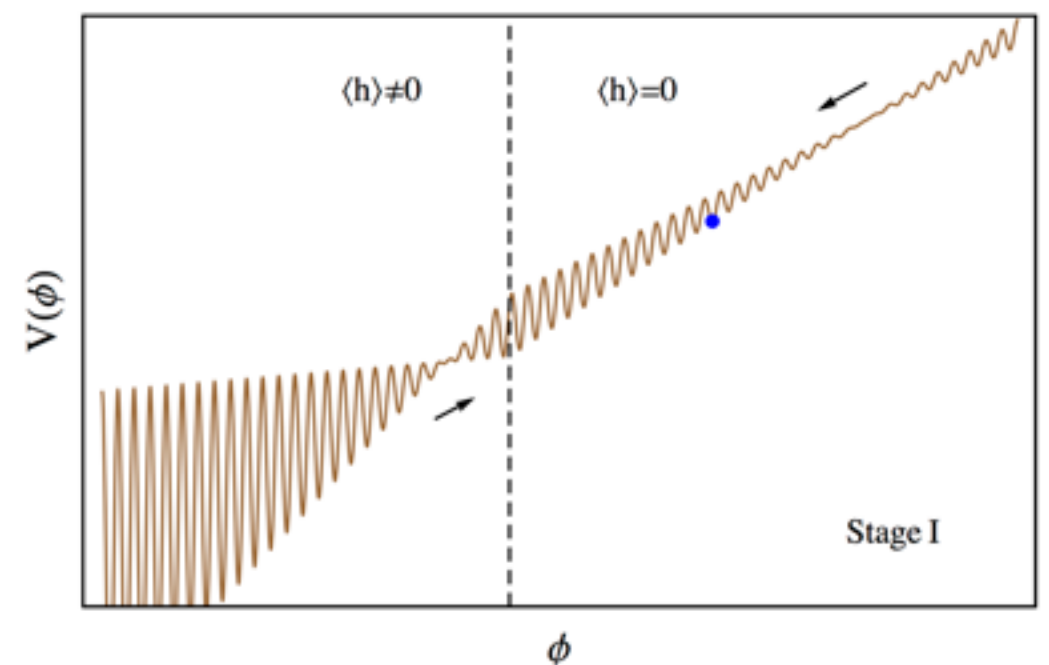
## ► Phase I:

potential barrier  $\sim O(\epsilon\Lambda^4)$

$$A(\phi, \sigma, H) \equiv \epsilon\Lambda^4 \left( \beta + c_\phi \frac{g\phi}{\Lambda} - c_\sigma \frac{g_\sigma \sigma}{\Lambda} + \frac{|H|^2}{\Lambda^2} \right)$$



- $\sigma$  starts with large value
- high potential barrier
- $\phi$  is stuck in some deep minimum
- $\sigma$  slowly rolls down  $\sigma(t) = \sigma_0 - g_\sigma \Lambda^3 t / (3H_I)$
- inflation prevents it from accelerating





# Cosmological Higgs-Axion Interplay dynamics

## ► Phase II:

potential barrier "<" linear potential  
(steepness of the oscillatory potential  
is smaller than  
steepness of the linear potential)

$$A(\phi, \sigma, h)/f < g\Lambda^3$$

$$\phi \in \left( \phi_c + \frac{c_\sigma g_\sigma}{c_\phi g} (\sigma - \sigma_c) \pm \frac{f}{c_\phi \epsilon} \right)$$

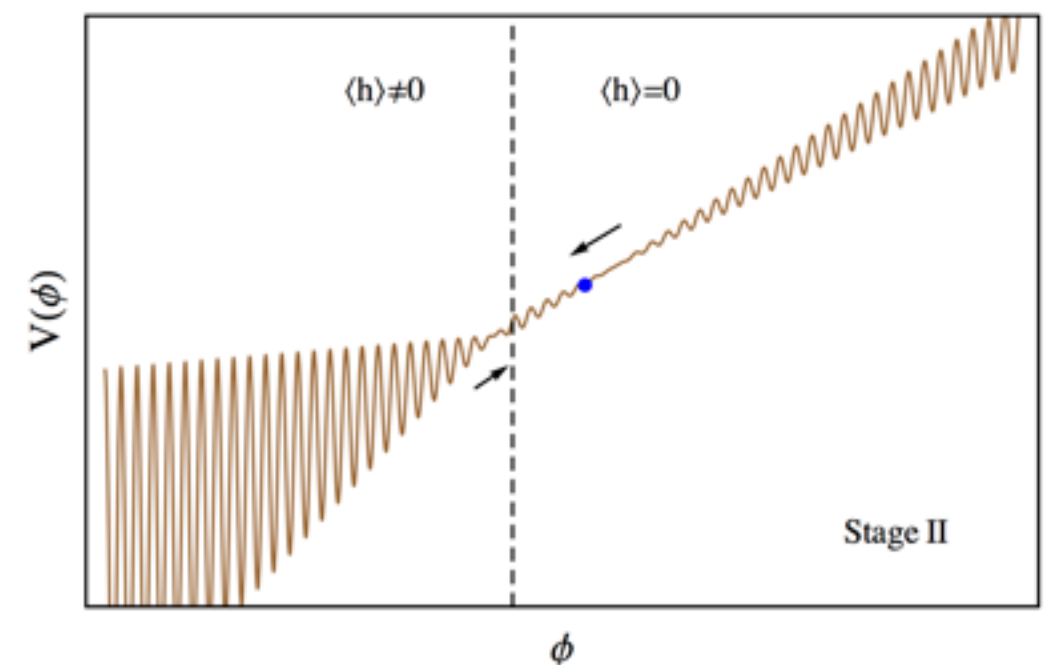
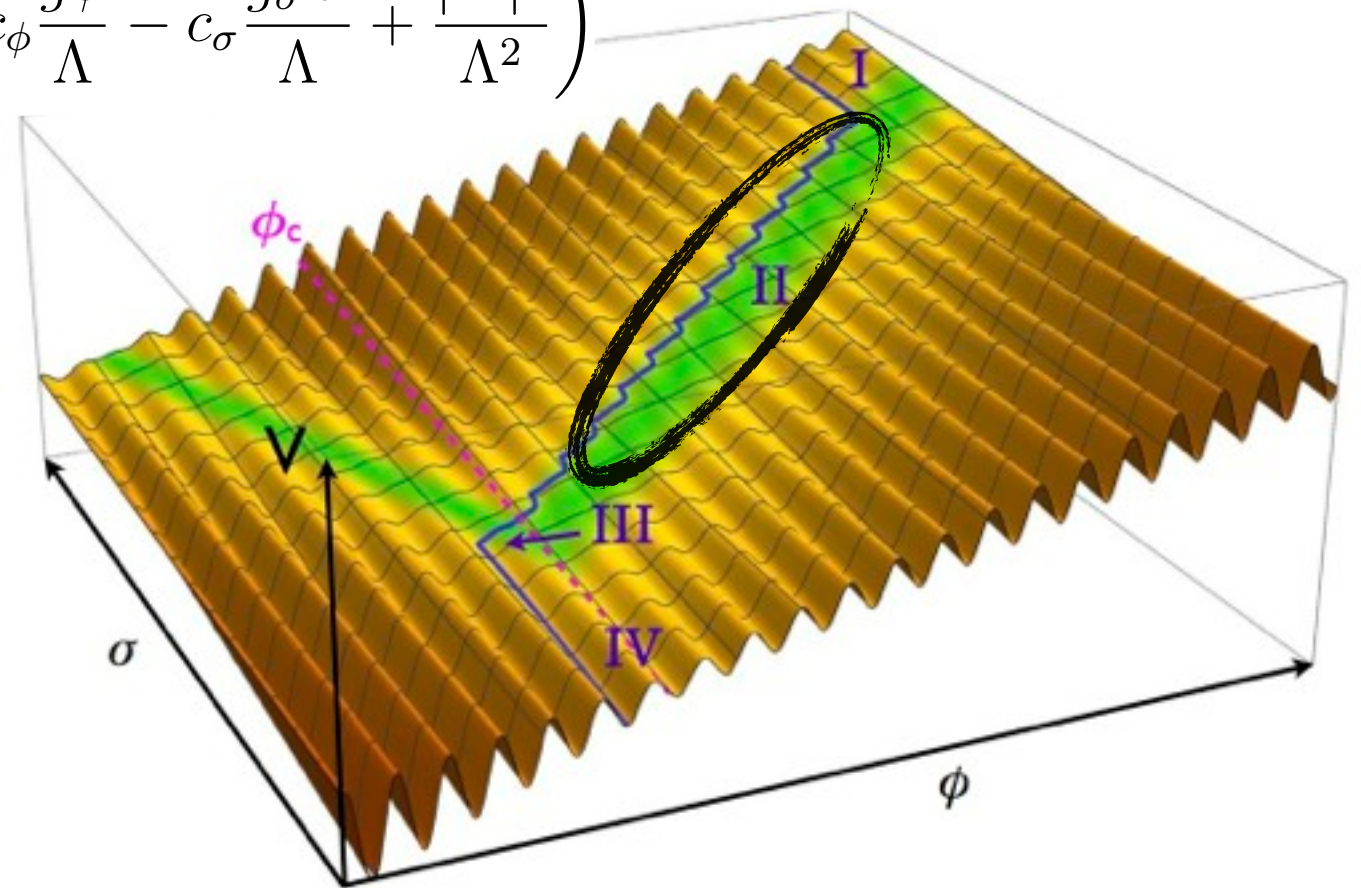
- $\sigma$  reaches a value where the barrier vanishes
- $\phi$  can start rolling down  $\phi(t) = \phi_0 - g\Lambda^3 t / (3H_I)$
- $\sigma$  and  $\phi$  roll down along the barrier-free valley

$\phi$  stays in the barrier-free region iff  
the gradient of the dynamical trajectory

$$(d\phi(t)/dt)/(d\sigma(t)/dt) = g/g_\sigma$$

is larger than the gradient  $d\phi/d\sigma$  of the valley

$$c_\phi g^2 > c_\sigma g_\sigma^2$$



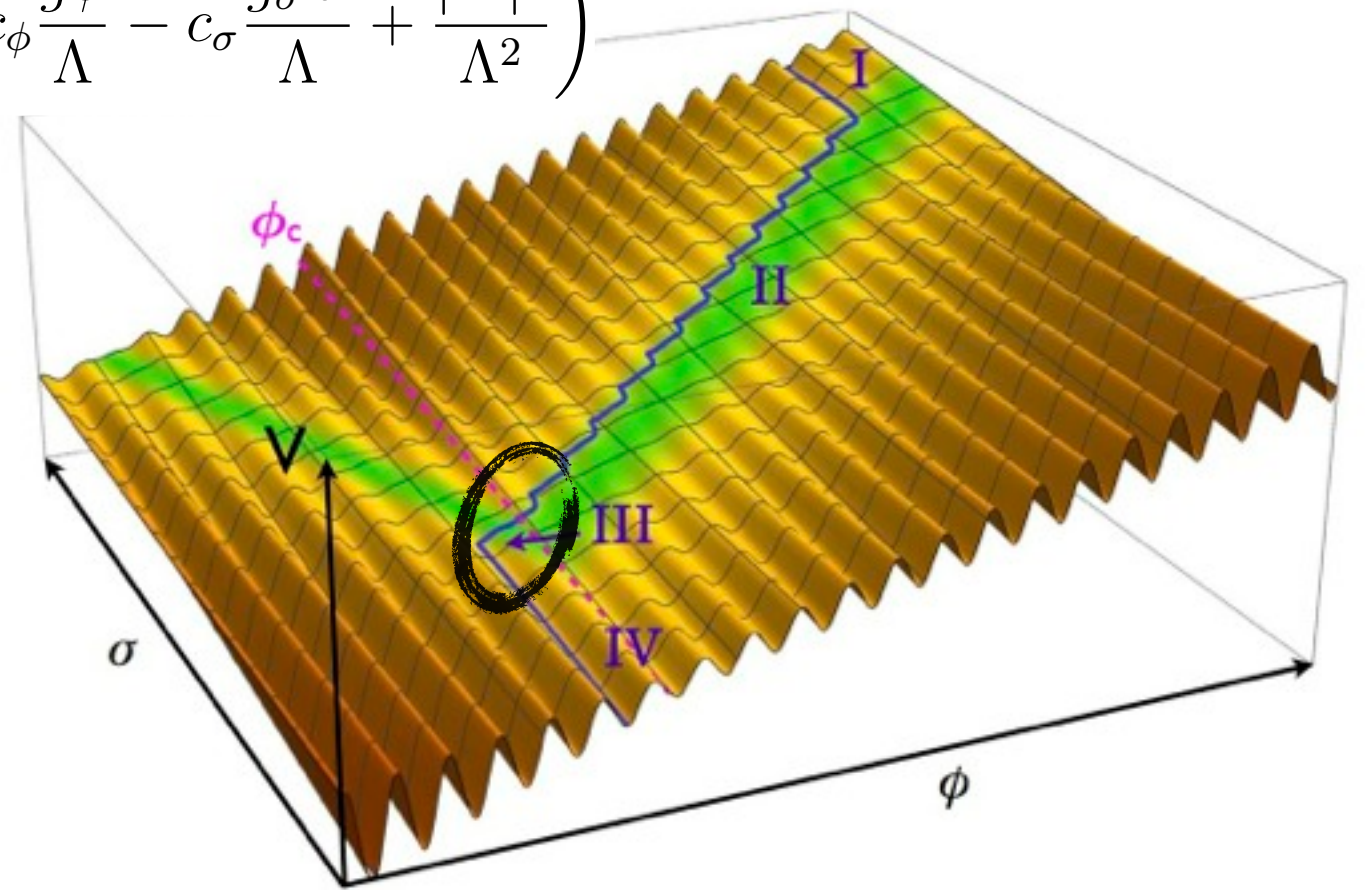
# Cosmological Higgs-Axion Interplay dynamics

## ► Phase III:

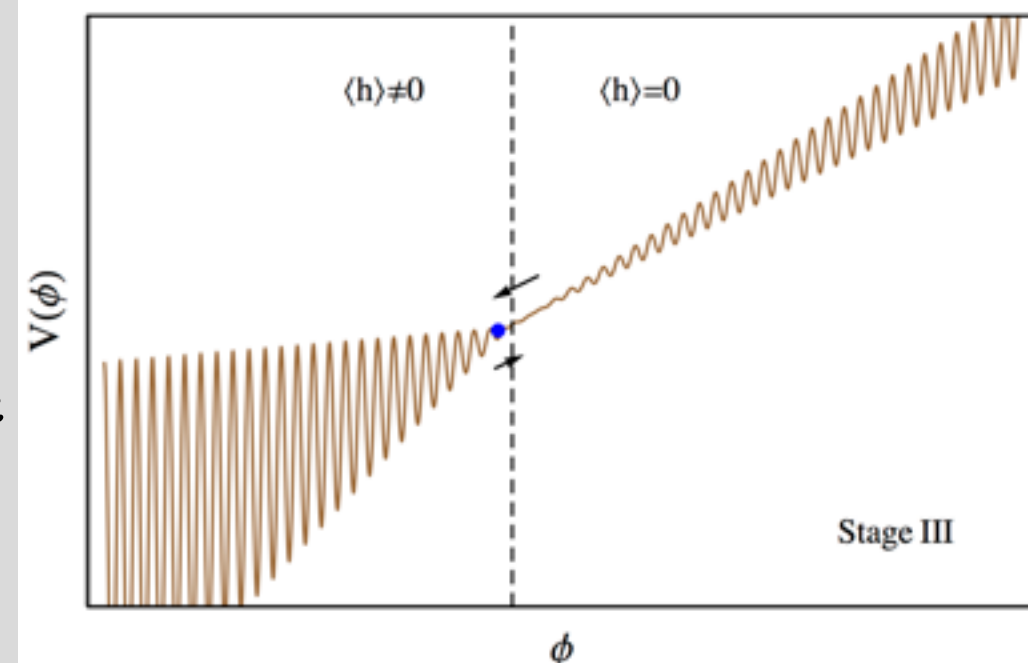
potential barrier "<" linear potential  
(steepness of the oscillatory potential  
is smaller than  
steepness of the linear potential)

$$A(\phi, \sigma, h)/f < g\Lambda^3$$

$$\phi \in \left( \phi_c + \frac{c_\sigma g_\sigma}{c'_\phi g} (\sigma - \sigma_c) \pm \frac{f}{c'_\phi \epsilon} \right)$$



- $\phi$  reaches the critical value
- the Higgs mass becomes tachyonic\*
- a  $\phi$ -dependent Higgs vev turns on
- new contribution to the barrier  $c'_\phi = c_\phi - 1/(2\lambda)$
- change in the slope of the tracking valley in the  $\sigma$ - $\phi$  plane
- $\phi$  crosses/exits the barrier-free valley:  $c'_\phi g^2 < c_\sigma g_\sigma^2$
- the size of barrier start increasing



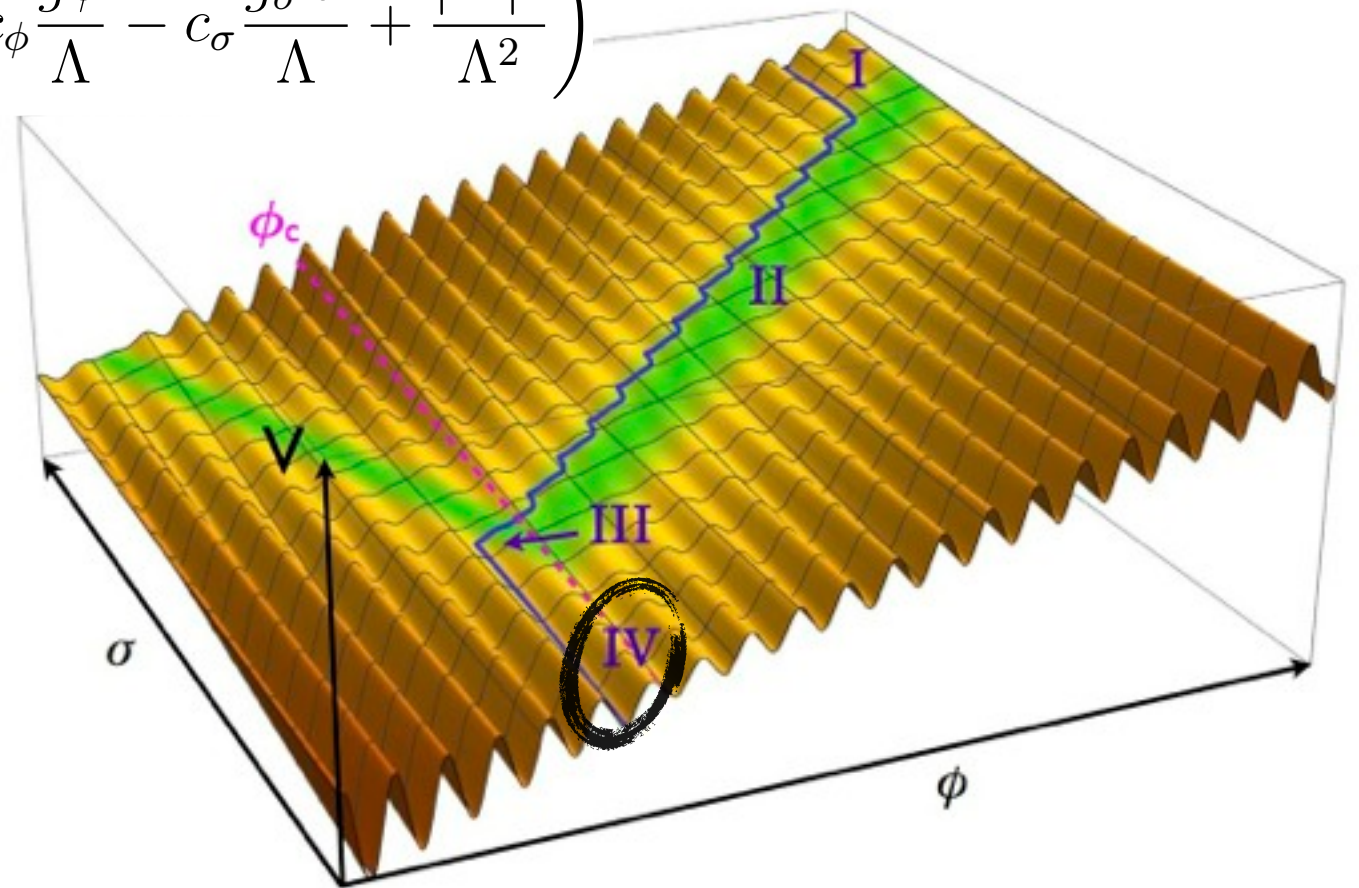


# Cosmological $H_{\text{iggs}}\text{-}A_{\text{xion}}$ $I\!N_{\text{terplay}}$ dynamics

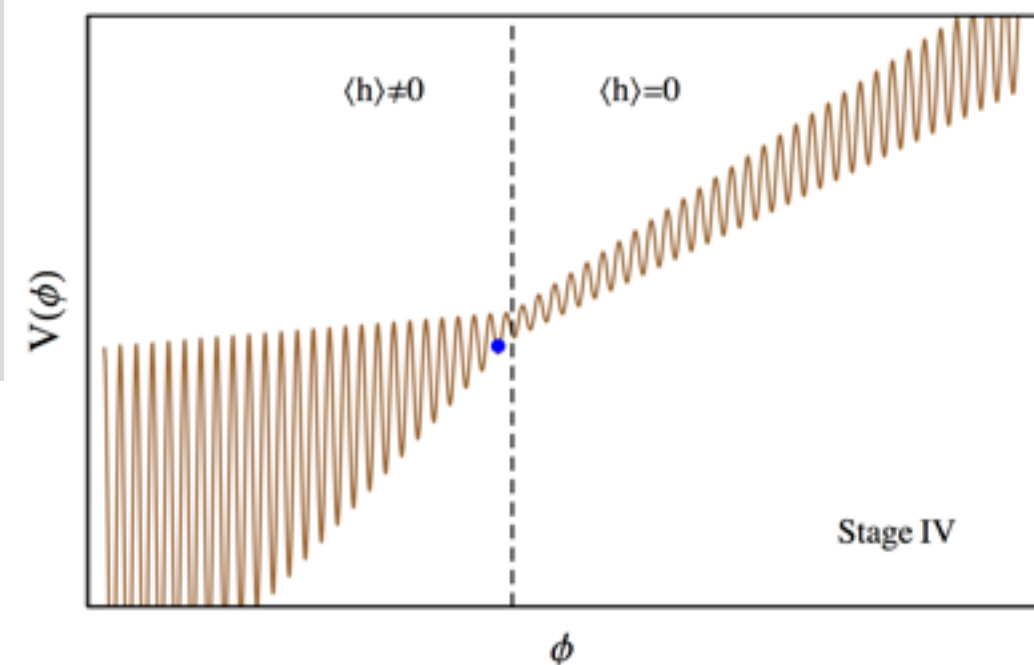
## ► Phase IV:

potential barrier  $\sim O(\epsilon\Lambda^4)$

$$A(\phi, \sigma, H) \equiv \epsilon\Lambda^4 \left( \beta + c_\phi \frac{g\phi}{\Lambda} - c_\sigma \frac{g_\sigma \sigma}{\Lambda} + \frac{|H|^2}{\Lambda^2} \right)$$



- potential barrier is large
- $\phi$  is stuck in a deep minimum
- $\sigma$  continues rolling down until it reaches its own minimum
- all the fields are frozen



# Same problem, same solution?

EX SCALE AS COSMOLOGICAL ERRATIC

courtesy to JR Espinosa



Okotoks glacial erratic,  
Alberta, Canada

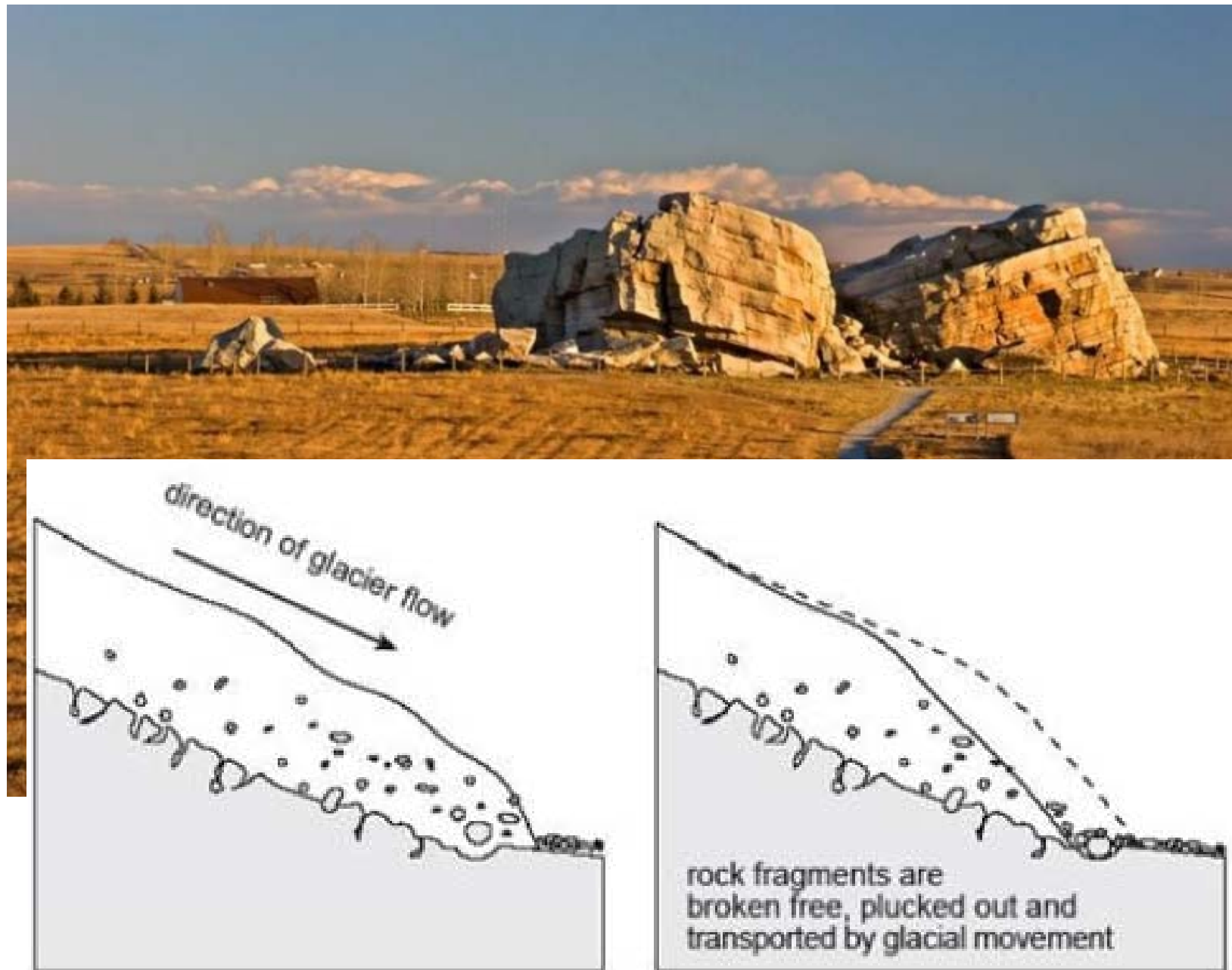
**Unnatural** large rocks differing in composition from the typical surrounding ones



# Same problem, same solution?

EX SCALE AS COSMOLOGICAL ERRATIC

courtesy to JR Espinosa



Unnatural large rocks differing in composition from the typical surrounding ones

Standard geological history:

they were transported by ancient glaciers over hundreds of kilometers

# Consistency conditions

► Quantum stability of the potential  $\epsilon \lesssim v^2/\Lambda^2$

ensures that terms  $\epsilon^2 \Lambda^4 \cos^2(\phi/f)$  don't affect the tracking solution

Ex.   $\epsilon^2 \Lambda^4 \cos^2(\phi/f)$

should be subleading compared to  $\epsilon \Lambda^2 h^2 \cos(\phi/f)$

Requires  $\epsilon \lesssim v^2/\Lambda^2$

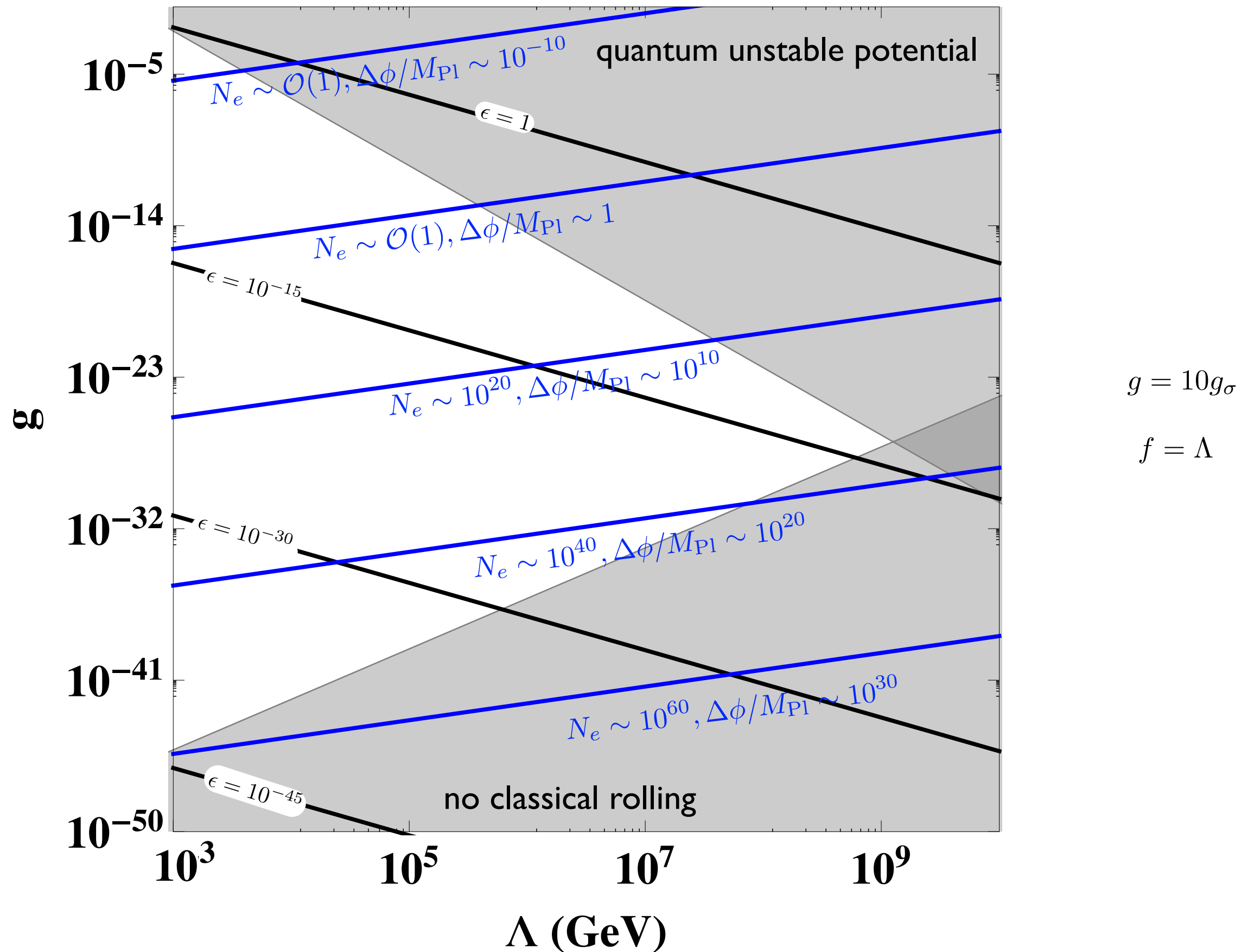
courtesy to JR Espinosa

# Consistency conditions

- ▶ Quantum stability of the potential  $\epsilon \lesssim v^2/\Lambda^2$   
ensures that terms  $\epsilon^2 \Lambda^4 \cos^2(\phi/f)$  don't affect the tracking solution
- ▶ Higgs vev stops cosmological rolling  $\frac{\epsilon \Lambda^2 v^2}{f} \sim \frac{\partial}{\partial \phi} (\Lambda^4 V(g\phi/\Lambda)) \simeq g\Lambda^3$
- ▶ Slow rolling:  $H_I > \frac{\Lambda^2}{M_P}$  ensures that the energy density stored in  $\sigma$  and  $\phi$  does not affect inflation
- ▶ Classical rolling:  $H_I^3 < g\Lambda^3$
- ▶  $\phi$  tracks  $\sigma$  in the barrier-free valley before EWSB:  $c_\phi g^2 > c_\sigma g_\sigma^2$
- ▶  $\phi$  exits the barrier-free valley after EWSB:  $(c_\phi - \frac{1}{2\lambda})g^2 < c_\sigma g_\sigma^2$
- ▶ large field excursions:  $\Delta\phi, \Delta\sigma > \Lambda/g$  to ensure that the Higgs mass scans values  $\Lambda$  from to the weak scale

$$\boxed{\frac{\Lambda^3}{M_{\text{Pl}}^3} \lesssim g_\sigma \lesssim g \lesssim \frac{v^4}{f\Lambda^3}} \quad \Rightarrow \quad \boxed{\Lambda \lesssim (v^4 M_{\text{Pl}}^3)^{1/7} \simeq 2 \times 10^9 \text{ GeV}}$$

# Consistency conditions





# UV completion: where CHAIN can come from?

Strong sector à la QCD with vector-like elementary quarks

+

axion-like field  $\frac{\phi}{f} G'_{\mu\nu} \tilde{G}'^{\mu\nu}$

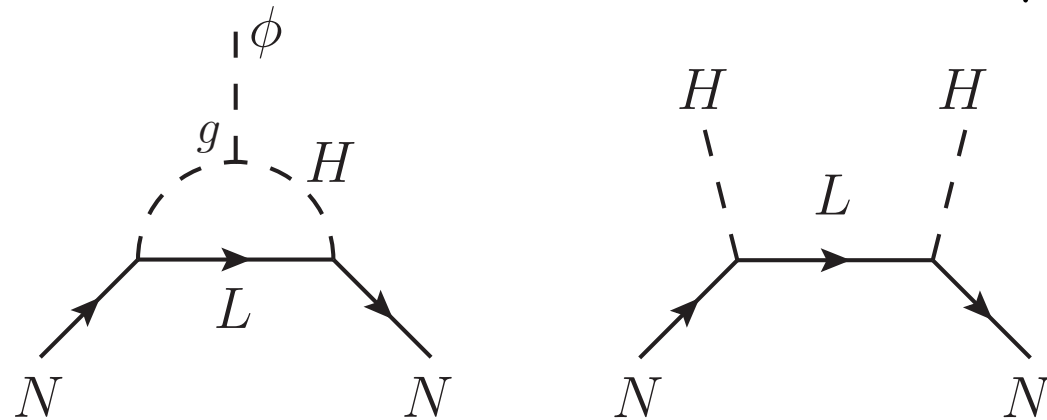
L  $SU(2)_L$  Dirac doublet  
N  $SU(2)_L$  Dirac singlet

$$\mathcal{L}_{\text{mass}} = \underbrace{\Lambda \bar{L} L + \epsilon \Lambda \bar{N} N}_{\text{Dirac masses}} \quad \mathcal{L}_{\text{Yuk}} = \underbrace{\sqrt{\epsilon} \bar{L} H N + h.c.}_{\text{Higgs interactions}}$$

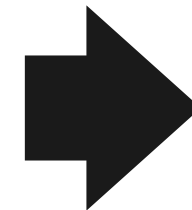
Dirac masses

Higgs interactions

$\epsilon \rightarrow 0$ , additional chiral symmetry (broken by axial anomaly)



$$\mathcal{L}_N = \epsilon g \phi \bar{N} N + \epsilon g_\sigma \sigma \bar{N} N$$



$$m_N \simeq \epsilon \left( \Lambda + g_\sigma \sigma + g \phi - \frac{|H|^2}{\Lambda} \right)$$

$$\langle N \bar{N} \rangle \sim \Lambda^3 \quad \Rightarrow \quad V = \Lambda^3 m_N \cos \frac{\phi}{f}$$

composite baryons and mesons @  $\Lambda$  but no light meson since axial U(1) is anomalous

# Phenomenological signatures

Nothing to be discovered at the LHC/ILC/CLIC/CepC/SppC/FCC!



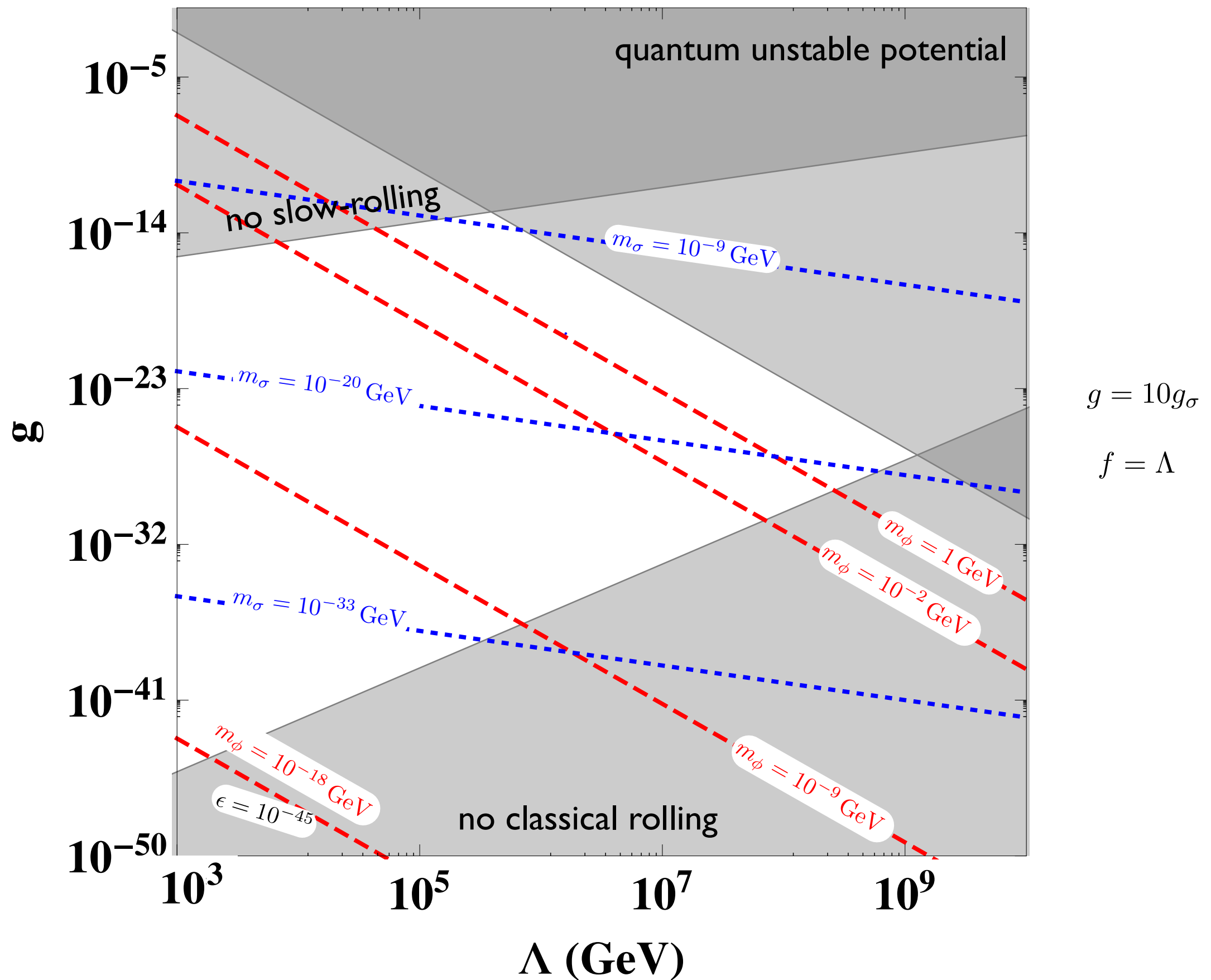
only BSM physics below  $\Lambda$

two (very) light and very weakly coupled axion-like scalar fields

$$m_\phi \sim (10^{-20} - 10^2) \text{ GeV}$$

$$m_\sigma \sim (10^{-45} - 10^{-2}) \text{ GeV}$$

# Phenomenological signatures



# Phenomenological signatures

interesting signatures in cosmology and possibly at SHiP



$\phi$  and  $\sigma$  couple to SM matter via their mixing with the Higgs

$$\theta_{\phi h} \sim \frac{g\Lambda v}{m_h^2}, \quad \theta_{\sigma\phi} \sim \frac{g_\sigma f v^2}{\Lambda^3}, \quad \theta_{\sigma h} \sim \text{Max} \left\{ \theta_{\sigma\phi}\theta_{\phi h}, \frac{g^2}{16\pi^2} \frac{g_\sigma \Lambda^7}{f^2 v^3 m_h^2} \right\}$$

from oscillatory potential

tree-level

quantum mixing  
from  $\phi$ -loop

unsuppressed quartic interaction with the Higgs:  $\phi\phi hh : \epsilon\Lambda^2/f^2$

$\phi$  and  $\sigma$  decay to SM particles

(mostly photons in a large region of parameter space)

$$\Gamma_\phi \sim \theta_{\phi h}^2 \Gamma_h(m_\phi), \quad \Gamma_\sigma \sim \theta_{\sigma h}^2 \Gamma_h(m_\sigma)$$

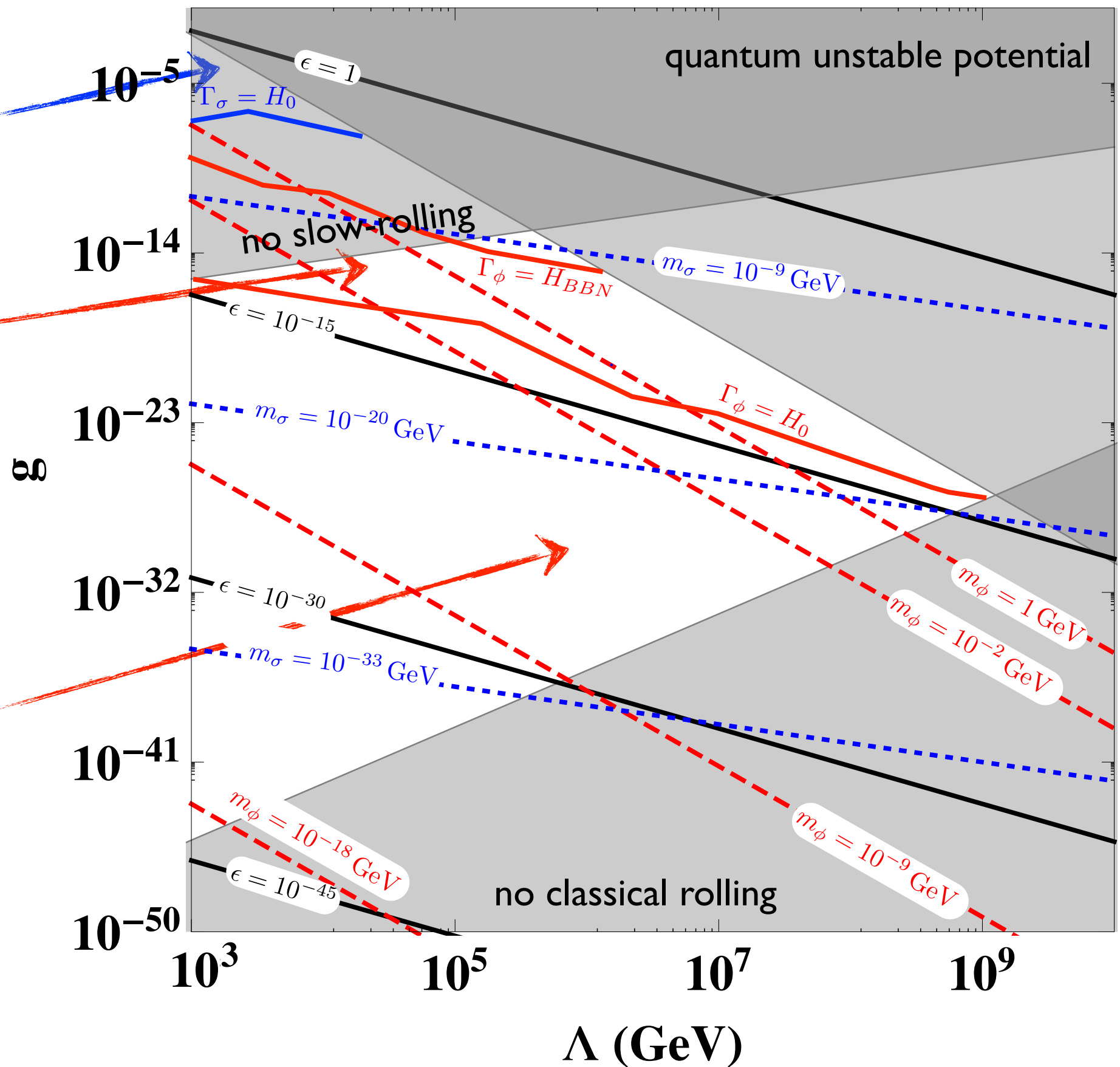


# Phenomenological signatures

$\sigma$  decays within the age of the Universe

$\phi$  decays after BBN

$\phi$  cosmologically stable



# Phenomenological signatures

vacuum misalignment: (after reheating)

quantum spreading makes the scalars oscillate around their minima

$$\Delta\sigma \sim \Delta\phi \sim \sqrt{N_e} H_I$$

the energy stored in these field oscillations behave like cold DM

$$\rho_{\text{ini}}^\sigma \sim m_\sigma^2 (\Delta\sigma)_{\text{ini}}^2 \sim H_I^4$$

$$\rho_{\text{ini}}^\phi \sim H_I^4$$

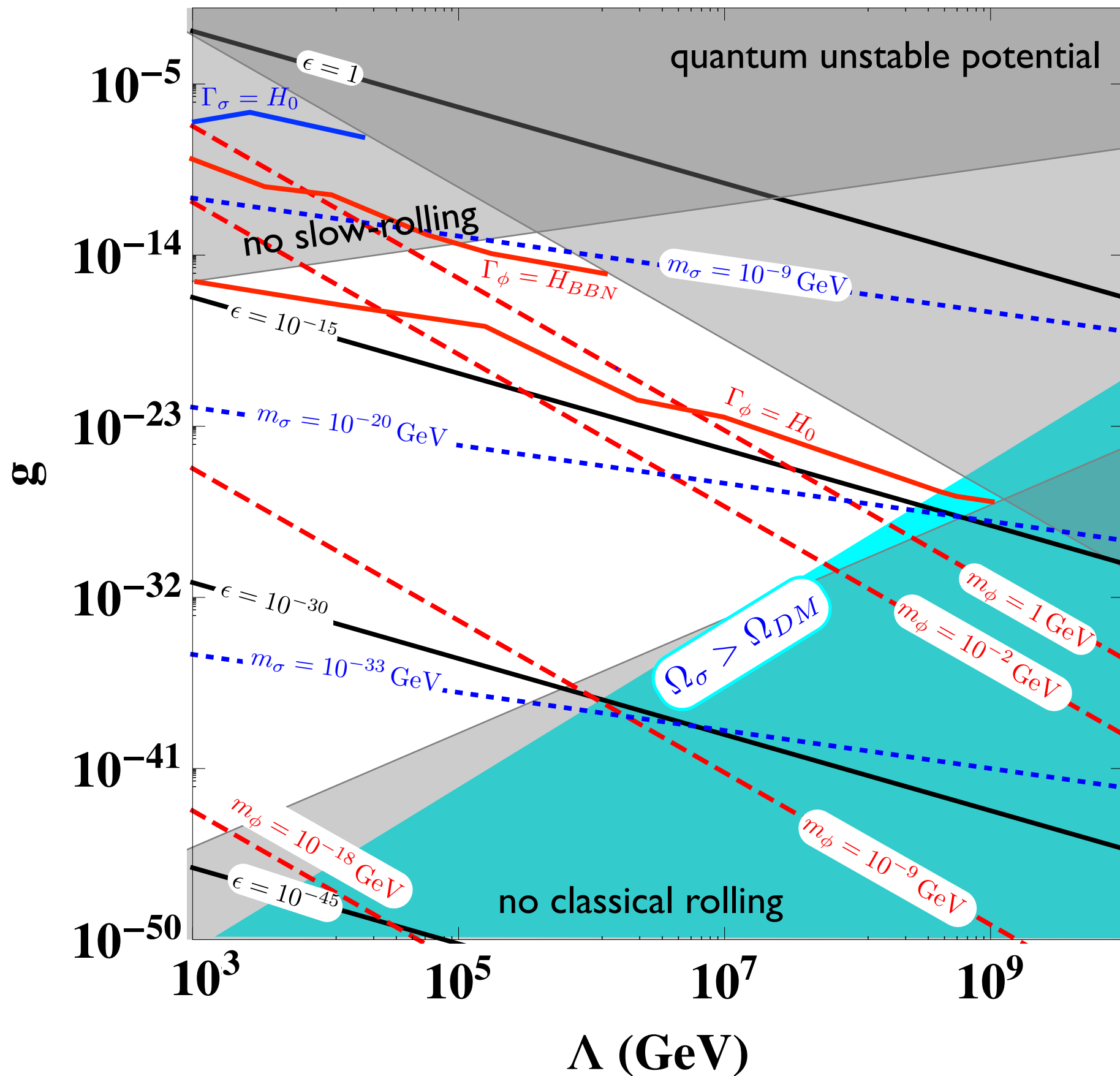
the oscillations start when  $H \sim m_i$  i.e.  $T_{\text{osc}}^i \sim \sqrt{m_i M_{\text{Pl}}}$

the energy density is then redshifted till today

$$\Omega_\sigma \sim \left( \frac{4 \times 10^{-27}}{g_\sigma} \right)^{3/2} \left( \frac{\Lambda}{10^8} \text{ GeV} \right)^{13/2}$$

$\Omega_\phi$  always very small since  $m_\phi \gg m_\sigma$  i.e.  $T_{\text{osc}}^\phi \gg T_{\text{osc}}^\sigma$

# Phenomenological signatures



# Phenomenological signatures

$\phi$  thermal production via interaction with the Higgs

$$h + h \rightarrow \phi + \phi \quad \text{or} \quad SM + SM \rightarrow h^{(*)} \rightarrow \phi + \phi$$

single production is subdominant since linear interactions are suppressed by small mixing angle

$\phi$  almost never in thermal equilibrium (except above  $\Gamma_{\text{BBN}}$  line)

number density is obtained from Boltzmann equation

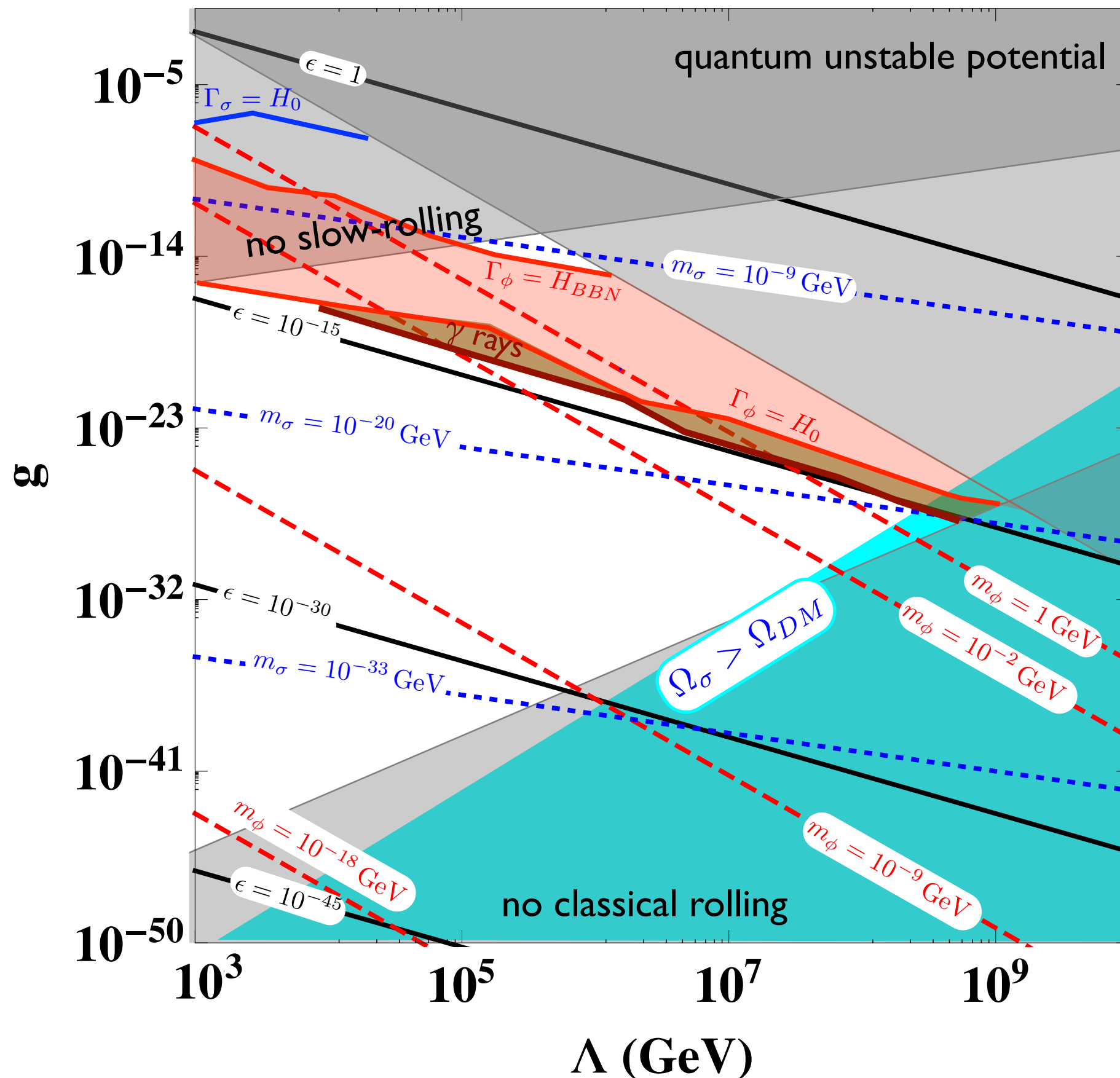
$$\frac{dn_\phi}{dt} + 3Hn_\phi = -\langle\sigma_{AV}\rangle(n_\phi^2 - n_{\phi,eq}^2) \quad \Rightarrow \quad Y_\phi \sim 10^{-4} \epsilon^2 \frac{\Lambda^4}{f^4} \frac{M_{\text{Pl}}}{m_h}$$

© BBN constraints

© distortions in galactic and extra galactic diffuse X-ray and  $\gamma$ -ray backgrounds



# Phenomenological signatures



# Some open questions

- ▶ large field excursions  $\hookrightarrow$  monodromy? McAllister, Schwaller, Servant, Westphal 'in progress  
Ibanez, Montero, Uranga, Valenzuela '15
- ▶ non-periodic potential for an axion-like particle?  
Gupta, Komargodski, Perez, Ubaldi '15
- ▶ hierarchy of decay constants:  $F \gg f$  is  $\sim$  to non-periodic potential  
Choi, Im '15      Kaplan, Rattazzi '15
- ▶ eternal inflation vs classical evolution?  
Arvanitaki, Dimopoulos, Villadoro 'private communication
- ▶ long period of inflation?  
Riotto et al 'in progress
- ▶ other source of friction to prevent over-shooting the EW scale?  
Hardy '15
- ▶ UV completion?
- ▶ can other scales be relaxed too? SUSY breaking scale?  
Batell, Giudice, McCullough '15
- ▶ signatures in atomic physics?

**A new playground for model builders  
at the cross-road between exp/cosmo/pheno/strings  
Joined forces needed**

# Conclusions

- Existence proof that technical naturalness doesn't require new physics at the weak scale.
- Is technical naturalness the right criterion?
- The energy frontier might be different than what we thought for many years!

let us think further and be prepared to be surprised