

Multi-Higgs Doublet Models and symmetries

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Multi-Higgs Doublet Models

Multi-Higgs Doublet Models: add more doublets.

Well motivated Beyond Standard Model scenario:

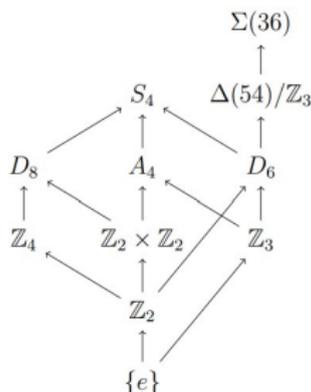
- Baryogenesis
- Dark Matter candidates
- Spontaneous CP violation (SCPV)

Multi-Higgs and symmetries

2HDM: Nishi (2006), Ivanov (2006, 2007)

3HDM list of realizable discrete symmetries: Ivanov, Vdovin

<https://arxiv.org/abs/1210.6553>



Recognizing symmetries in 3HDM in basis-independent way:
IdMV, Ivanov

<https://arxiv.org/abs/1903.11110>

My talk today

FCNC-free multi-Higgs-doublet models from broken family symmetries: IdMV, Talbert

<https://arxiv.org/abs/1908.10979>

Exploring multi-Higgs models with softly broken large discrete symmetry groups: IdMV, Ivanov, Levy

<https://arxiv.org/abs/2107.08227>

Softly-broken A_4 or S_4 3HDMs with stable states: IdMV, Ivo

<https://arxiv.org/abs/2202.00681>

Residual symmetries

$$A \rightarrow T_A A, \quad \text{with } A \in \{u_L, u_R, d_L, d_R, l_L, e_R\},$$

$$T_A = \text{diag}(e^{i\alpha_A}, e^{i\beta_A}, e^{i\gamma_A}).$$

Diagonal

$$T_A \rightarrow T_{AU} = U_A T_A U_A^\dagger,$$

← U_A , fermion mixing

$$m_{AU} = T_{AU}^\dagger m_{AU} T_{AU}.$$

← Symm. of mass matrix
(not of L)

UV Origin

e.g. Higgs VEVs

$$T_A \langle \phi \rangle_A = \langle \phi \rangle_A$$

FCNC in MHDM

In MHDMs

$$\mathcal{L}^Y = - \sum_{k=1}^N \left\{ \bar{Q}'_L \left(Y_k^{d,\prime} H'_k d'_R + Y_k^{u,\prime} \tilde{H}'_k u'_R \right) + \bar{L}'_L Y_k^{e,\prime} H'_k e'_R + h.c. \right\}.$$

Higgs Basis

$$H_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2} G^+ \\ v + S_1^0 + iG^0 \end{pmatrix}, \quad H_{k>1} = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2} S_k^+ \\ S_k^0 + iP_k^0 \end{pmatrix},$$

$$\begin{aligned} \mathcal{L}^Y = & - \left(1 + \frac{S_1^0}{v} \right) (\bar{d}_L m_d d_R + \bar{u}_L m_u u_R + \bar{l}_L m_e e_R) \\ & - \frac{1}{v} \sum_{k=2}^N (S_k^0 + iP_k^0) (\bar{d}_L Y_k^d d_R + \bar{u}_L Y_k^u u_R + \bar{l}_L Y_k^e e_R) \quad \text{FCNC} \\ & - \frac{\sqrt{2}}{v} \sum_{k=2}^N S_k^+ \left(\bar{u}_L V Y_k^d d_R - \bar{u}_R Y_k^{u,\dagger} V d_L + \bar{\nu}_L Y_k^e e_R \right) \\ & + h.c., \end{aligned}$$

Yukawa alignment from residual symmetries

$$Y_k^A \stackrel{!}{=} T_A Y_k^A T_A^\dagger, \quad \text{Alignment}$$

$$\begin{pmatrix} Y_{11} & e^{i(\alpha_1 - \beta_1)} Y_{12} & e^{i(\alpha_1 - \gamma_1)} Y_{13} \\ e^{i(\beta_1 - \alpha_1)} Y_{21} & Y_{22} & e^{i(\beta_1 - \gamma_1)} Y_{23} \\ e^{i(\gamma_1 - \alpha_1)} Y_{31} & e^{i(\gamma_1 - \beta_1)} Y_{32} & Y_{33} \end{pmatrix} \stackrel{!}{=} \begin{pmatrix} Y_{11} & Y_{12} & Y_{13} \\ Y_{21} & Y_{22} & Y_{23} \\ Y_{31} & Y_{32} & Y_{33} \end{pmatrix}$$

"Realistic Toy Model"

A_4 with $\langle \phi_i \rangle \propto \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$

$$\frac{H_1'}{\Lambda} (y_e^1 [\bar{L}_L \phi] e_R + y_\mu^1 [\bar{L}_L \phi]' \mu_R + y_\tau^1 [\bar{L}_L \phi]'' \tau_R) +$$

$$\frac{H_2'}{\Lambda} (y_e^2 [\bar{L}_L \phi] e_R + y_\mu^2 [\bar{L}_L \phi]' \mu_R + y_\tau^2 [\bar{L}_L \phi]'' \tau_R).$$

$$Y_1^{e,\prime} = \frac{v_l}{\Lambda} \begin{pmatrix} y_e^1 & 0 & 0 \\ 0 & y_\mu^1 & 0 \\ 0 & 0 & y_\tau^1 \end{pmatrix}, \quad Y_2^{e,\prime} = \frac{v_l}{\Lambda} \begin{pmatrix} y_e^2 & 0 & 0 \\ 0 & y_\mu^2 & 0 \\ 0 & 0 & y_\tau^2 \end{pmatrix}.$$

Soft breaking terms

$$A_1, S_4, \Delta(S_4), \Sigma(36)$$

$$V_0 = -m^2(\phi_1^\dagger\phi_1 + \phi_2^\dagger\phi_2 + \phi_3^\dagger\phi_3) + V_4,$$

Soft terms

$$V_{\text{soft}} = m_{11}^2\phi_1^\dagger\phi_1 + m_{22}^2\phi_2^\dagger\phi_2 + m_{33}^2\phi_3^\dagger\phi_3 + \left(m_{12}^2\phi_1^\dagger\phi_2 + m_{23}^2\phi_2^\dagger\phi_3 + m_{31}^2\phi_3^\dagger\phi_1 + h.c.\right)$$

$\Sigma(36)$

$$\begin{aligned}
V_0 = & -m^2 \left[\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 + \phi_3^\dagger \phi_3 \right] + \lambda_1 \left[\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 + \phi_3^\dagger \phi_3 \right]^2 \\
& - \lambda_2 \left[|\phi_1^\dagger \phi_2|^2 + |\phi_2^\dagger \phi_3|^2 + |\phi_3^\dagger \phi_1|^2 - (\phi_1^\dagger \phi_1)(\phi_2^\dagger \phi_2) - (\phi_2^\dagger \phi_2)(\phi_3^\dagger \phi_3) - (\phi_3^\dagger \phi_3)(\phi_1^\dagger \phi_1) \right] \\
& + \lambda_3 \left(|\phi_1^\dagger \phi_2 - \phi_2^\dagger \phi_3|^2 + |\phi_2^\dagger \phi_3 - \phi_3^\dagger \phi_1|^2 + |\phi_3^\dagger \phi_1 - \phi_1^\dagger \phi_2|^2 \right).
\end{aligned}$$

alignment A : $A_1 = (\omega, 1, 1)$, $A_2 = (1, \omega, 1)$, $A_3 = (1, 1, \omega)$

alignment A' : $A'_1 = (\omega^2, 1, 1)$, $A'_2 = (1, \omega^2, 1)$, $A'_3 = (1, 1, \omega^2)$

alignment B : $B_1 = (1, 0, 0)$, $B_2 = (0, 1, 0)$, $B_3 = (0, 0, 1)$

alignment C : $C_1 = (1, 1, 1)$, $C_2 = (1, \omega, \omega^2)$, $C_3 = (1, \omega^2, \omega)$

$\Sigma(36)$ masses

A, A'

$$m_{h_{SM}}^2 = 2\lambda_1 v^2 = 2m^2,$$

$$m_{H^\pm}^2 = \frac{1}{2}\lambda_2 v^2 \quad (\text{double degenerate}),$$

$$m_h^2 = \frac{1}{2}\lambda_3 v^2 \quad (\text{double degenerate}),$$

$$m_H^2 = 3m_h^2 = \frac{3}{2}\lambda_3 v^2 \quad (\text{double degenerate}).$$

B, C

$$m_{h_{SM}}^2 = 2(\lambda_1 + \lambda_3)v^2 = 2m^2,$$

$$m_{H^\pm}^2 = \frac{1}{2}(\lambda_2 - 3\lambda_3)v^2 \quad (\text{double degenerate}),$$

$$m_h^2 = -\frac{1}{2}\lambda_3 v^2 \quad (\text{double degenerate}),$$

$$m_H^2 = 3m_h^2 = -\frac{3}{2}\lambda_3 v^2 \quad (\text{double degenerate}).$$

NO SCPV

Alignment preserving soft breaking

$$\frac{\partial V_0}{\partial \phi_i^*} = -m^2 \phi_i + \frac{\partial V_4}{\partial \phi_i^*} = 0.$$

Add

$$V_{\text{soft}} = \phi_i^\dagger M_{ij} \phi_j, \quad M_{ij} = \begin{pmatrix} m_{11}^2 & m_{12}^2 & (m_{31}^2)^* \\ (m_{12}^2)^* & m_{22}^2 & m_{23}^2 \\ m_{31}^2 & (m_{23}^2)^* & m_{33}^2 \end{pmatrix},$$

$$\frac{\partial V}{\partial \phi_i^*} = M_{ij} \phi_j - m^2 \phi_i + \frac{\partial V_4}{\partial \phi_i^*} = 0.$$

Reprint

$$v|_V \text{ extremum} = \zeta \cdot v|_{V_0 \text{ extremum}}.$$

$$\left. \frac{\partial V_4}{\partial \phi_i^*} \right|_{V_0 \text{ extremum}} = m^2 \phi_i|_{V_0 \text{ extremum}}.$$

$$\left. \frac{\partial V_4}{\partial \phi_i^*} \right|_{V \text{ extremum}} = \zeta^2 \cdot m^2 \phi_i|_{V \text{ extremum}}.$$

$$M_{ij} \phi_j = (1 - \zeta^2) m^2 \phi_i. \quad \text{Eigenvector of } M$$

Example (1, 1, 1)

$$M_{ij} = \mu_1 n_{1i} n_{1j}^* + \mu_2 n_{2i} n_{2j}^* + \mu_3 n_{3i} n_{3j}^* .$$

$$\begin{array}{ccc} (1, 1, 1) & (0, 1, -1) & (2, -1, -1) \\ \ell_1 & \ell_2 & \ell_3 \end{array}$$

$$\vec{n}_i = U_{ij} \vec{e}_j, \quad i, j = 2, 3, \quad \text{where } U = \begin{pmatrix} \cos \theta & e^{i\xi} \sin \theta \\ -e^{-i\xi} \sin \theta & \cos \theta \end{pmatrix} .$$

$$\Sigma = \mu_2 + \mu_3, \quad \delta = \mu_2 - \mu_3, \quad \theta, \quad \xi .$$

Universal results

Universal Correlations (A, A', B, C)

$$\Delta m_{H_1^\pm}^2 = \mu_2 = \frac{\Sigma + \delta}{2}, \quad \Delta m_{H_2^\pm}^2 = \mu_3 = \frac{\Sigma - \delta}{2}.$$

$$m_{h_1}^2 = \frac{1}{2} \left(2|\lambda_3|v^2 + \Sigma - \sqrt{(\lambda_3 v^2)^2 + \delta^2 + 2x|\lambda_3||\delta|v^2} \right),$$

$$m_{h_2}^2 = \frac{1}{2} \left(2|\lambda_3|v^2 + \Sigma - \sqrt{(\lambda_3 v^2)^2 + \delta^2 - 2x|\lambda_3||\delta|v^2} \right),$$

$$m_{H_1}^2 = \frac{1}{2} \left(2|\lambda_3|v^2 + \Sigma + \sqrt{(\lambda_3 v^2)^2 + \delta^2 - 2x|\lambda_3||\delta|v^2} \right),$$

$$m_{H_2}^2 = \frac{1}{2} \left(2|\lambda_3|v^2 + \Sigma + \sqrt{(\lambda_3 v^2)^2 + \delta^2 + 2x|\lambda_3||\delta|v^2} \right),$$

$$x = \sqrt{1 - (\sin 2\theta \sin \xi)^2}.$$

All non-SM Higgses decay
(some decays suppressed...
displaced vertices?)

A_4, S_4 : Dark Matter candidates

$A_4, S_4, (1, 0, 0)$ Alignment

$$\rho = \begin{pmatrix} 1 & & \\ & -1 & \\ & & -1 \end{pmatrix},$$

The AP SBP_2 preserve ρ !

Dark Matter Candidates

Not a general feature!

$\Delta(S_4), (1, 0, 0)$ Alignment

$$\rho_{23} = \begin{pmatrix} 1 & & \\ & 1 & \\ & & 1 \end{pmatrix}, \text{ but } M_{22} \neq M_{33}$$

Conclusions

- Multi-Higgs with symmetries are well motivated.
- Symmetries control flavour changing processes.
- Softly broken symmetries - interesting phenomenology.
- Multi-Higgs with softly broken A_4 , S_4 - Dark Matter.