# All about the Triplet

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- Introduction
- The Georgi-Machacek model, solving the Tparameter: susy breaking
- ..... and dark matter.
- Conclusions

# Introduction

- The run I of the LHC has discovered the last building block of the SM: The Higgs
- Its properties are compatible with the ones of the SM: minimal realization of the Higgs potential
- But an extended Higgs sector is not yet excluded and can have different and interesting signatures.

- One way of deviating from the minimal model is to include extra Higgses:
  - More doublets
  - singlets
  - Triplets
  - •

## These extra fields could appear in a UV completion of the model and could remain light by the same mechanism as the SM Higgs

- Pseudo-goldstone boson
- Supersymmetry

- In this talk I am going to explore different aspects of an extended Higgs sector with triplets in a supersymmetric model.
- Triplets can help on the Higgs mass without destabilizing the EW scale
- In general there are two possible triplets that could mix with the Higgs, Y=0 and Y=1.
- The kind of couplings are:

 $W = \lambda_0 H_u T_0 H_d + \lambda_{+1} H_d T_{+1} H_d + \lambda_{-1} H_u T_{-1} H_u$ 

## In principle the triplet will get a vev once EWSB occurs and can lead to a dangerous contribution to the T-parameter

• I will make sure that the vev of the triplet component by implementing a symmetry to avoid contributions to the T-parameter.

# The (SUSY)-Georgi-Machacek model: SUSY breaking and DM

- The GM model was proposed to ensure a custodial structure with triplets.
- It was supersymmetrized in arXiv:1308.4025

$$\bar{H} = \begin{pmatrix} H_1 \\ H_2 \end{pmatrix}, \quad \bar{\Delta} = \begin{pmatrix} -\frac{\Sigma_0}{\sqrt{2}} & -\Sigma_{-1} \\ -\Sigma_1 & \frac{\Sigma_0}{\sqrt{2}} \end{pmatrix}$$

$$W_0 = \lambda \bar{H} \cdot \bar{\Delta} \bar{H} + \frac{\lambda_3}{3} \operatorname{tr} \bar{\Delta}^3 + \frac{\mu}{2} \bar{H} \cdot \bar{H} + \frac{\mu_{\Delta}}{2} \operatorname{tr} \bar{\Delta}^2 + h_t \overline{Q}_L \cdot H_2 t_R + h_b \overline{Q}_L \cdot H_1 b_R$$

 RGE evolution, yukawa and U(1) break the custodial limit and generate a non-zero ρ so in general:

$$W = -\lambda_a H_1 \cdot \Sigma_1 H_1 + \lambda_b H_2 \cdot \Sigma_{-1} H_2 + \sqrt{2}\lambda_c H_1 \cdot \Sigma_0 H_2 + \sqrt{2}\lambda_3 \operatorname{tr} \Sigma_1 \Sigma_0 \Sigma_{-1} - \mu H_1 \cdot H_2 + \frac{\mu_{\Delta_a}}{2} \operatorname{tr} \Sigma_0^2 + \mu_{\Delta_b} \operatorname{tr} \Sigma_1 \Sigma_{-1} + h_t \overline{Q}_L \cdot H_2 t_R + h_b \overline{Q}_L \cdot H_1 b_R$$

$$V_{\text{SOFT}} = m_{H_1}^2 H_1^{\dagger} H_1 + m_{H_2}^2 H_2^{\dagger} H_2 + m_{\Sigma_0}^2 \Sigma_0^{\dagger} \Sigma_0 + m_{\Sigma_1}^2 \Sigma_1^{\dagger} \Sigma_1 + m_{\Sigma_{-1}}^2 \Sigma_{-1}^{\dagger} \Sigma_{-1} - m_3^2 H_1 \cdot H_2$$
  
+  $\left\{ \frac{B_{\Delta_a}}{2} \text{tr} \Sigma_0^2 + B_{\Delta_b} \text{tr} \Sigma_1 \Sigma_{-1} - A_{\lambda_a} H_1 \cdot \Sigma_1 H_1 + A_{\lambda_b} H_2 \cdot \Sigma_{-2} H_2 \right.$   
+  $\sqrt{2} A_{\lambda_c} H_1 \cdot \Sigma_0 H_2 + \sqrt{2} A_{\lambda_3} \text{tr} \Sigma_1 \Sigma_0 \Sigma_{-1} + a_t \tilde{Q}_L \cdot H_2 \tilde{t}_R + a_b \tilde{Q}_L \cdot H_1 \tilde{b}_R + h.c. \right\}$ 

 I am going to embed this model into a predictive scenario of SUSY breaking like GMSB to see the deviation from the custodial limit:

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$$W = \left(\tilde{\lambda}_8^{ij}X + \mathcal{M}_8^{ij}\right)\Phi_{8i}\Phi_{8j} + \left(\tilde{\lambda}_3^{ij}X + \mathcal{M}_3^{ij}\right)\Phi_{3i}\Phi_{3j} + \left(\tilde{\lambda}_1^{ij}X + \mathcal{M}_1^{ij}\right)\bar{\Phi}_{1i}\Phi_{1j}$$

$$M_{3} = \frac{\alpha_{3}(\mathcal{M})}{4\pi} 3n_{8}g(\Lambda_{8}/\mathcal{M})\Lambda_{8},$$
  

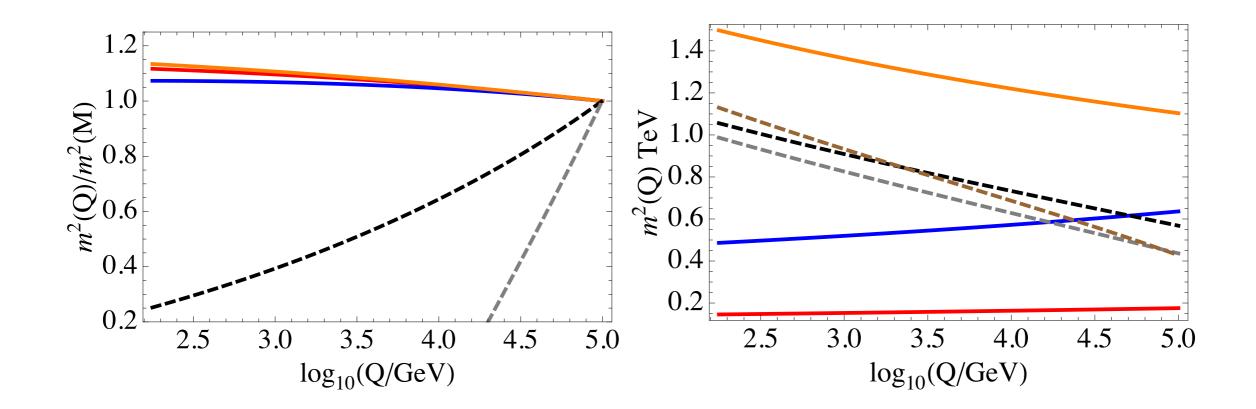
$$M_{2} = \frac{\alpha_{2}(\mathcal{M})}{4\pi} 2n_{3}g(\Lambda_{3}/\mathcal{M})\Lambda_{3},$$
  

$$M_{1} = \frac{\alpha_{1}(\mathcal{M})}{4\pi} \frac{6}{5}n_{1}g(\Lambda_{1}/\mathcal{M})\Lambda_{1},$$

$$m_{\tilde{f}}^2 = 2\left[C_3^f \left(\frac{\alpha_3(\mathcal{M})}{4\pi}\right)^2 3n_8 f(\Lambda_8/\mathcal{M})\Lambda_8^2 + C_2^f \left(\frac{\alpha_2(\mathcal{M})}{4\pi}\right)^2 2n_3 f(\Lambda_3/\mathcal{M})\Lambda_3^2 + C_1^f \left(\frac{\alpha_1(\mathcal{M})}{4\pi}\right)^2 \frac{1}{2} \left(\frac{6}{5}\right)^2 n_1 f(\Lambda_1/\mathcal{M})\Lambda_1^2\right].$$

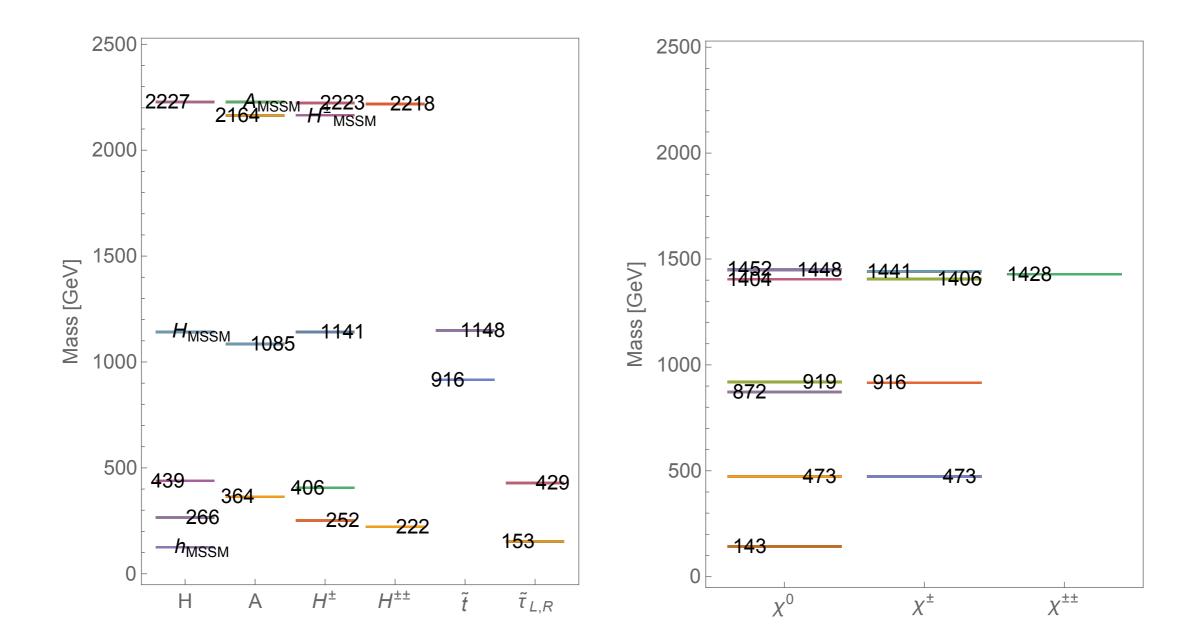
- The spectrum is generated at a high scale M
- Run down to the EW scale
- EWSB it is imposed
- m<sub>h</sub>=125 GeV
- All experimental constrains (direct & indirect) are satisfied

$$n_1 = 1, n_3 = 2, n_8 = 6$$
 and  $\tilde{\lambda}_1 = 0.9, \tilde{\lambda}_3 = 0.5, \tilde{\lambda}_8 = 0.1$ .



### Higgses

squarks & gauginos



 $\tan \beta = 1.38$ 

• There are deviations on Higgs properties and exotic decays like:

$$H^{\pm} \to W^{\pm}Z$$

Scenario #1	WW	ZZ	$b\overline{b}$	$t\overline{t}$	$\gamma\gamma$
$r_{hXX}$	1.05	1.04	1.01	1.01	1.22
$\begin{bmatrix} \mu_{hXX}^{(gF)}, \mu_{hXX}^{(htt)} \end{bmatrix}$	1.07	1.05	1	0.99	1.45
$\left[ \begin{array}{c} \mu_{hXX}^{(WF)}, \mu_{hXX}^{(Wh)} \end{array} \right]$	1.16	1.14	1.08	1.07	1.58
$\left[\begin{array}{c} \mu_{hXX}^{(ZF)}, \mu_{hXX}^{(Zh)} \end{array}\right]$	1.14	1.11	1.06	1.05	1.54

# How about DM?

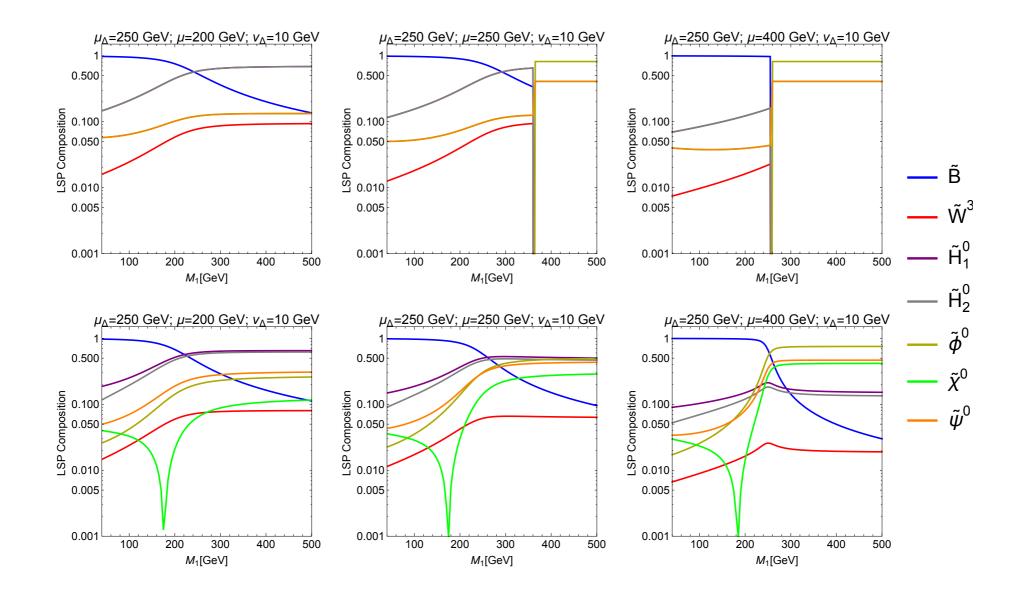
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- Having extra fermionic states can have an impact in the composition and properties of the LSP
- I am going to present different cases where the relic abundance is obtained
- NB: I am deviating from GMSB. i e the gravitino is not the LSP!!!!!!

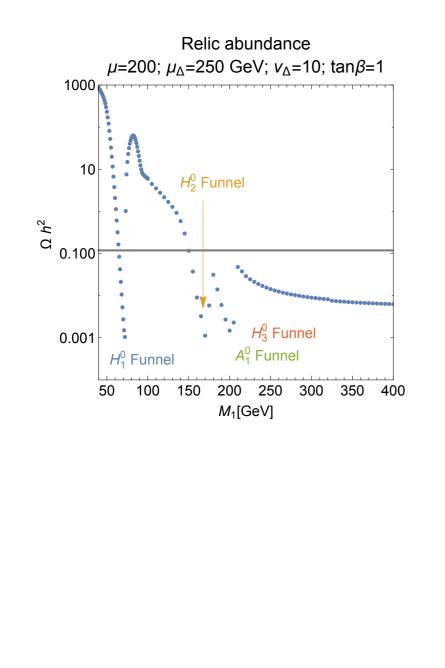
$$\mathbf{M} = \begin{pmatrix} M_{1} & 0 & -\frac{1}{\sqrt{2}}g'c_{\beta}v_{H} & \frac{\sqrt{2}}{2}g's_{\beta}v_{H} & 0 & -g'v_{\Delta} & g'v_{\Delta} \\ 0 & M_{2} & \frac{\sqrt{2}}{2}g_{2}c_{\beta}v_{H} & -\frac{1}{\sqrt{2}}g_{2}s_{\beta}v_{H} & 0 & g_{2}v_{\Delta} & -g_{2}v_{\Delta} \\ -\frac{1}{\sqrt{2}}g'c_{\beta}v_{H} & \frac{1}{\sqrt{2}}g_{2}c_{\beta}v_{H} & -\sqrt{2}\lambda v_{\Delta} & -\frac{1}{\sqrt{2}}\lambda v_{\Delta} - \mu & -\lambda s_{\beta}v_{H} & 0 & -2\lambda c_{\beta}v_{H} \\ \frac{\sqrt{2}}{2}g's_{\beta}v_{H} & -\frac{1}{\sqrt{2}}g_{2}s_{\beta}v_{H} & -\frac{1}{\sqrt{2}}\lambda v_{\Delta} - \mu & -\sqrt{2}\lambda v_{\Delta} & -\lambda c_{\beta}v_{H} & 0 & -2\lambda s_{\beta}v_{H} \\ 0 & 0 & -\lambda s_{\beta}v_{H} & \lambda c_{\beta}v_{H} & \mu_{\Delta} & -\frac{1}{\sqrt{2}}\lambda_{3}v_{\Delta} & -\frac{1}{\sqrt{2}}\lambda_{3}v_{\Delta} \\ g'v_{\Delta} & g_{2}v_{\Delta} & 0 & -2\lambda s_{\beta}v_{H} & -\frac{1}{\sqrt{2}}\lambda_{3}v_{\Delta} & 0 & \mu_{\Delta} - \frac{1}{\sqrt{2}}\lambda_{v_{\Delta}} \\ g'v_{\Delta} & -g_{2}v_{\Delta} & -2\lambda c_{\beta}v_{H} & 0 & -\frac{1}{\sqrt{2}}\lambda_{3}v_{\Delta} & 0 \end{pmatrix}$$

### Neutralino Mass-Matrix

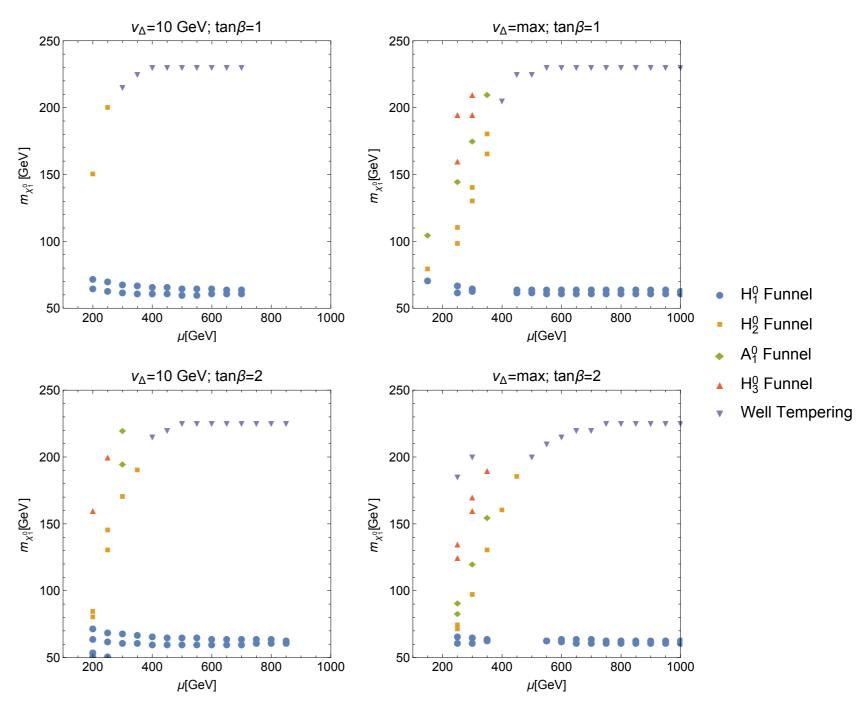
#### Composition of the LSP



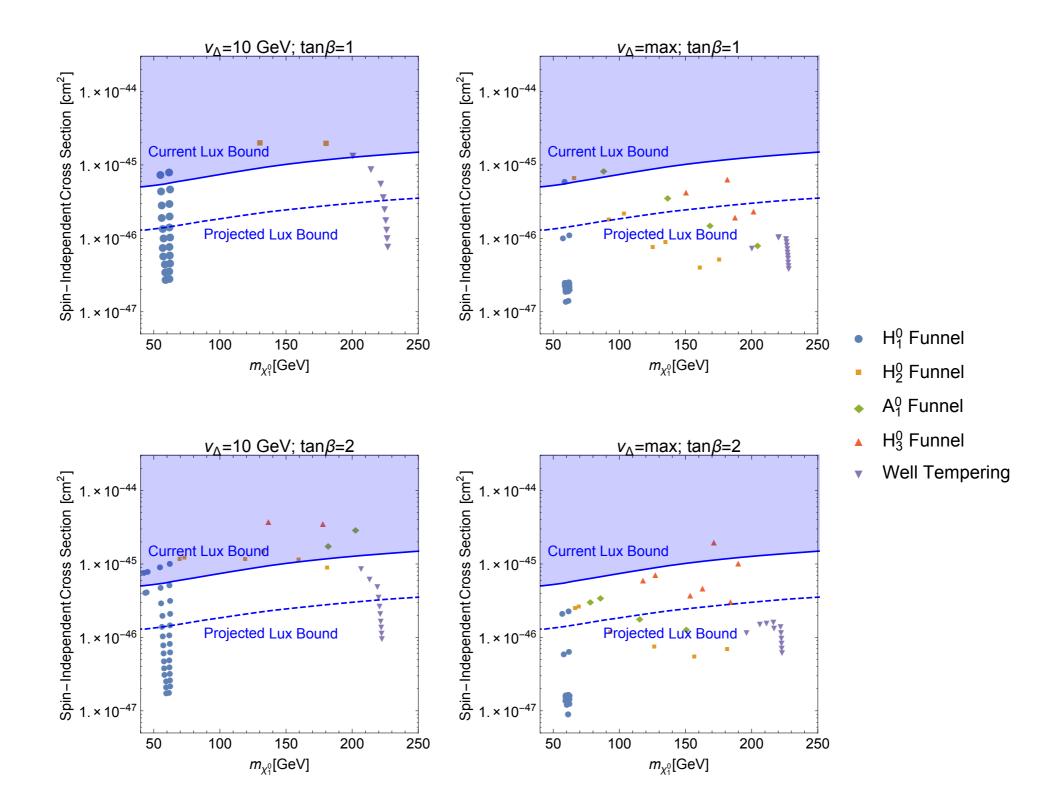
M<sub>2</sub>=1 TeV



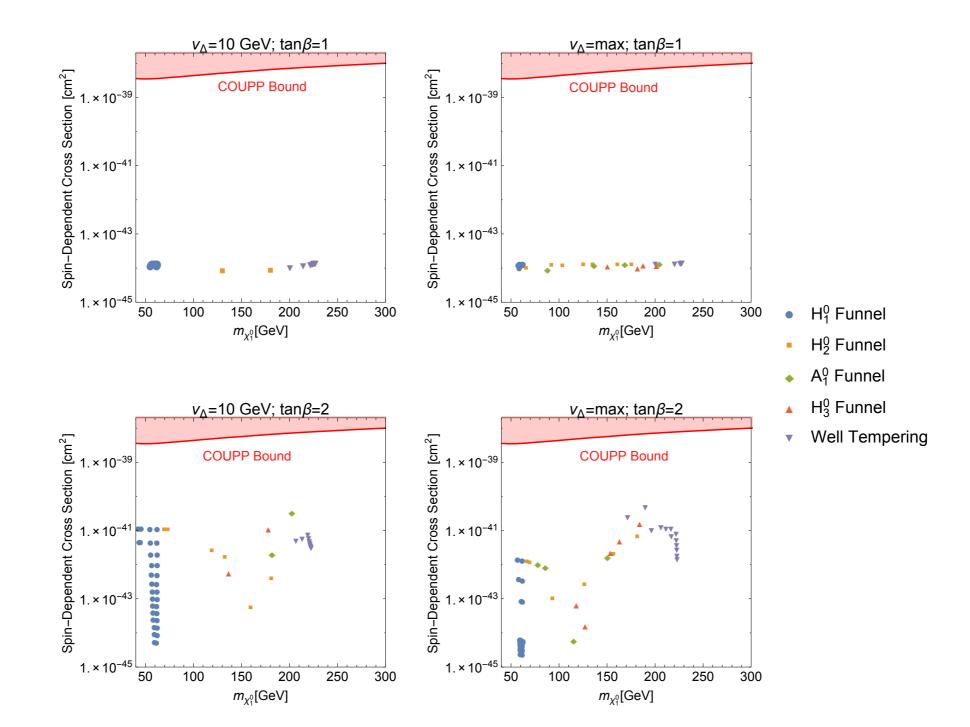
#### **Relic Abudance**



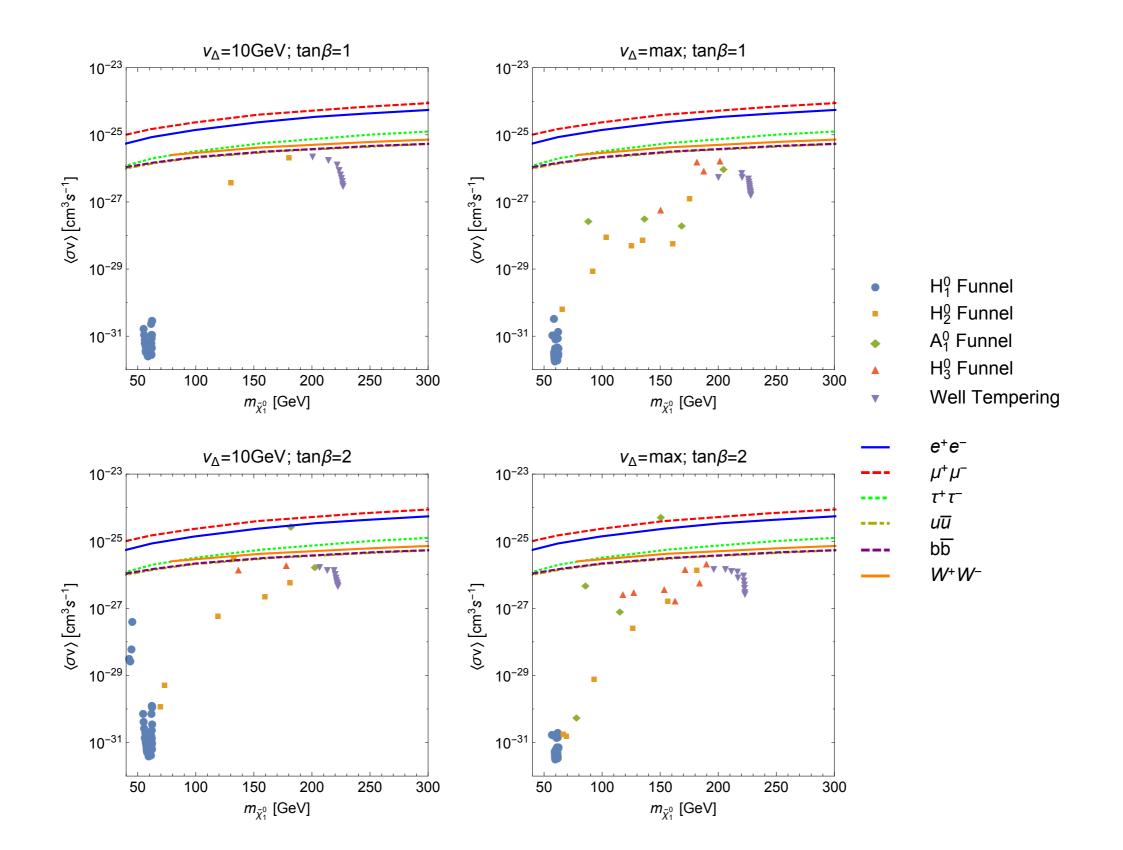
#### Spin independent xsec



#### Spin dependent xsec



#### Indirect detection



# Conclusions

- One of the multiple possibilities for physics beyond the SM is an extended Higgs sector
- Extra scalars appear naturally in UV theories attempting to explain the Hierarchy problem
- In this talk I have supposed that supersymmetry is the explanation of the EW scale and moreover that there are triplets coupled to the usual Higgses.

- One way to automatically have the T-parameter under control is the GM model.
- I have introduced the SCTM and study how can it be embedded into GMSB
- It naturally leads to a low messenger scale
- But one can successfully have a complete model with low tan  $\boldsymbol{\beta}$

### Finally I have analyzed the implications that having an extended sector of neutralinos on DM.

- New regions appear that can have very interesting implications for direct and indirect detection.
- The SCTM has very exotic decays for the Higgs sectors that I am currently studying.