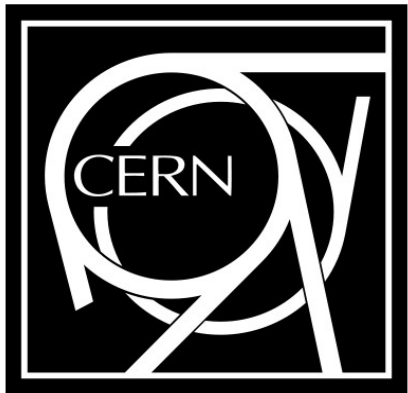
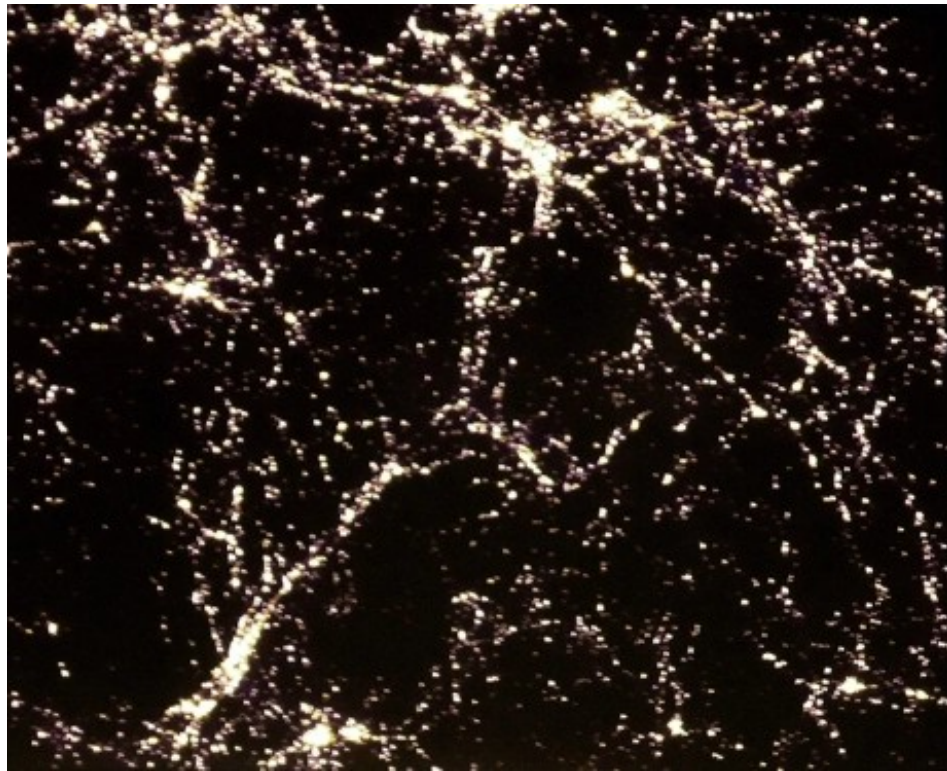
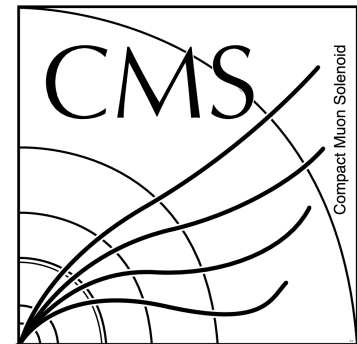


Dark Matter Mediators @ Colliders

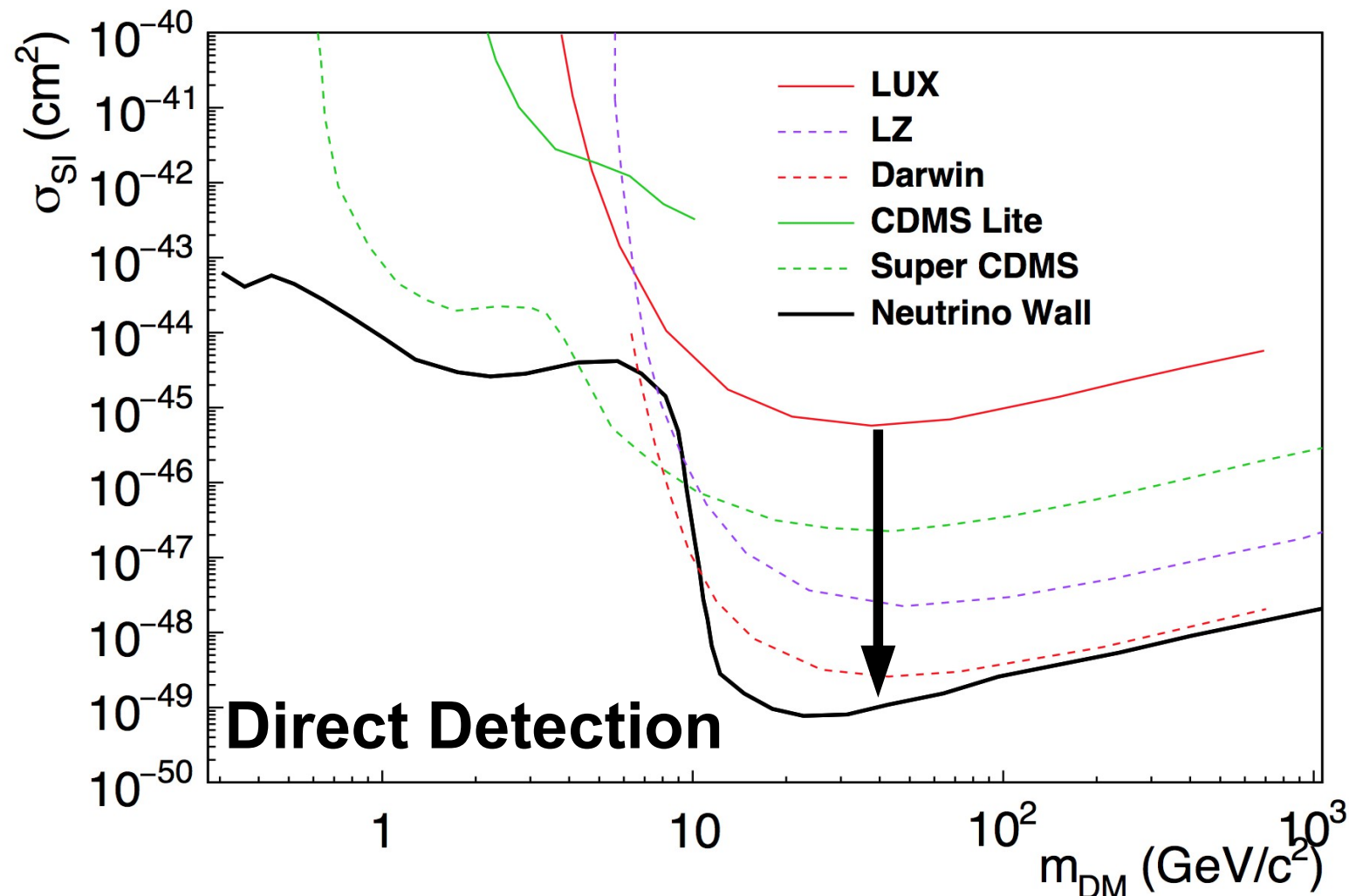


Phil Harris
(CERN)



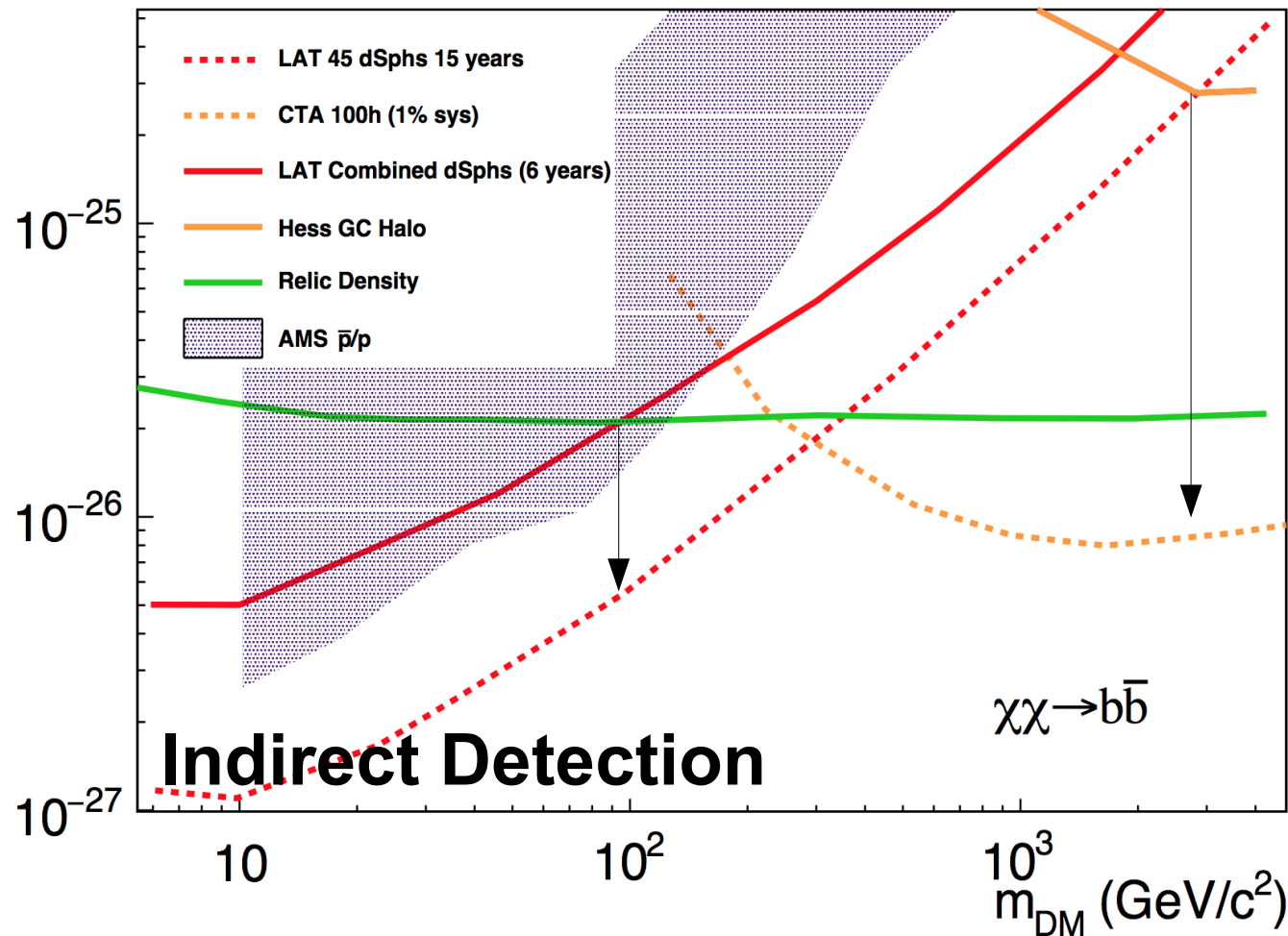
Dark Matter searches not @ collider

Dark matter searches not at colliders have **clear benchmarks**



Goal: **get to the Neutrino background wall**

Dark Matter searches not @ collider



Goal: get to the Relic density

Question:

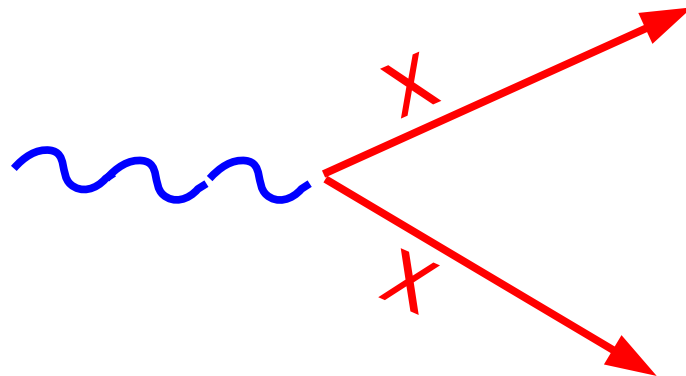
- Whats the simplest way to present LHC results in the context of Dark Matter?

Question:

- Whats the simplest way to present LHC results in the context of Dark Matter?

- Answer:

– $\sigma_{\text{Invisible}}$



- Assumes dark matter coupling to standard model

– $\mathcal{L} = g_{\text{DM}} \chi \bar{\chi} Y \longrightarrow \text{Mediator} + \text{SM interactions}$

Dark Matter

Adding Dark Matter

- What drives dark matter interaction is production
 - Take the approach that this is defined by the mediator

• $\mathcal{L}' = g_{\text{DM}} \cancel{\chi\chi} Y$

Z'^μ Spin 1

Uniform coupling to SM

$$\mathcal{L}' = \mathcal{L}' + g_{\text{SM}} Z'_\mu \bar{q} \gamma^\mu q$$

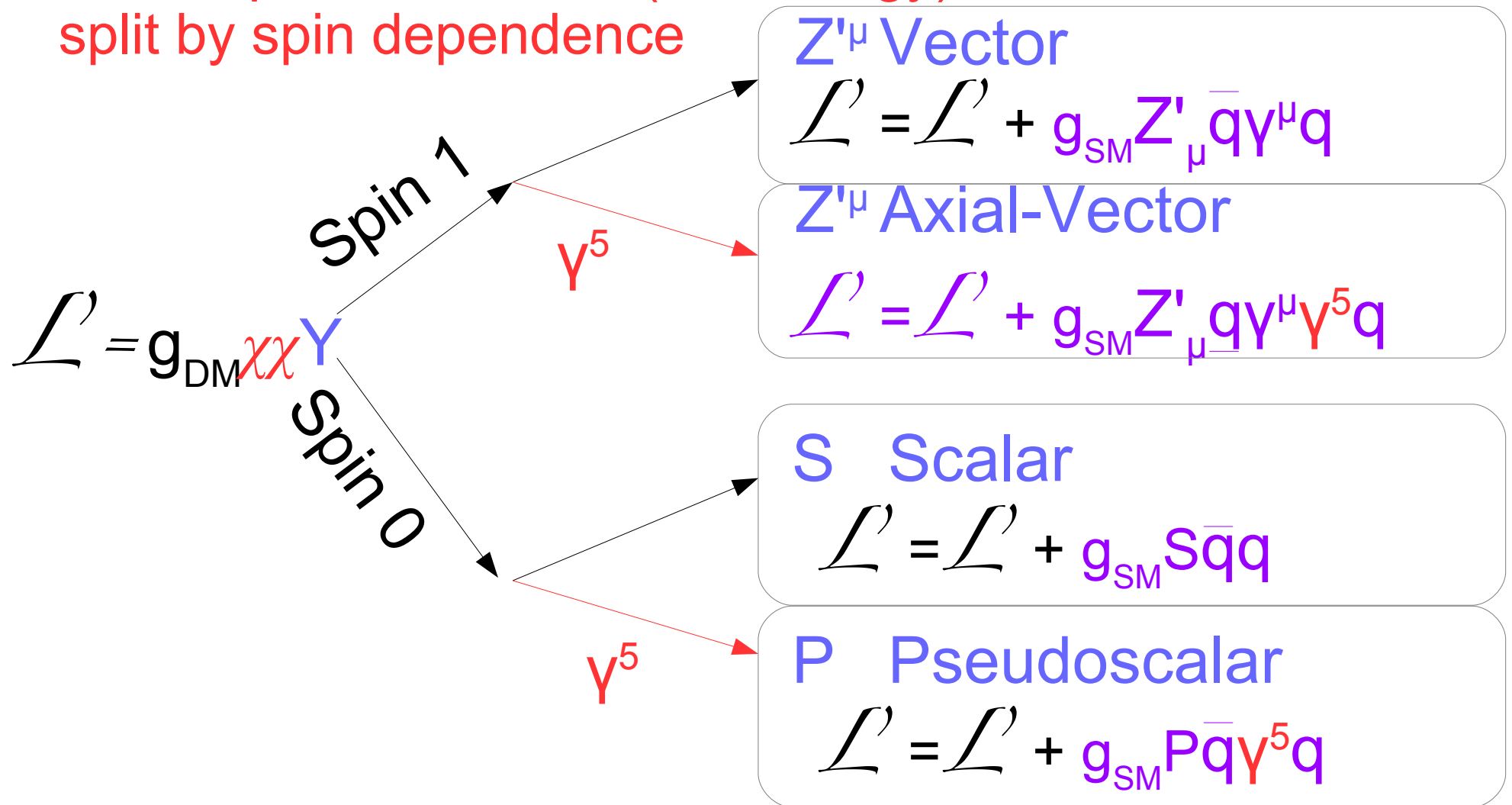
S Spin 0

Yukawa* couplings to SM

$$\mathcal{L}' = \mathcal{L}' + g_{\text{SM}} S \bar{q} q$$

Preserving Generality?

To compare with other (low energy) searches :
split by spin dependence



Strategy of searches in LHC does not change much

Interpretation against Direct Detection/Indirect Changes a lot

Simplified Models 101

Vector

$$g_{\text{DM}} Z'_\mu \bar{\chi} \gamma^\mu \chi$$

EWK style coupling
(equal to all quarks/leptons)

Axial vector

$$g_{\text{DM}} Z''_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi$$

EWK style coupling
(equal to all quarks/leptons)

Scalar

$$g_{\text{DM}} S \bar{\chi} \chi$$

Yukawa style coupling
(Mass based coupling)

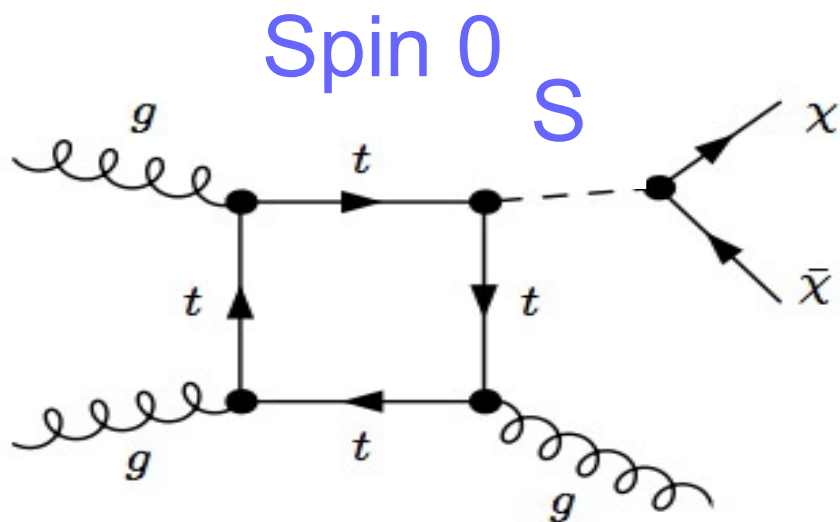
Pseudoscalar

$$g_{\text{DM}} P \bar{\chi} \gamma^5 \chi$$

Yukawa style coupling
(Mass based coupling)

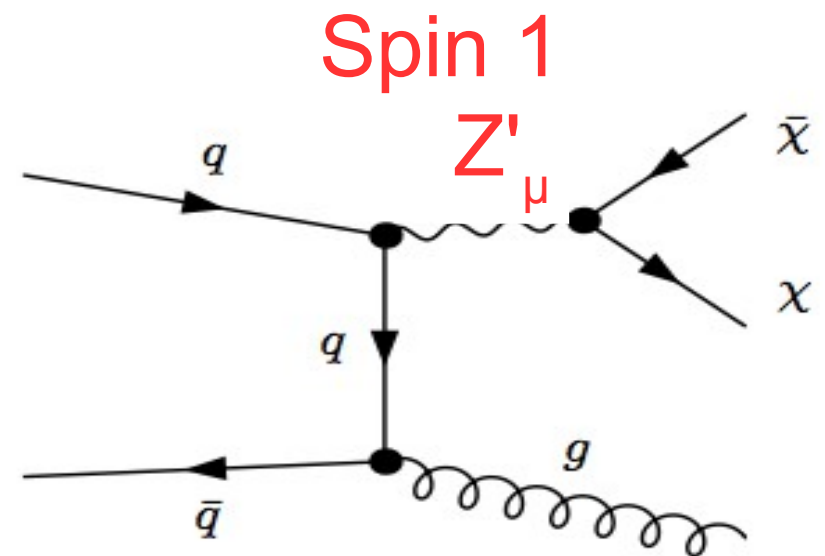
Impact of the spin in production

- Couplings to SM force two different scenarios



Yukawa coupling to quarks
Dominated by heavy quarks

Small cross sections :
Probe low masses or
Large couplings



Flavor universal to quarks
Dominated by light quark

Large cross sections :
Probe large masses or
Small couplings

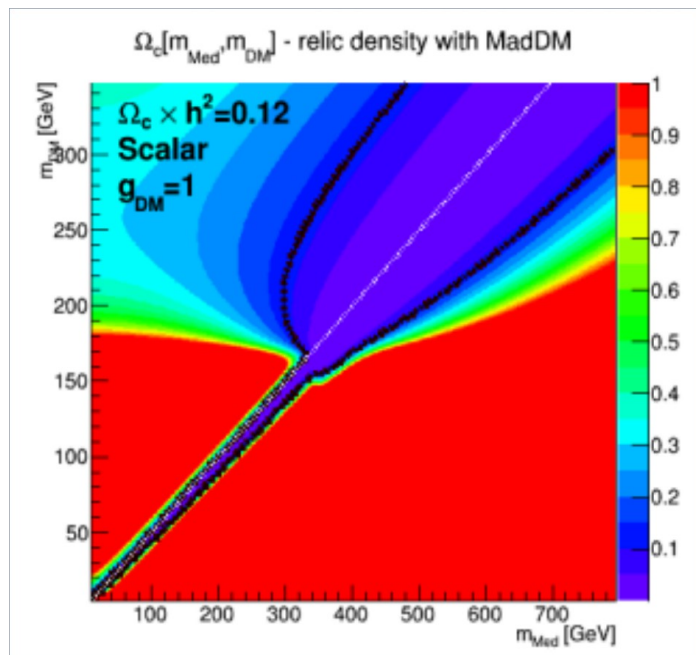
Establishing a
collider benchmark

Relic Density??

Relic Density on Simplified Models

- Lets calculate relic density on simplified models
 - Calculation is performed with MadDM (2.0)  MadDM

Acknowledge these models are simplified
All of our results come with caveats



However it still serves as a
motivating benchmark

Note : I am experimentalist
Don't expect a strong theoretical motivation

Features of the relic density

- Considering form of the annihilation cross section

$$\sigma_{\text{ann},s}^A \cdot v = \sum_q \frac{N_c^q g_{\text{DM}}^2 g_q^2 \beta_q}{2\pi} \frac{m_q^2 (4m_{\text{DM}}^2 - M_{\text{med}}^2)^2}{M_{\text{Med}}^4 \left[(M_{\text{med}}^2 - 4m_{\text{DM}}^2)^2 + M_{\text{med}}^2 \Gamma_{\text{med}}^2 \right]}$$

$$\Gamma_{\text{vector}}^{\chi\bar{\chi}} = \frac{g_{\text{DM}}^2 M_{\text{med}}}{12\pi} (1 - 4z_{\text{DM}})^{1/2} (1 + 2z_{\text{DM}}),$$

$$z_i = m_i^2 / M_{\text{med}}^2$$

Width a function of :
 $g_{\text{SM}}^2, g_{\text{DM}}^2, m_{\text{DM}}$ and m_{MED}

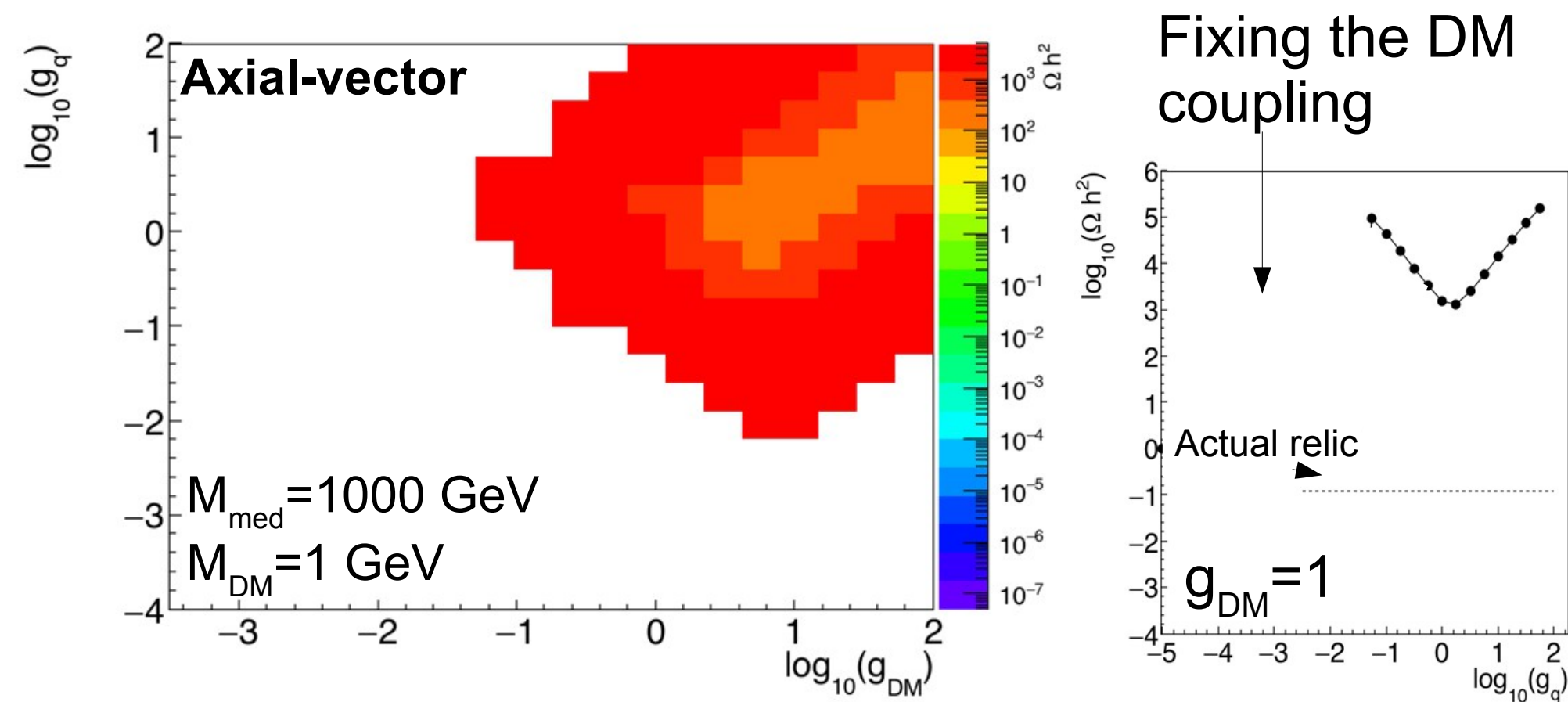
For constant $m_{\text{DM}}, m_{\text{MED}}, g_{\text{DM}}$ and $\langle \sigma v \rangle_{\text{relic}}$:

Above form is **quadratic** in g_{SM} for fixed relic

Yields two possible solutions for the relic density

Solving for the relic

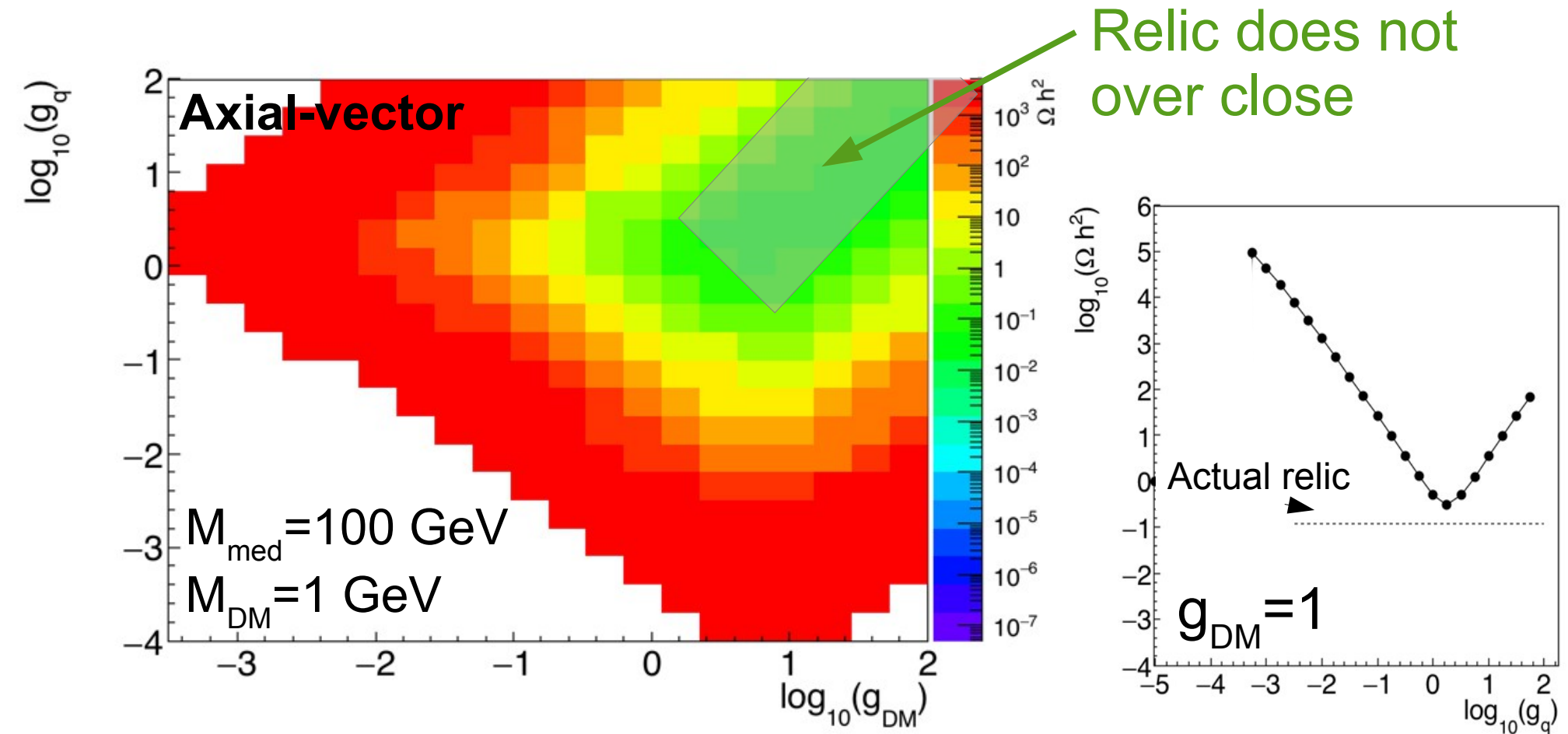
- We can numerically solve for the solution



For a large mediator we have no allowed solution

Solving for the relic

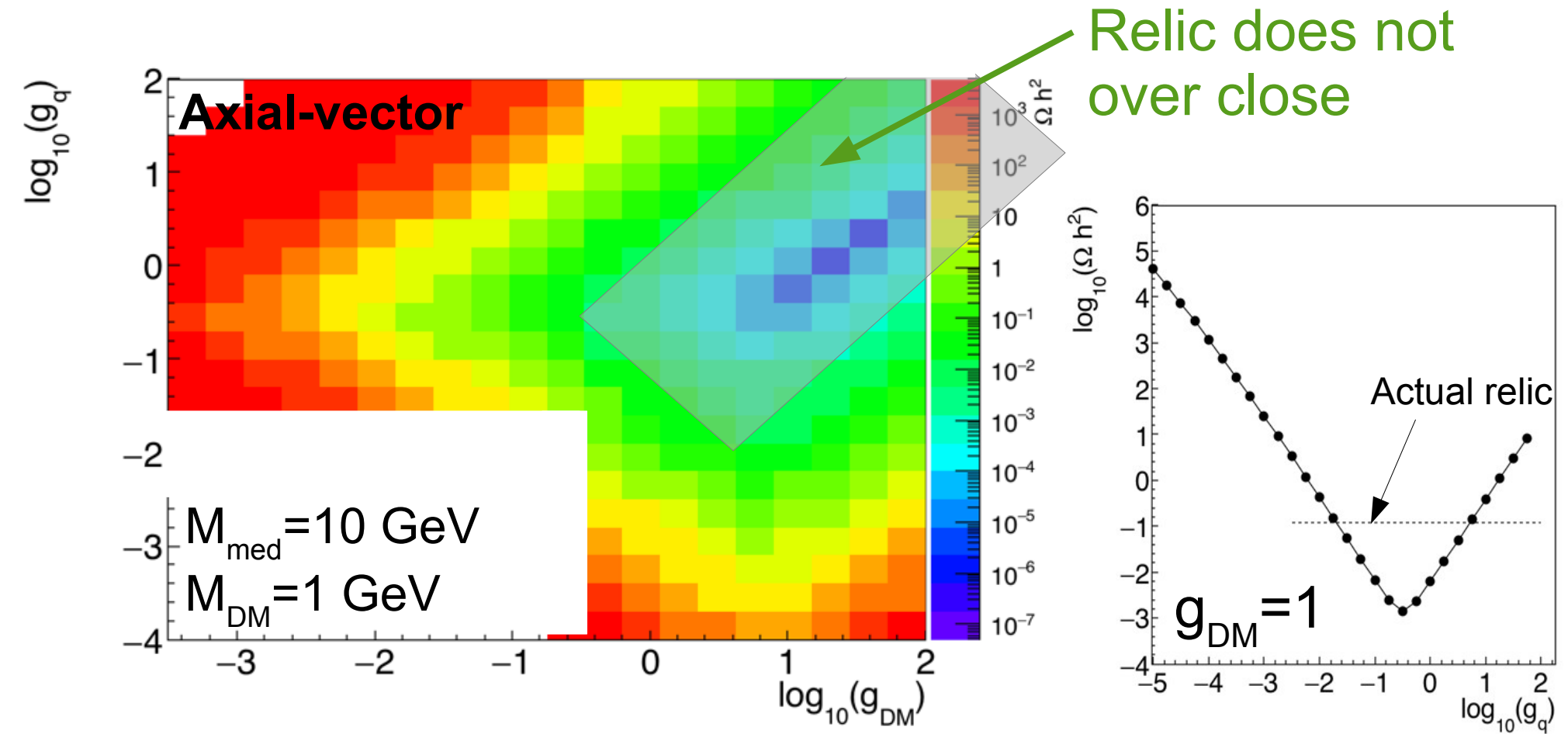
- We can numerically solve for the solution



For smaller mediators allowed solution

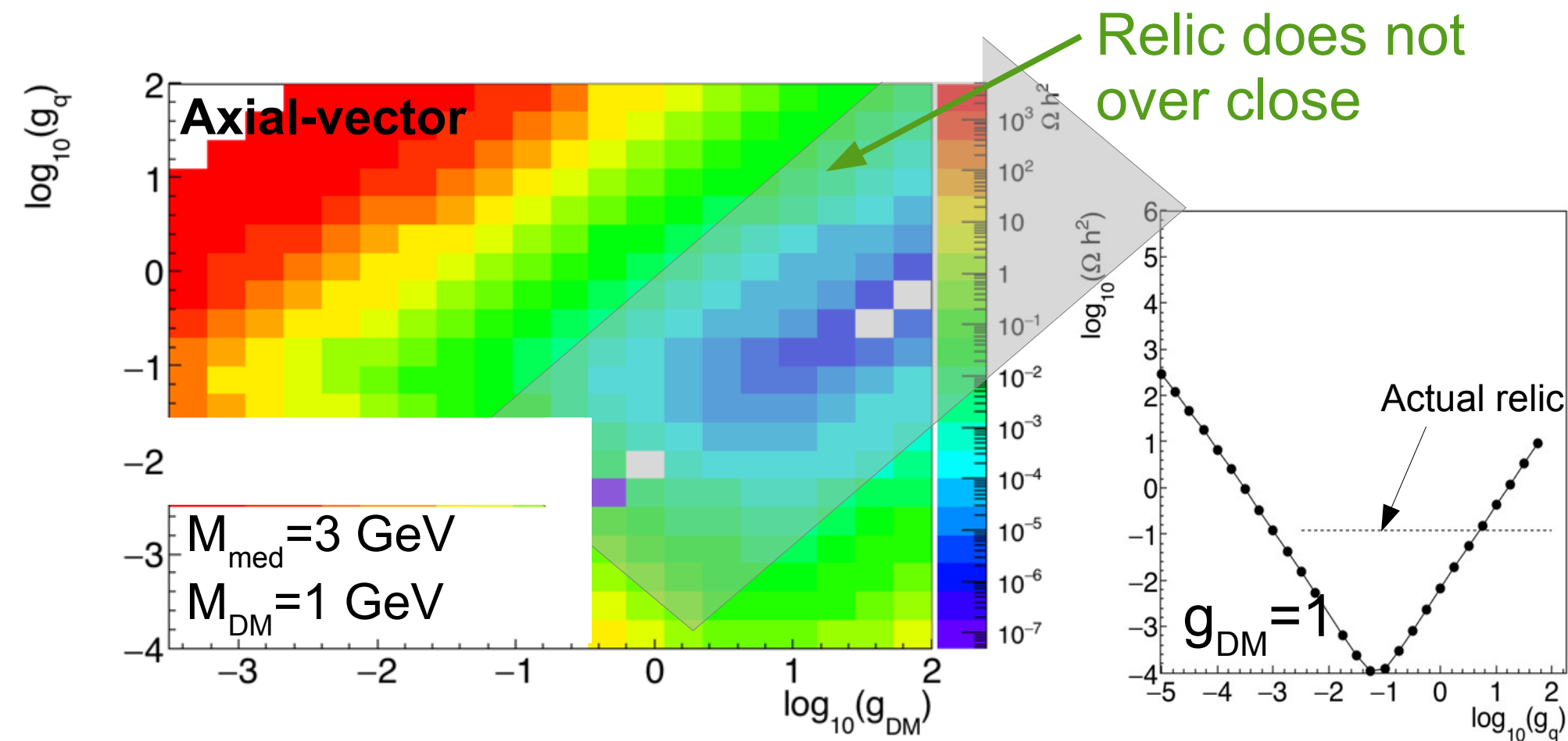
Solving for the relic

- We can numerically solve for the solution



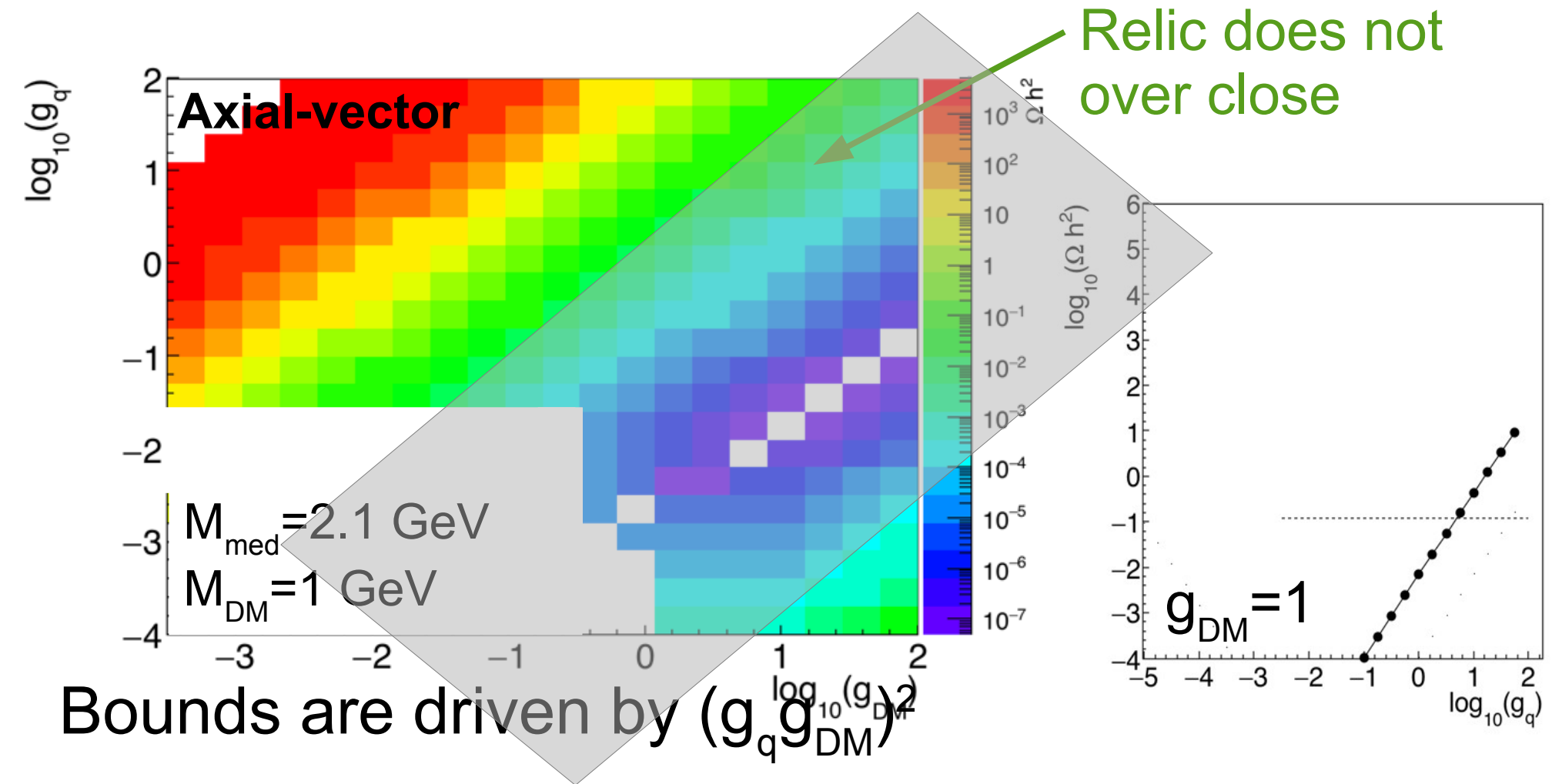
Solving for the relic

- We can numerically solve for the solution



Solving for the relic

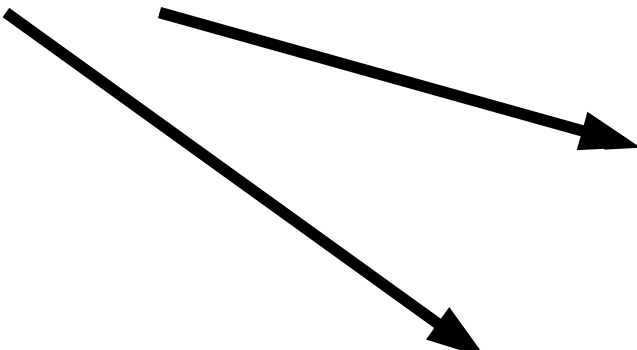
- We can numerically solve for the solution



A common theme of DM talks

- For a constant value of:

$$-(g_q g_{\text{DM}})^2 = C$$

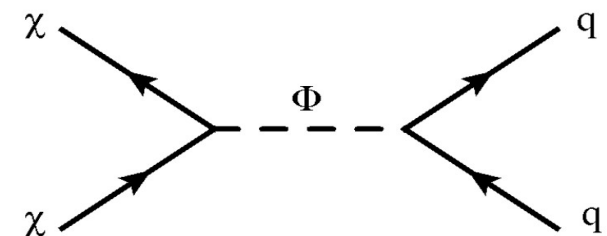
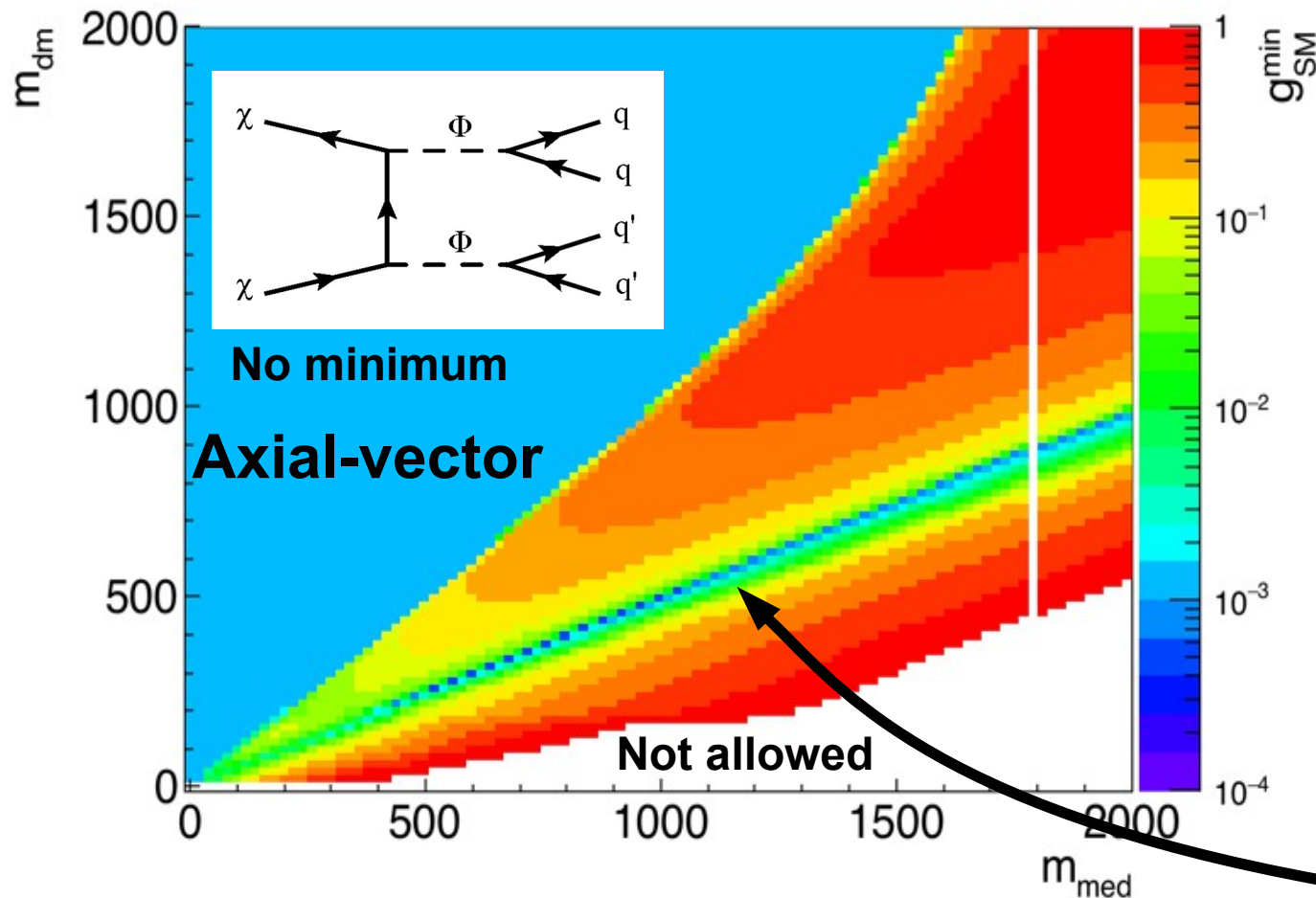


Set this to be **large**
still get right relic

Set this to be **small**
Weak coupling with the SM

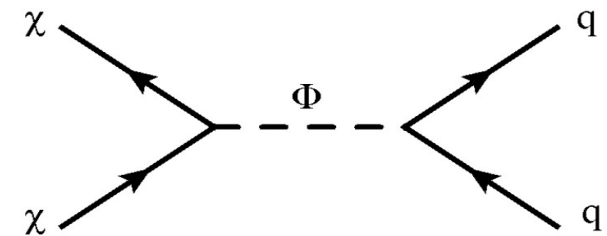
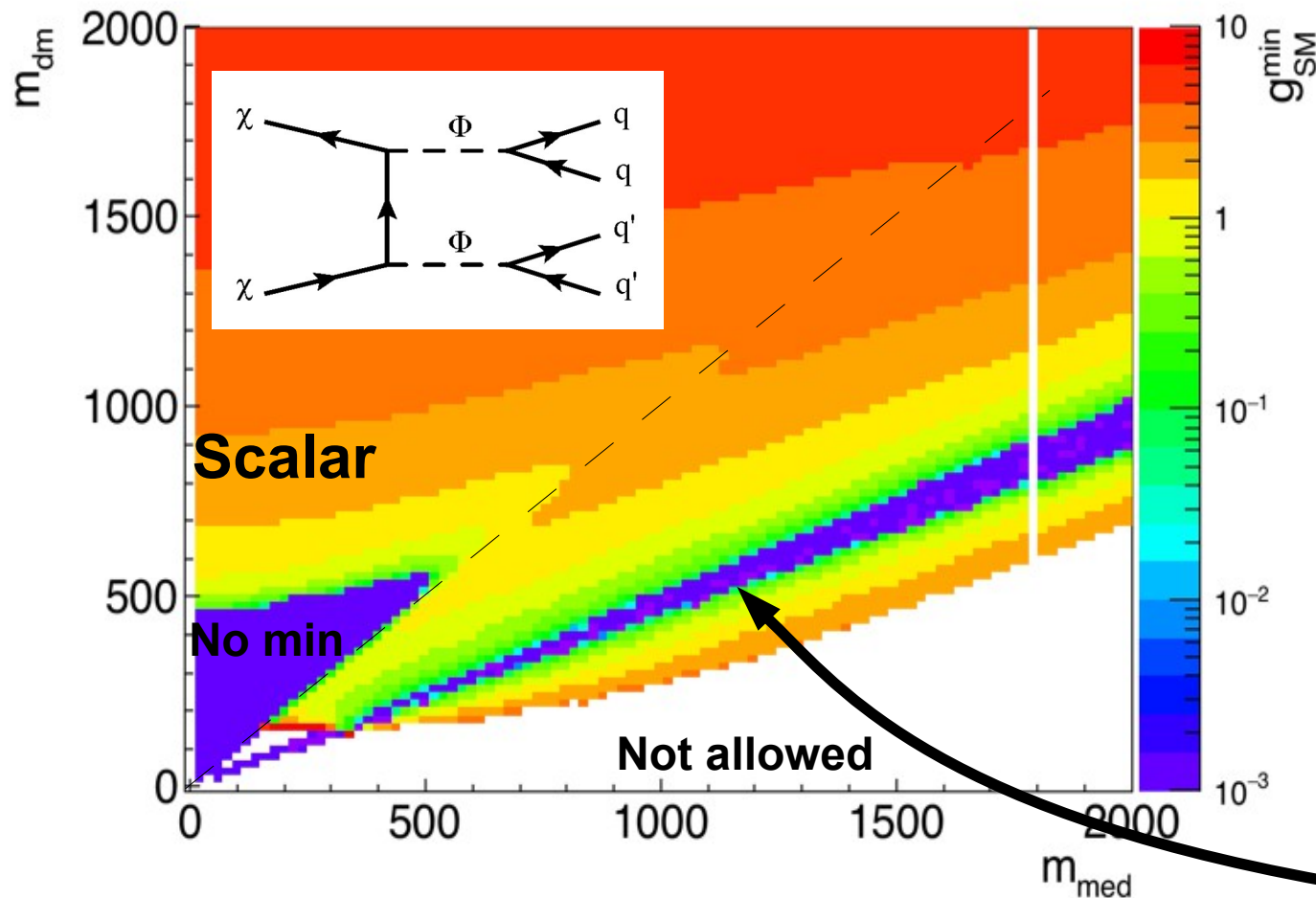
Min couplings

- Can split the solution to the max and min coupling
 - In this case we fix $g_{\text{DM}}=1$ (product $g_q g_{\text{DM}}$ defines bound)



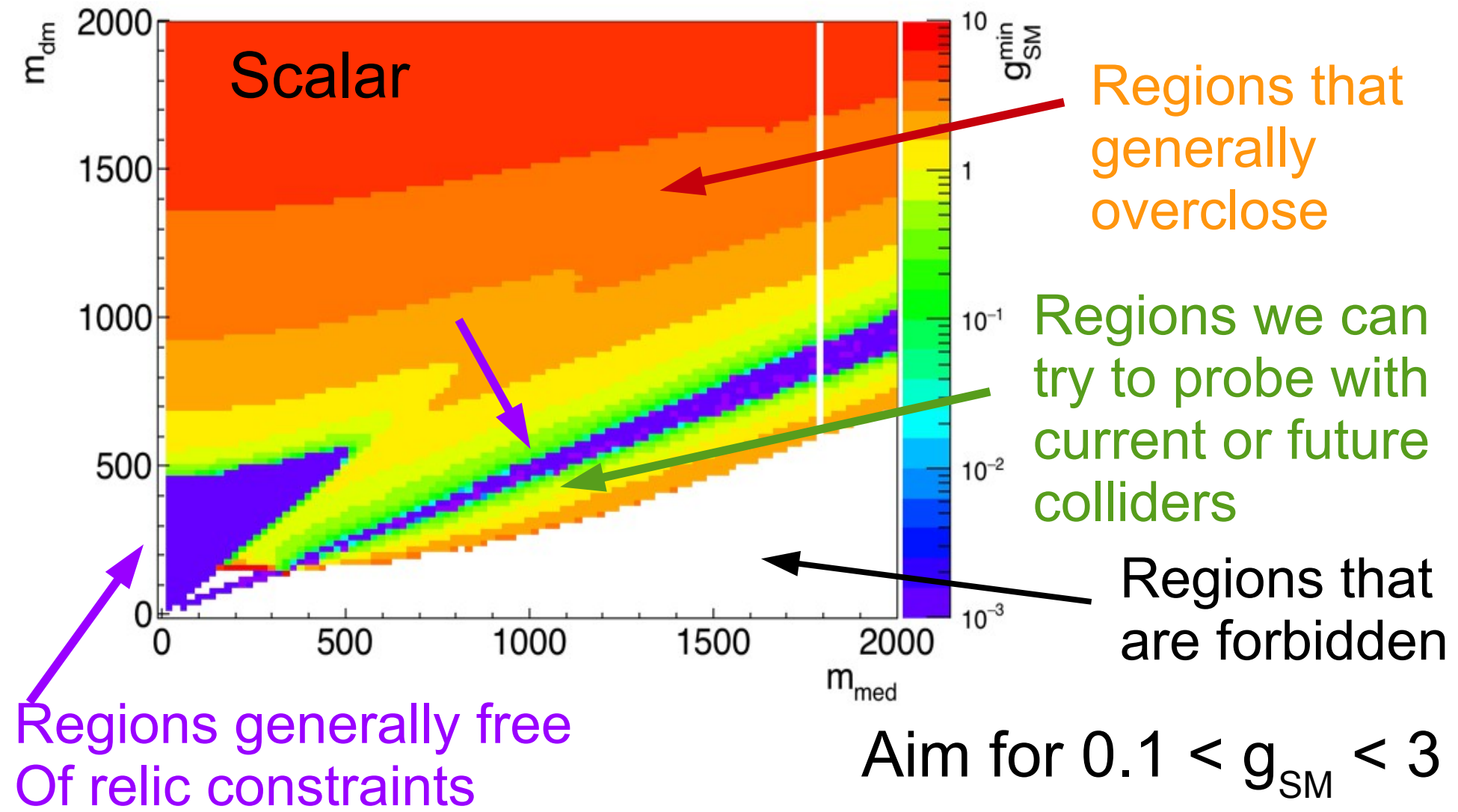
Min couplings

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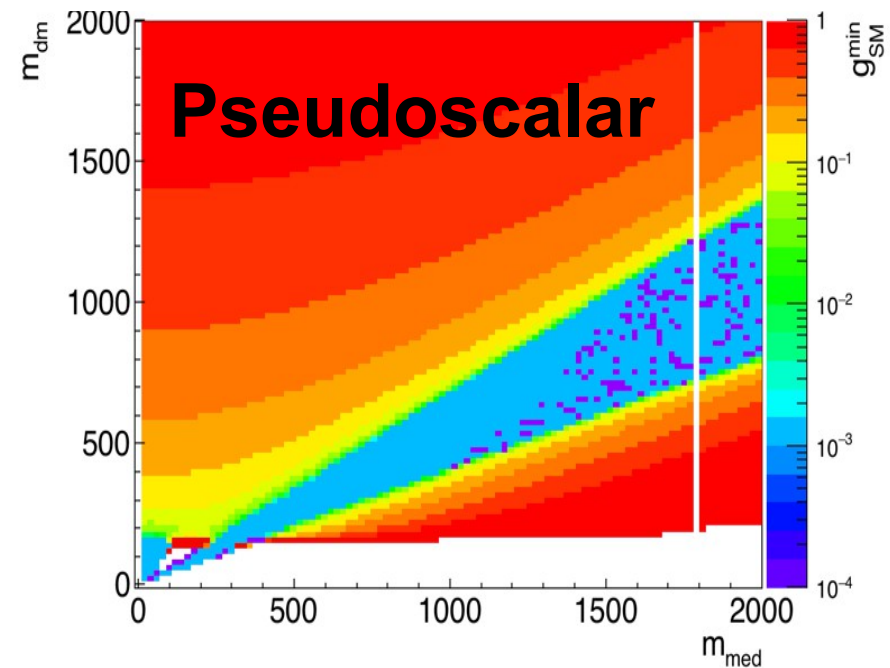
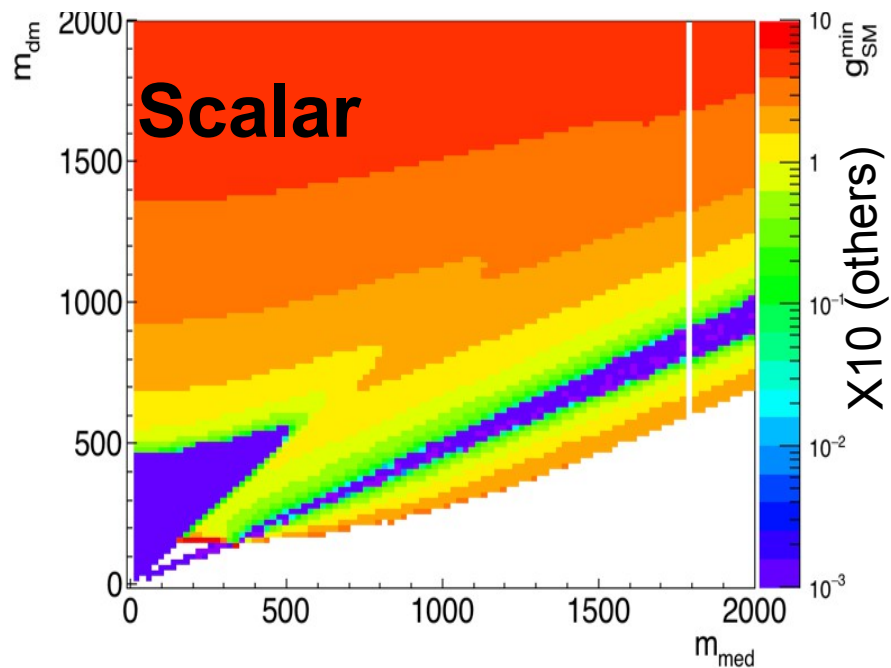
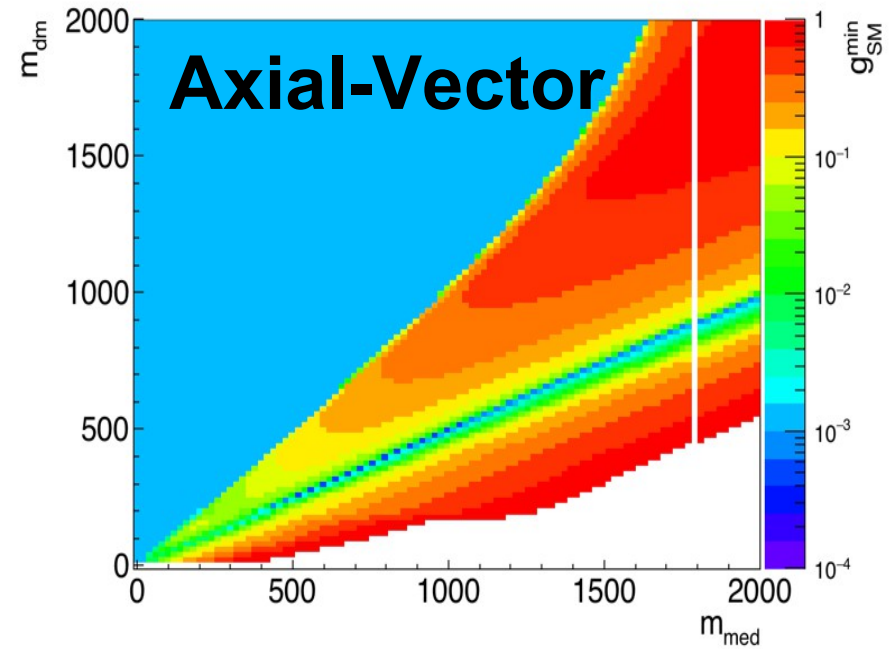
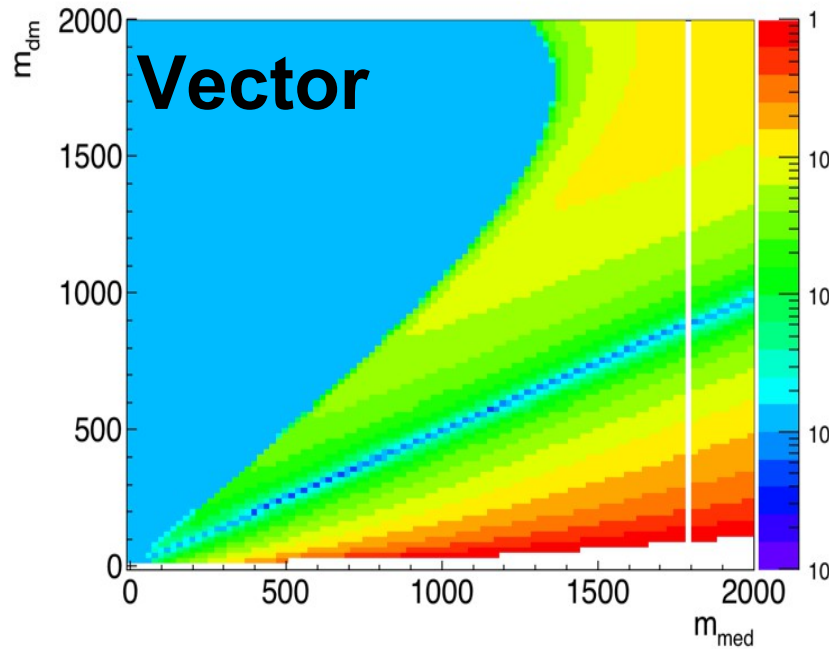


What do we conclude?

- What is driving the results is the coupling

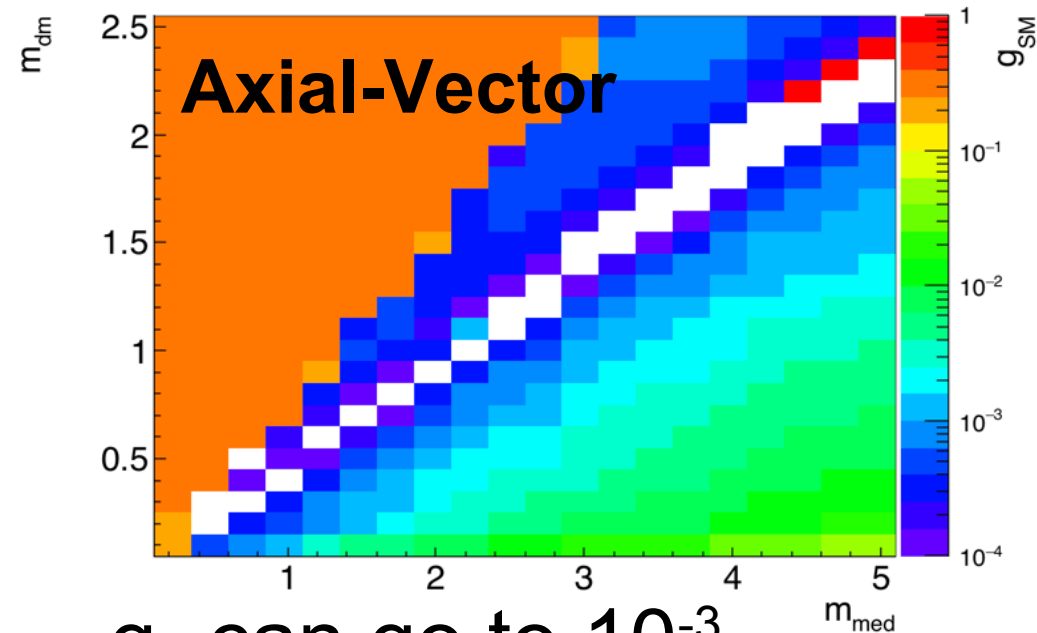
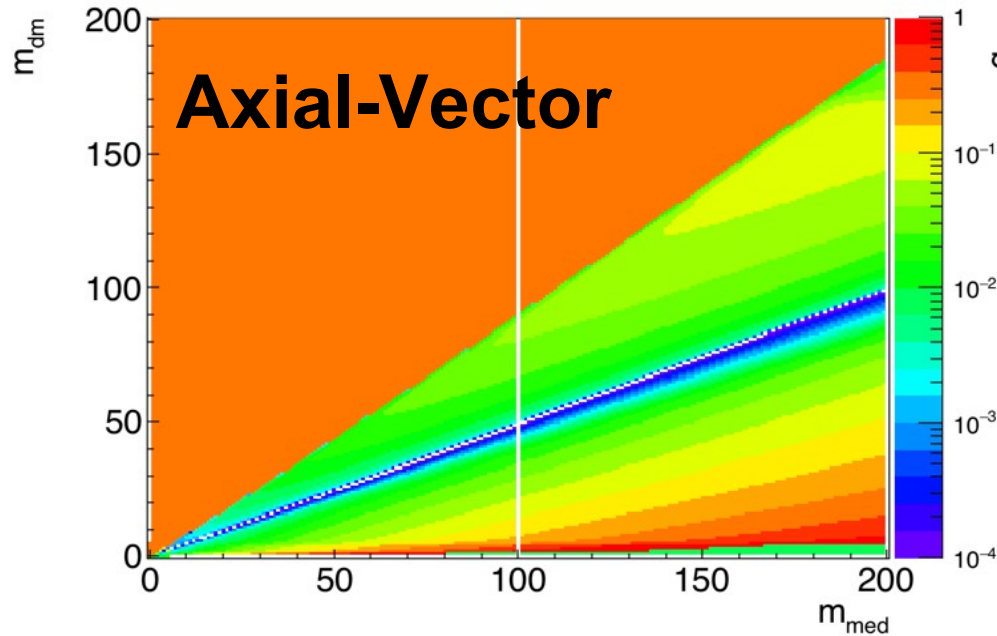


Min Couplings for all²²



What about at low mass?

- Coupling becomes a real challenge



g_q can go to 10^{-3}

At low masses we can have very small couplings

However we have more strategies

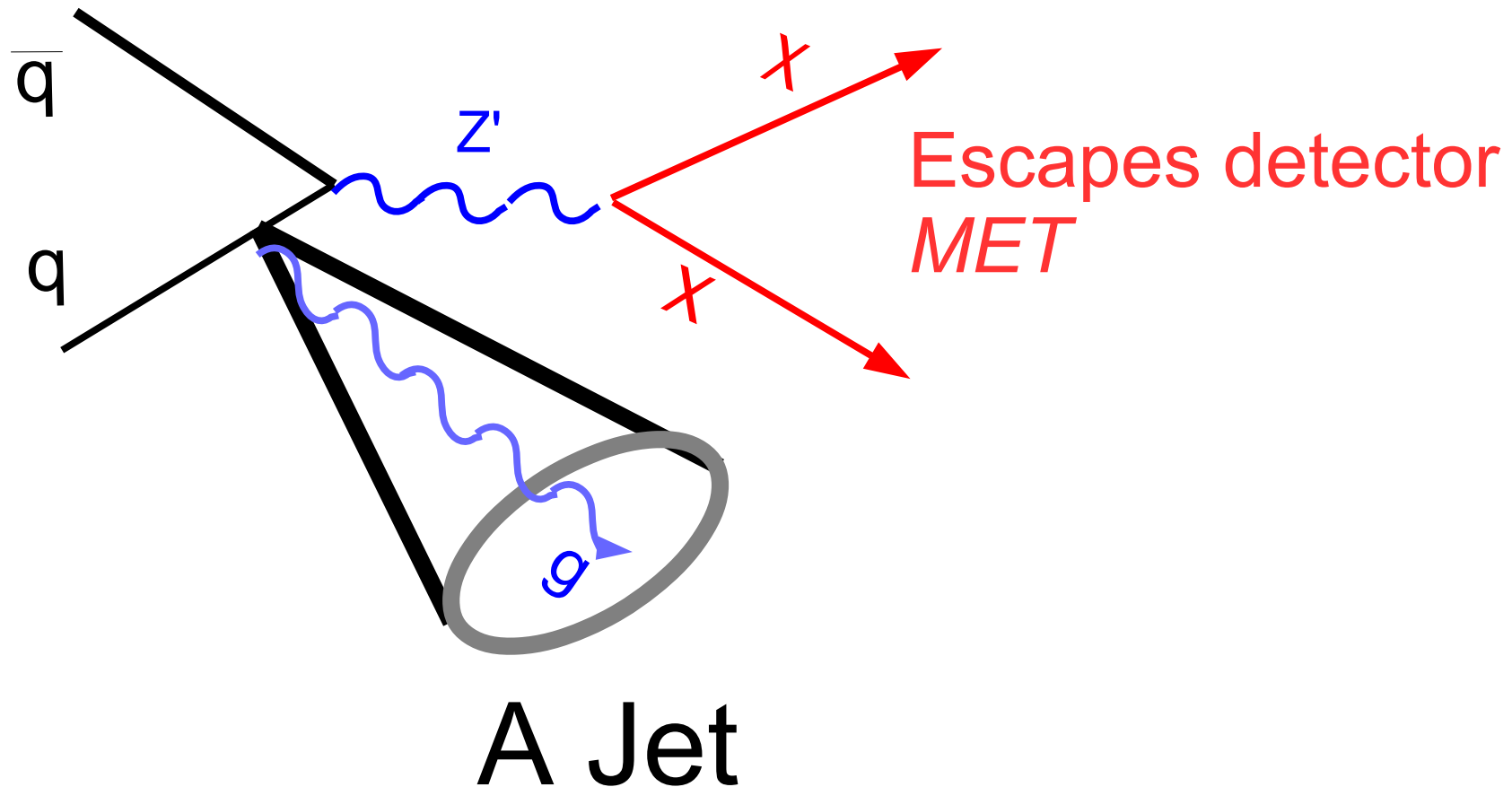
Summary Benchmarks

- Spin 1 :
 - Aim to probe couplings down 0.01 for $m_{\text{Med}} > 100 \text{ GeV}$
 - For $10 < m_{\text{Med}} < 100 \text{ GeV}$ aim to probe down to 10^{-3}
 - For $m_{\text{Med}} < 10 \text{ GeV}$ aim to probe coupling to 10^{-4}
- Spin 0 :
 - Aim to probe couplings down 0.1 for $m_{\text{Med}} > 300 \text{ GeV}$
 - Try to cover $m_{\text{Med}} < 300$ by any means possible
- Covers most of the phase space

Collider Searches

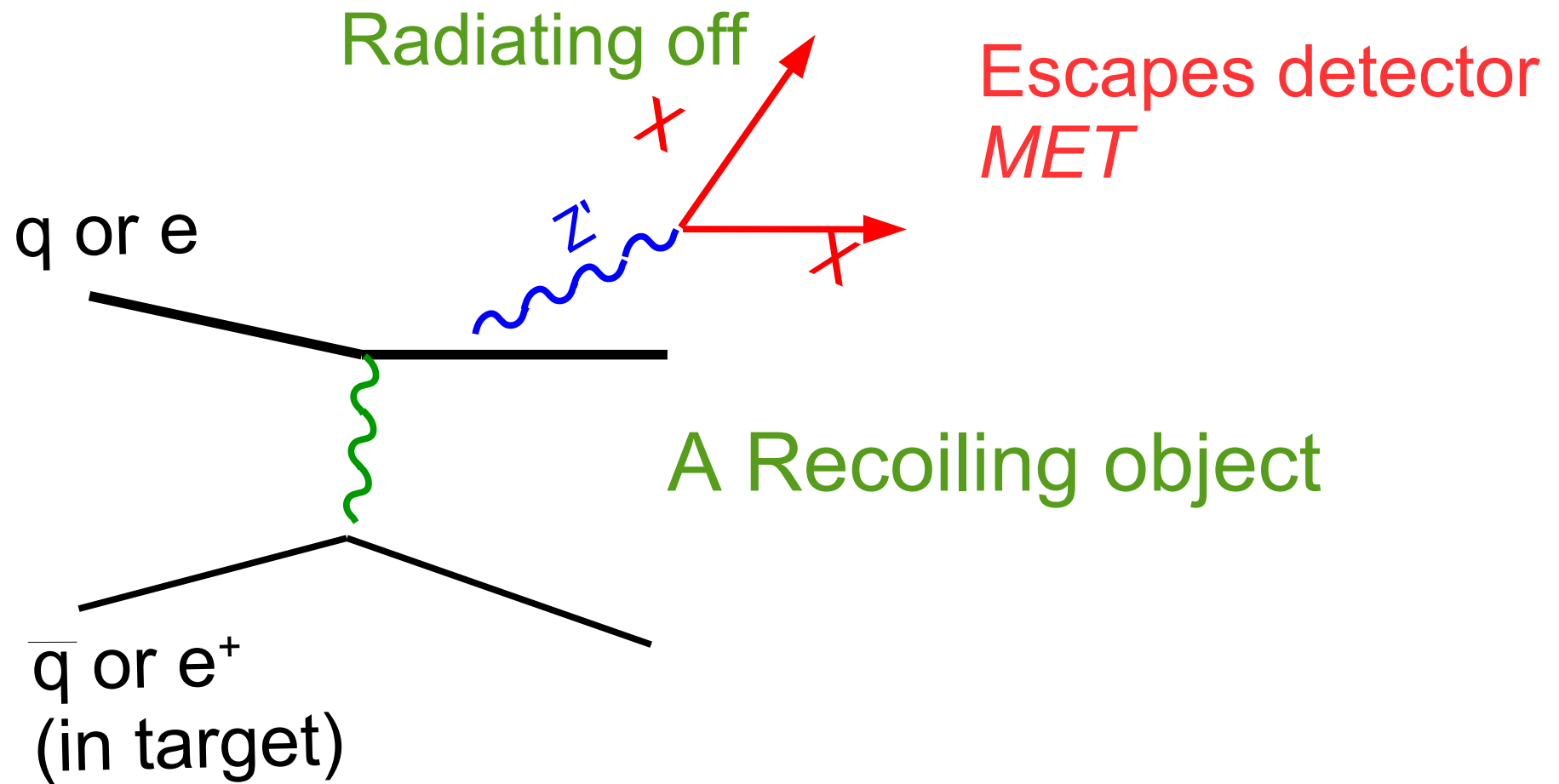
The Basic Monojet Search

Escaping detector gives us signatures of *MET*



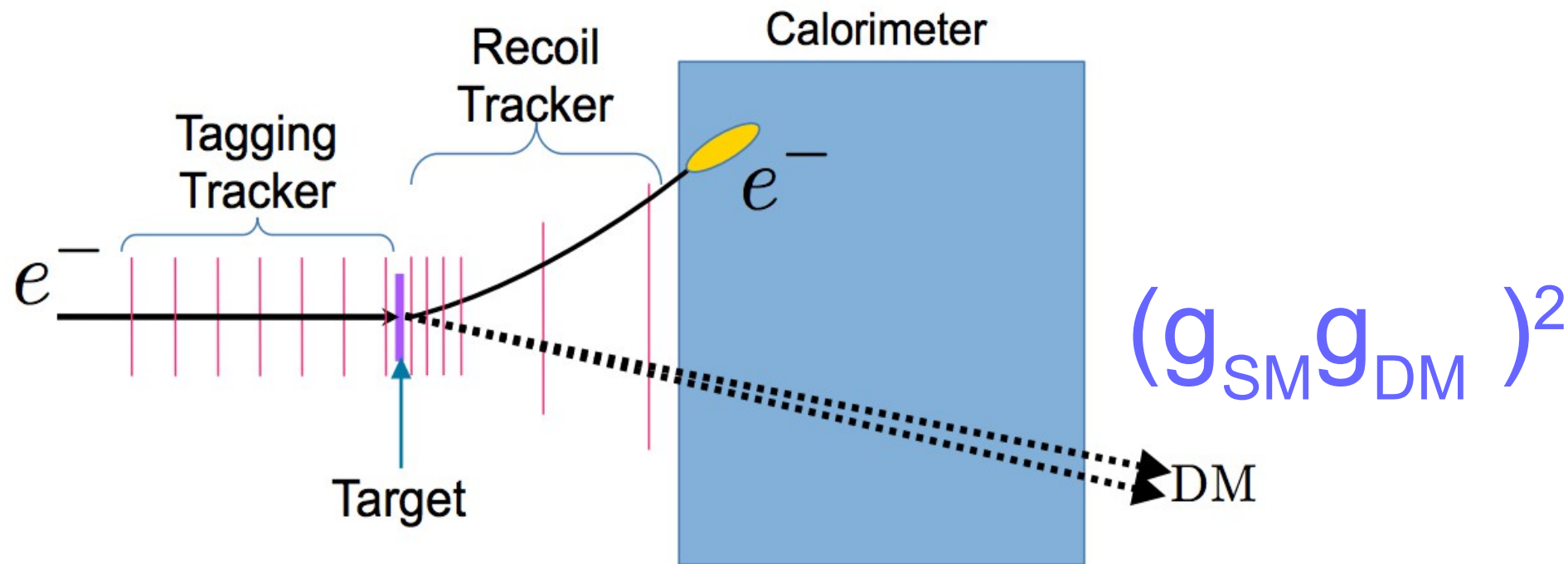
At Low masses

We still get *MET signatures but with other stuff*



@low mass aim is for low couplings

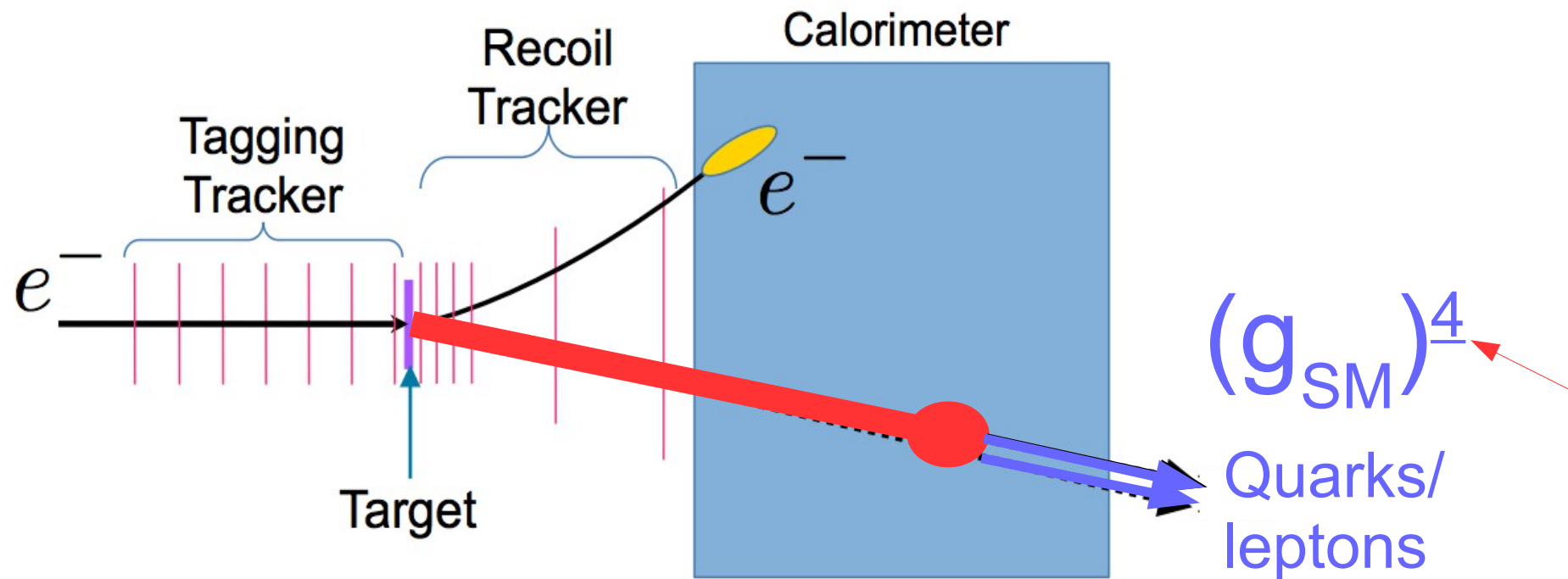
To probe small couplings you need very high rates
 This is something we can do with a beam dump



Now many experiments being made or proposed

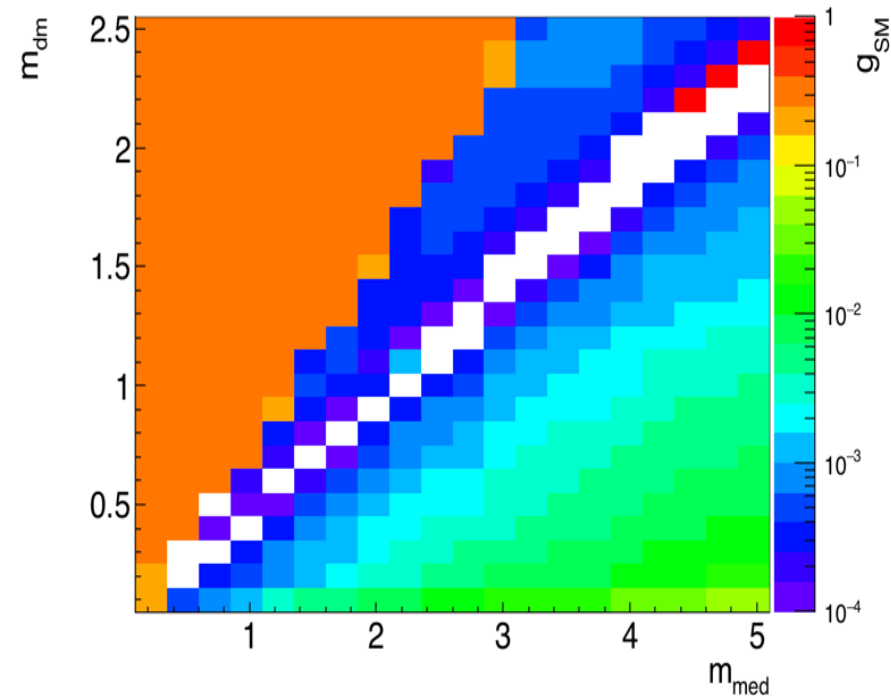
@low mass aim is for low couplings

To probe small couplings you need very high rates
 This is something we can do with a beam dump



Alternative approach is to look for decays from dm
 Suppressed if SM coupling is small

@low mass aim is for low couplings

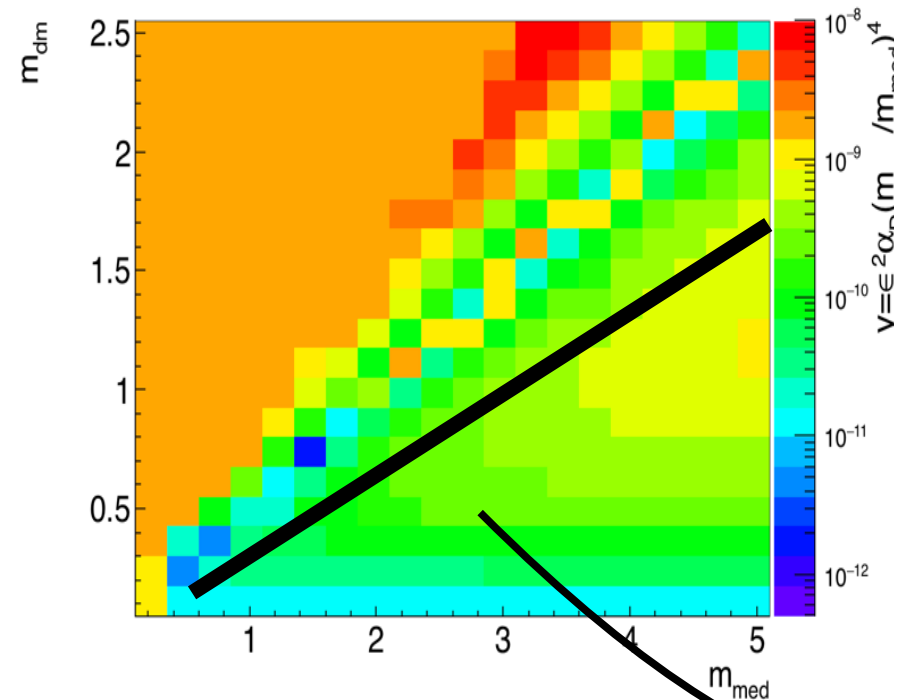


At low mass we want to probe the $g_{SM}=10^{-4}$ region

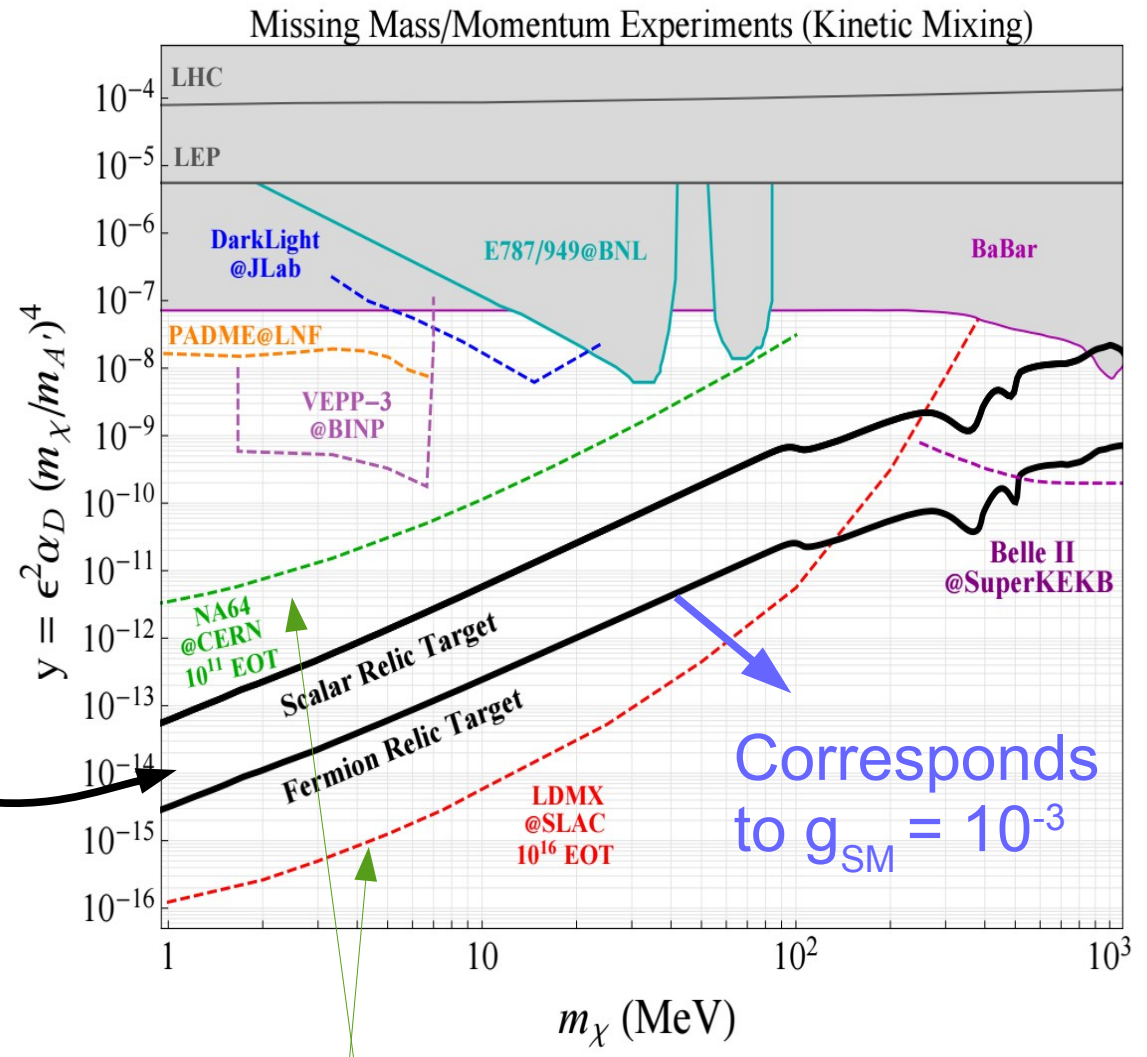
However, just to confuse everybody

Coupling units are changed : $\epsilon^2 \alpha_D (m_{DM}/m_{MED})^4$

@low mass aim is for low couplings



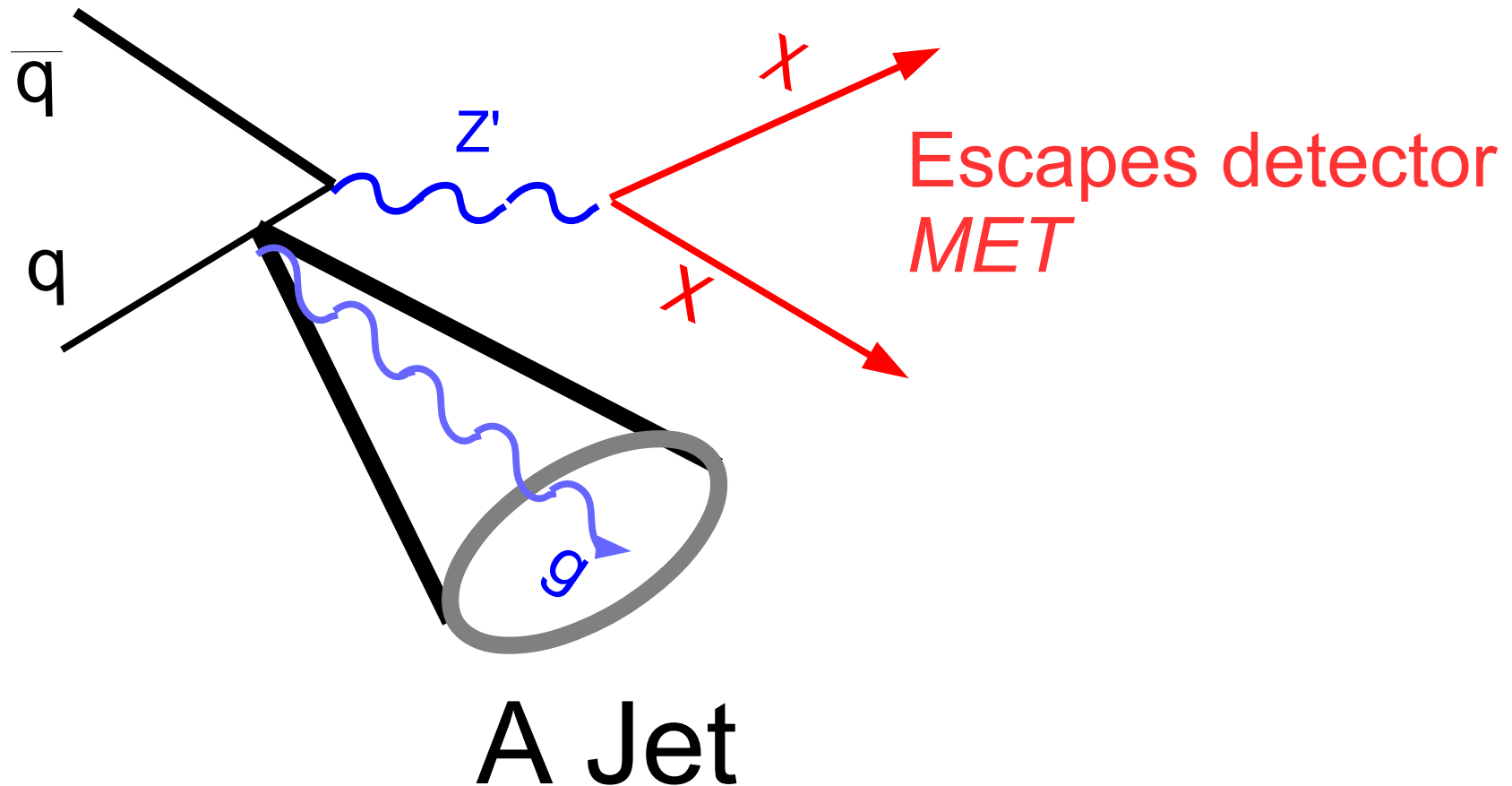
Consider a diagonal projection $m_{DM}/m_{MED} = 1/3$



Reach of a few proposed beam dump experiments
Can probe the interesting region

The Basic Monojet Search

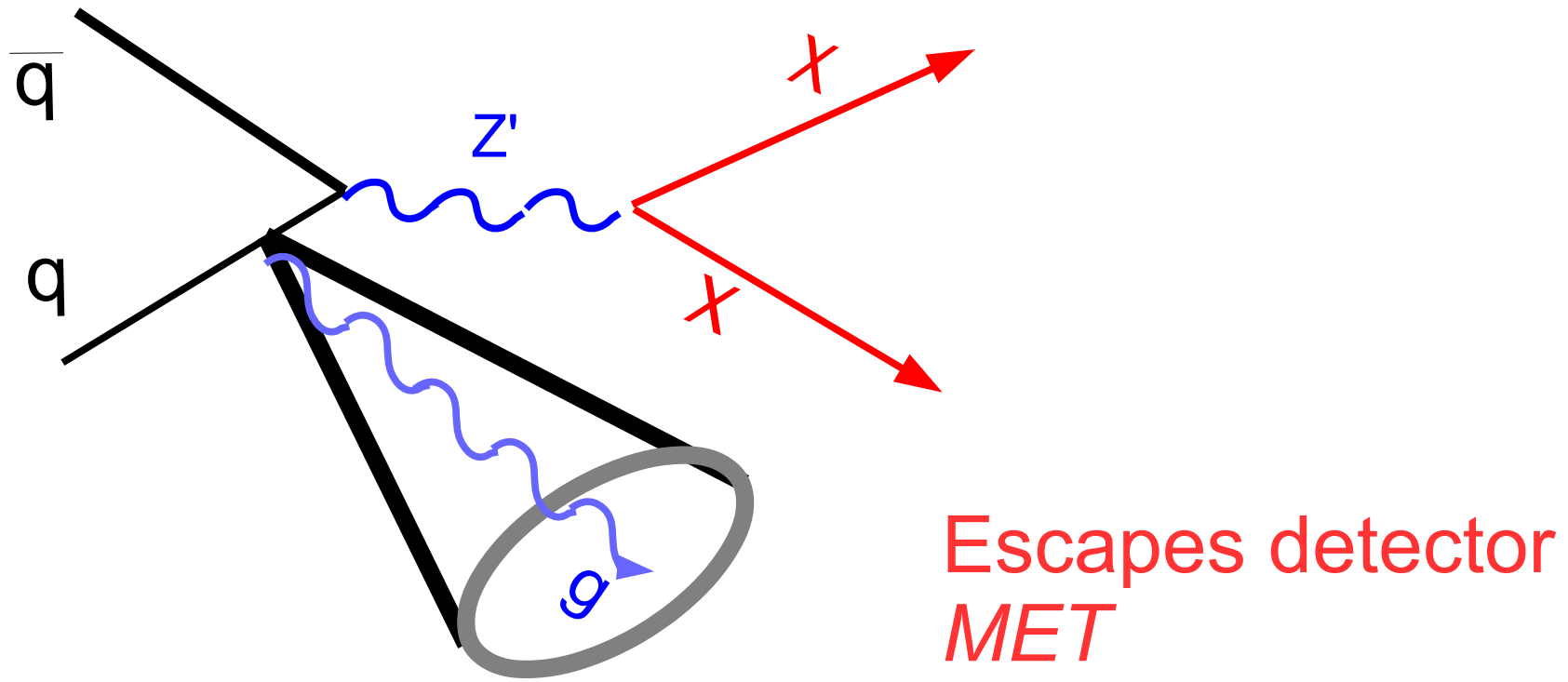
Escaping detector gives us signatures of *MET*





Monojet search
Straddling SM and BSM

Searching for *MET*

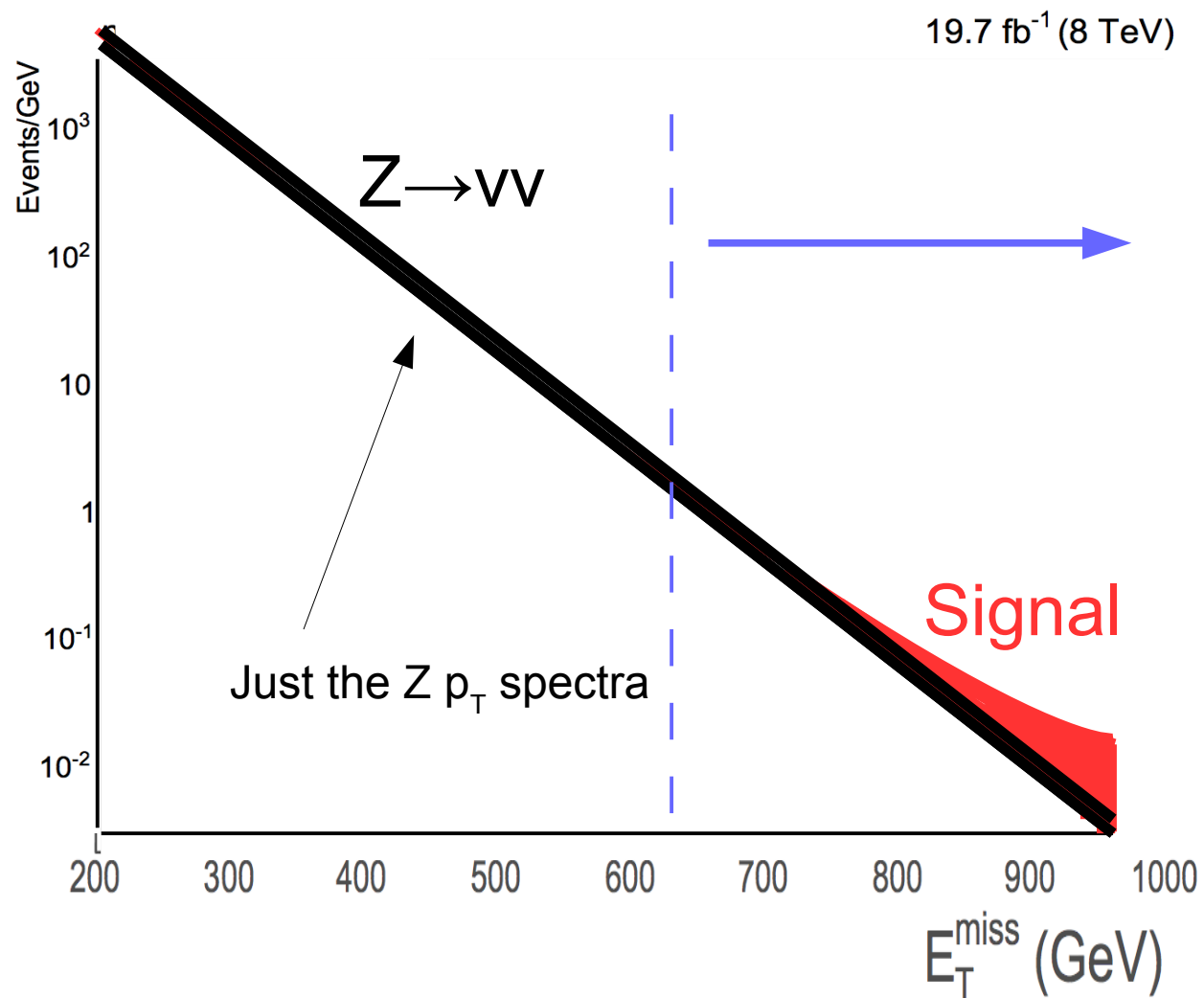


“To find nothing you have to reconstruct everything”[1]

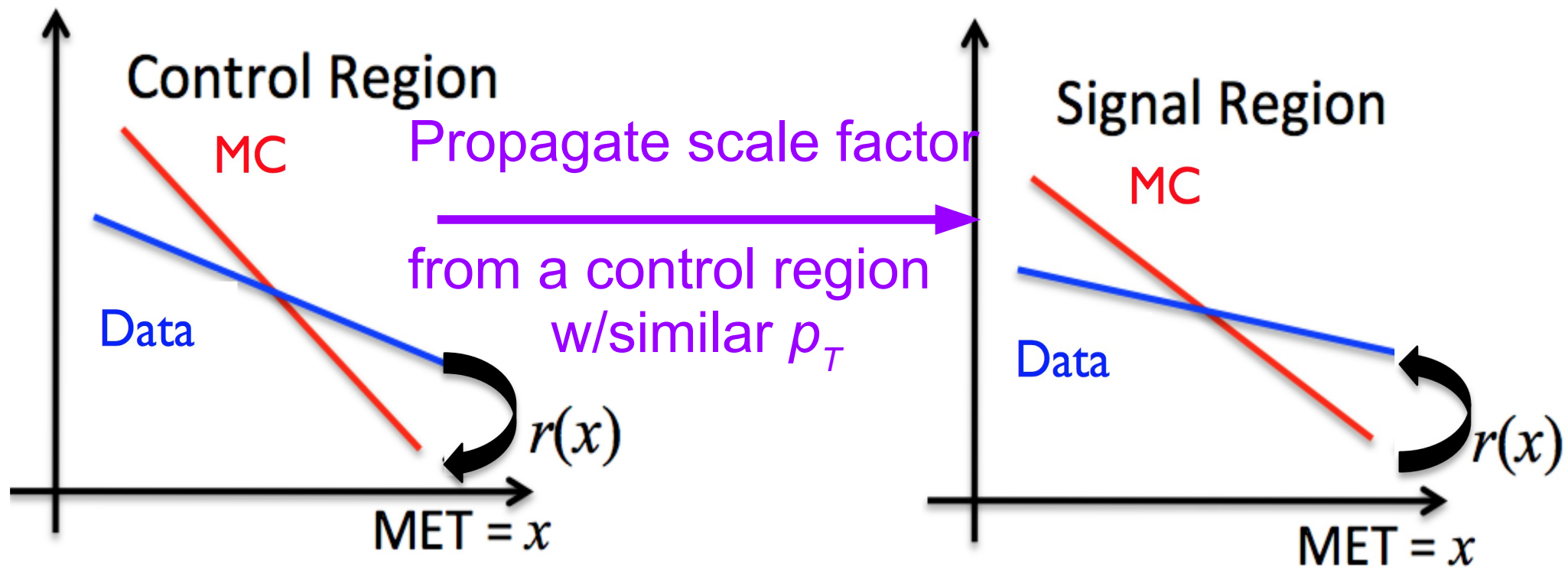
$$-\sum_{\text{All particles}} p_T = \textcolor{red}{MET} \text{ } (E_T^{\text{Miss}})$$

$$\text{-Boson } p_T = \textcolor{red}{MET} \text{ } (E_T^{\text{Miss}})$$

How do we search?



Strategy to fix agreement



Control: another decay of a Z boson

~~$Z \rightarrow \mu\mu$~~ $\rightarrow$ $Z \rightarrow \nu\nu$
 Remove

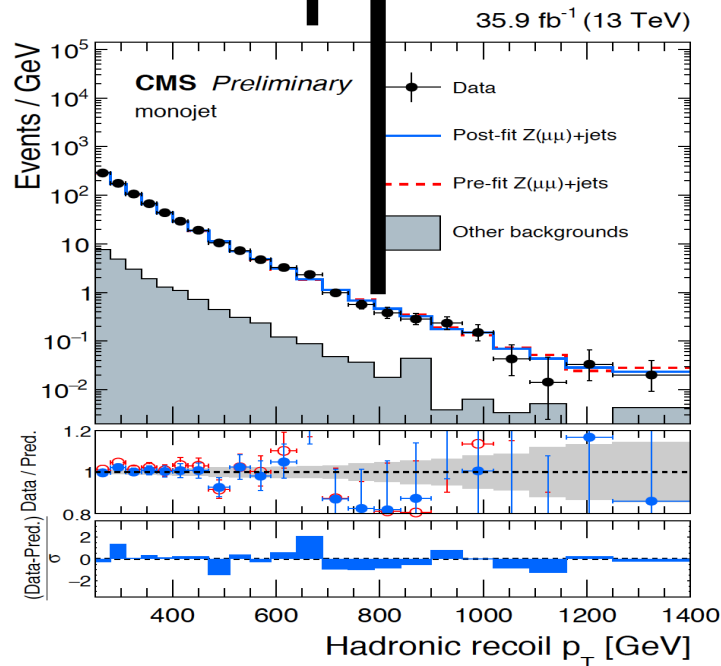
hadronic recoil : Transverse sum of all particles in event excluding leptons/photons

CMS-EXO-16-037
 CMS-EXO-16-010
 CMS-EXO-12-055

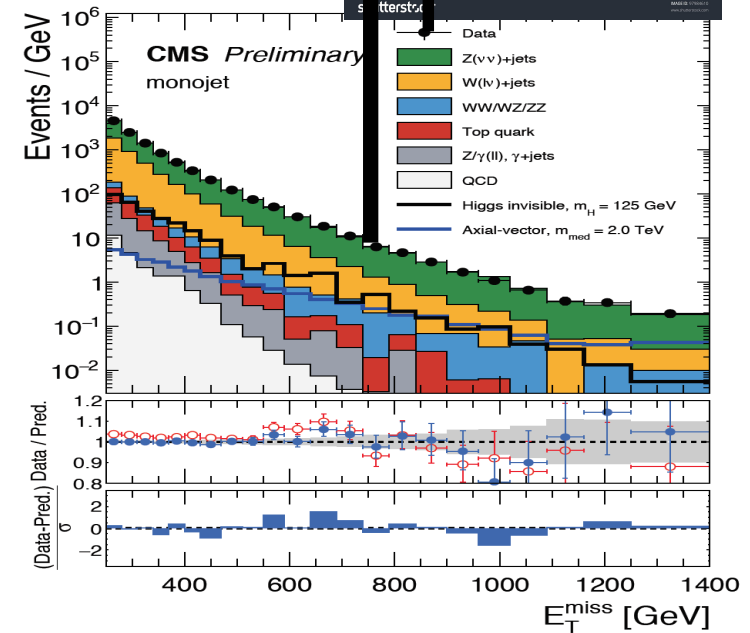
What is the transfer factor?

Propagate the data/MC agreement of the hadronic recoil
From a control region to a signal region

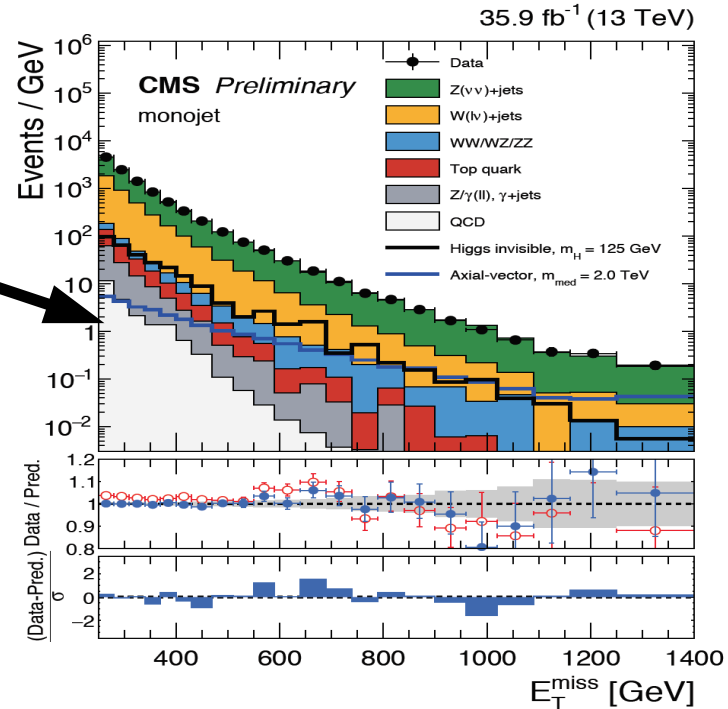
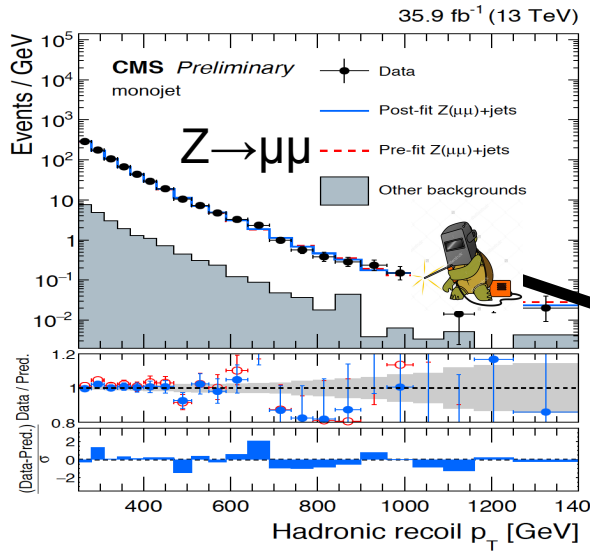
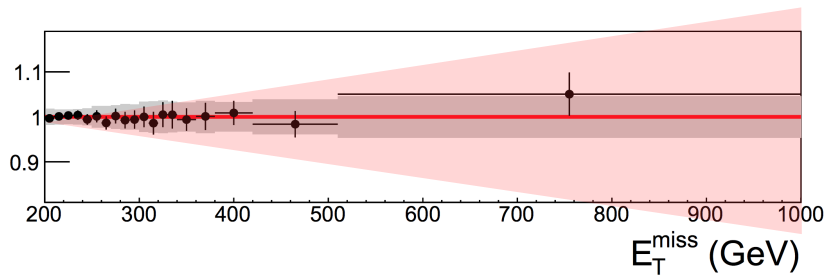
Control



Signal



MC/data

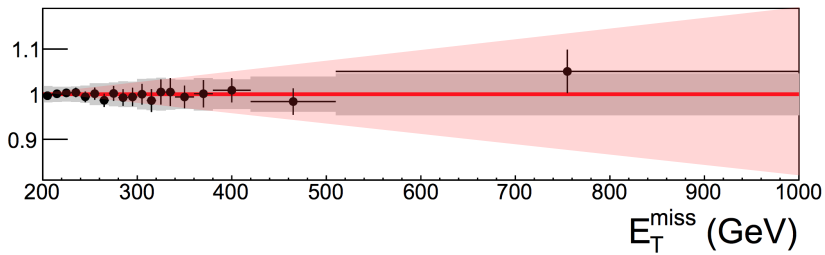


Control regions have less events than signal

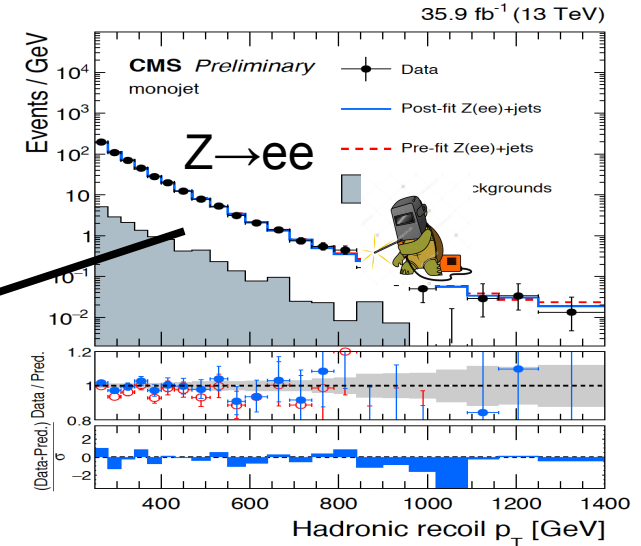
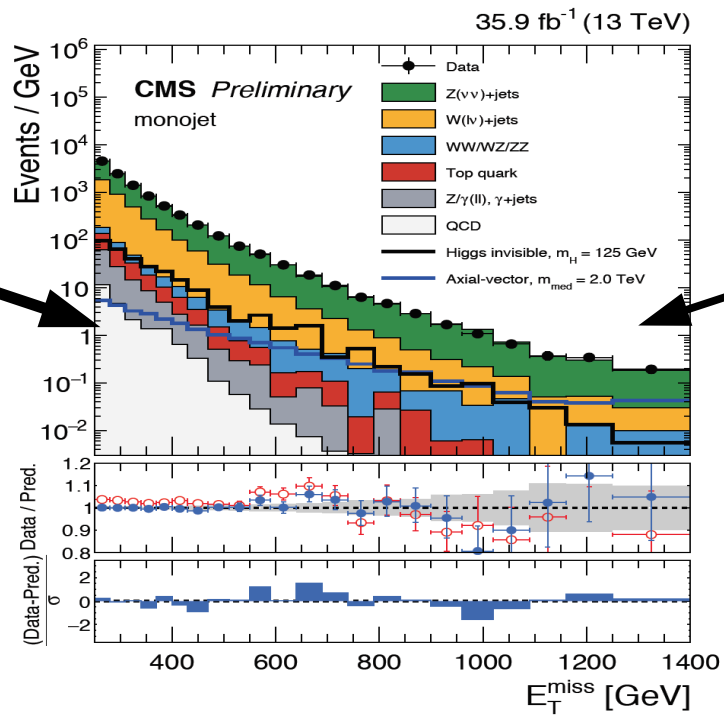
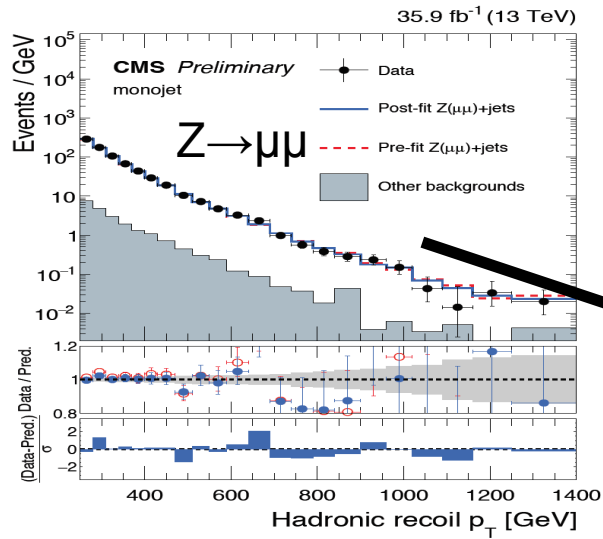
$$\sigma_{\mu\mu} = 0.1 \sigma_{\nu\nu}$$

Statistical precision is 4x worse
Not good enough!

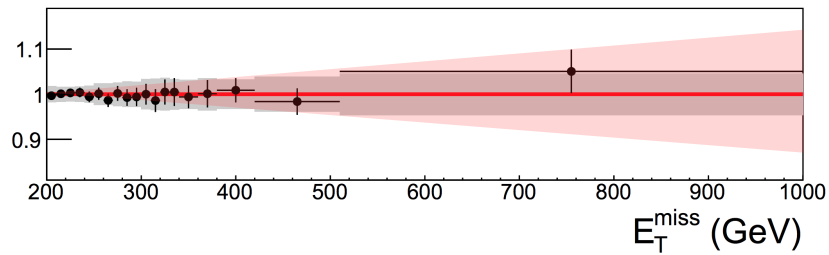
MC/data



2 Control regions
60% uncertainty @ 1 TeV

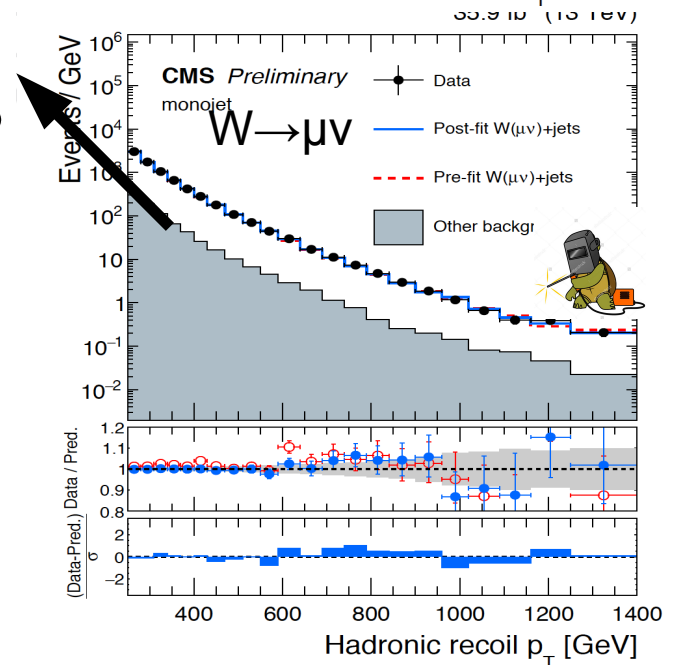
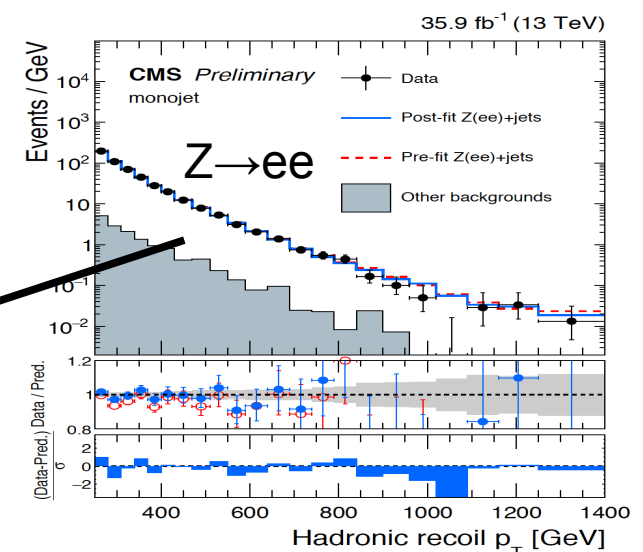
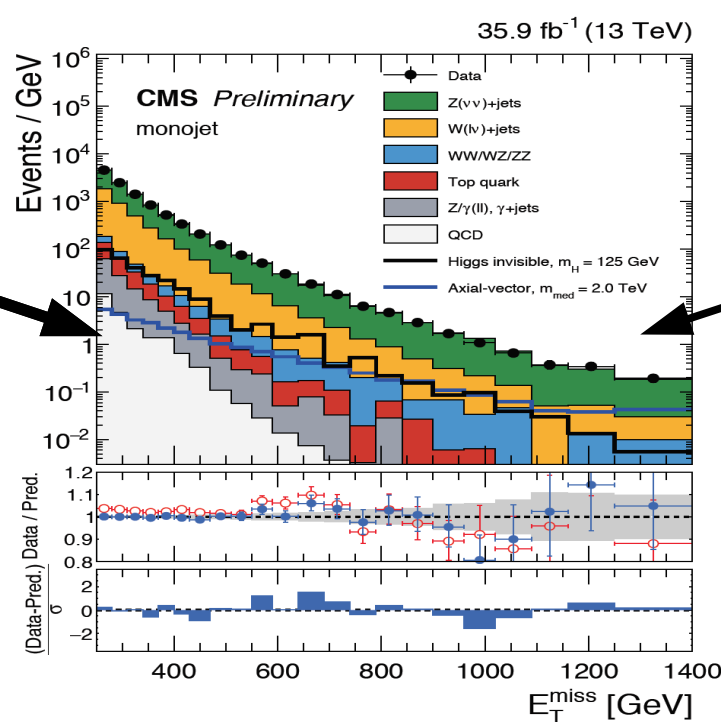
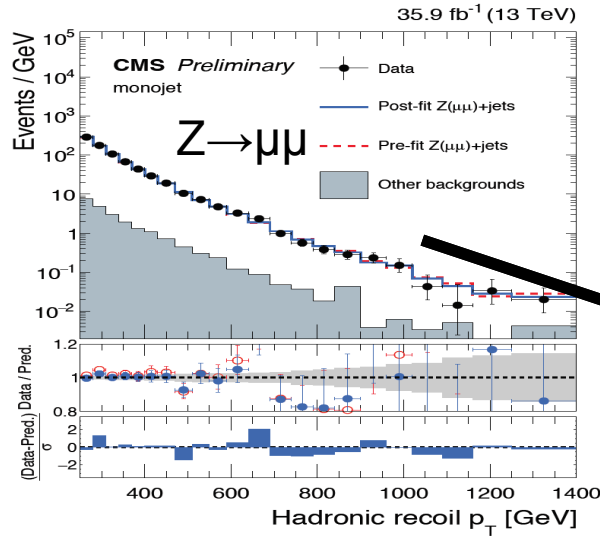


MC/data

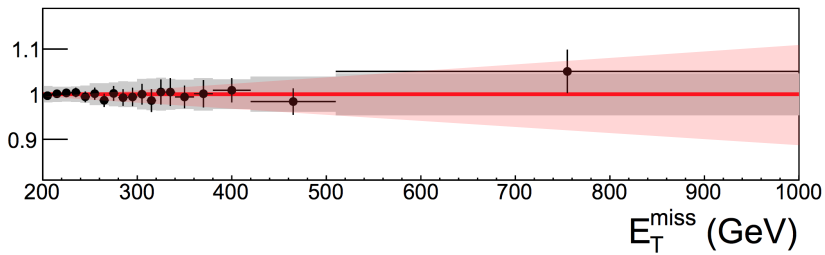


3 Control regions

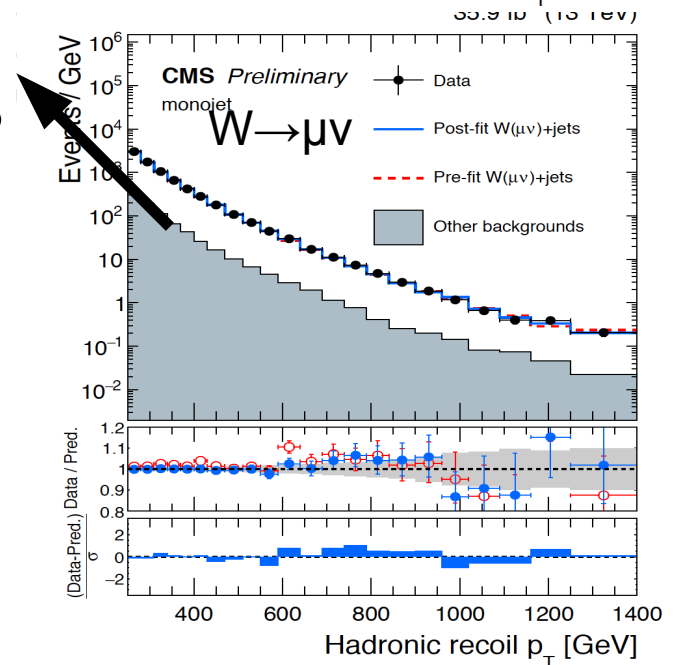
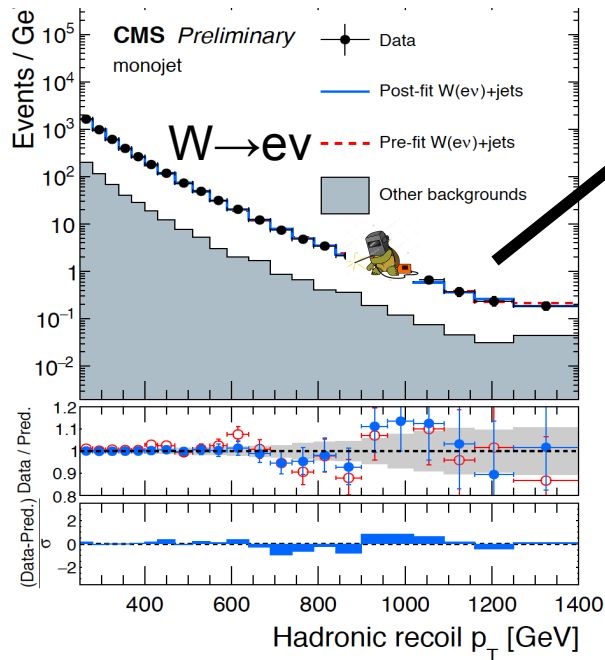
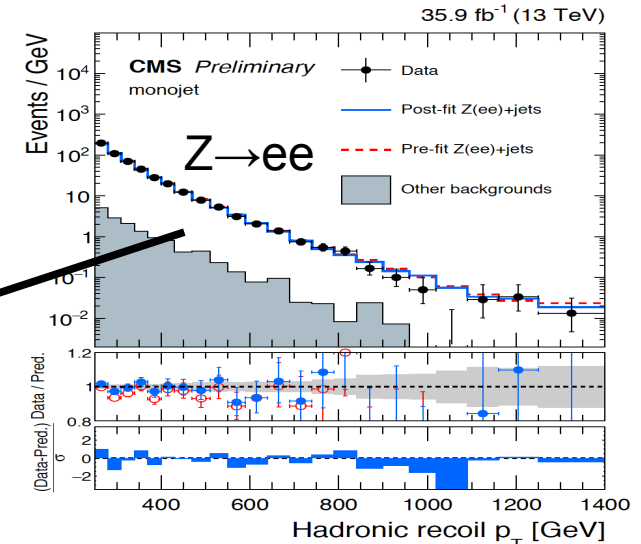
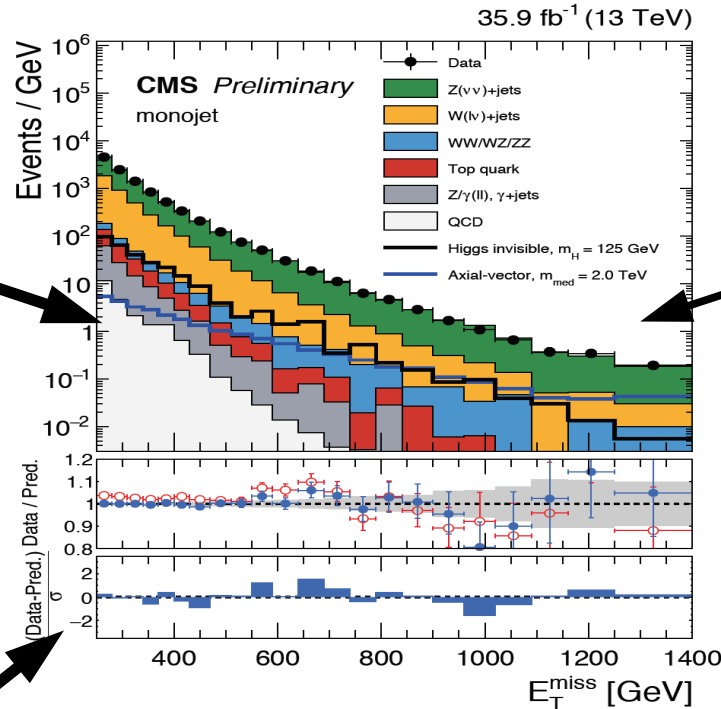
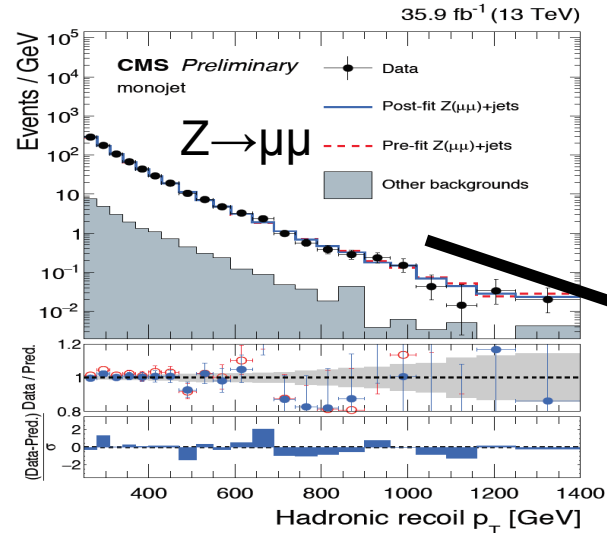
40% uncertainty @ 1 TeV



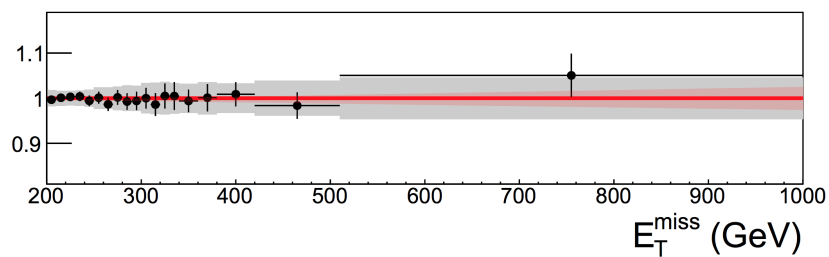
MC/data



4 Control regions
30% uncertainty @ 1 TeV

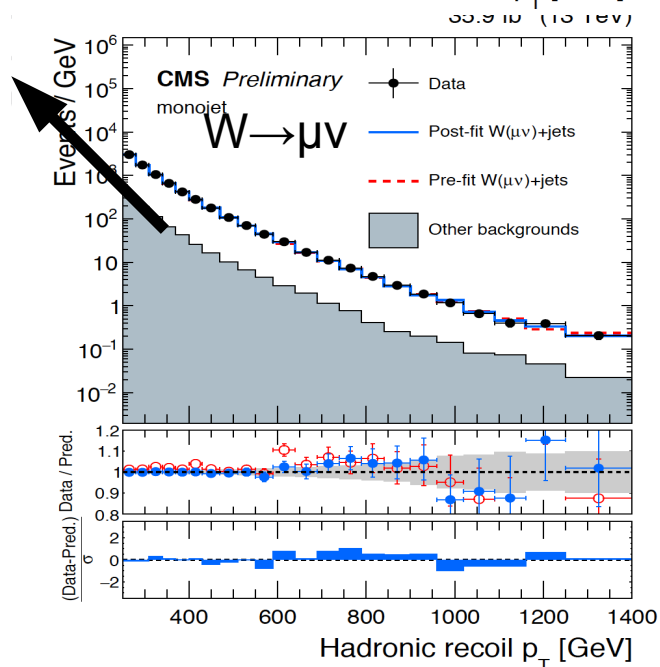
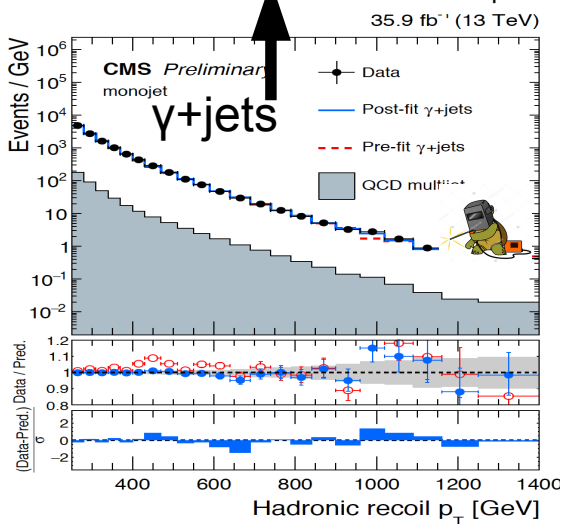
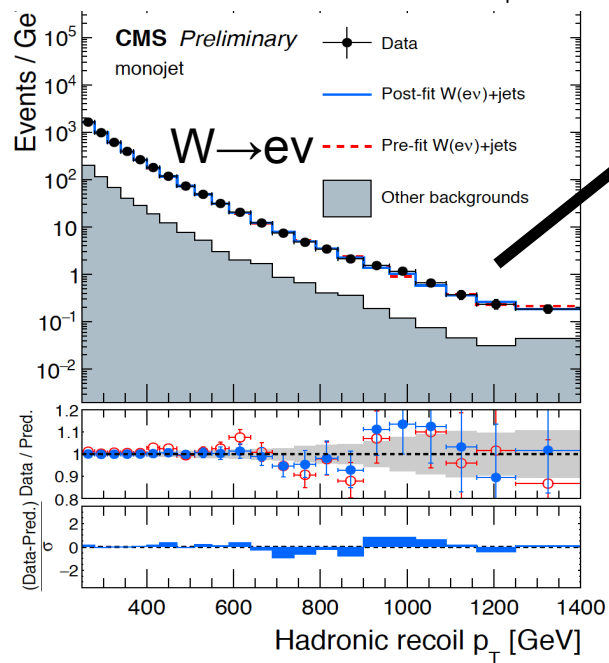
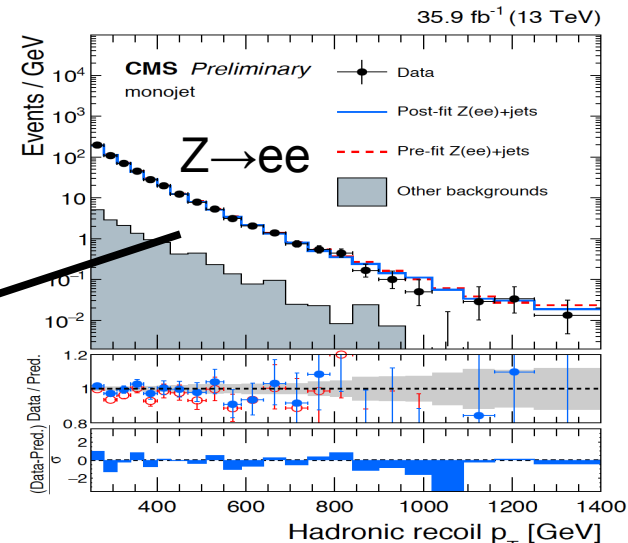
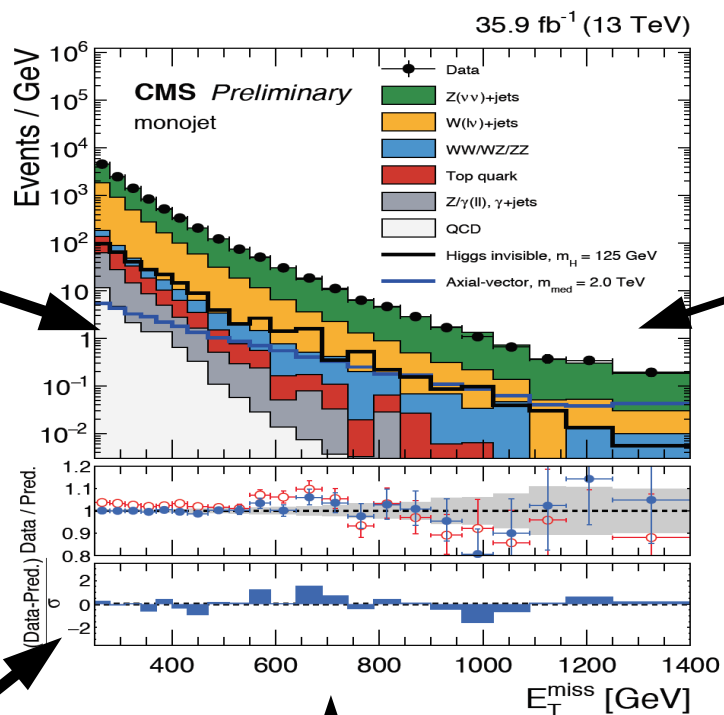
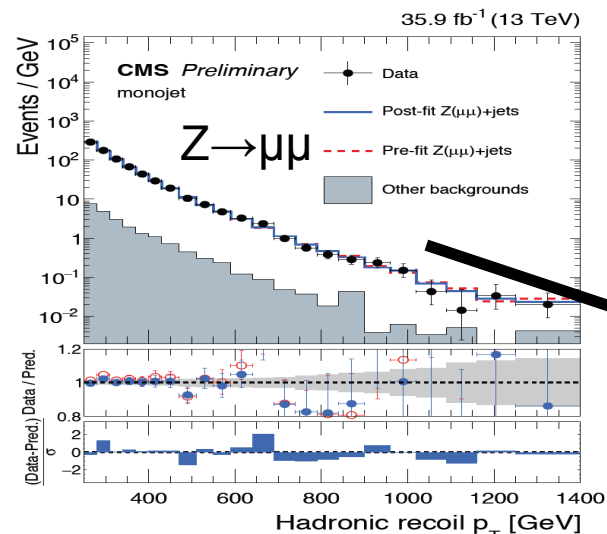


MC/data



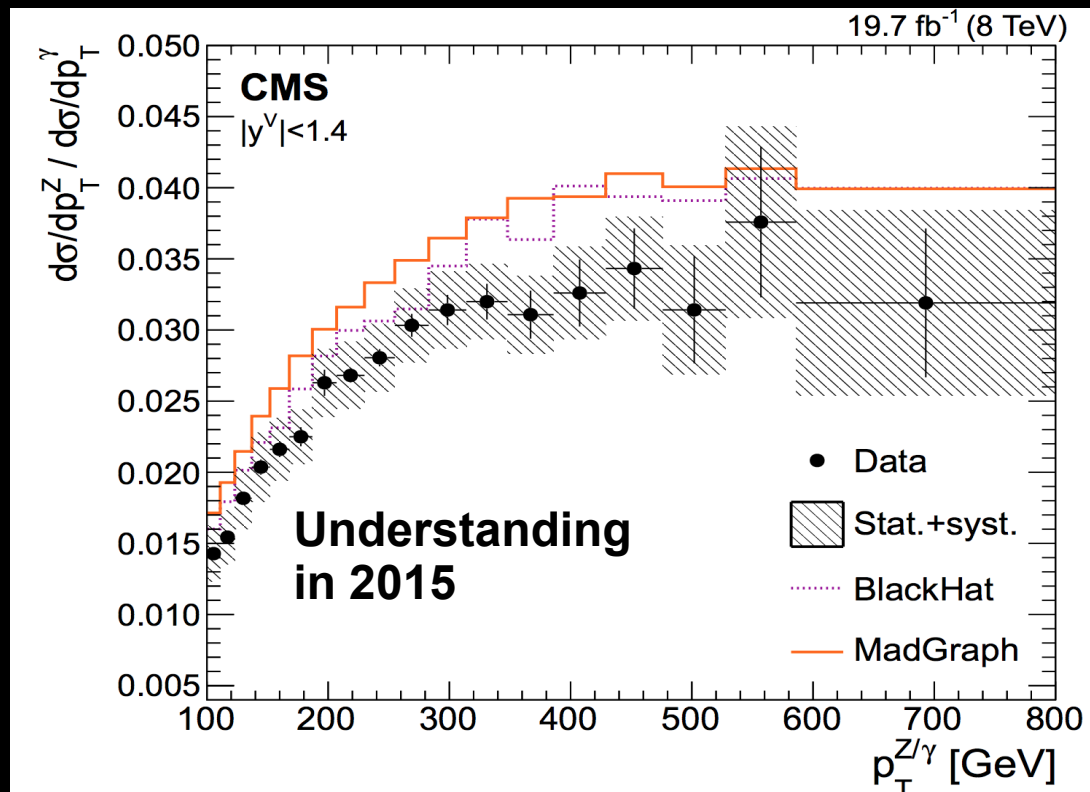
5 Control regions

15% uncertainty @ 1 TeV



Can we really use
all these regions?

Answer: Sort of.....



However we still have a problem!

Going from γ or $W \rightarrow Z$

Unc. $\longrightarrow \frac{d\sigma^{\gamma(W)}}{dp_T} / \frac{d\sigma^Z}{dp_T}$

Need to know the uncertainty on the ratios
 @NNLO QCD @NLO EWK
This is not a light statement!

However we still have a problem!

Going from γ or $W \rightarrow Z$

Unc. $\longrightarrow \frac{d\sigma^{\gamma(W)}}{dp_T} / \frac{d\sigma^Z}{dp_T}$

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Out last Monday

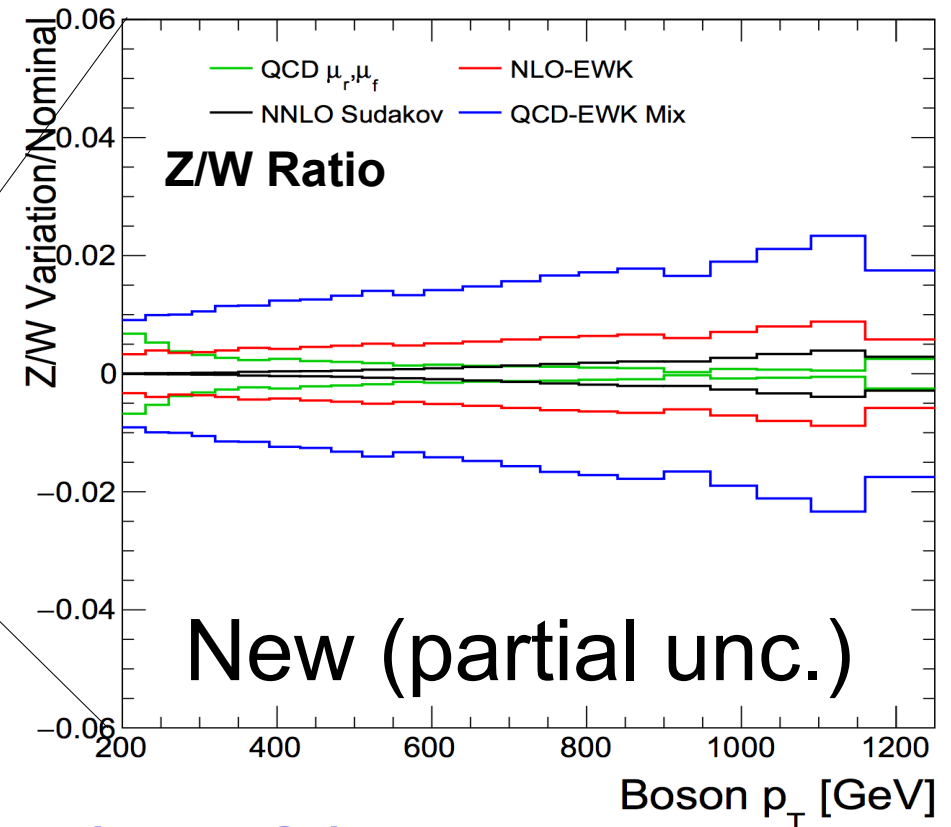
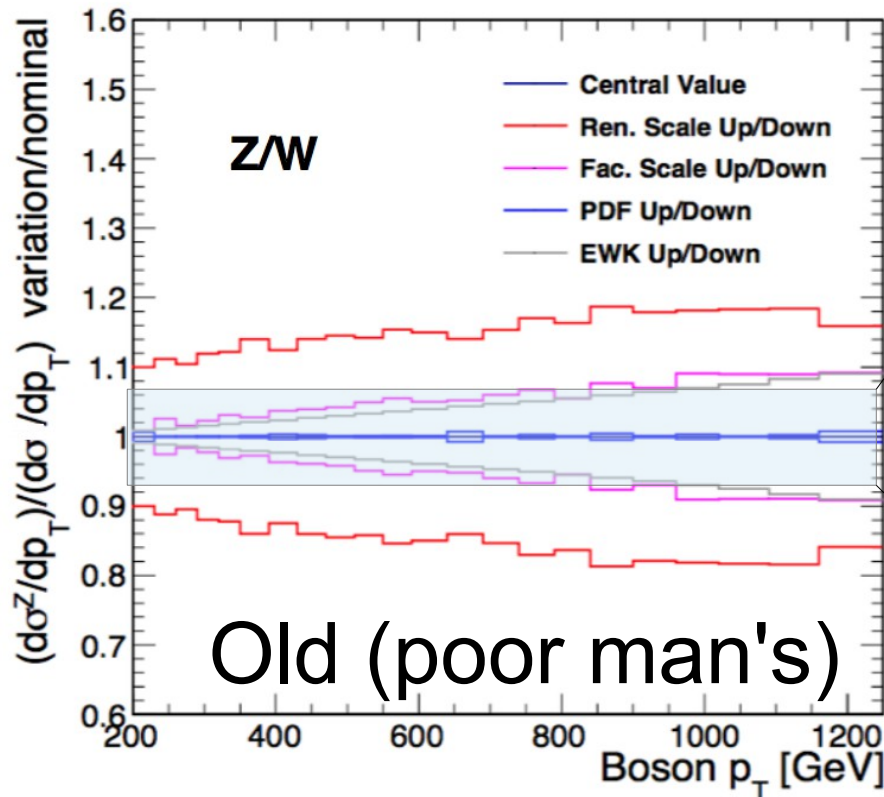
Arxiv:1705.04664

Precise predictions for V +jets dark matter backgrounds

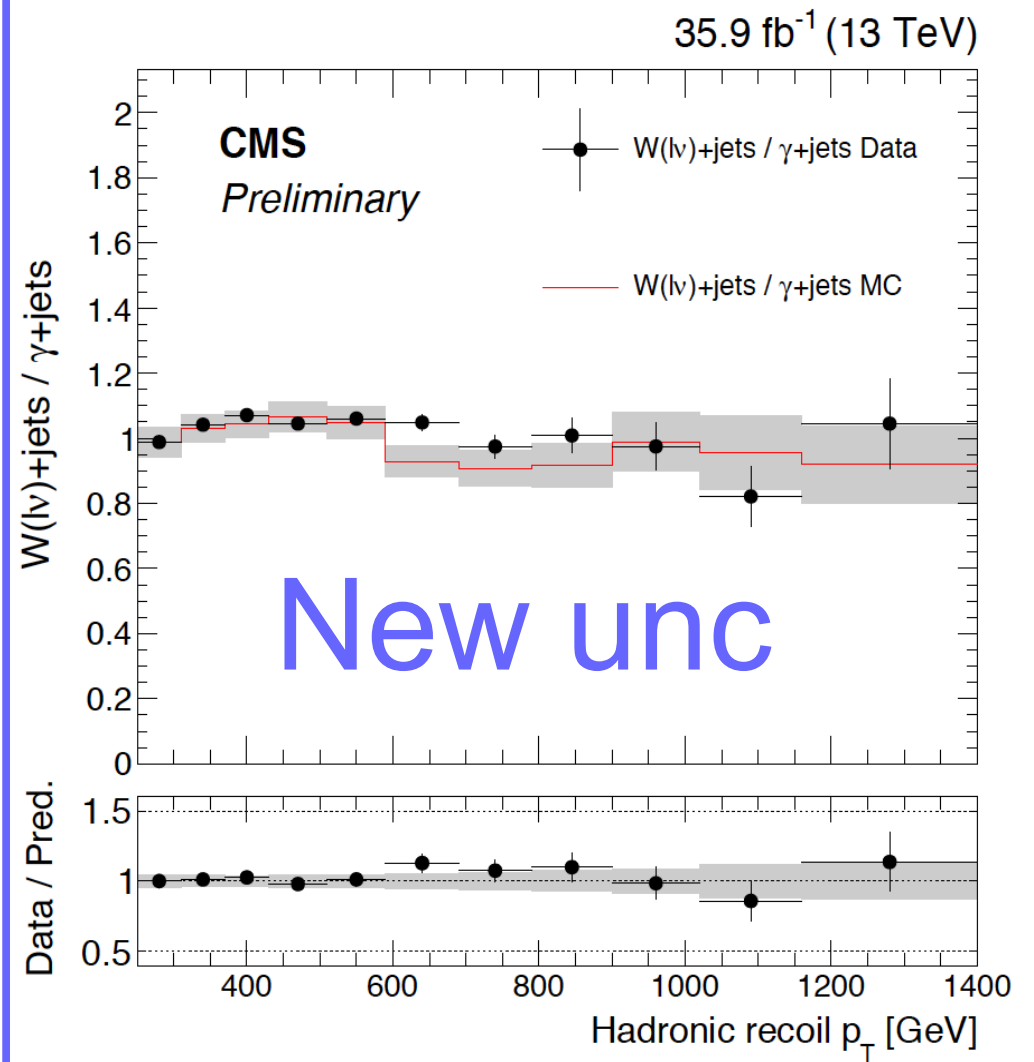
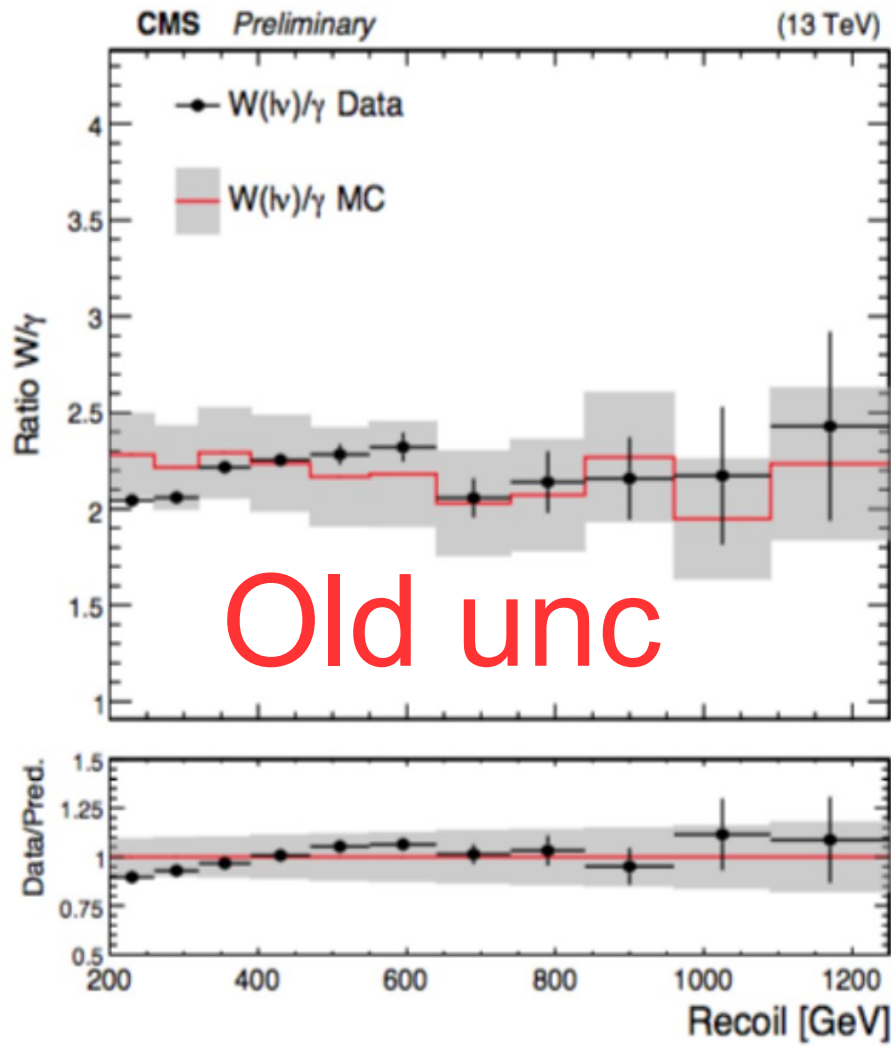
J. M. Lindert¹, S. Pozzorini², R. Boughezal³, J. M. Campbell⁴, A. Denner⁵,
 S. Dittmaier⁶, A. Gehrmann-De Ridder^{2,7}, T. Gehrmann², N. Glover¹, A. Huss⁷,
 S. Kallweit⁸, P. Maierhöfer⁶, M. L. Mangano⁸, T.A. Morgan¹, A. Mück⁹,
 F. Petriello^{3,10}, G. P. Salam^{*8}, M. Schönherr², and C. Williams¹¹

What did we get out of this?

- Large reduction in theory uncertainties
 - Experimental effects of the same order
- Can now correlate unc. across boson p_T bins

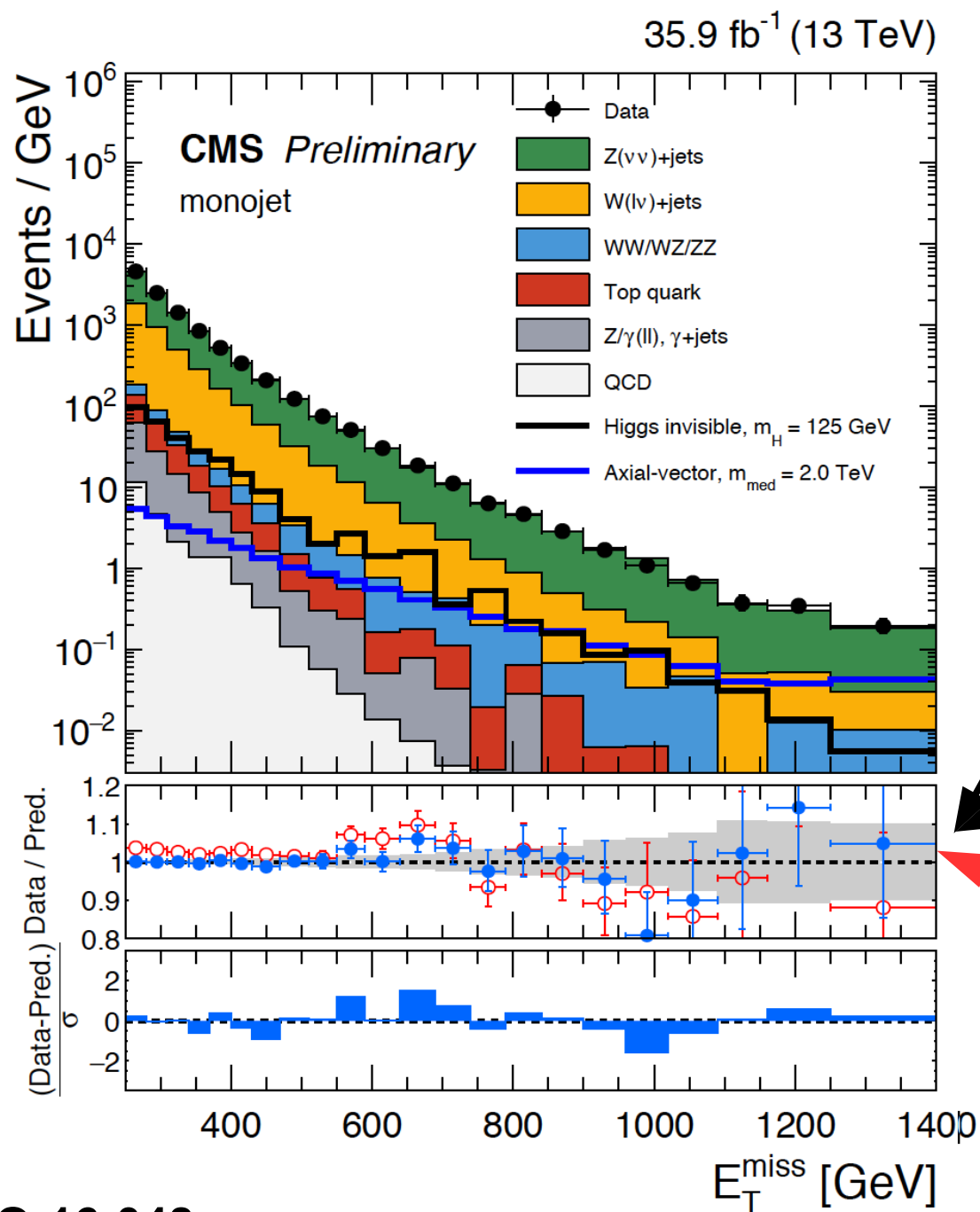


Reducing these unc. by a 5x can give x2 improvement



Uncertainties are considerably smaller
data constraints still better

Current Monojet Sensitivity⁴⁹



Uncertainty
Now at 1% level

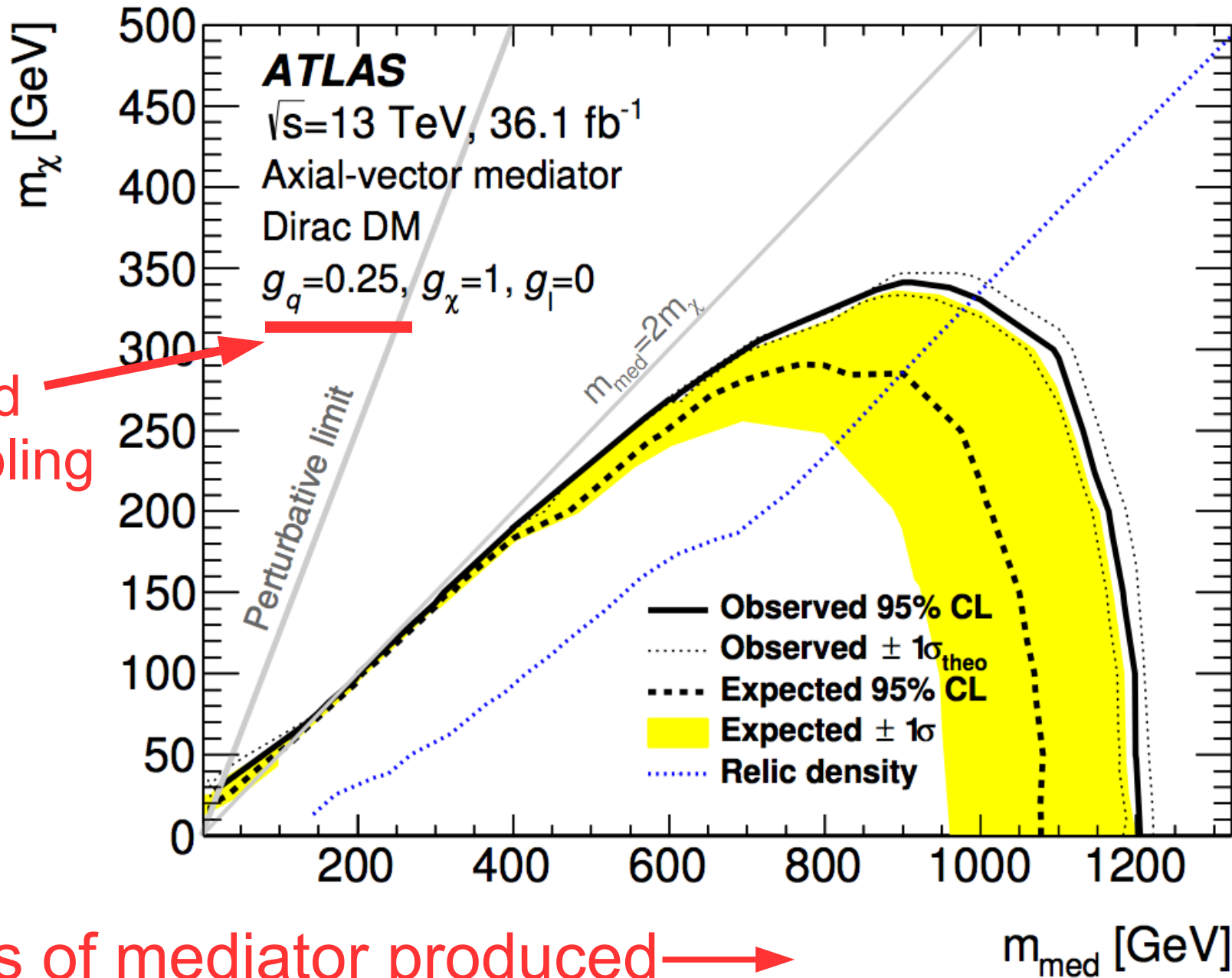
No excess

Almost **no**
systematic wall

Dark Matter
Mass

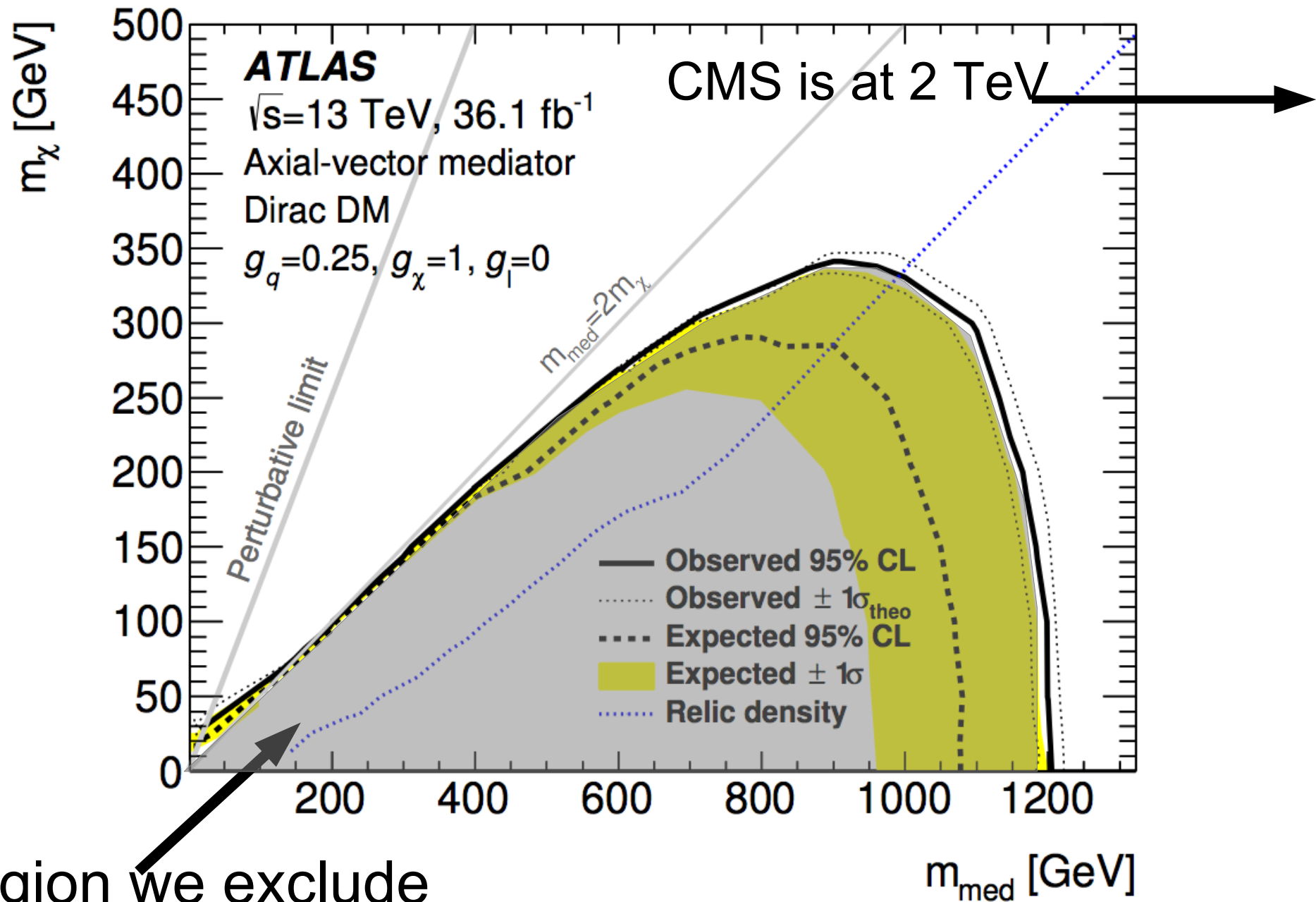
Pick a Model

Fixed
coupling



Mass of mediator produced

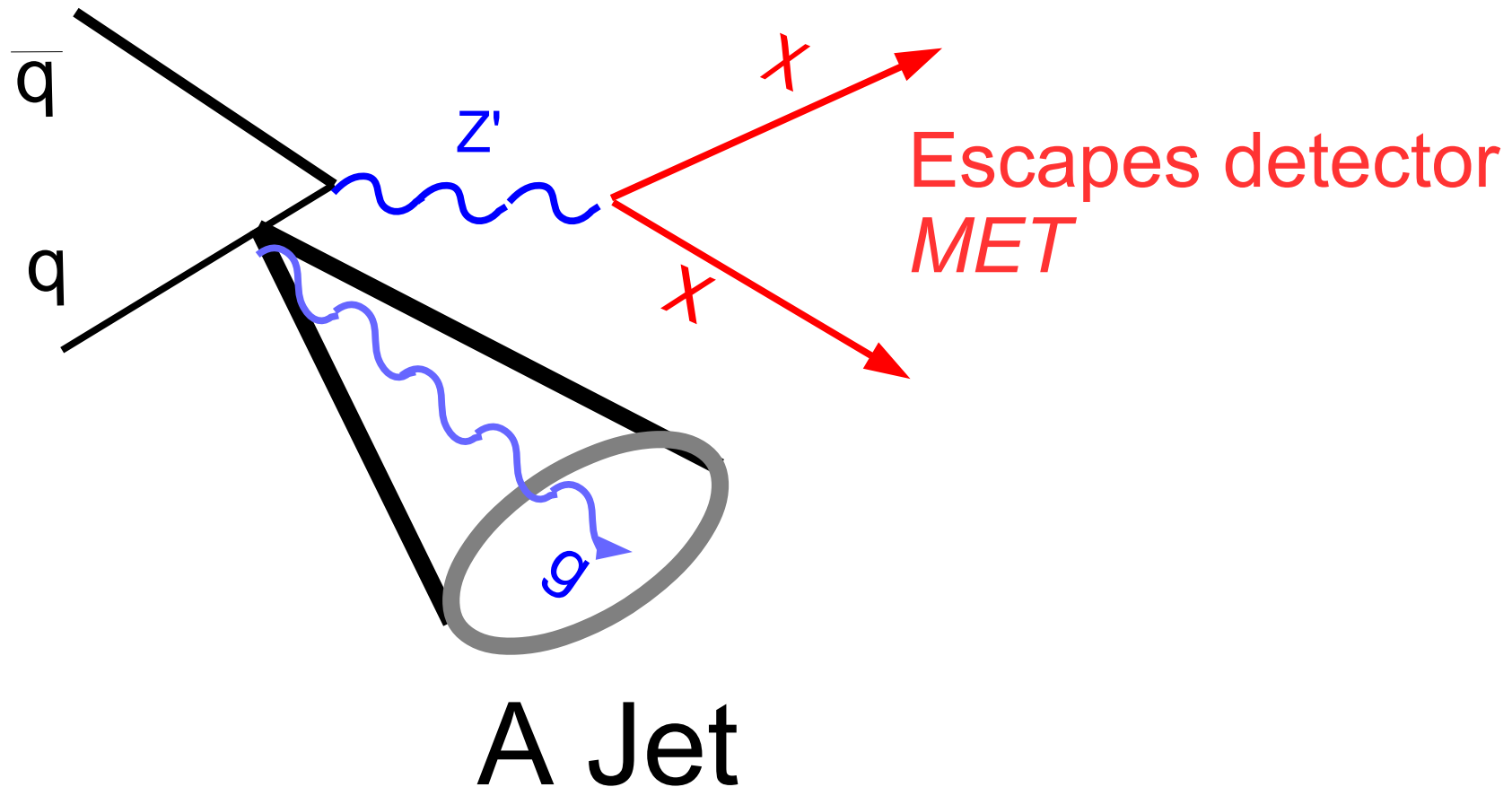
Understanding sensitivity



Beyond Monojet

The Basic Monojet Search

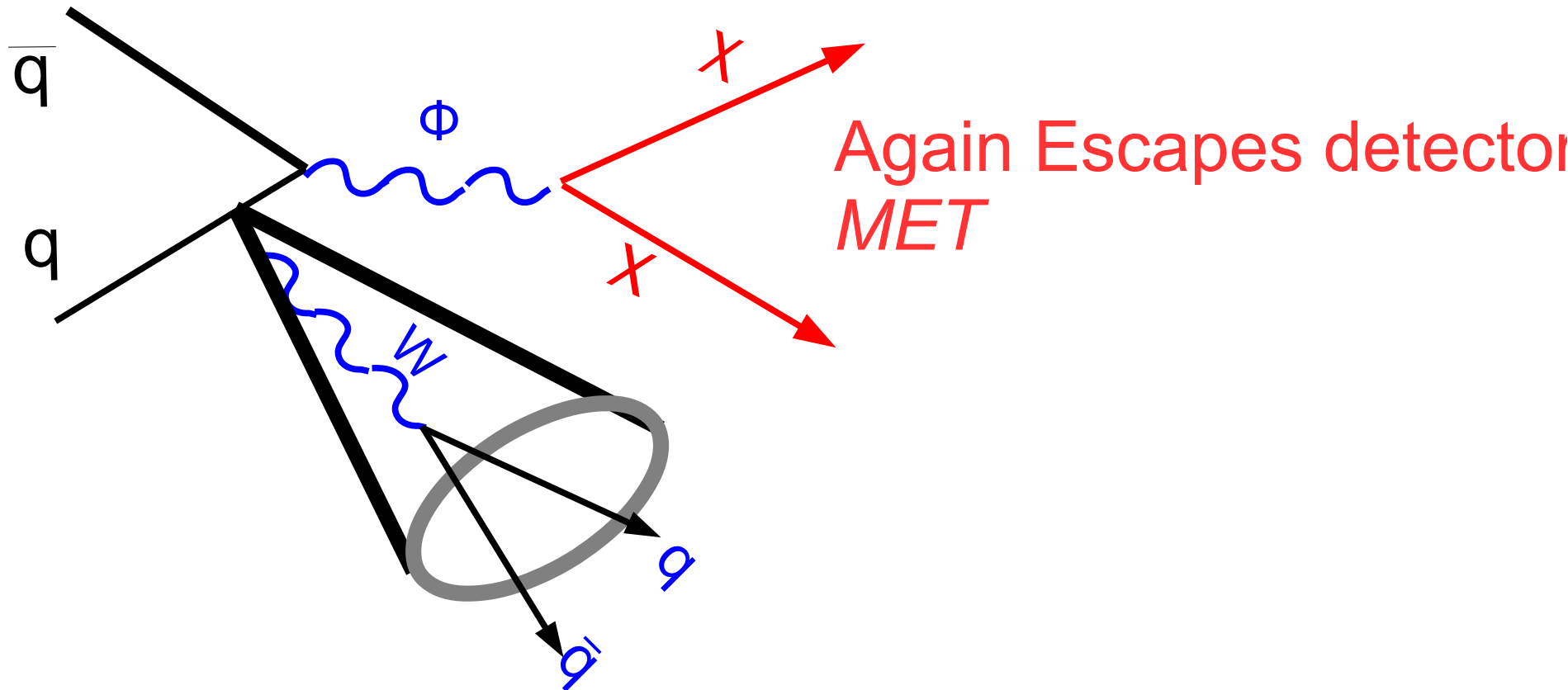
Escaping detector gives us signatures of *MET*



How Do We Discriminate Models?

Mono-W

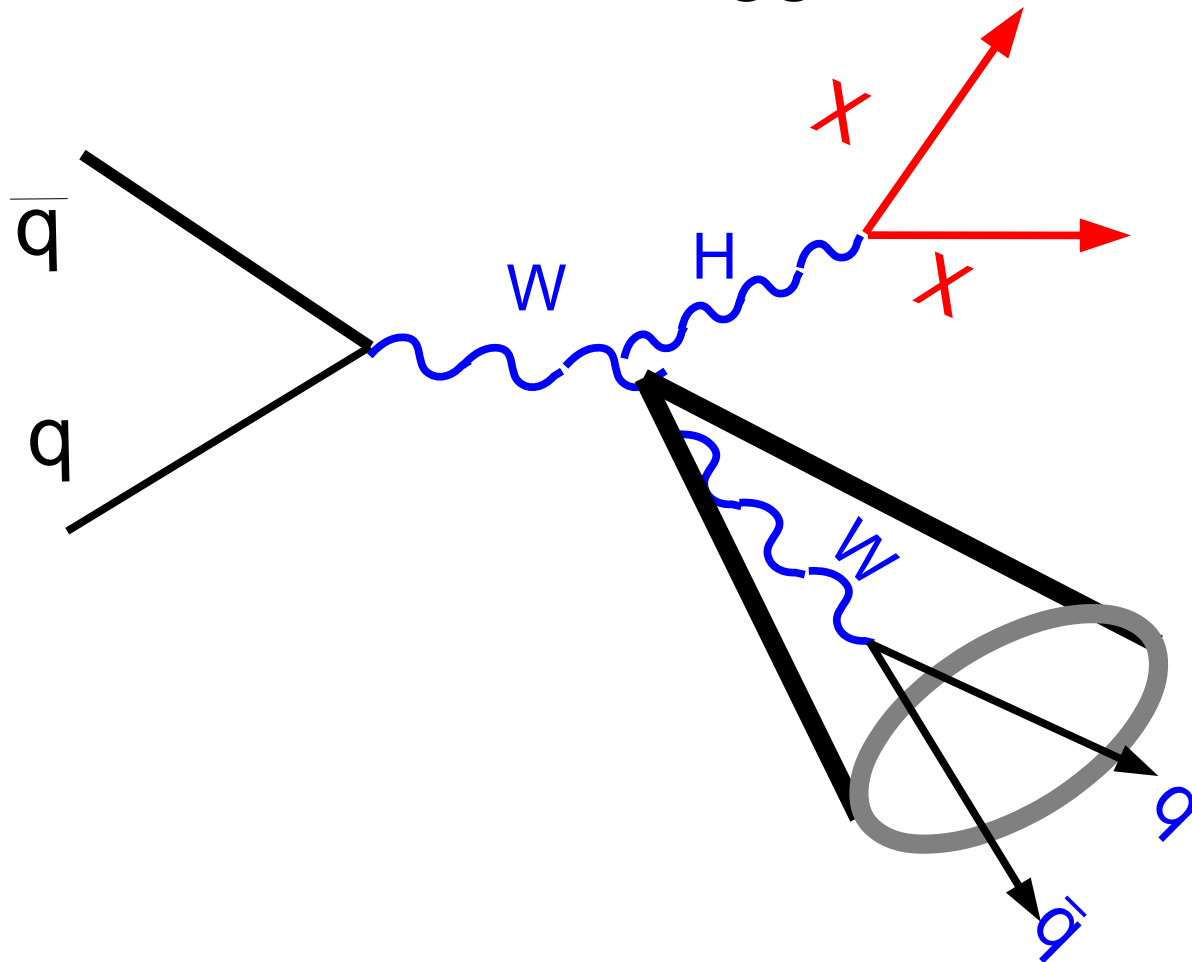
Monos : Vector, Axial, Higgs



How Do We Discriminate Models?

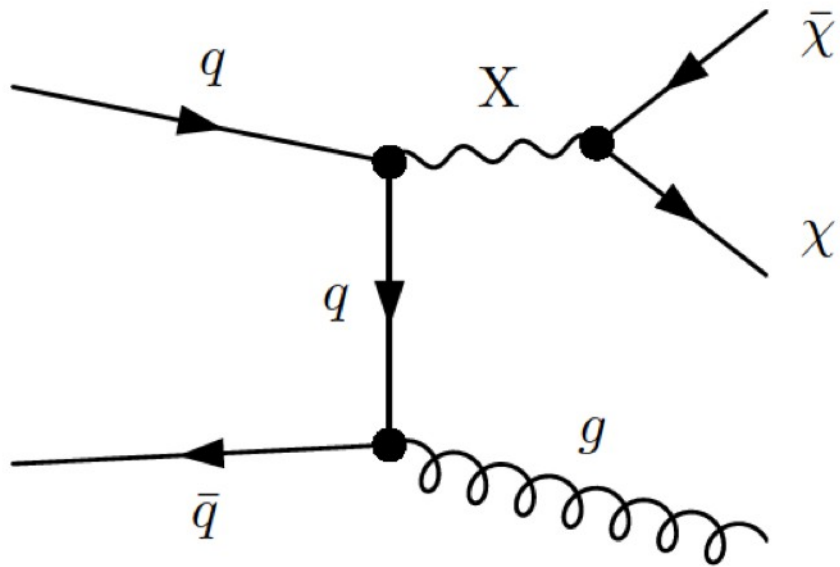
Mono-W

Monos : Vector, Axial, Higgs



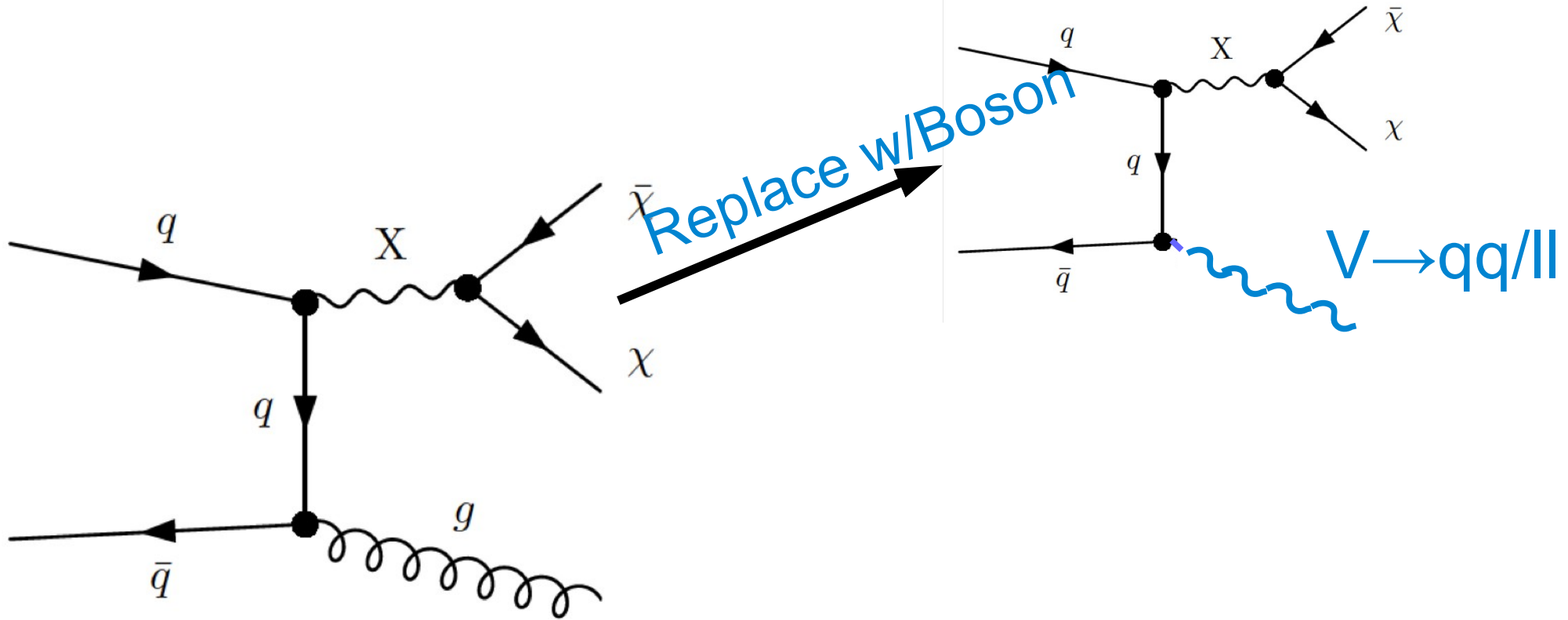
Spin 1 DM Searches

Spin 1 production on SM couplings for final state
Easily extend this to other final states



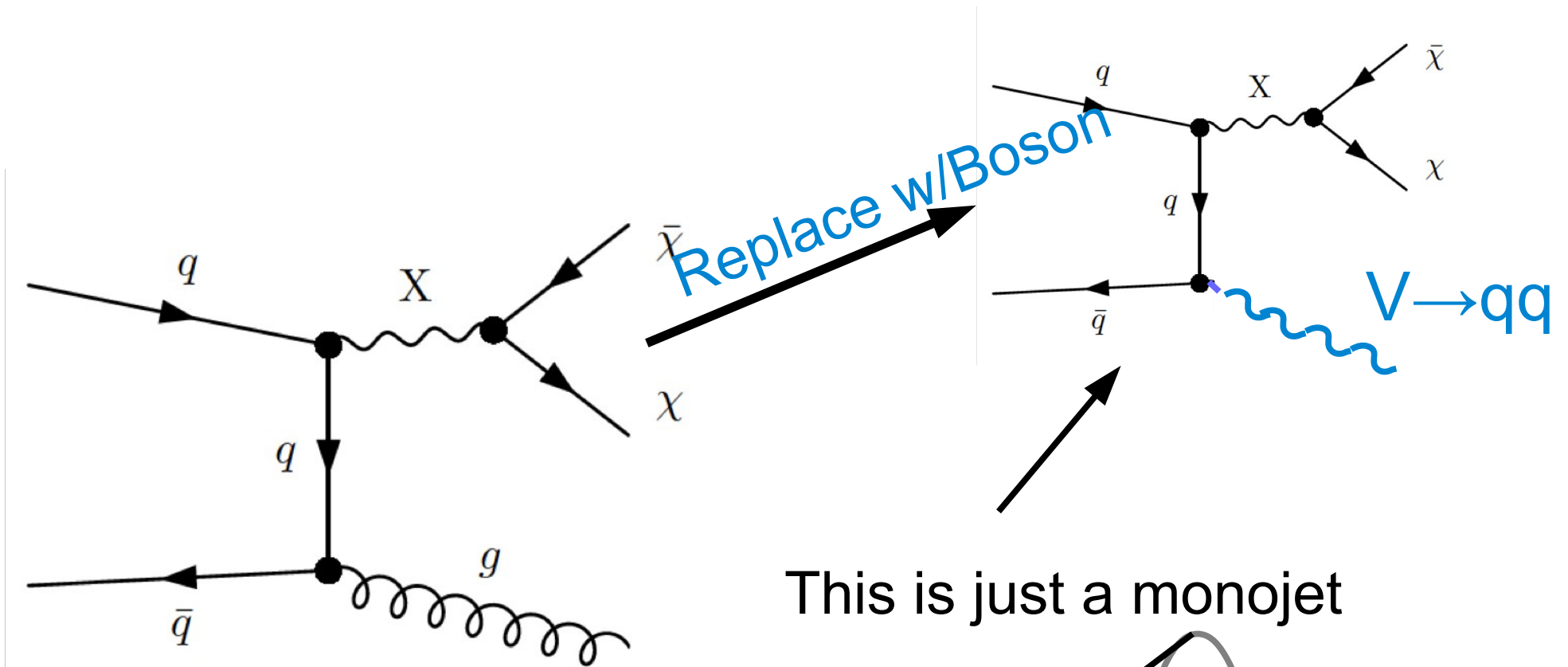
Spin 1 DM Searches

Can look for a Vector boson+*MET* as well



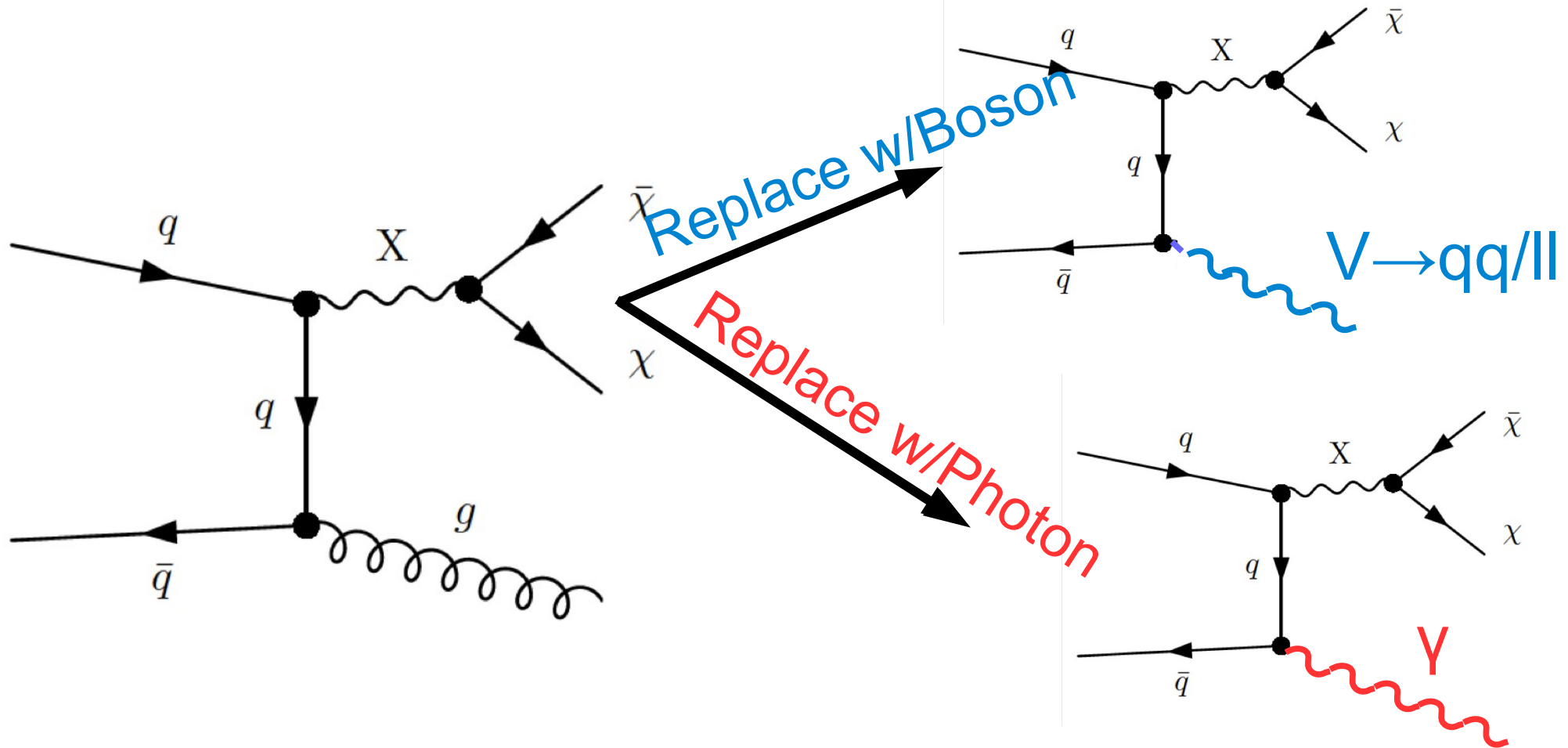
The split in simplified model terms

- With spin 1 can generate other final states :



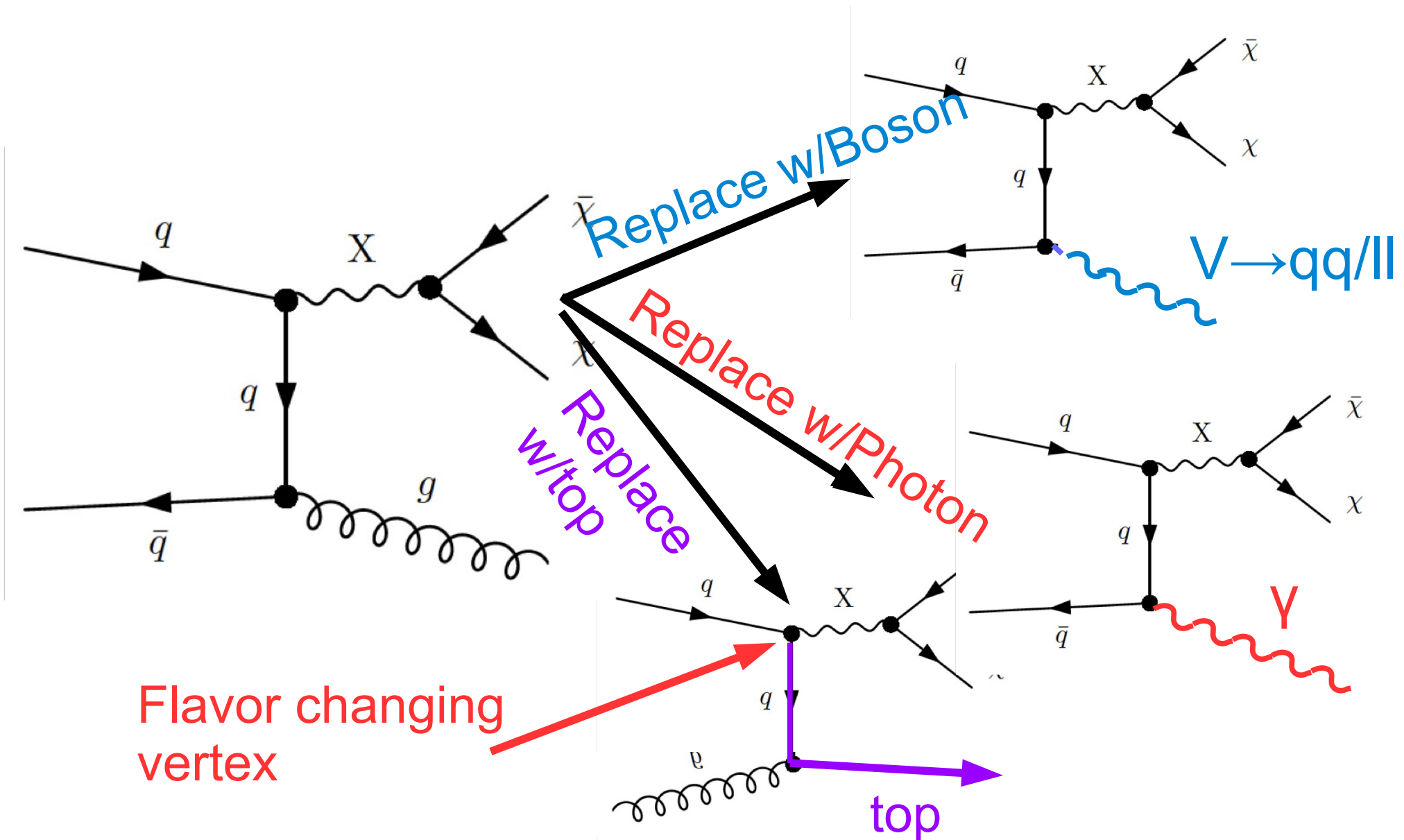
Spin 1 DM Searches

Can look for a Photon+*MET* as well



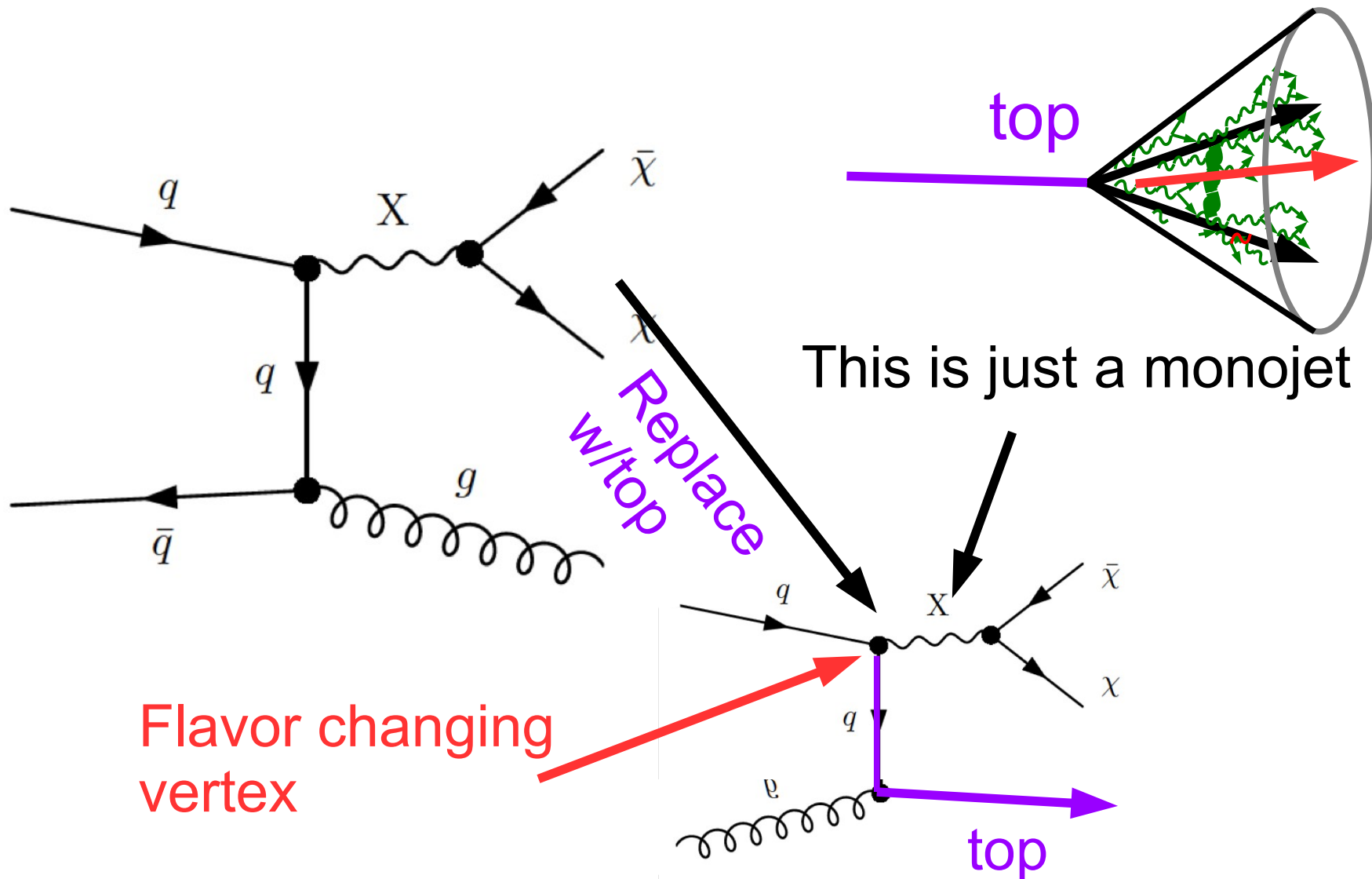
Spin 1 DM Searches

If vertex is flavor changing



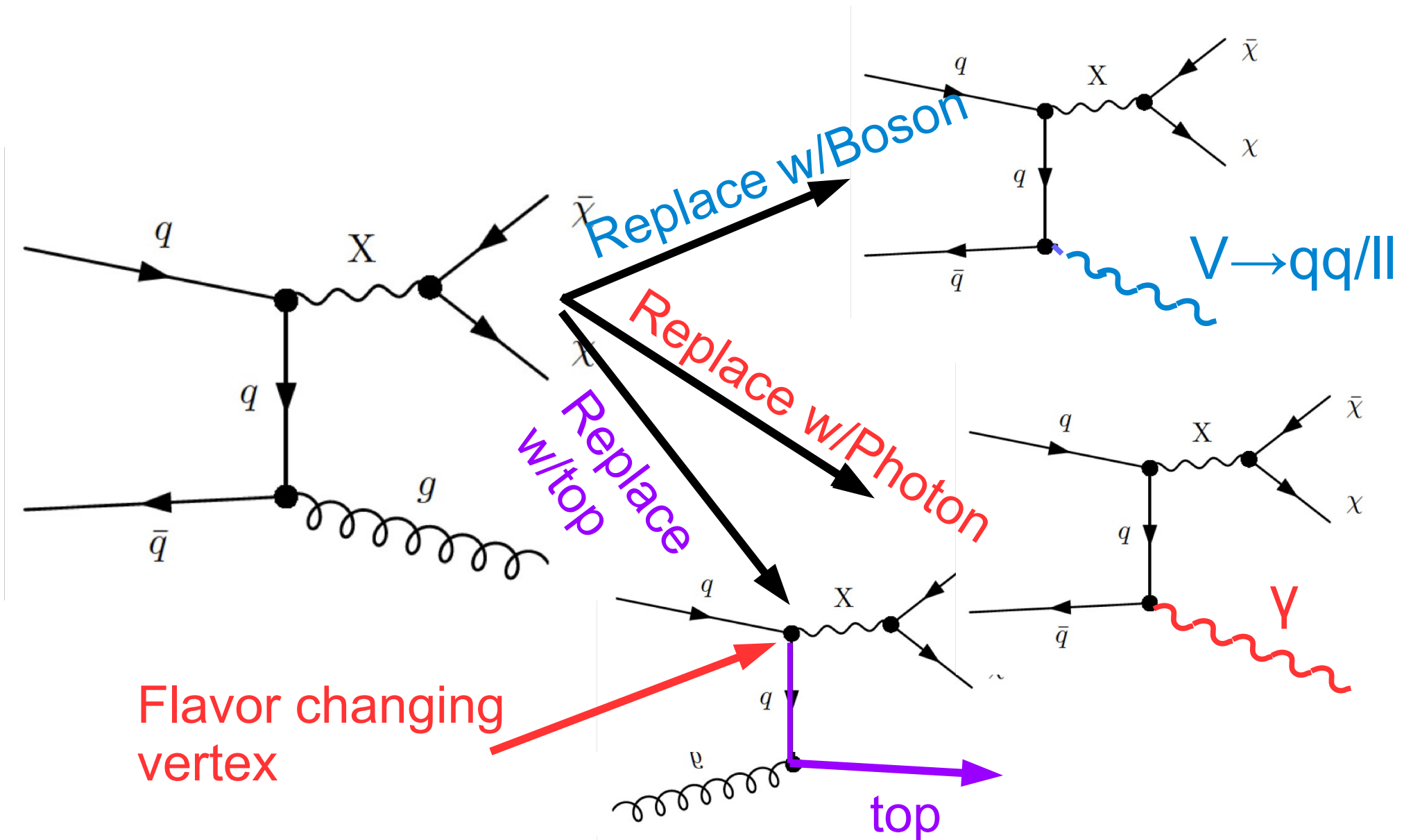
The split in simplified model terms

- With spin 1 can generate other final states :



Spin 1 DM Searches

If vertex is flavor changing



Spin 1 DM Searches

ATLAS-CONF-2017-028
arXiv:1703.05236



Modify the model

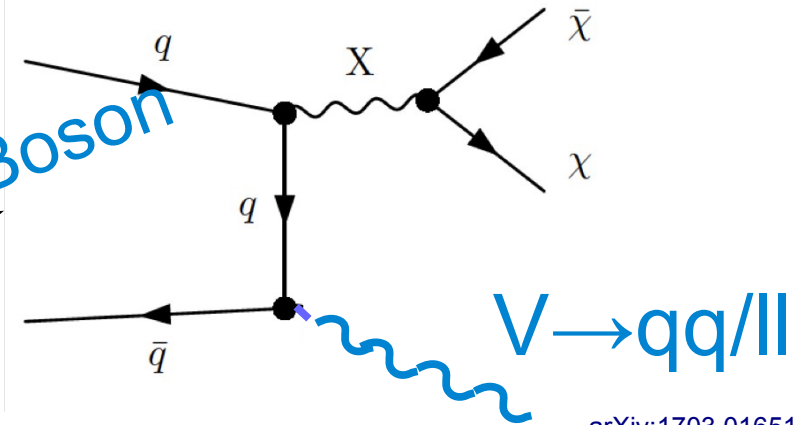
Replace w/Higgs

Replace w/Boson

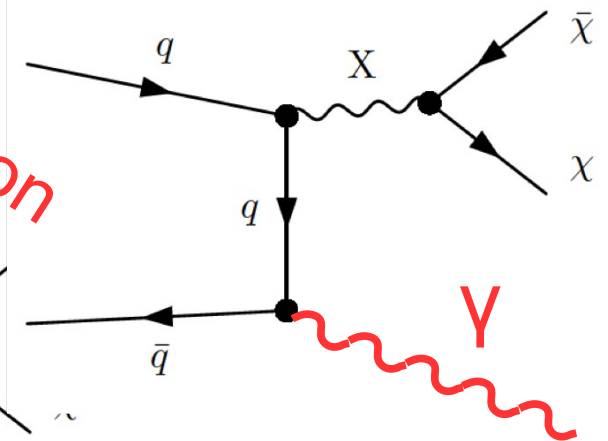
Replace w/Photon

Replace w/top

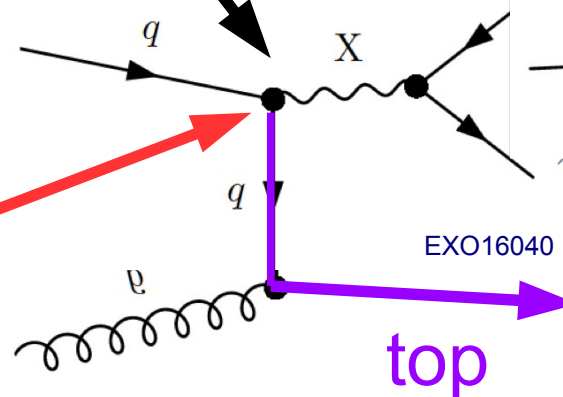
Flavor changing vertex



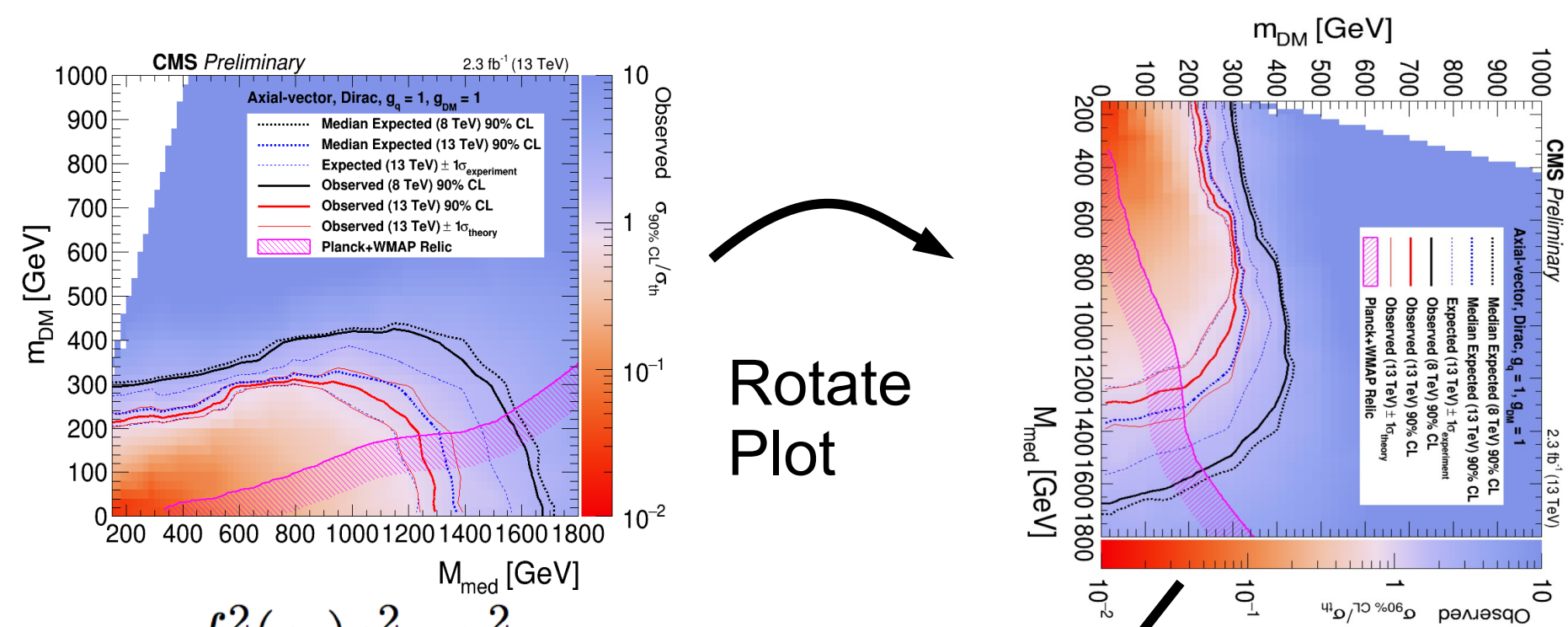
arXiv:1703.01651



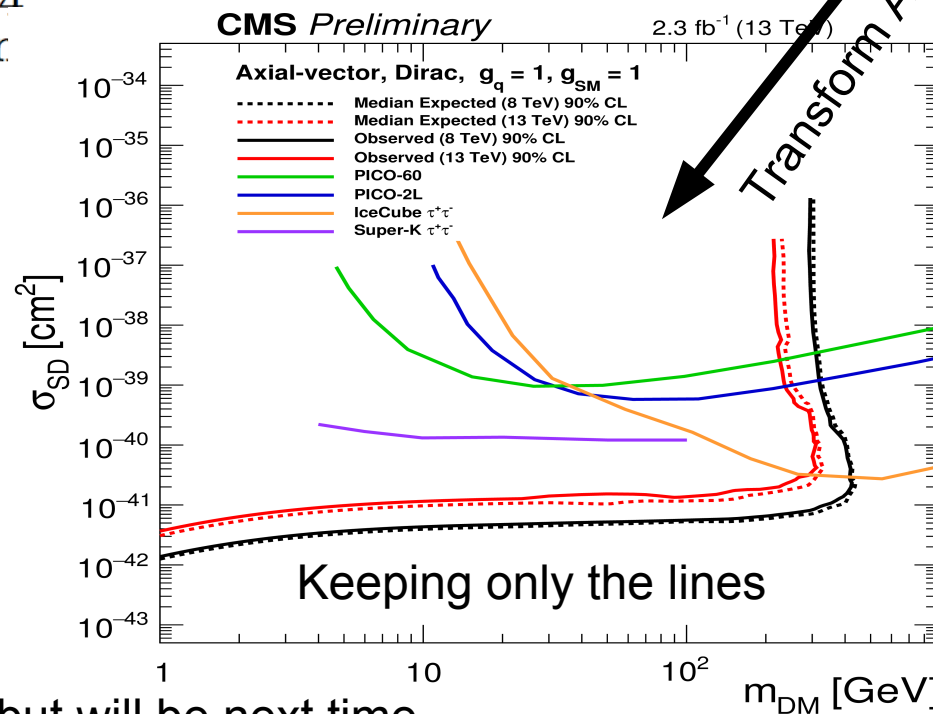
arxiv:1704.03848



EXO16040



$$\sigma_{\text{SI}} = \frac{f^2(g_q)g_{\text{DM}}^2\mu_{n\chi}^2}{\pi M_r^4}$$



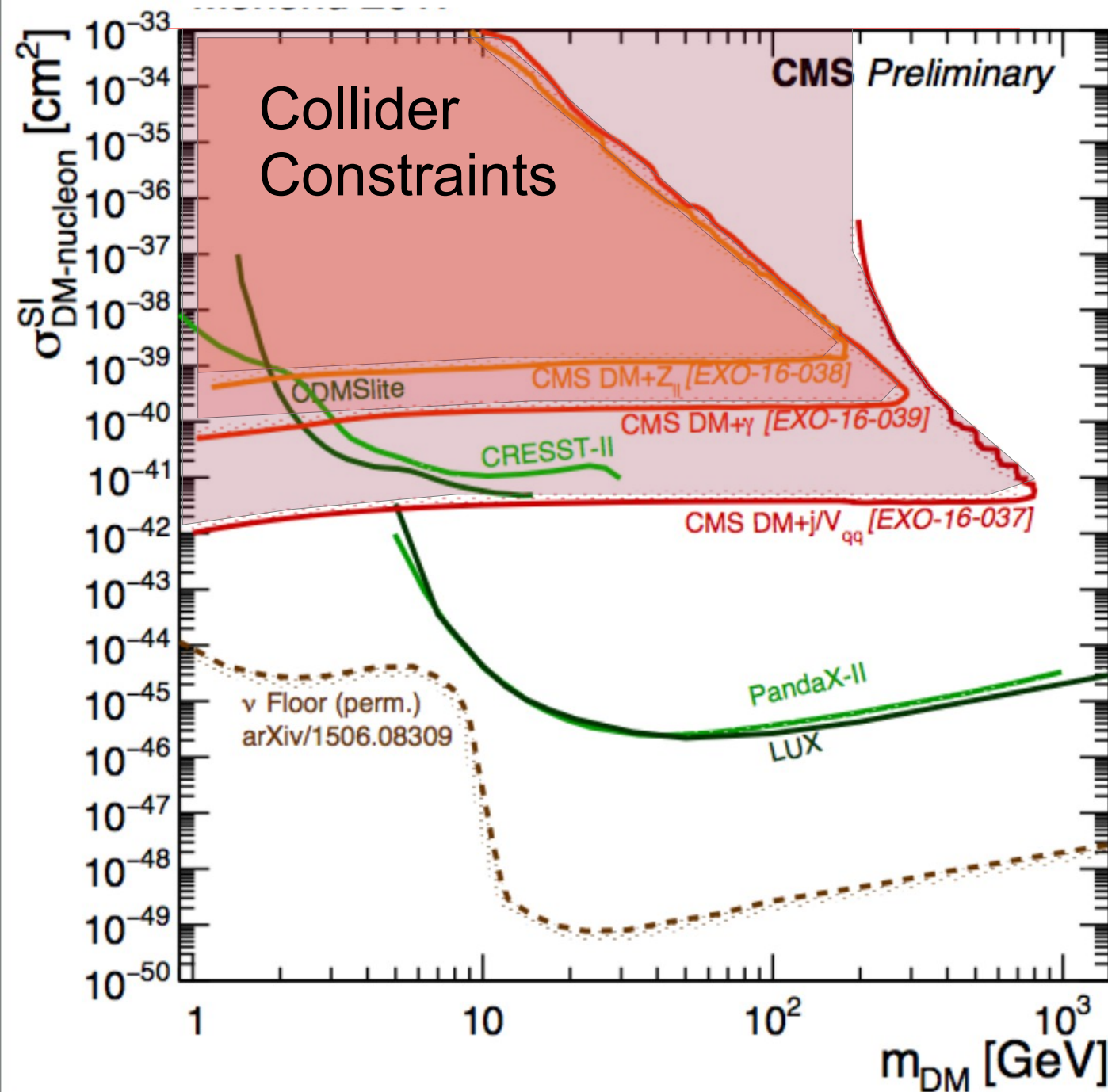
Now that search
is cast in terms
of mediator

No concerns in
the translation
Fixed couplings

1605.04917

Note: scale not run, but will be next time

Vector Mediator



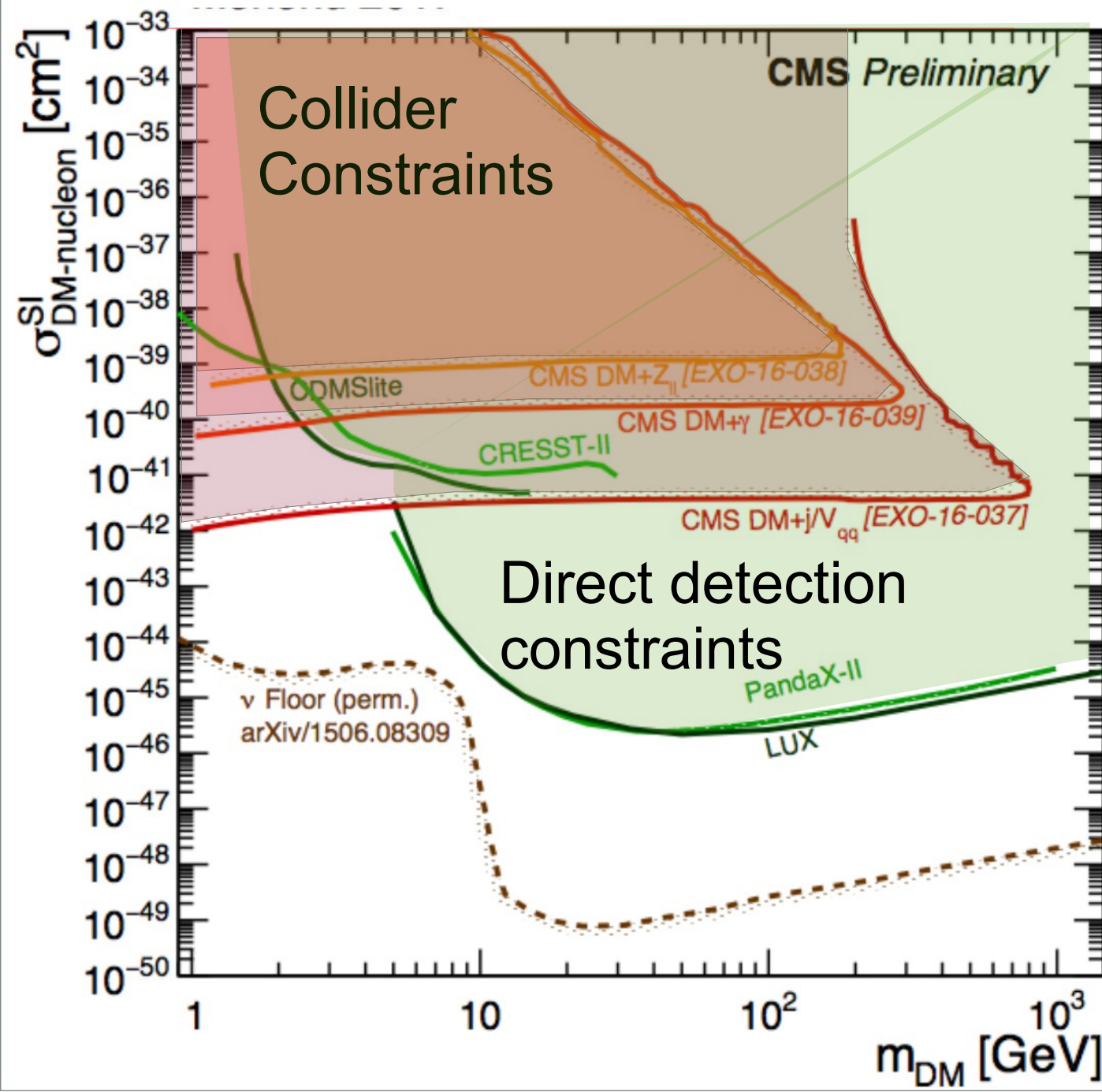
CMS observed exclusion 90% CL
Vector med., Dirac DM; $g_q = 0.25$, $g_{\text{DM}} = 1.0$

- ⋯ CMS DM+ j/V_{qq} [EXO-16-037]
- CMS DM+ γ [EXO-16-039]
- ⋯ CMS DM+ Z_{ll} [EXO-16-038]

DD observed exclusion 90% CL

- CRESST-II [arXiv:1509.01515]
- CDMSlite [arXiv:1509.02448]
- PandaX-II [arXiv:1607.07400]
- LUX [arXiv:1608.07648]

Vector Mediator



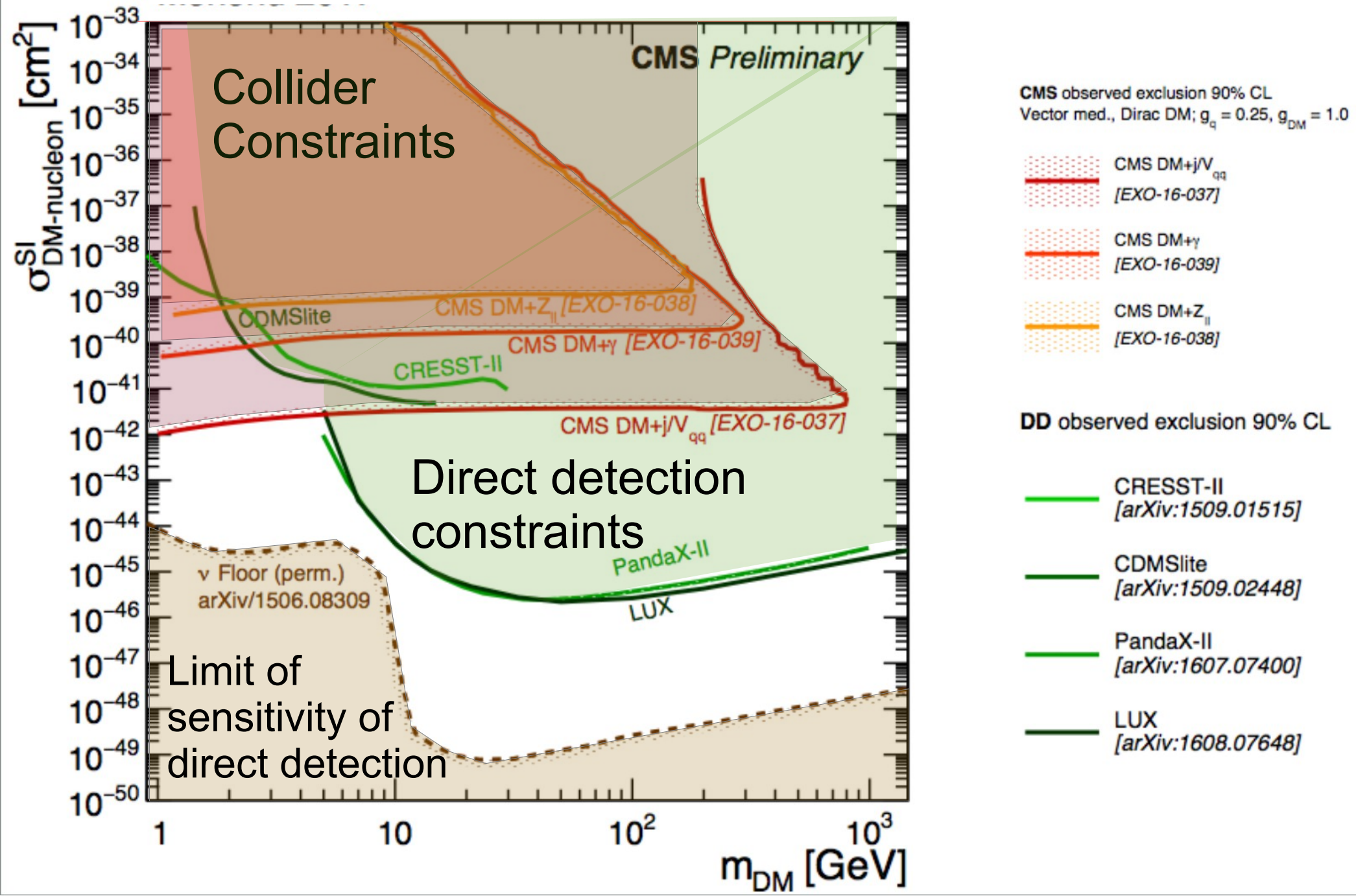
CMS observed exclusion 90% CL
Vector med., Dirac DM; $g_q = 0.25$, $g_{\text{DM}} = 1.0$

- ⋯ CMS DM+j/V_{qq} [EXO-16-037]
- ⋯ CMS DM+ γ [EXO-16-039]
- ⋯ CMS DM+Z_{ll} [EXO-16-038]

DD observed exclusion 90% CL

- CRESST-II [arXiv:1509.01515]
- CDMSlite [arXiv:1509.02448]
- PandaX-II [arXiv:1607.07400]
- LUX [arXiv:1608.07648]

Vector Mediator



With Direct Detection

Vector(SI)



Spin independent

Extremely good

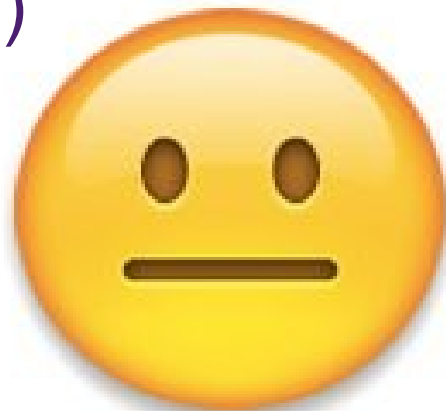
Axial(SD)



Spin dependent

Not so great

Scalar(SI)



So-so

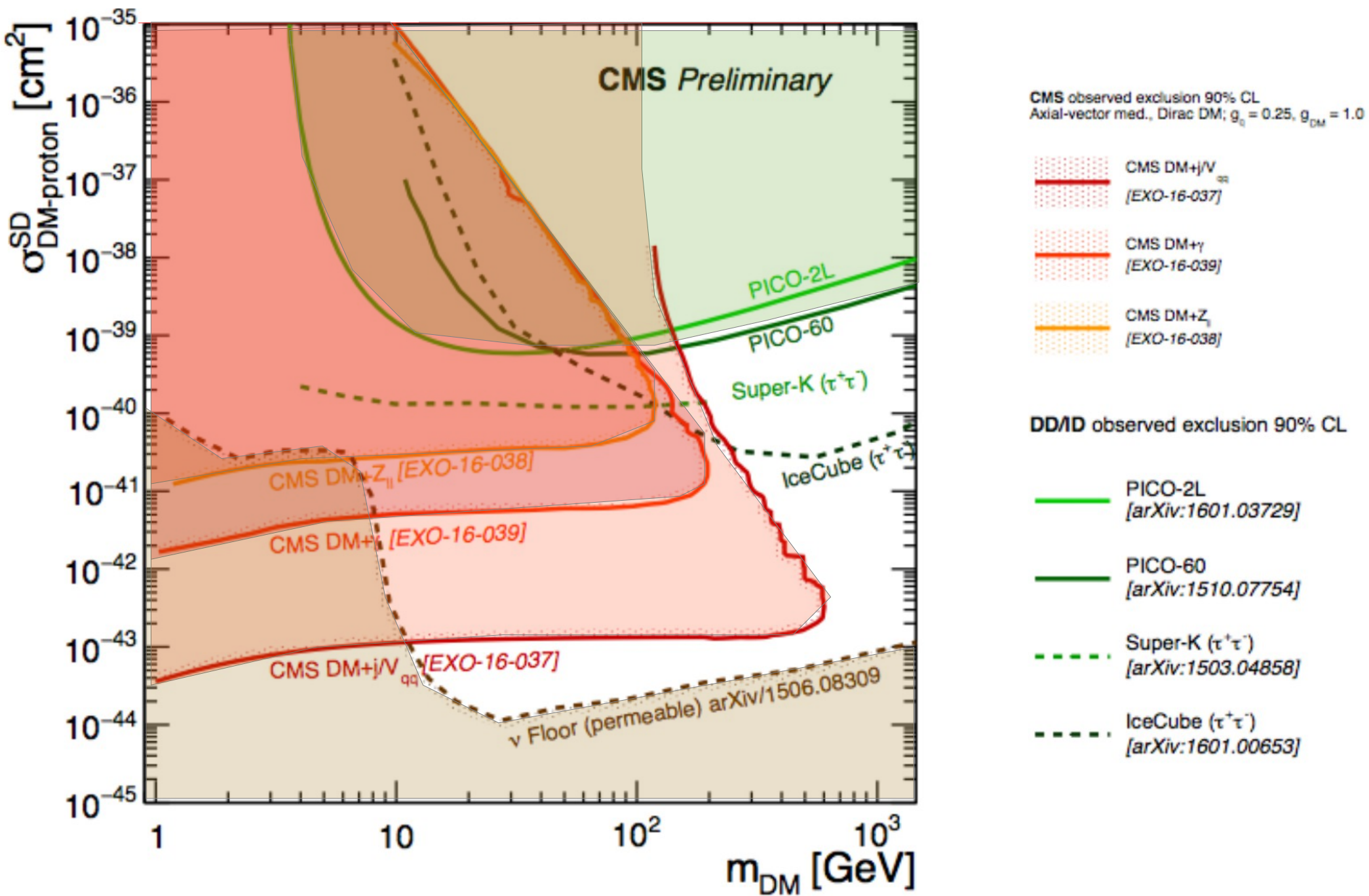
Spin independent

Pseudoscalar



Use indirect detection

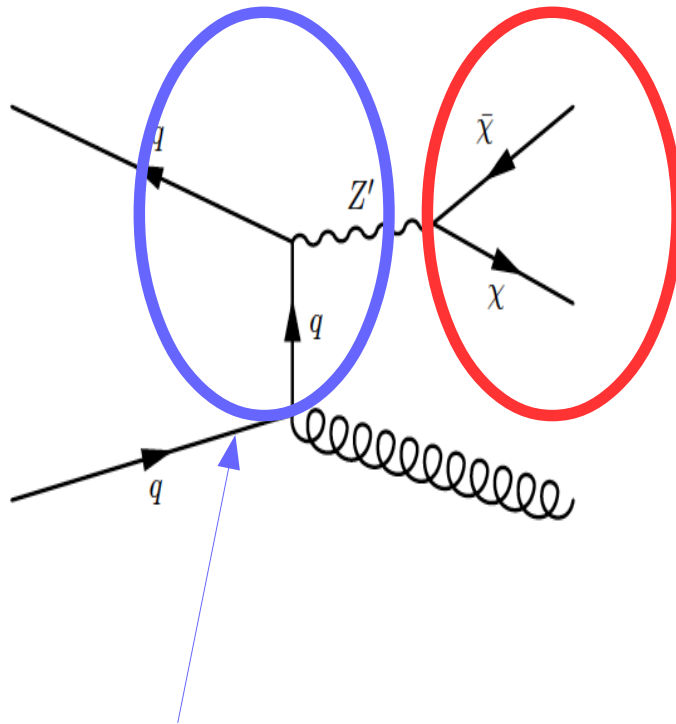
Axial Mediator



Beyond Invisible Searches

What else?

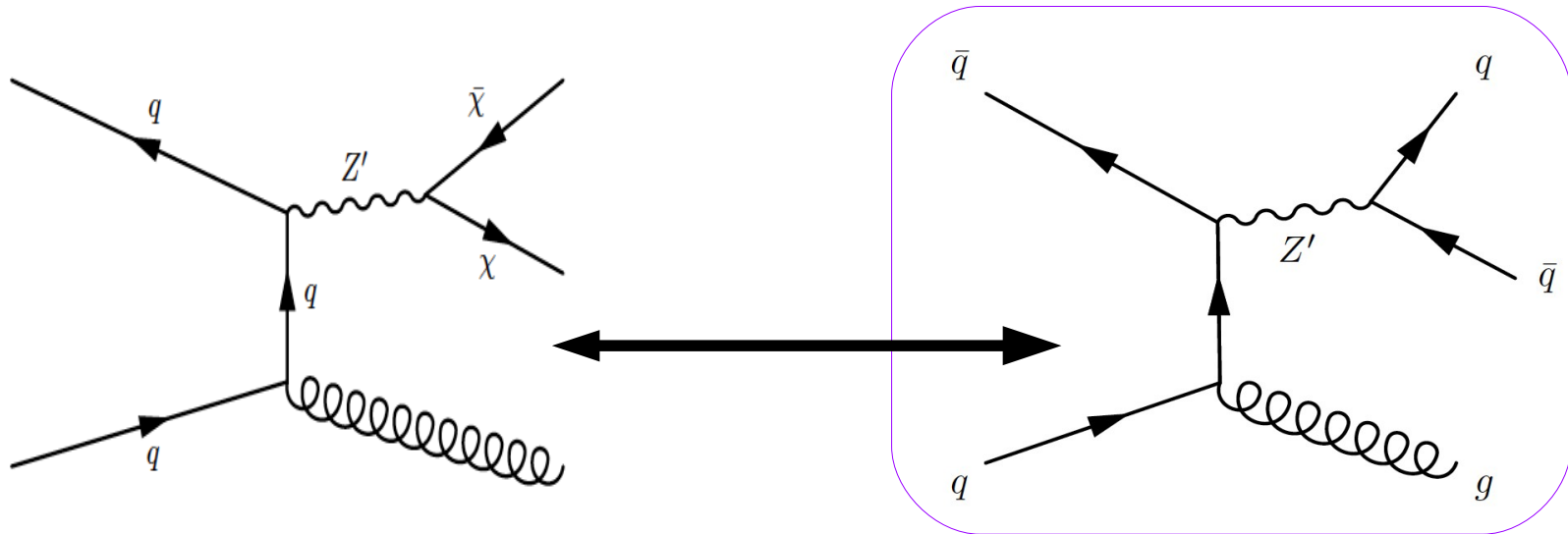
- Without loss of generality we also have dijets



Mediator is coupling to **quarks** and to **Dark matter**

What else?

- Without loss of generality we also have dijets

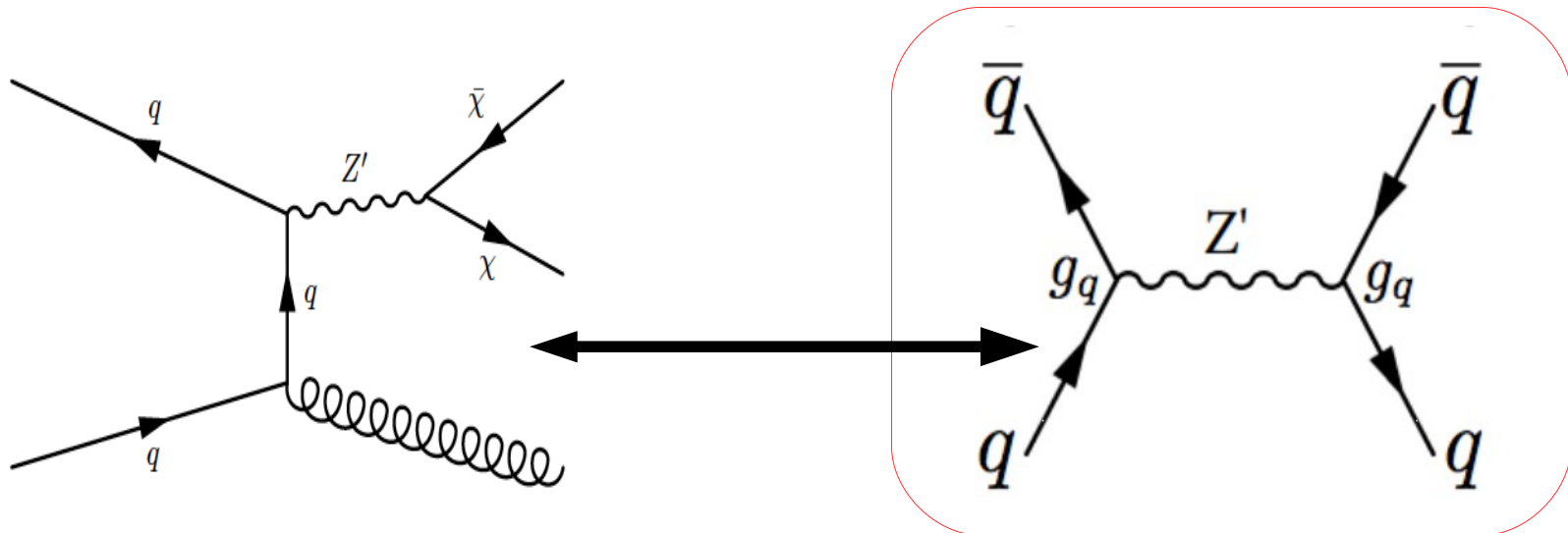


This is a dijet+ISR search

Mediator is coupling to quarks and to Dark matter
 Mediator can decay to quarks

What else?

- Without loss of generality we also have dijets



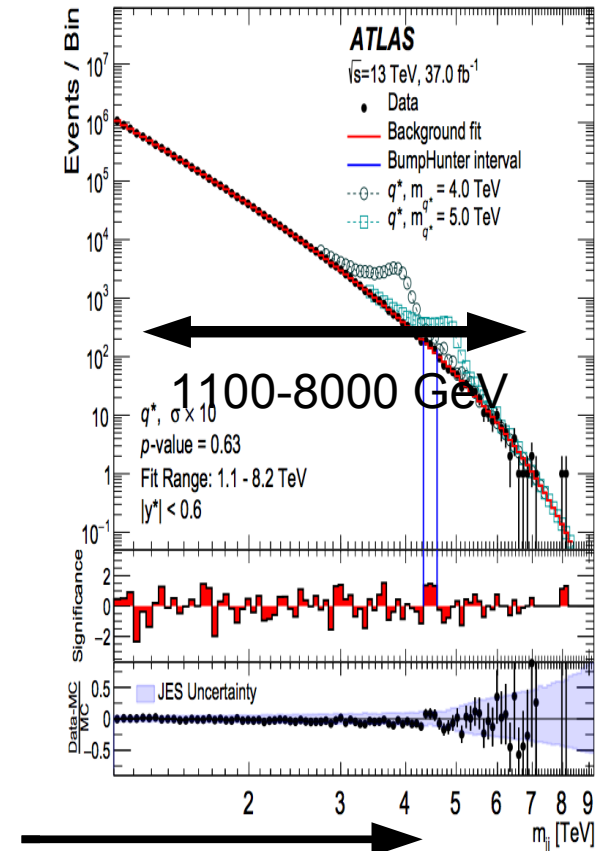
Can also just do a plain dijet search

When doing a dijet search **don't need additional jet**

$$\text{BR}(Z' \rightarrow qq) \approx 0.5 \text{BR}(Z' \rightarrow \quad)$$

Probing the Mass range

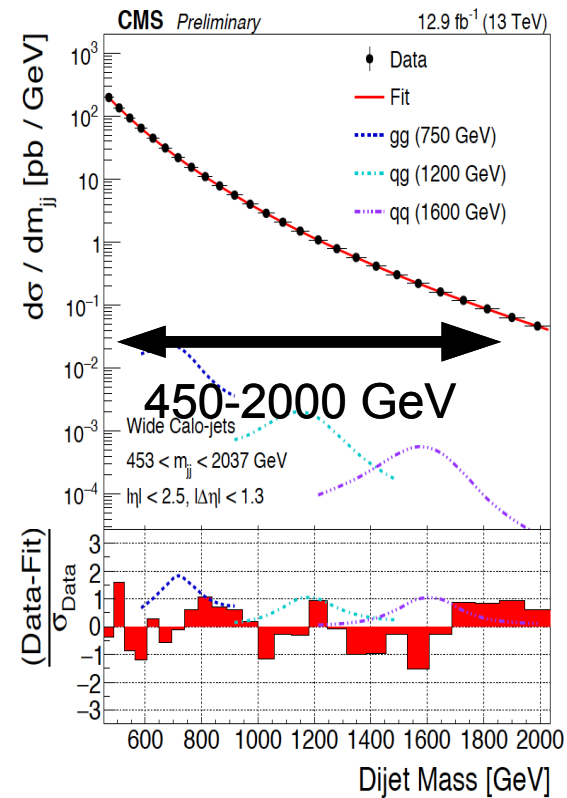
Like Monojet
we can expand to further regions
By tagging other objects



Standard jet
triggers

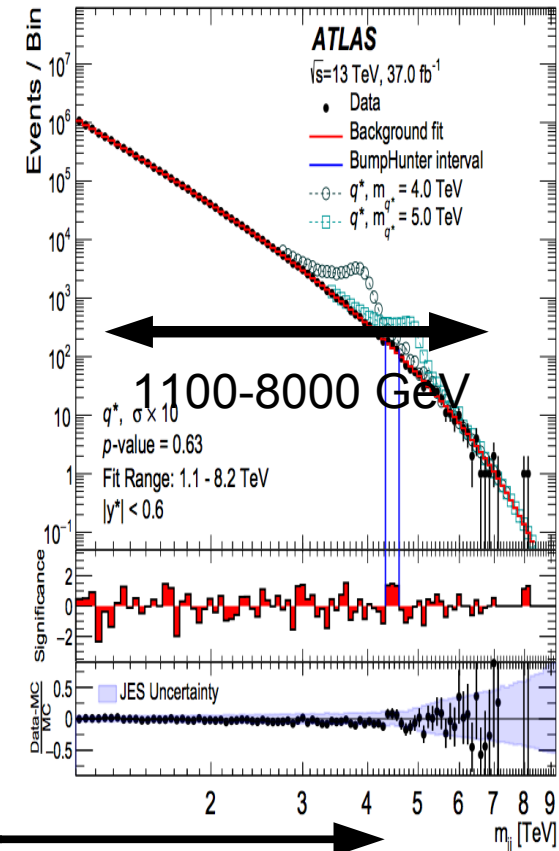
No tag

Probing the Mass range



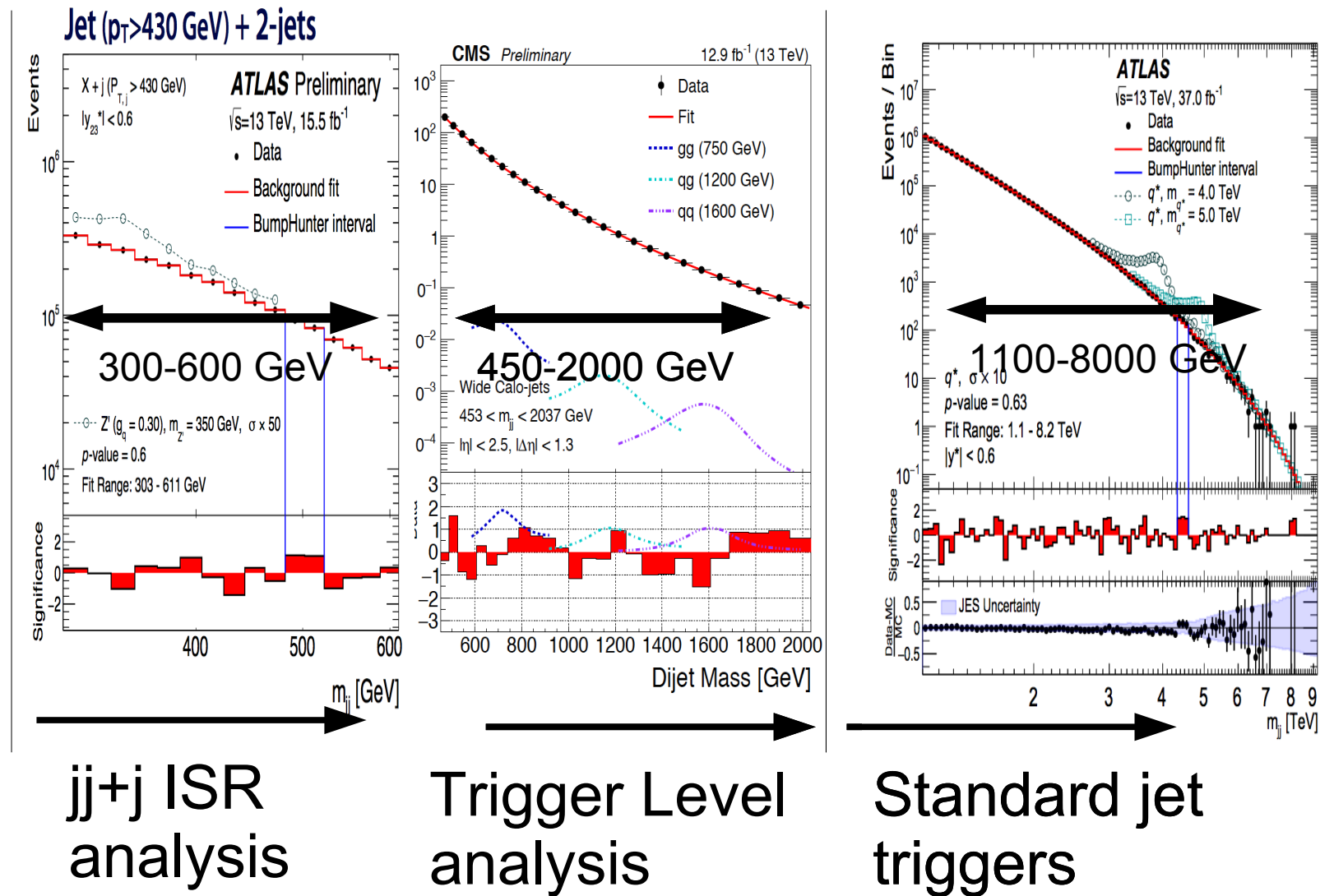
Trigger Level
analysis

Jets in trigger



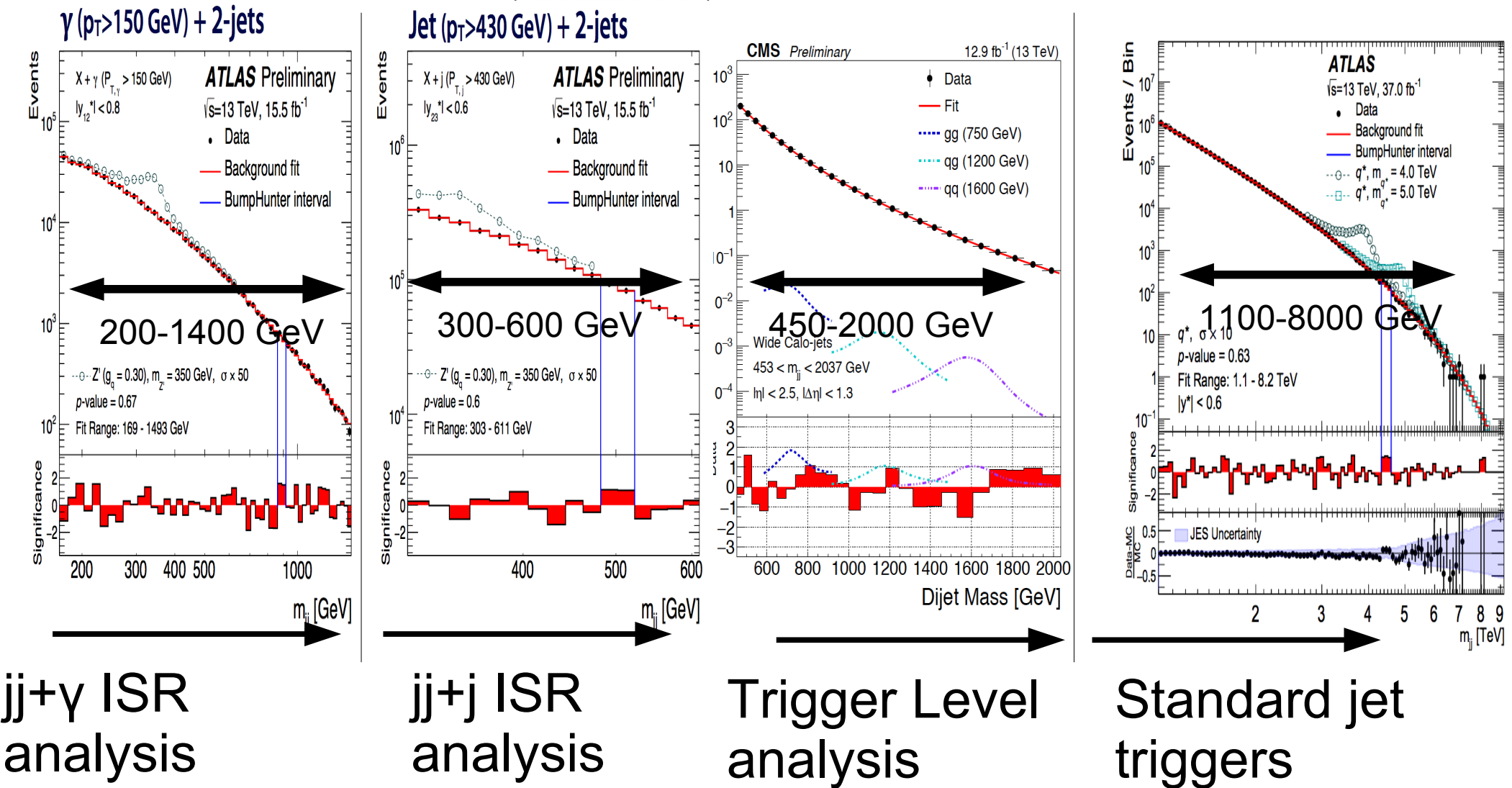
Standard jet
triggers

Probing the Mass range



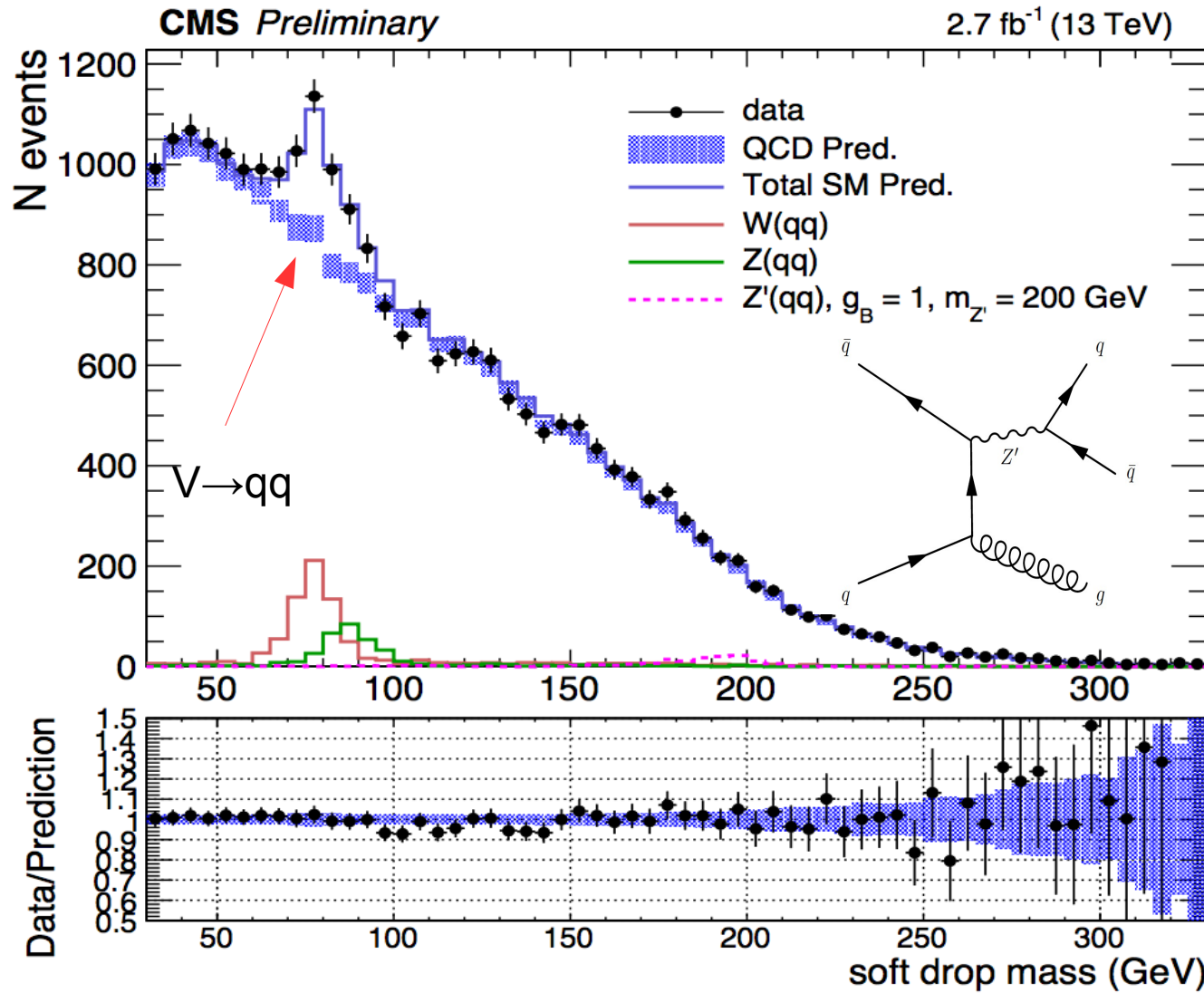
Additional jet (like monojet) **ATLAS**

Probing the Mass range

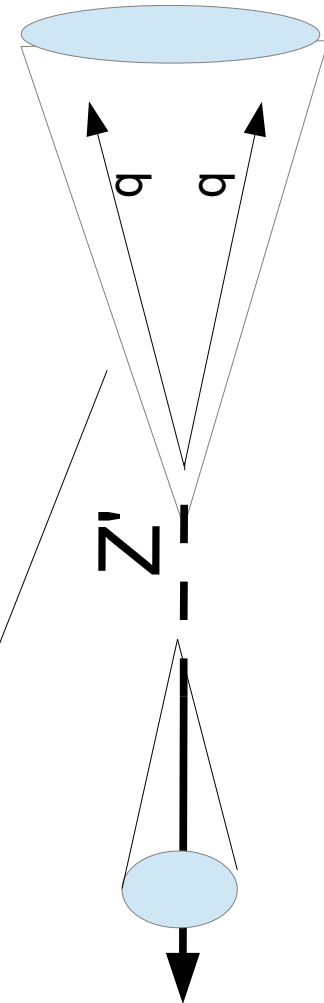


Additional photon (like monophoton) **ATLAS**

Going all the way down



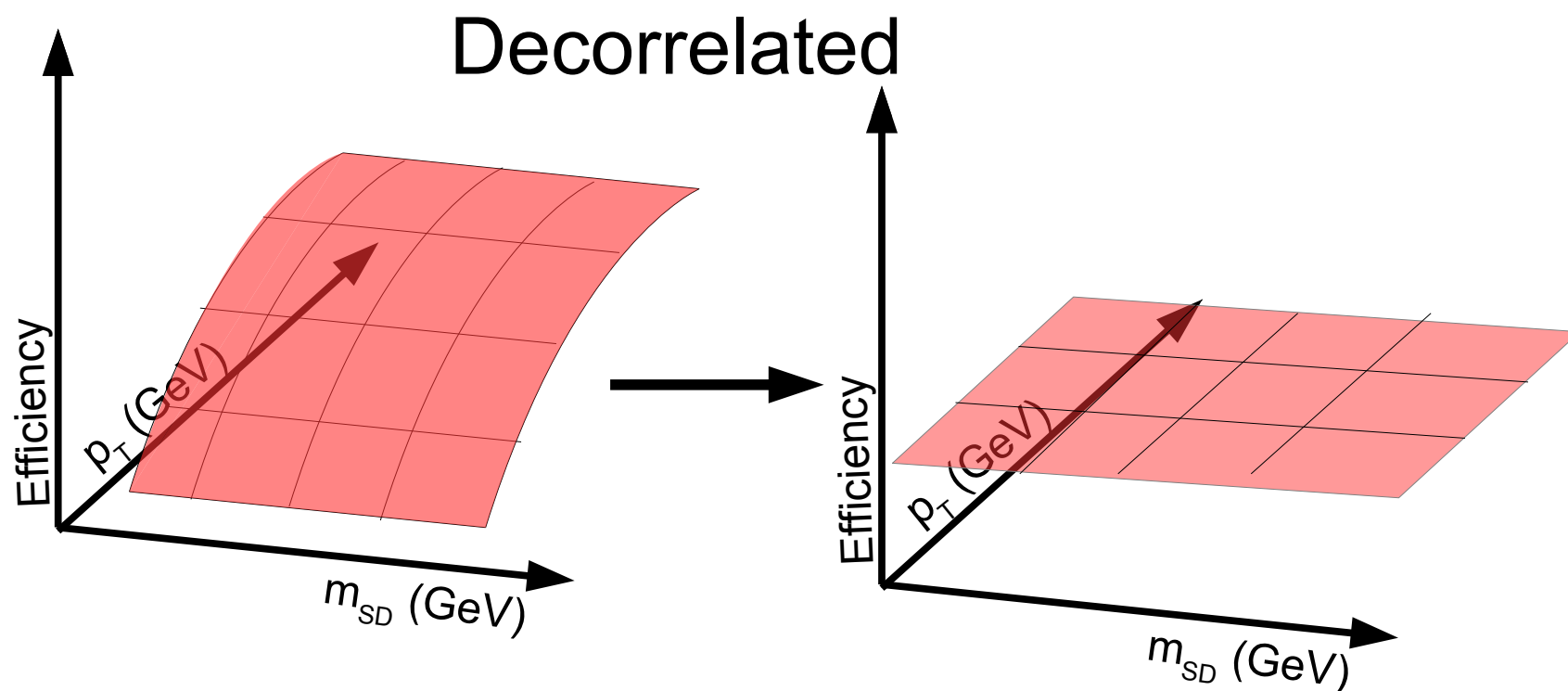
Fat jet + ISR



For this plot we invented a new substructure var ← [arXiv:1603.00027](https://arxiv.org/abs/1603.00027)

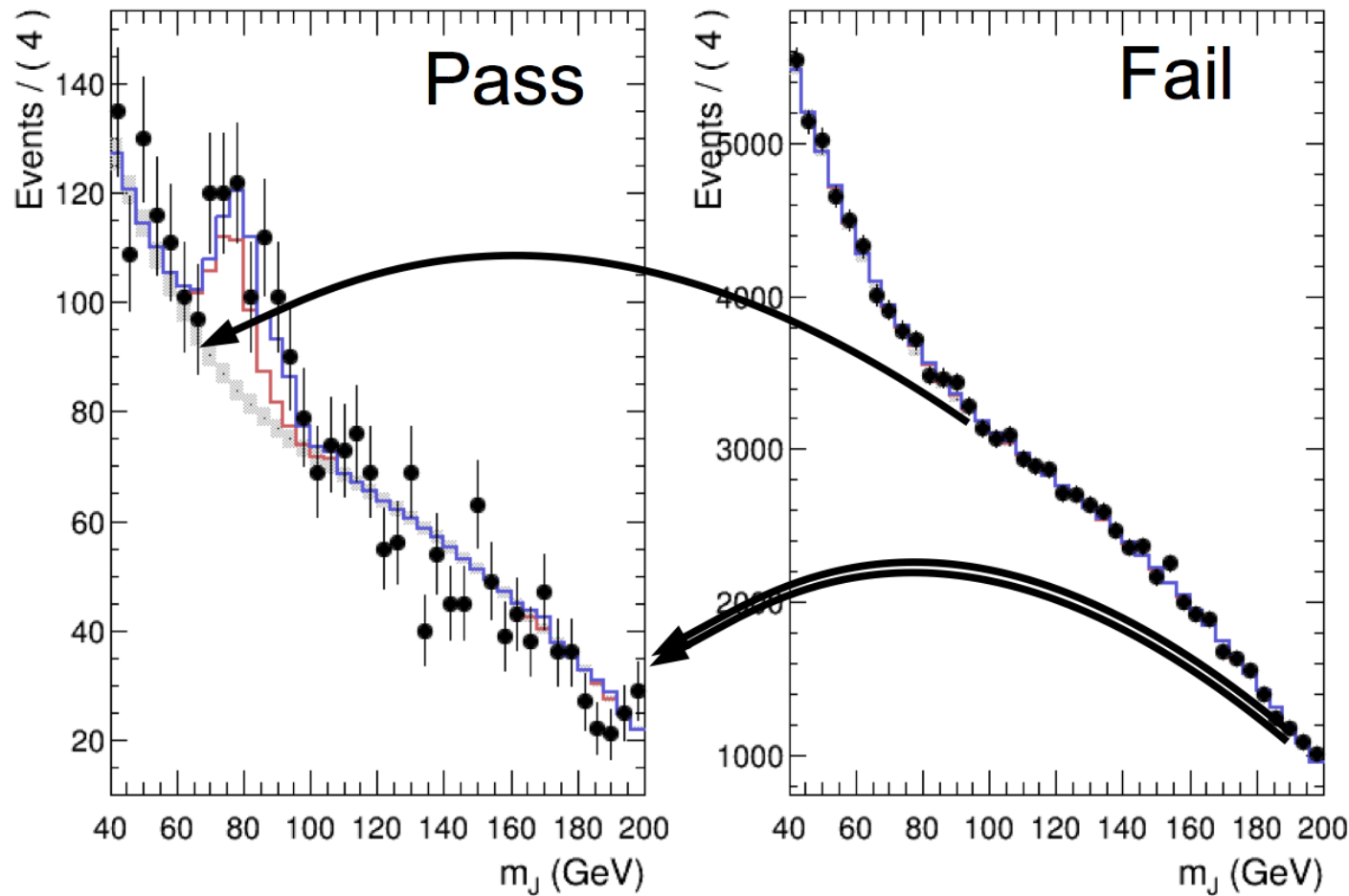
Design a transform to decorrelate against mass and p_T

Decorrelating avoid mass sculpting
allows us to cut tighter



Fit both Pass and Fail
Pass QCD = fail $\epsilon(\rho, p_T)$

General Idea

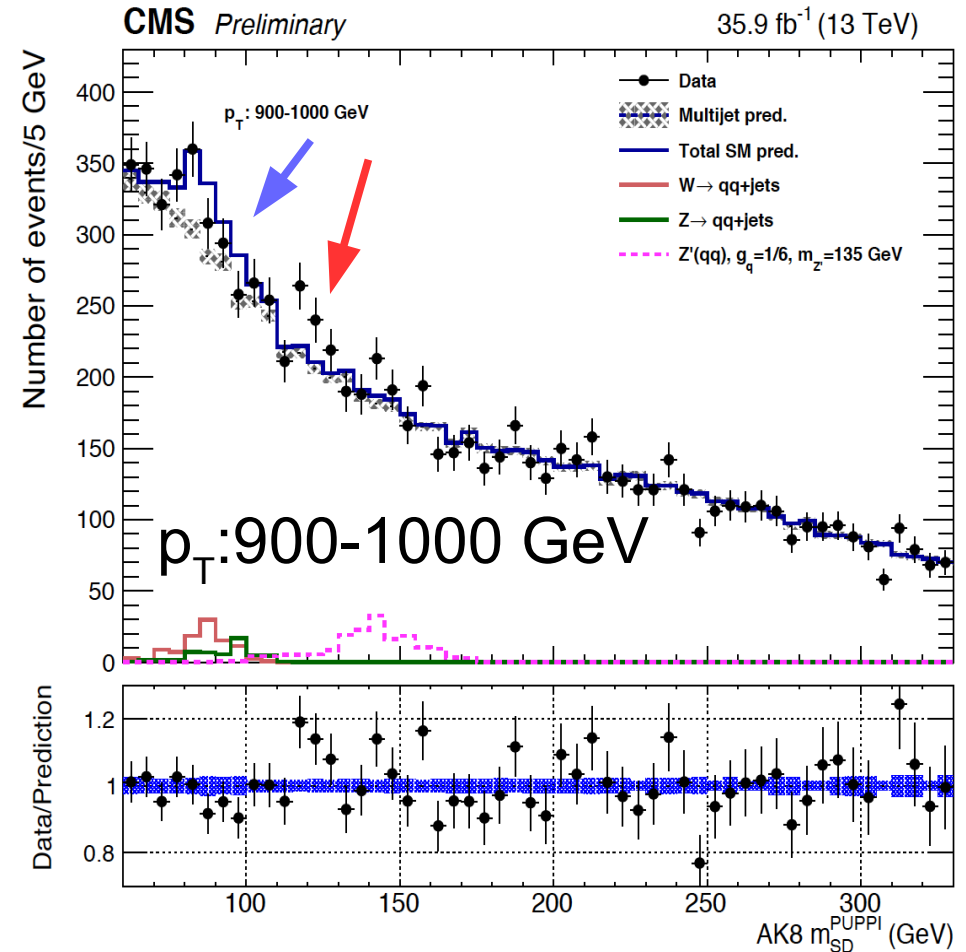
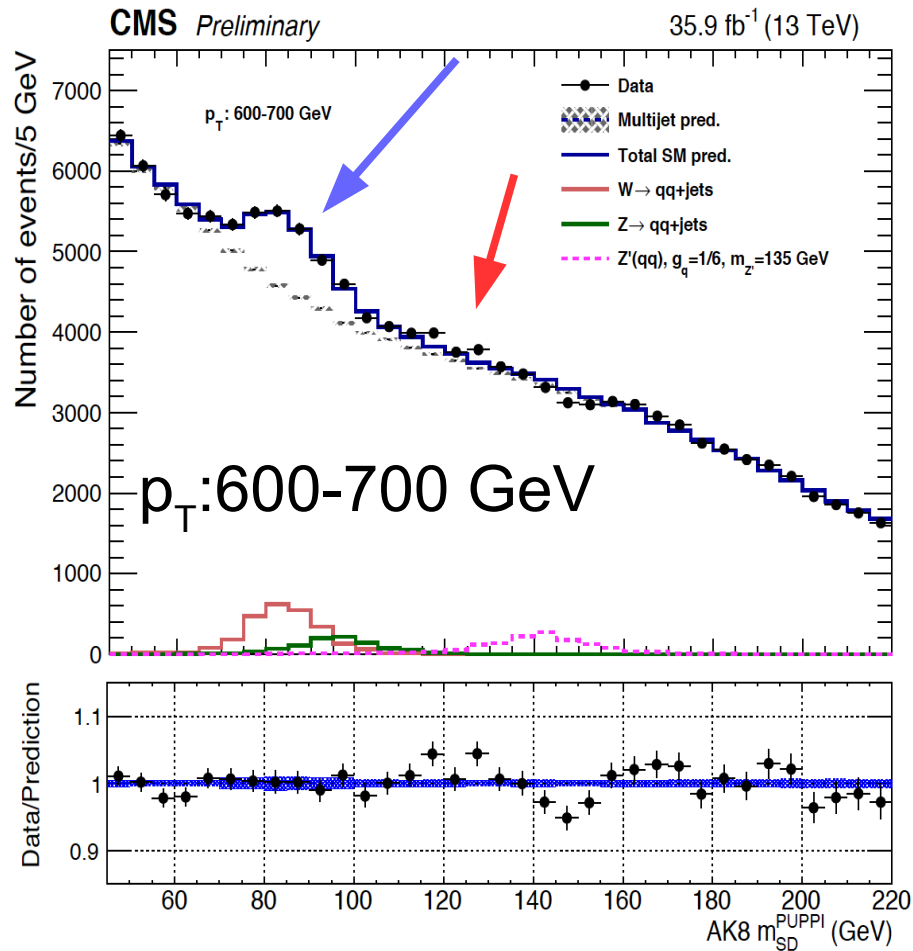


Fist iteration did this in two steps

Next iteration all in one fit (thanks to monojet technology)

$\text{Bin}_i^{\text{Pass}} = \epsilon_0(1 + a p_T + b p + c p^2 + \dots)$ $\text{Bin}_i^{\text{Fail}}$ (add orders til f-test ok)

What we see

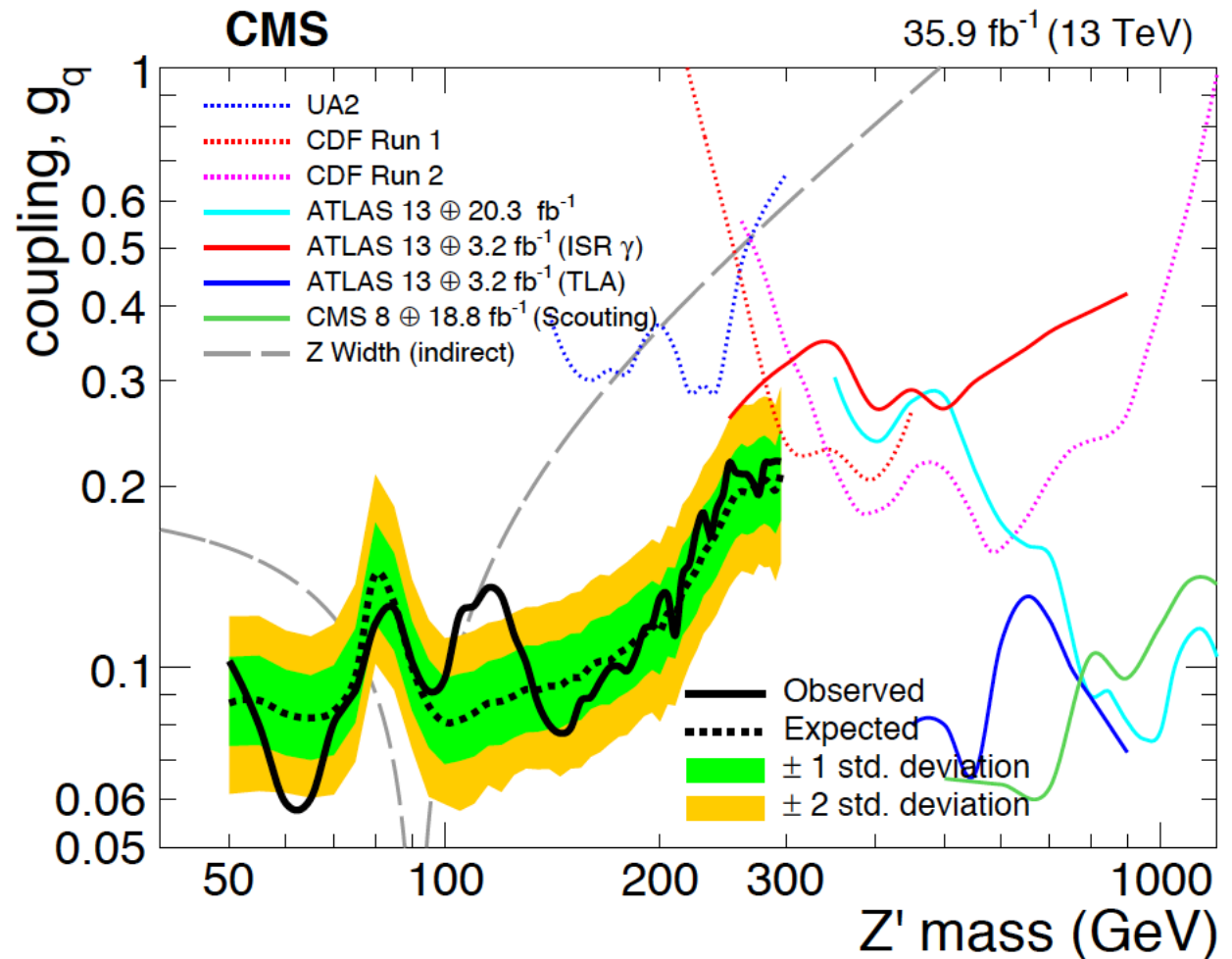
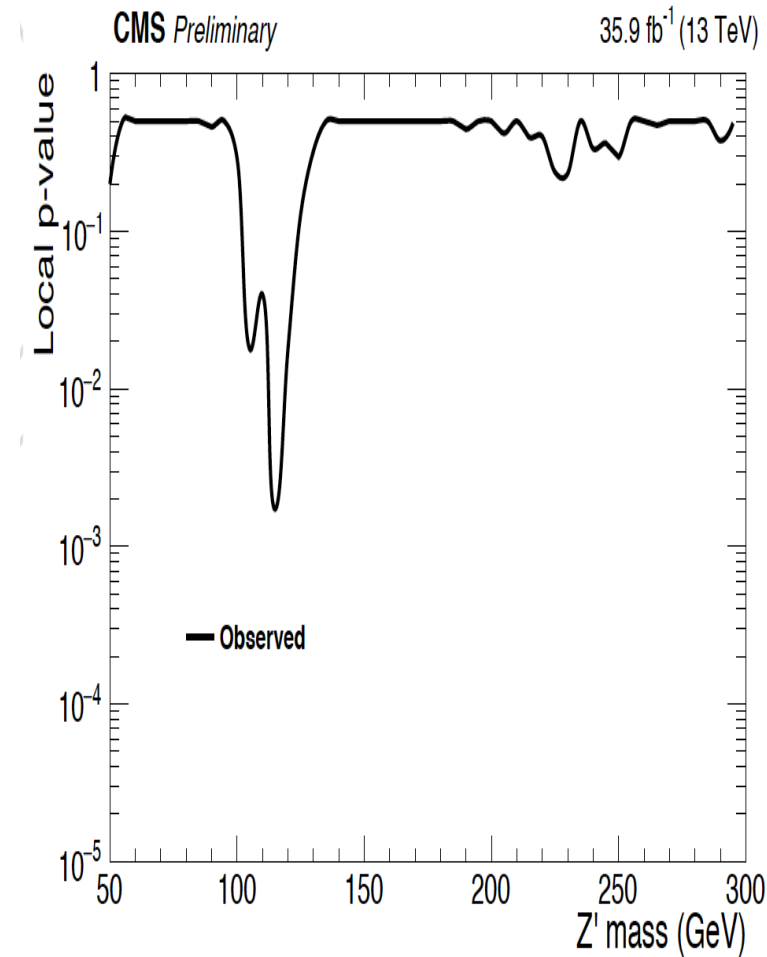


Nice $W/Z \rightarrow \text{Jet}$ peak

Excess present near 115-120 GeV

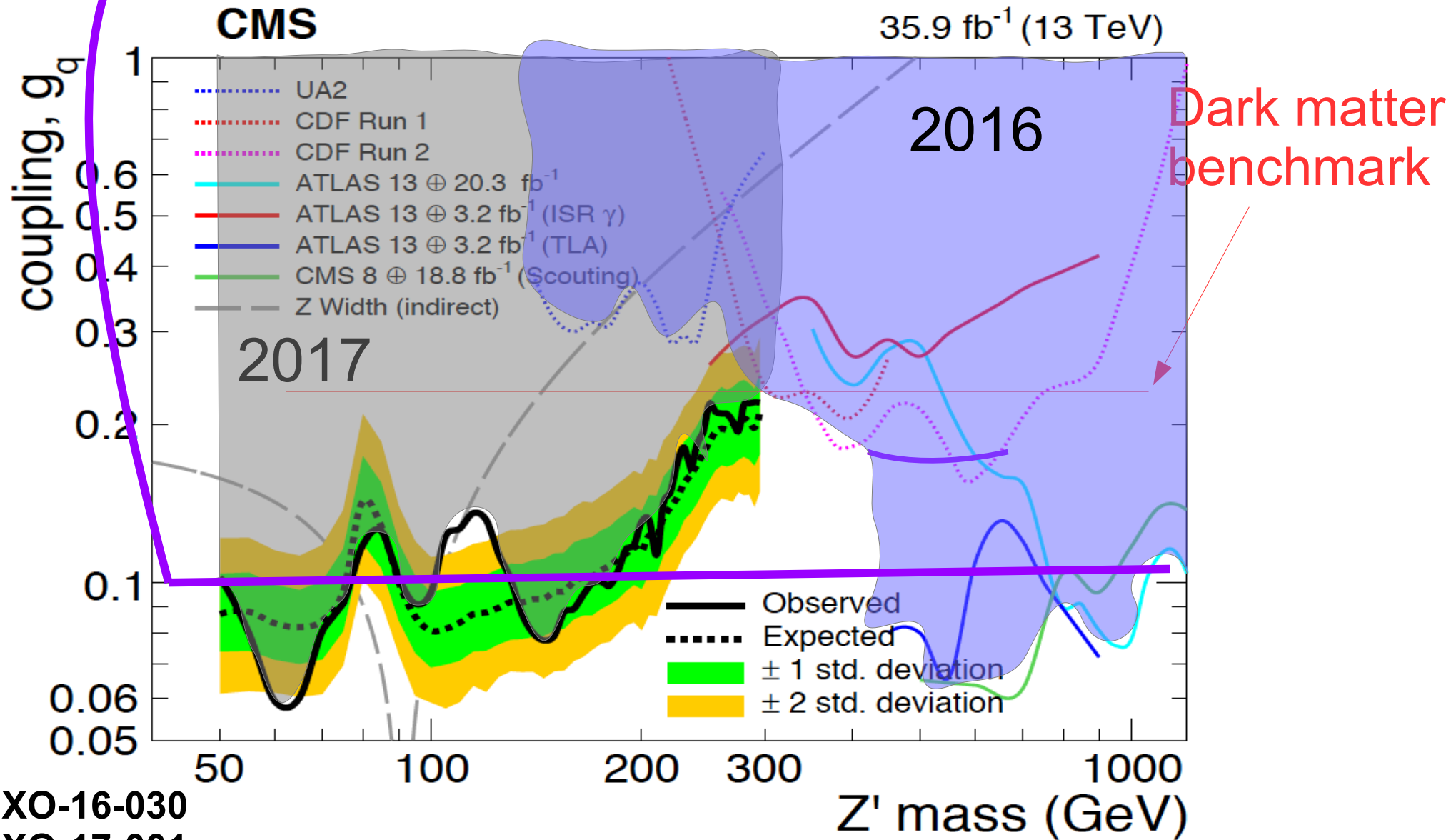
Precision level analysis

Combining



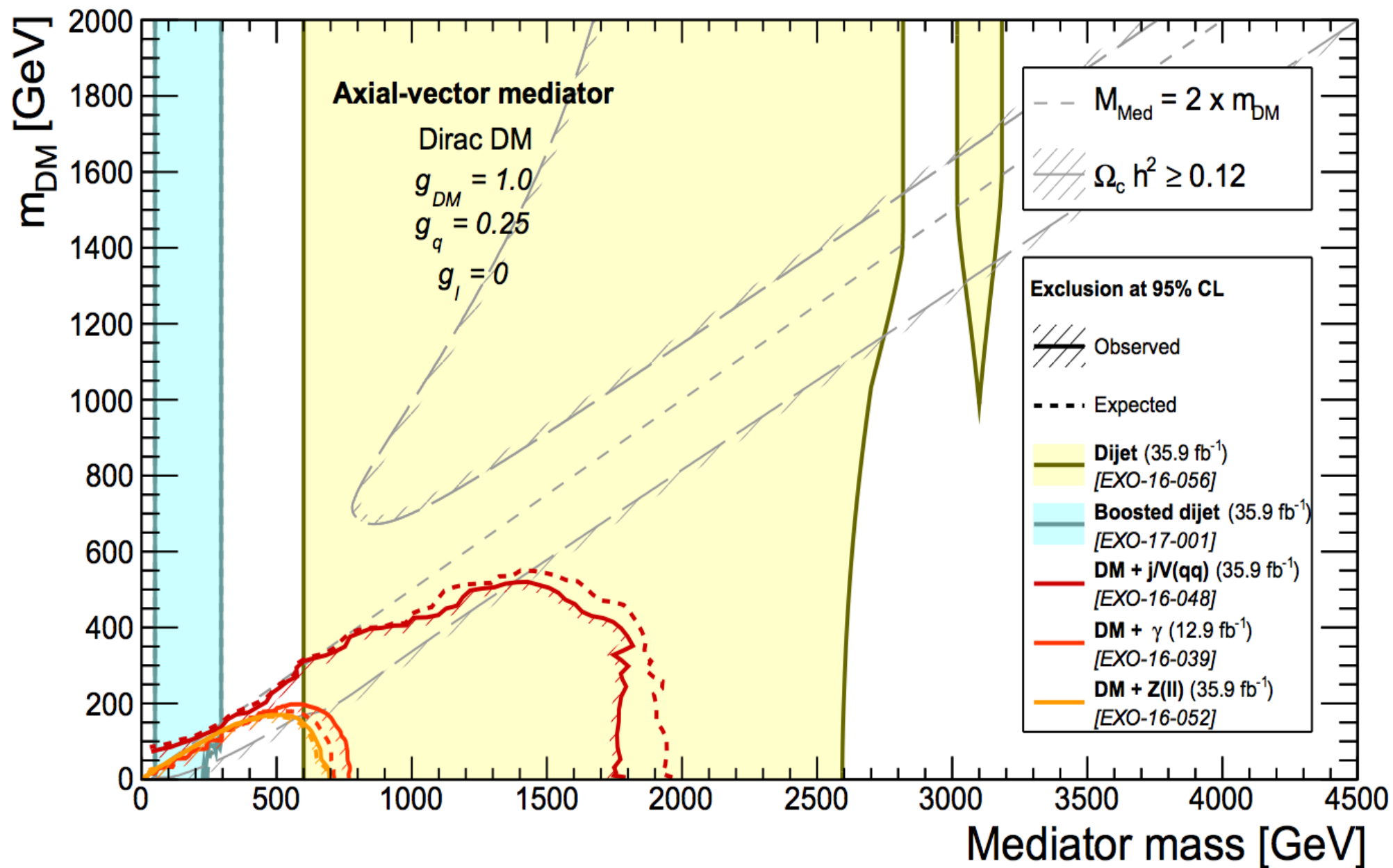
Excess with 2.9 (2.2 global) significance
 Visible in last bin

Interesting Region from earlier slides



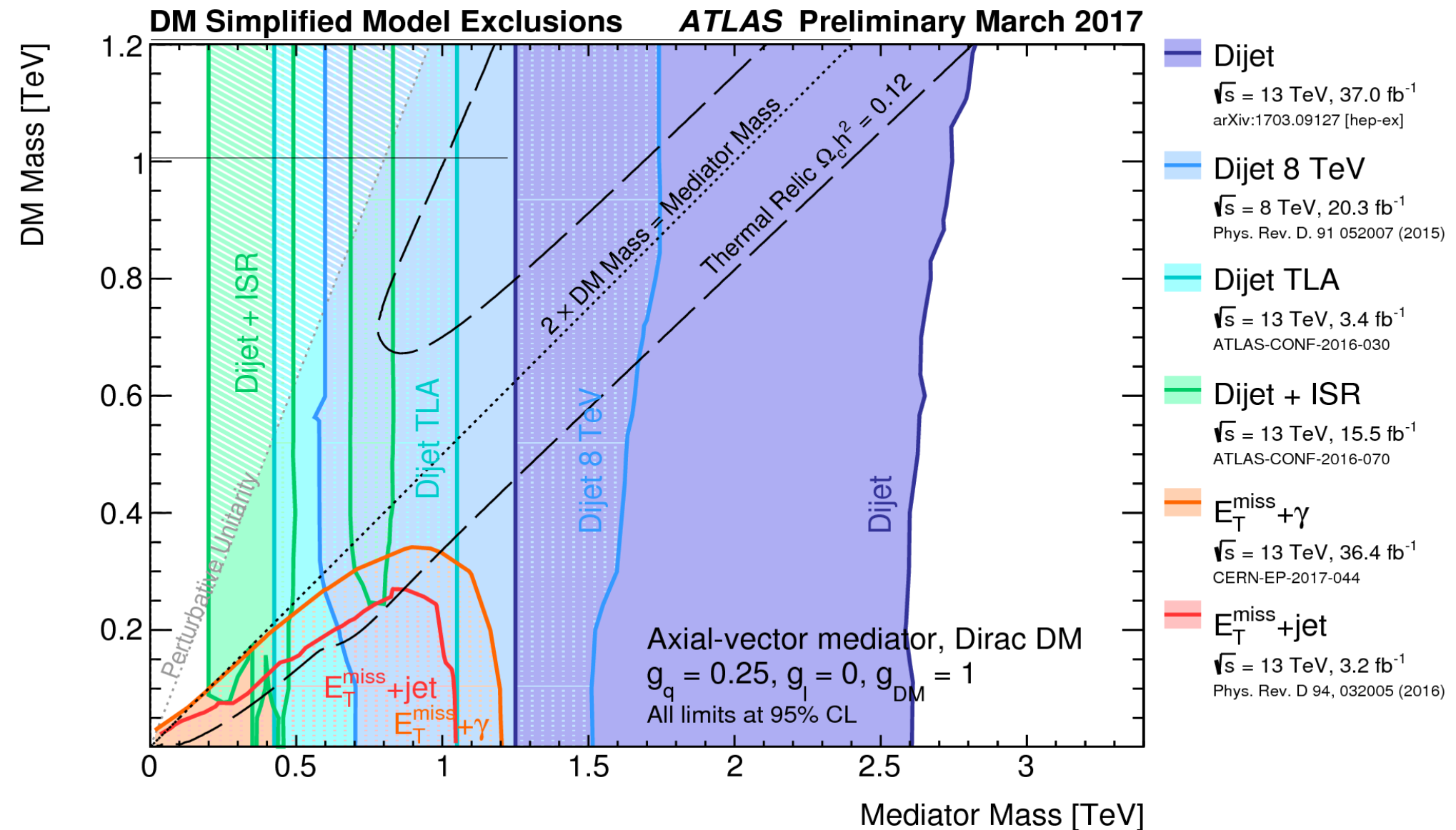
CMS Preliminary

LHCP 2017



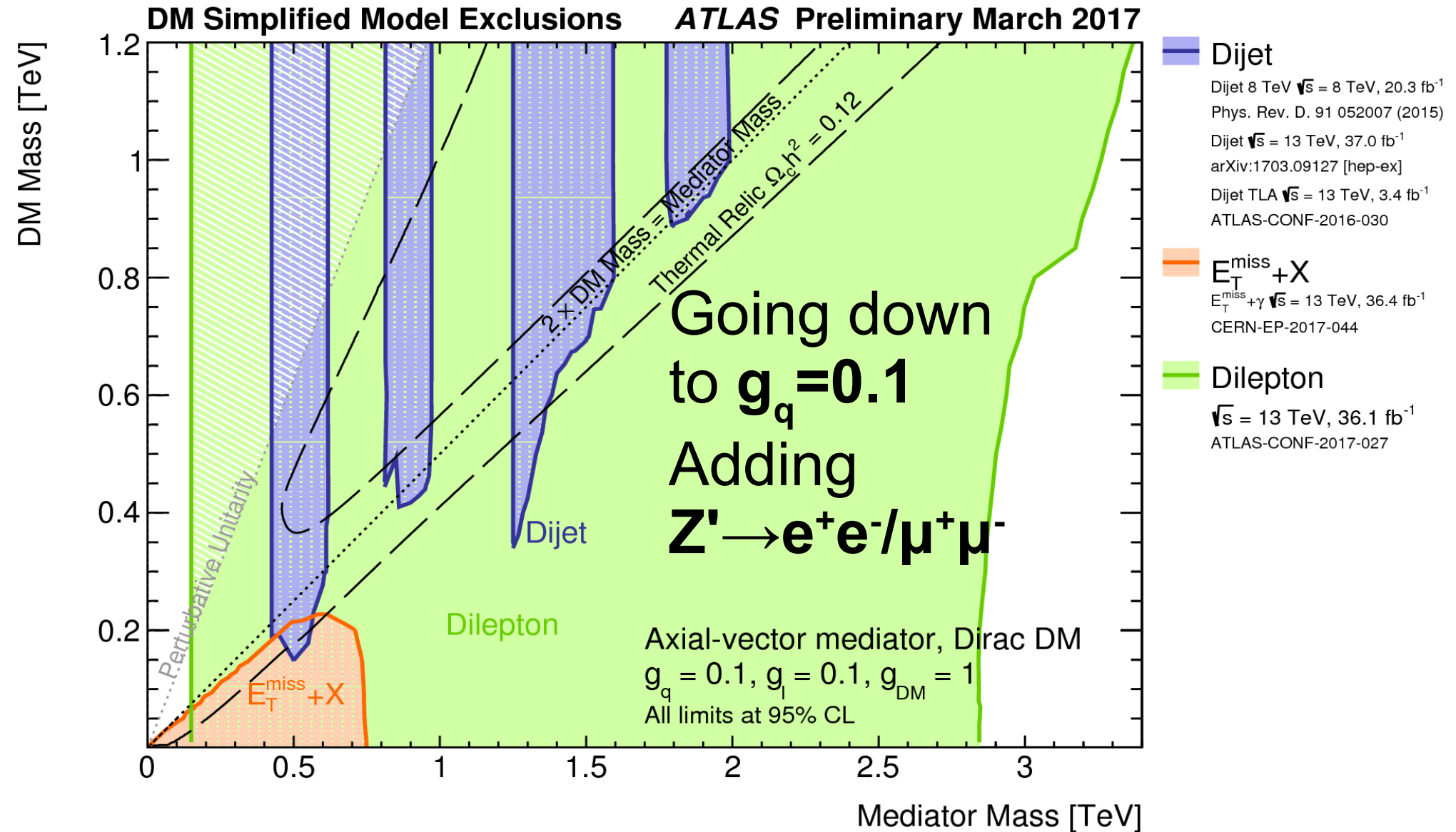
ATLAS Results : Split even further

<https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CombinedSummaryPlots/EXOTICS/index.html#>

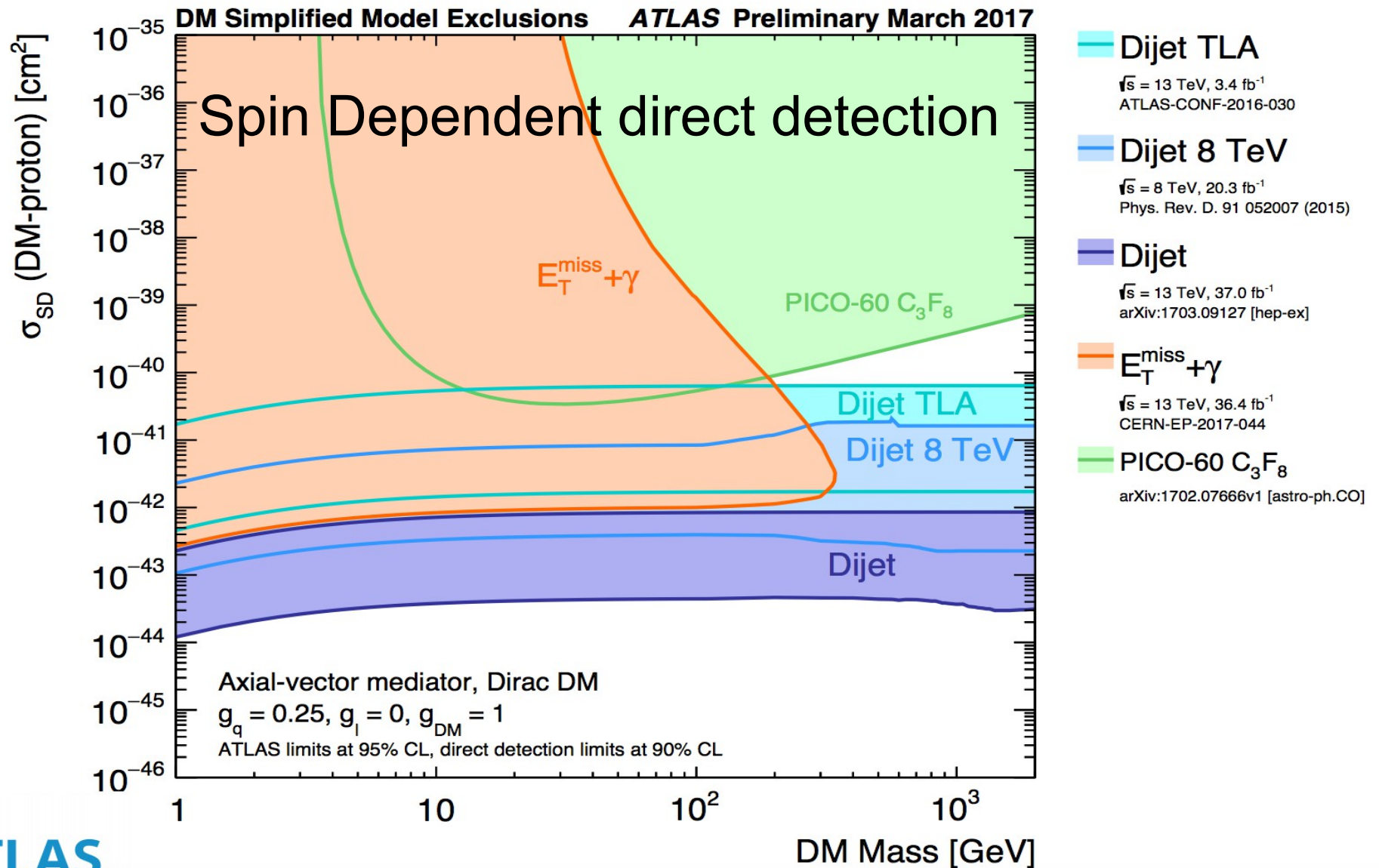


Now with adding the leptons

<https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CombinedSummaryPlots/EXOTICS/index.html#>



Direct Detection

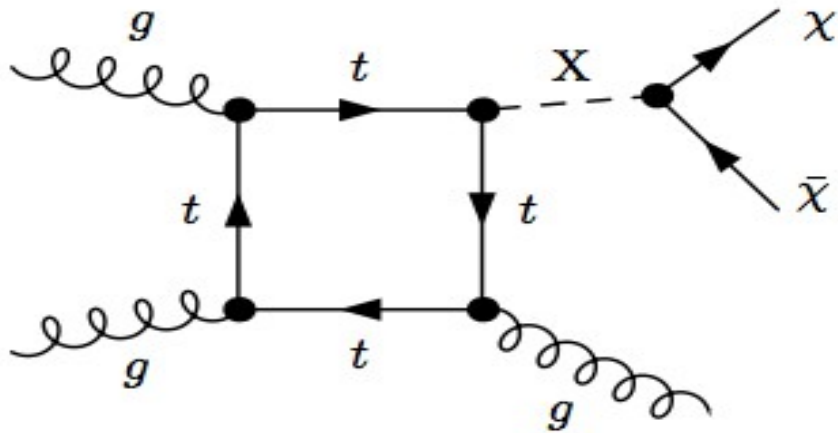


Spin 0

What can you do with Spin 0?

Basic production is gluon fusion

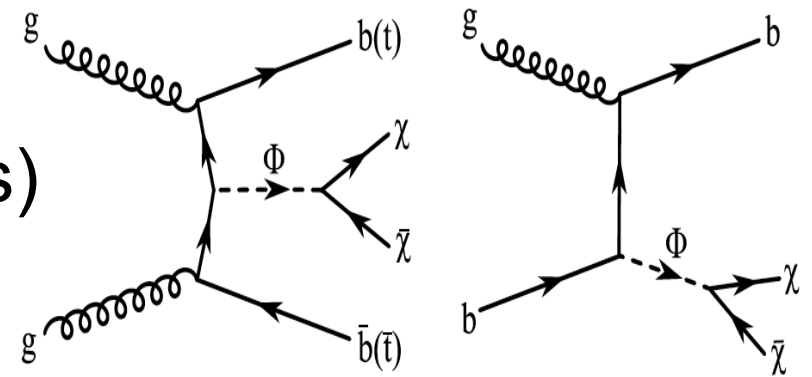
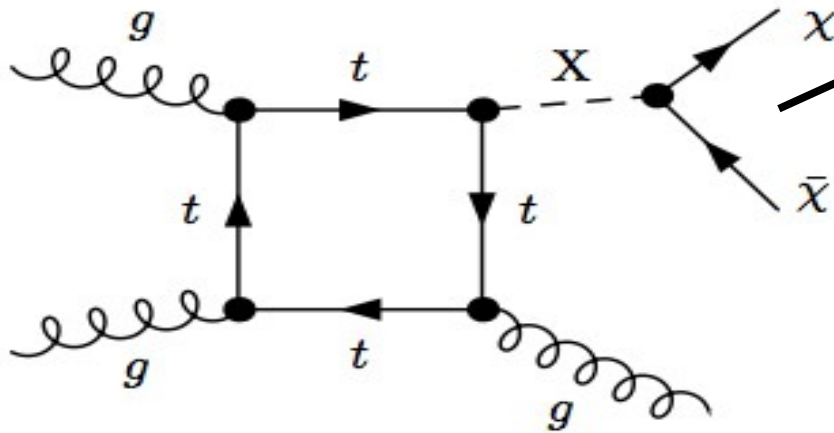
Amplitude is double for pseudoscalar mediator



What can you do with Spin 0?

Heavy flavor channels

Can be added with (same couplings)

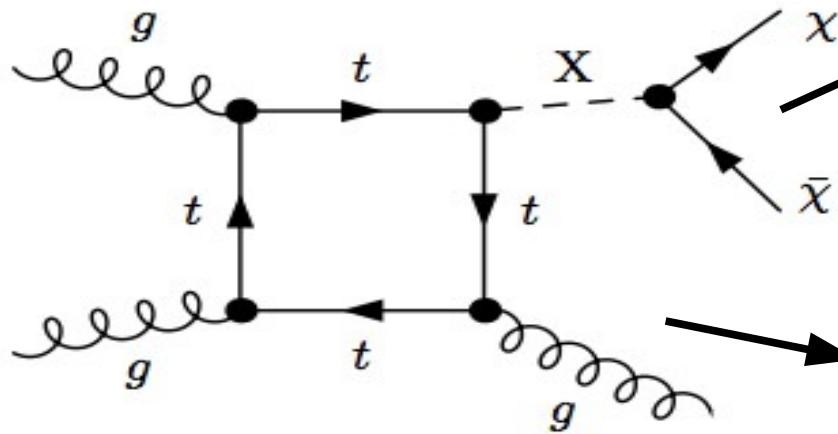


Heavy Flavor
final states

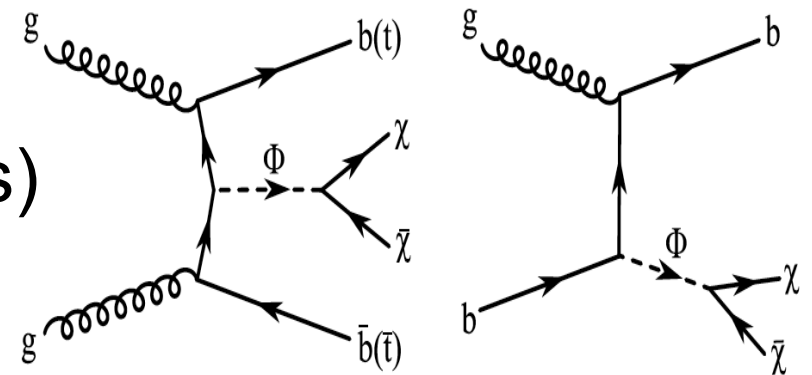
What can you do with Spin 0?

Heavy flavor channels

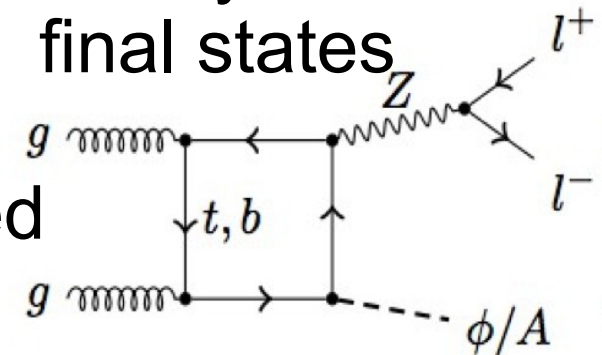
Can be added with (same couplings)



Loop induced
Z boson



Heavy Flavor
final states

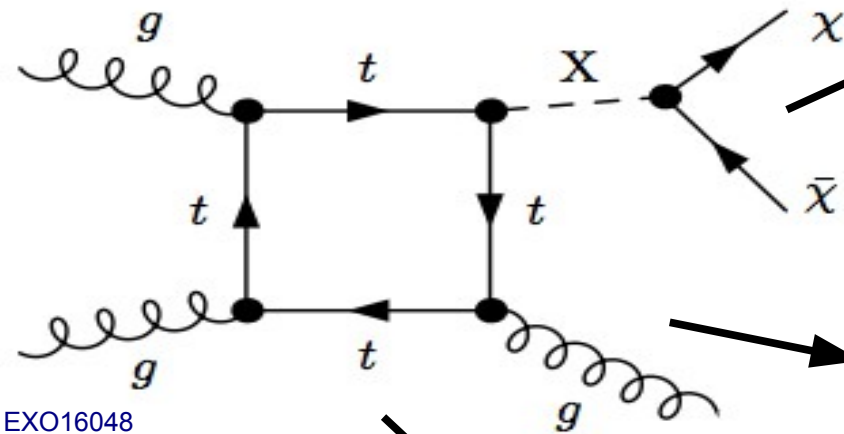


EXO16028

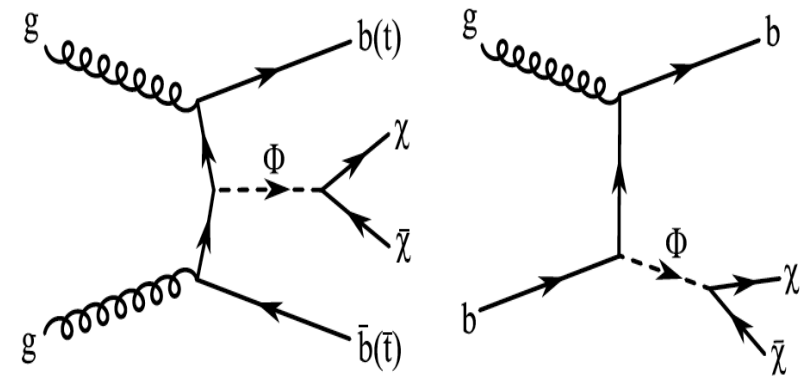
EXO16052

What can you do with Spin 0?

Big Assumption :
No mixing w/Higgs

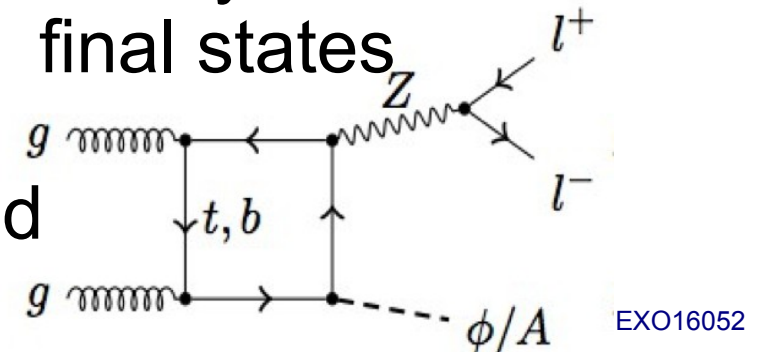


Loop induced
Z boson

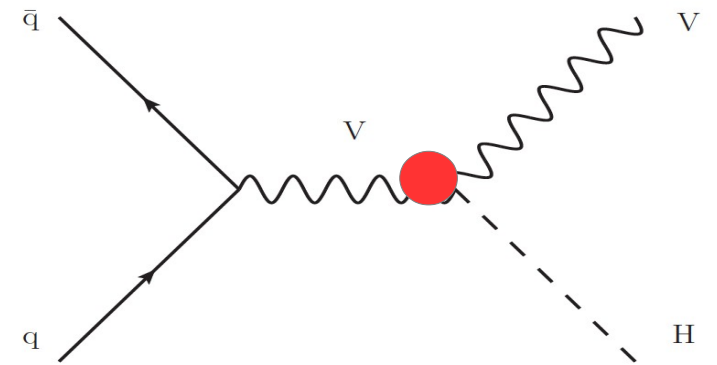
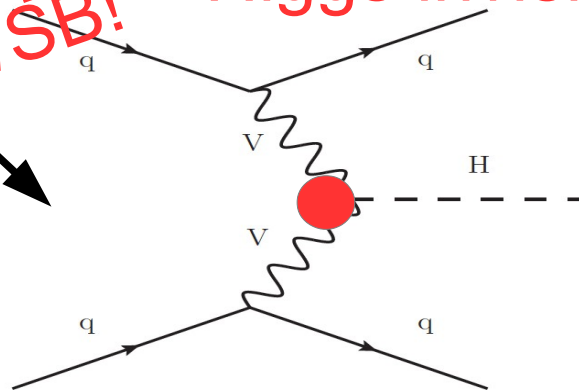


Heavy Flavor
final states

EXO16028

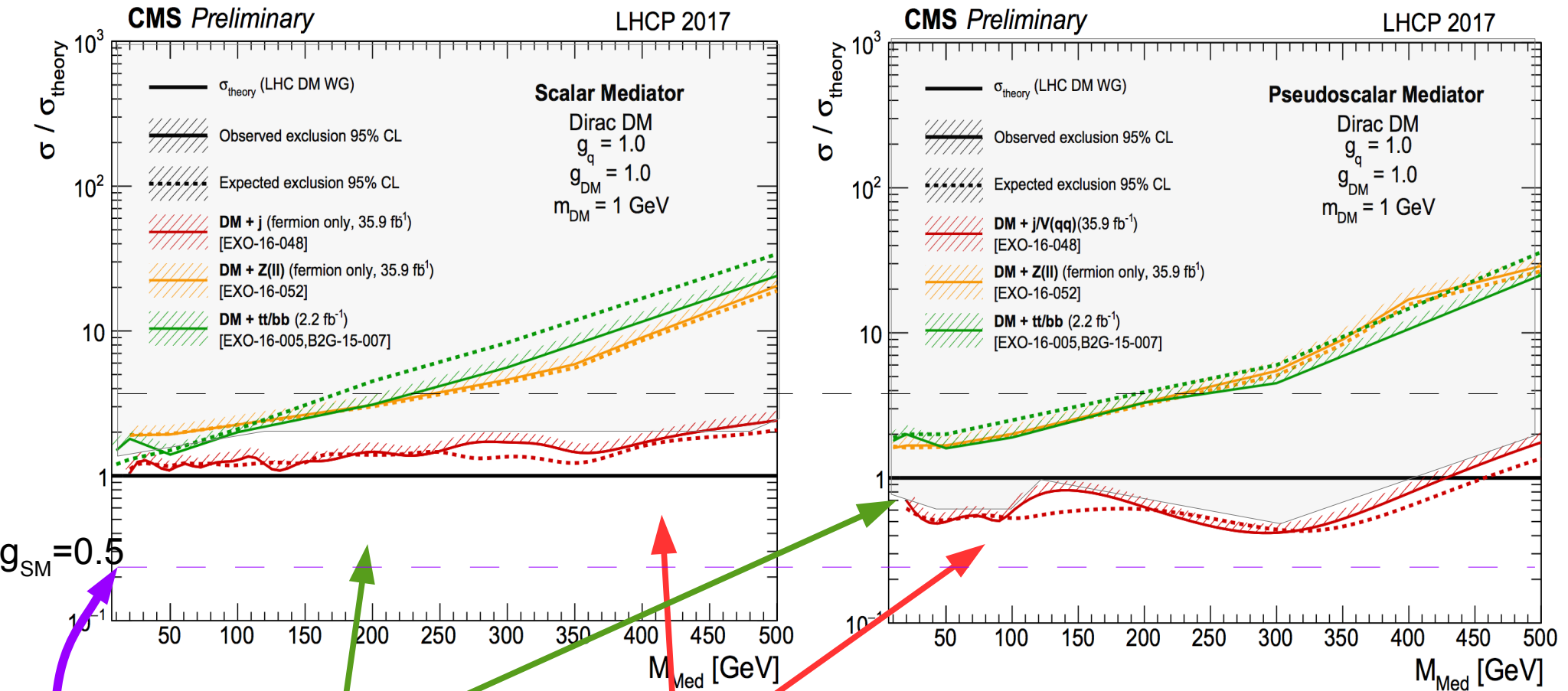


Applying EWSB! Higgs invisible or Scalar w/EWSB



No EWSB

Comparing all channels

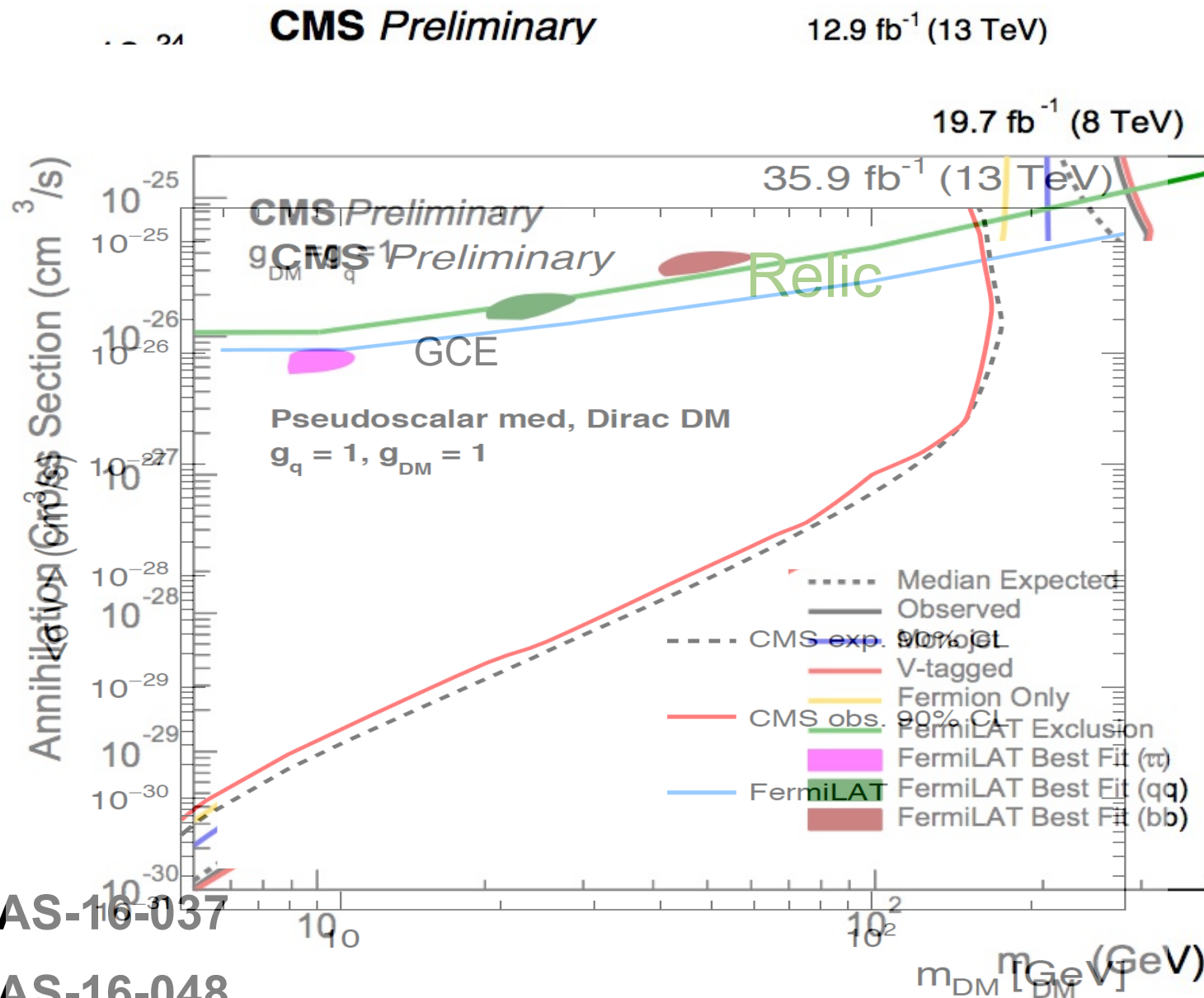


$tt + \text{DM}(2\text{fb}^{-1})$ and **monojet** drive the combination

Not far from an intermediate benchmark of $g_{\text{SM}} = 0.5$

Whats the impact?

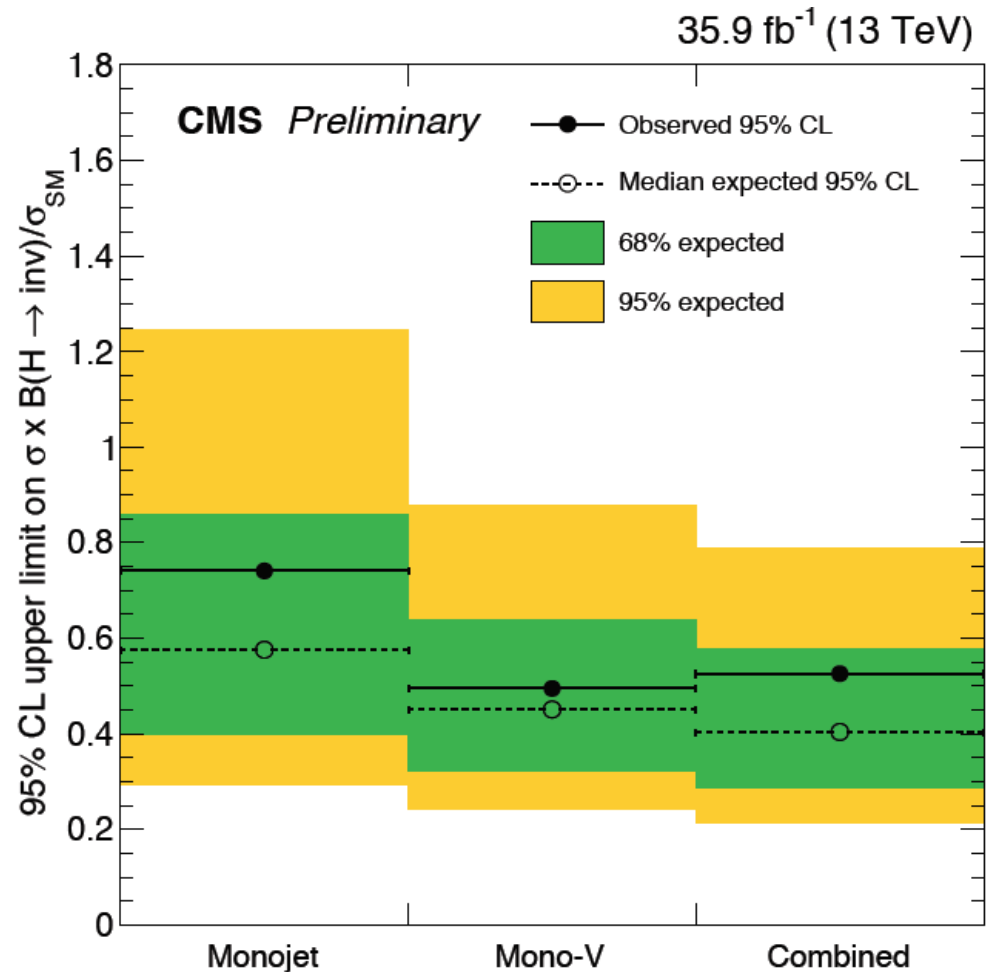
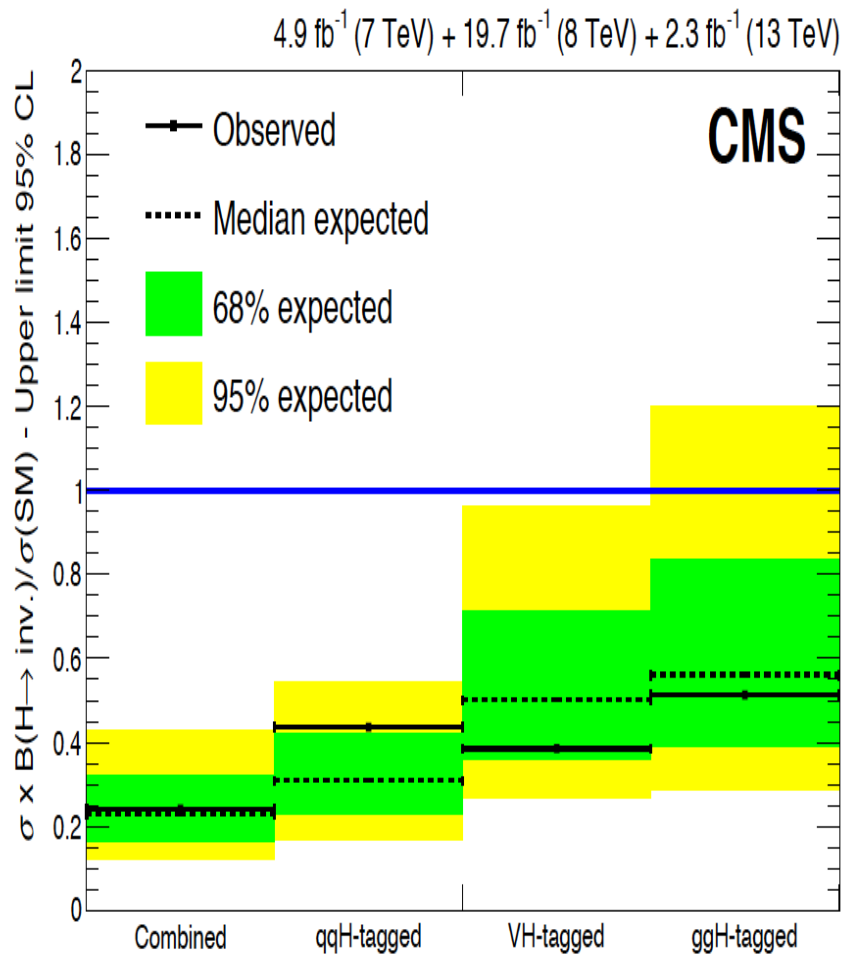
No EWSB



EWSB

At the Higgs mass

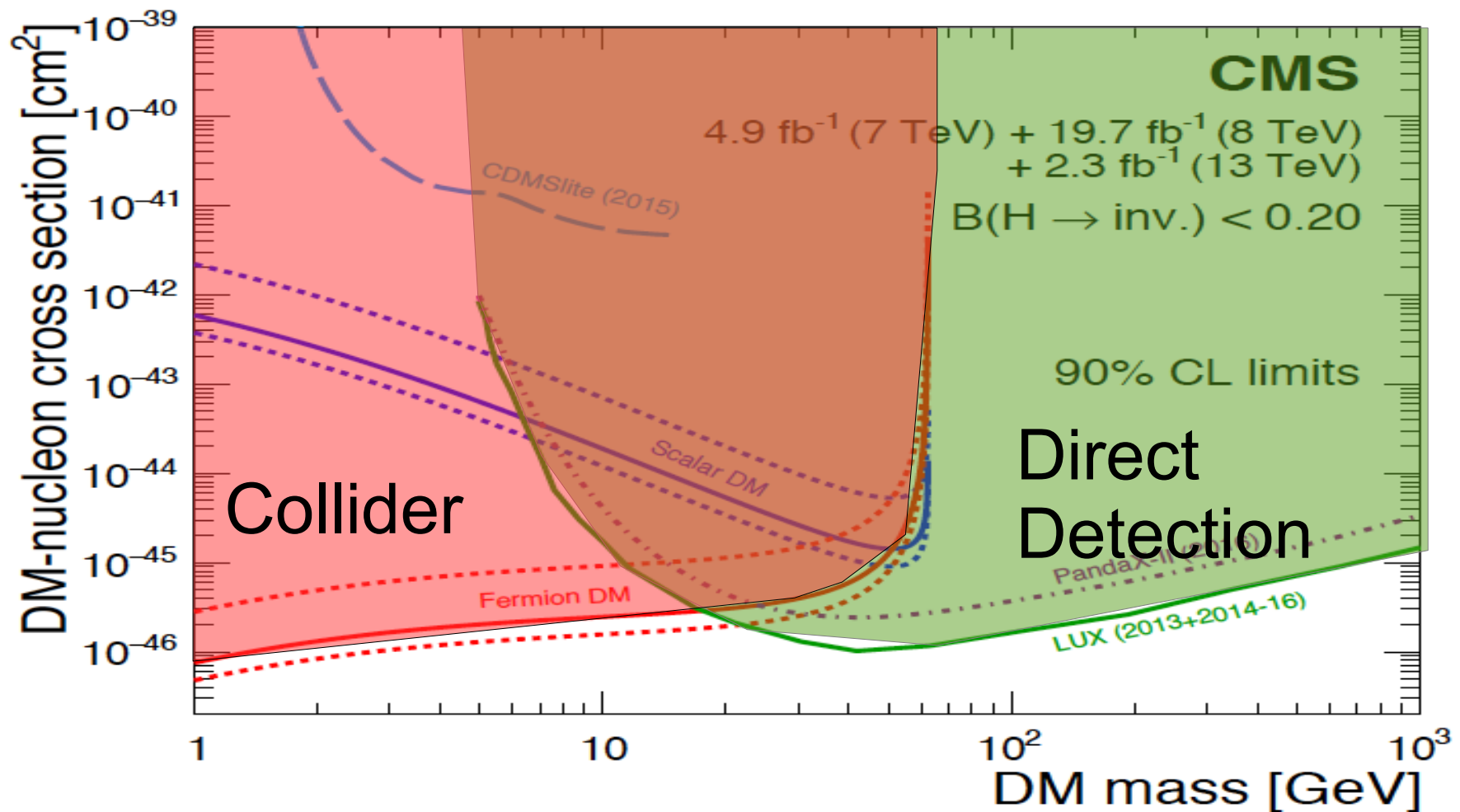
- This model is the same as Higgs invisible search



$\text{BR}(H \rightarrow \text{Inv}) < 24\% \text{ (CMS)} \quad 25\% \text{ (ATLAS)}$

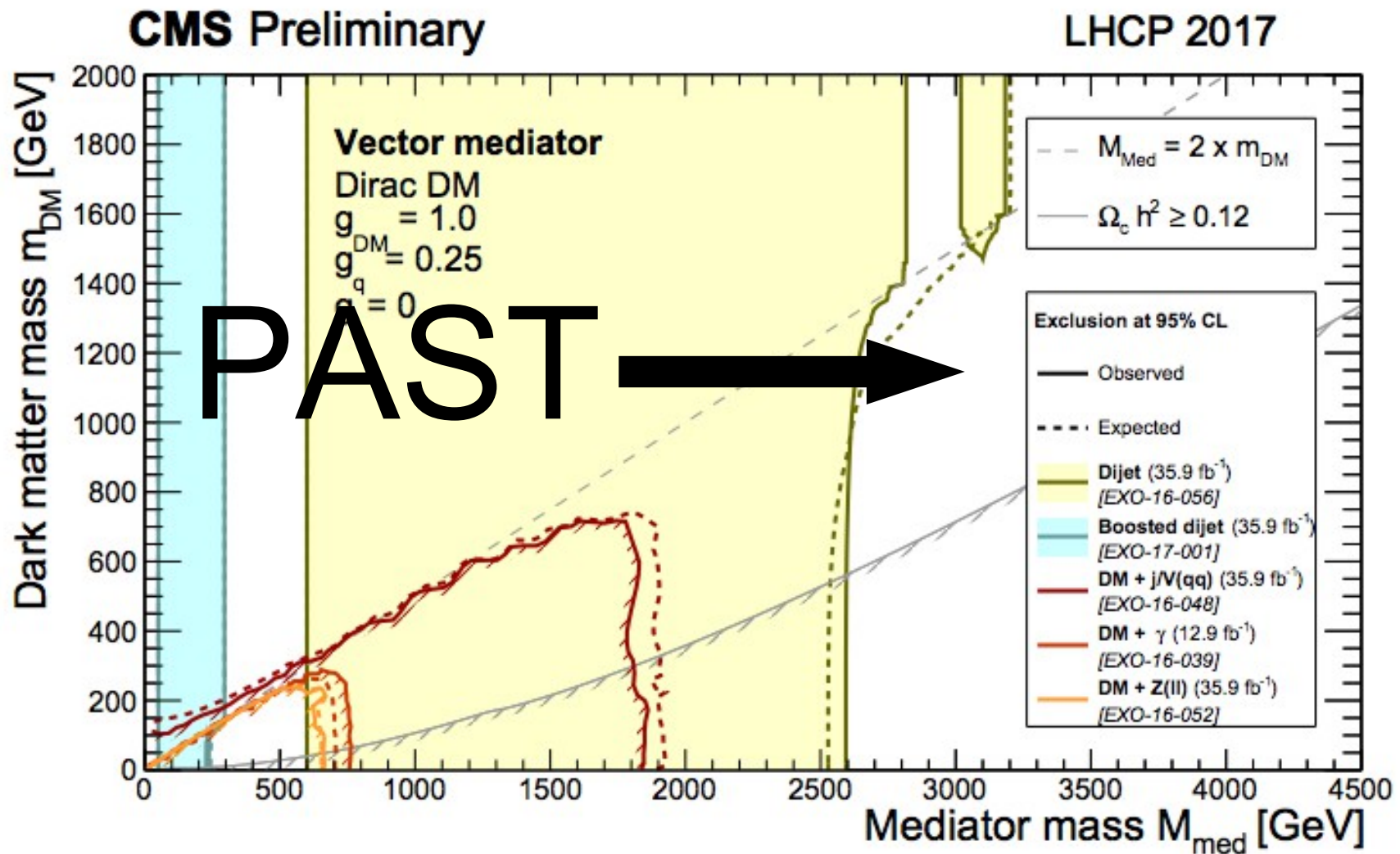
At the Higgs mass

- Higgs to invisible :
 - Direct detection and collider are head to head



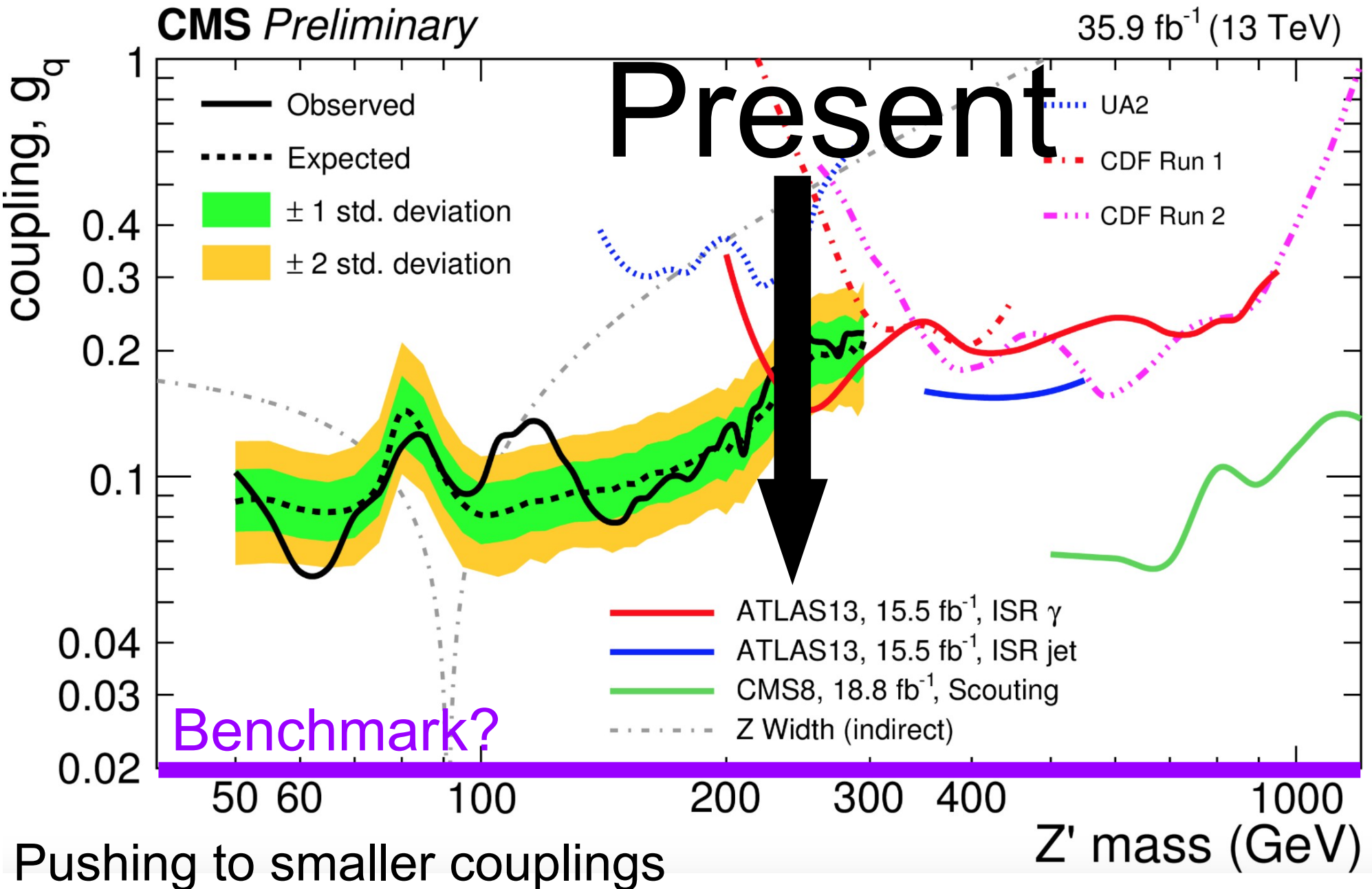
$\text{BR}(H \rightarrow \text{Inv}) < 24\% \text{ (CMS)} \quad 25\% \text{ (ATLAS)}$

Conclusions



Pushing to higher masses

Conclusions

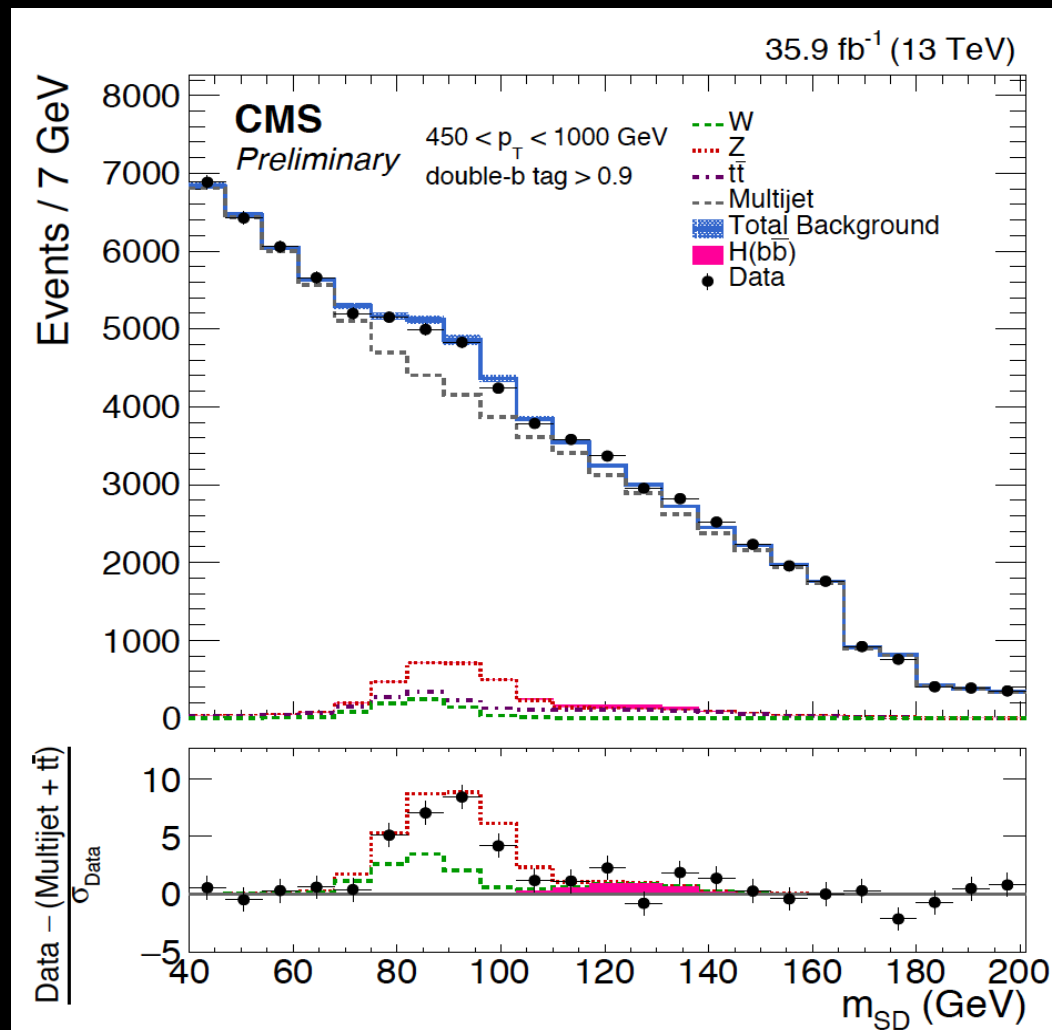




Forever

SM

BSM

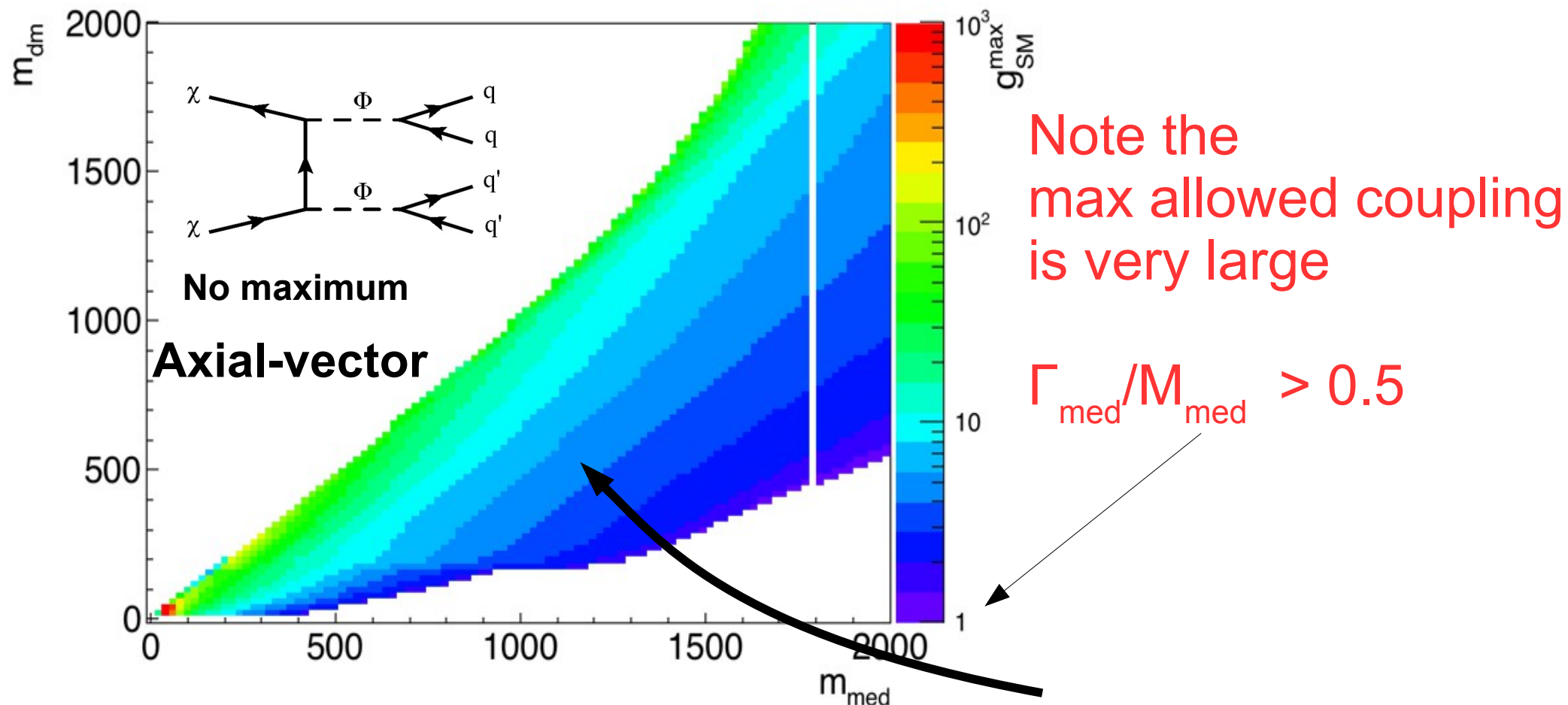


Thanks!

See Rainer Mankel's talk at Blois next week

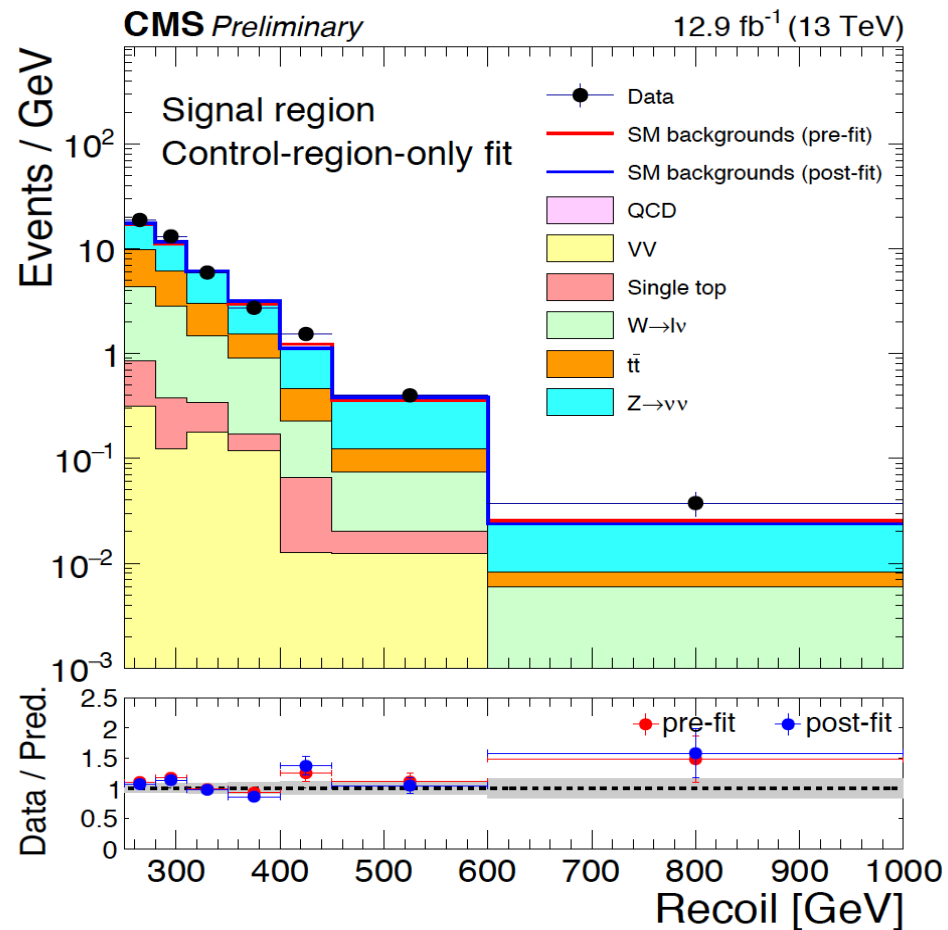
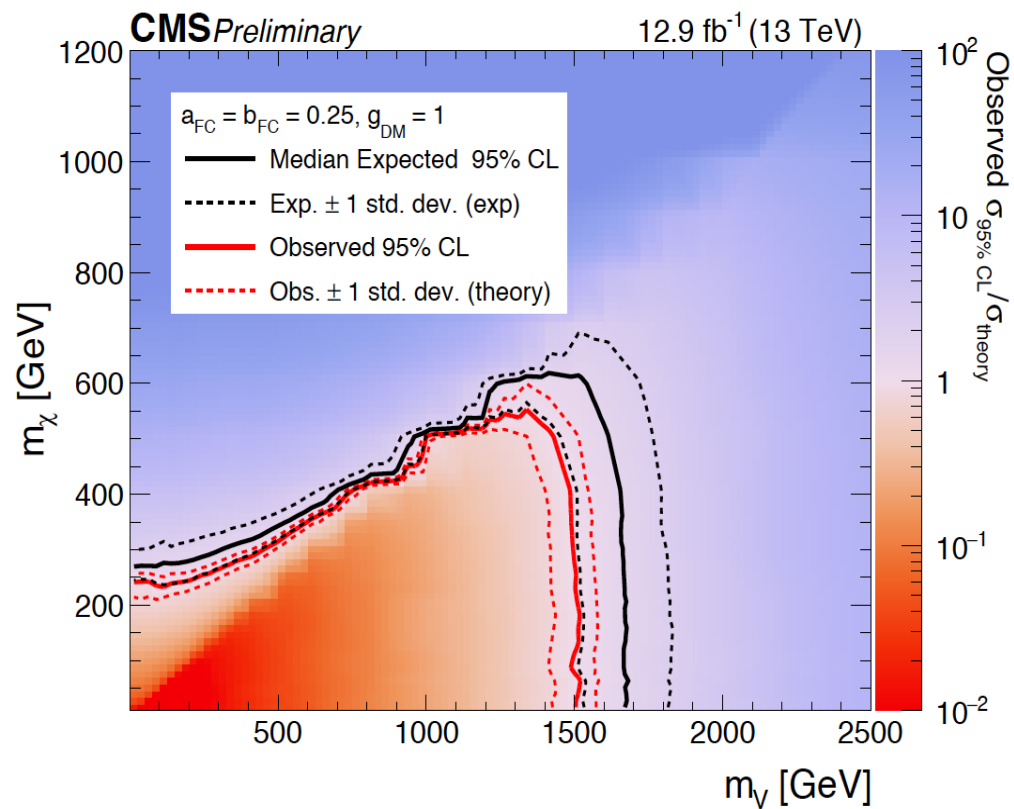
Max couplings

- Now consider the maximum coupling
 - In this case we fix $g_{\text{DM}}=1$ (product $g_q g_{\text{DM}}$ defines bound)



Given coupling so large we will not focus on this upper bound

CMS Monotop





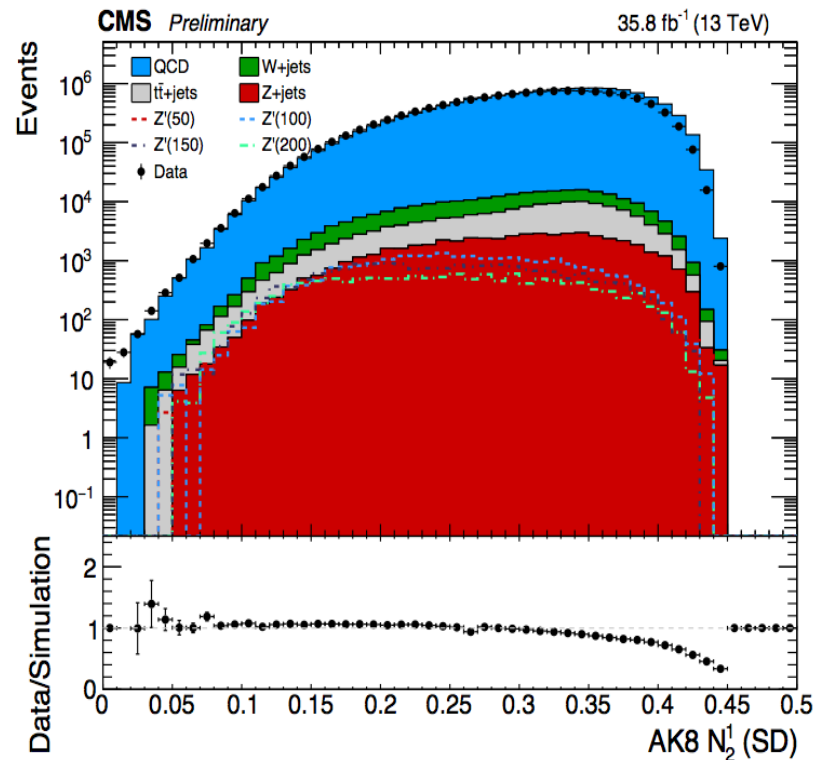
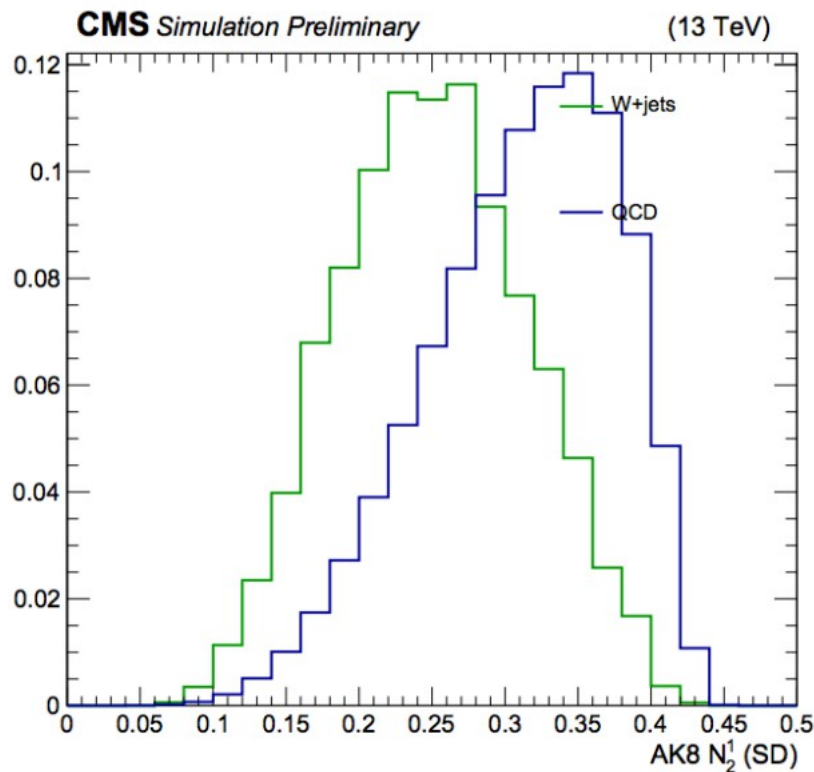
BRAND NEW

New Substructure Observables

$$N_2(\beta) = \frac{2e_3^\beta}{(e_2^\beta)^2} \quad 1e_2^\beta = e_2^\beta = \sum_{1 \leq i < j \leq n_l} z_i z_j \Delta R_{ij}^\beta \quad z_i \equiv \frac{p_{Ti}}{\sum_{j \in \text{jet}} p_{Tj}}$$

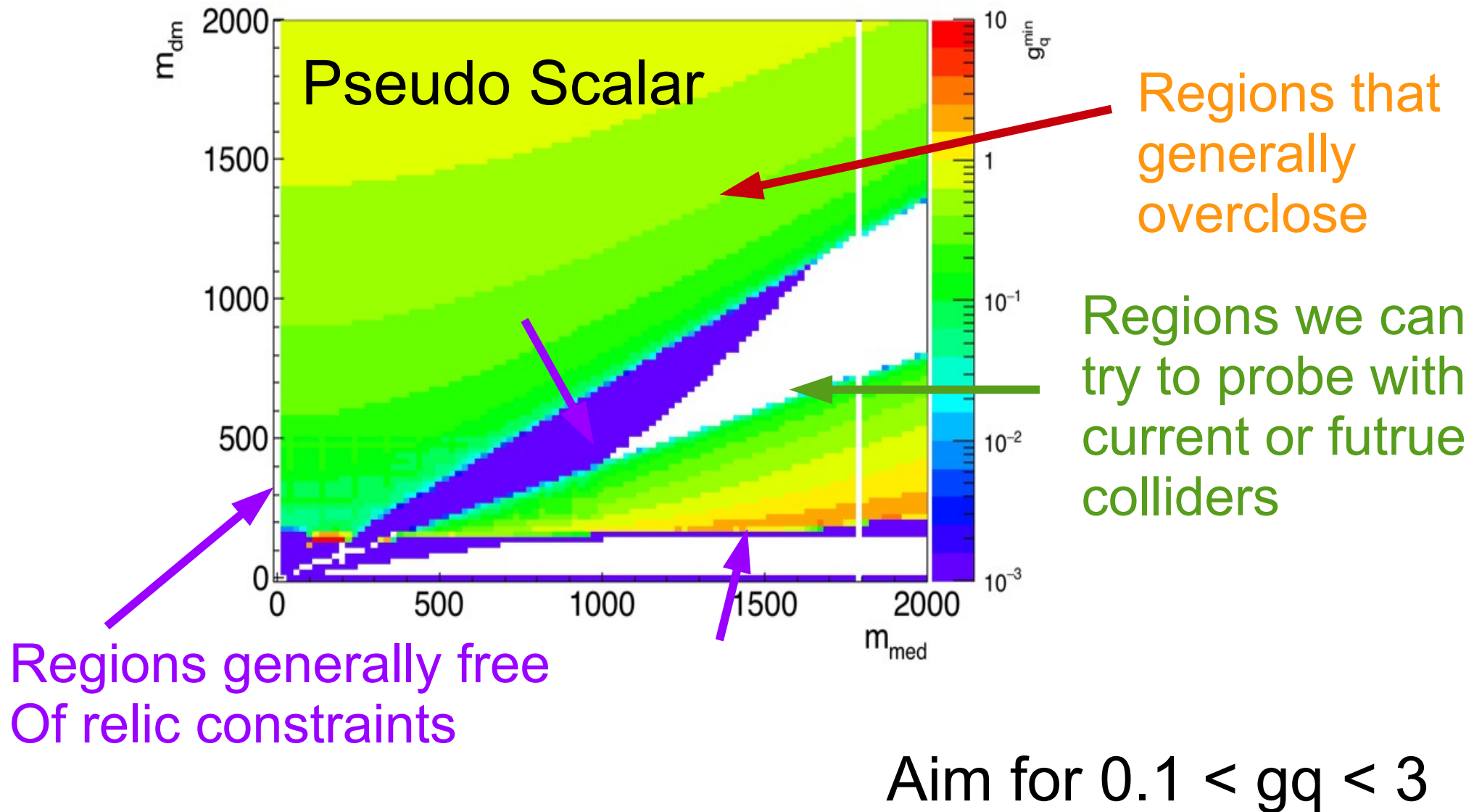
$$2e_3^\beta = \sum_{1 \leq i < j < k \leq n_l} z_i z_j z_k \min\{\Delta R_{ij}^\beta \Delta R_{ik}^\beta, \Delta R_{ij}^\beta \Delta R_{jk}^\beta, \Delta R_{ik}^\beta \Delta R_{jk}^\beta\}$$

Using AK8 PUPPI jets



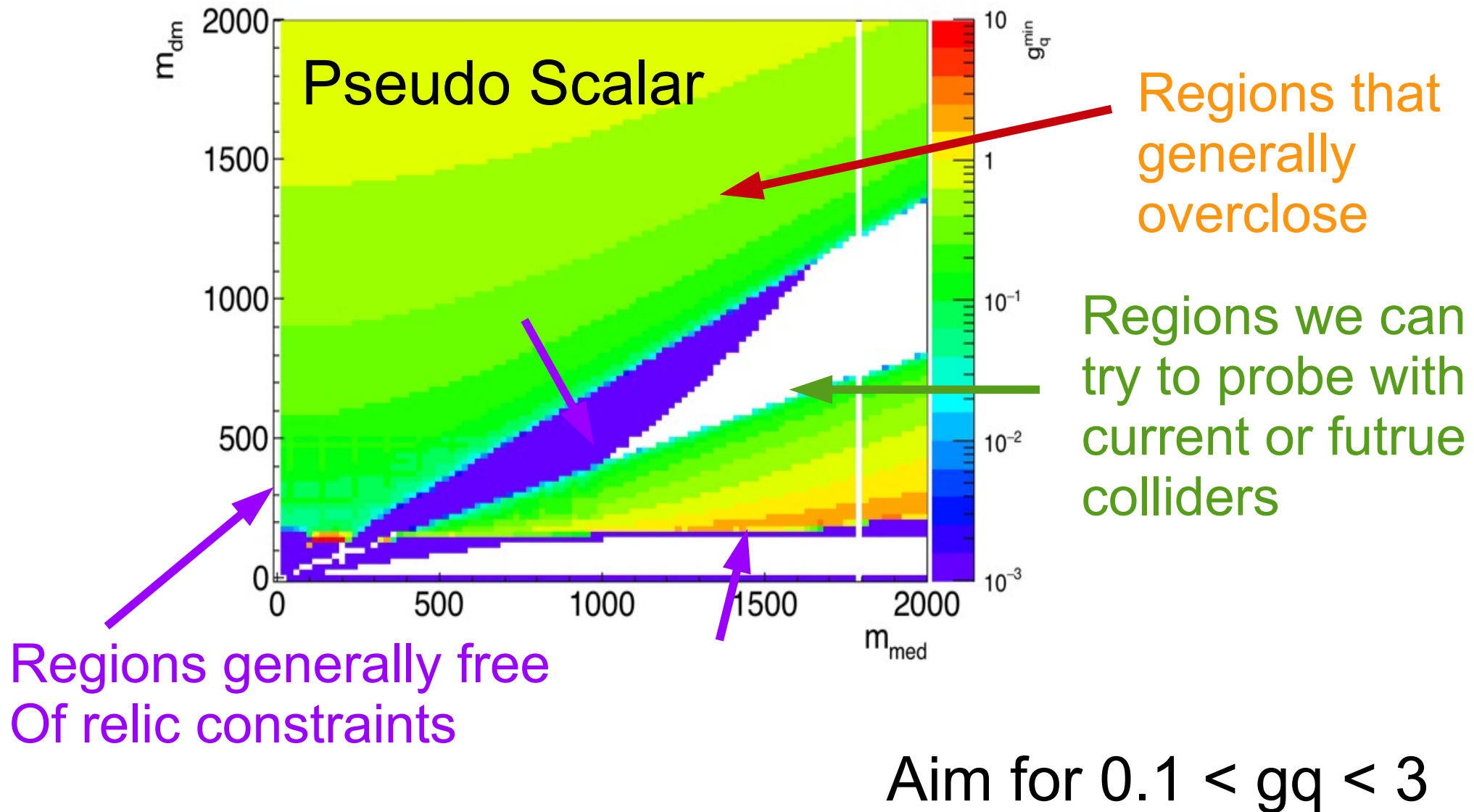
What do we conclude?

- For pseudoscalar we have looser constraints



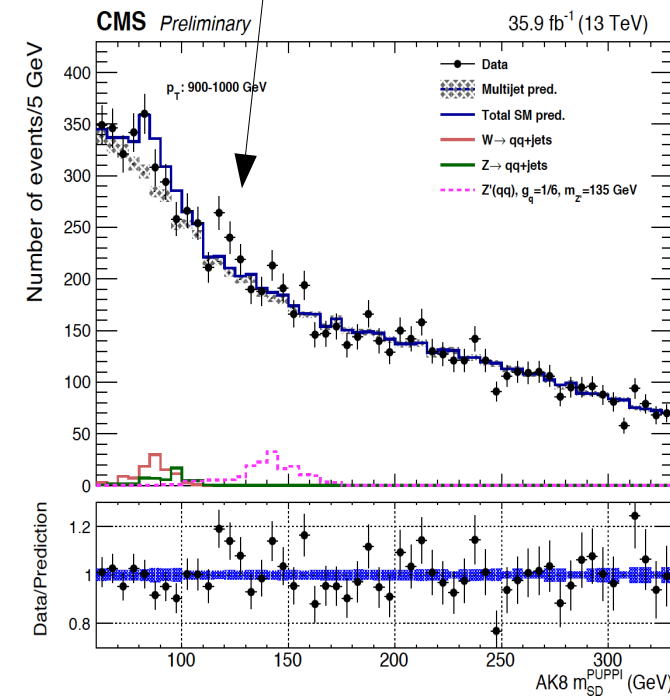
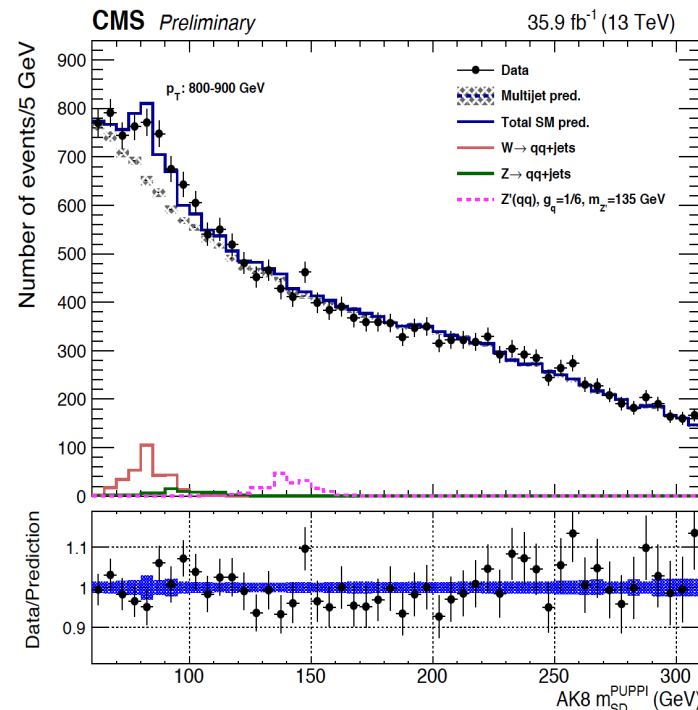
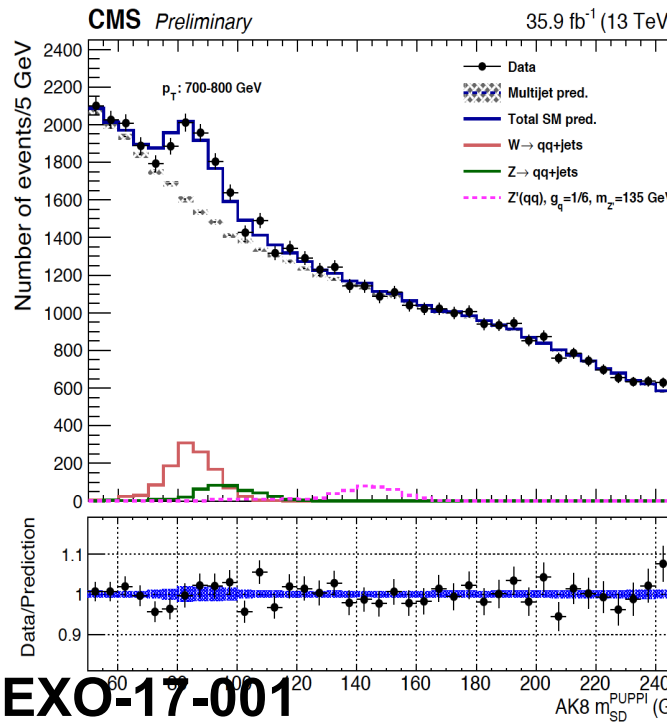
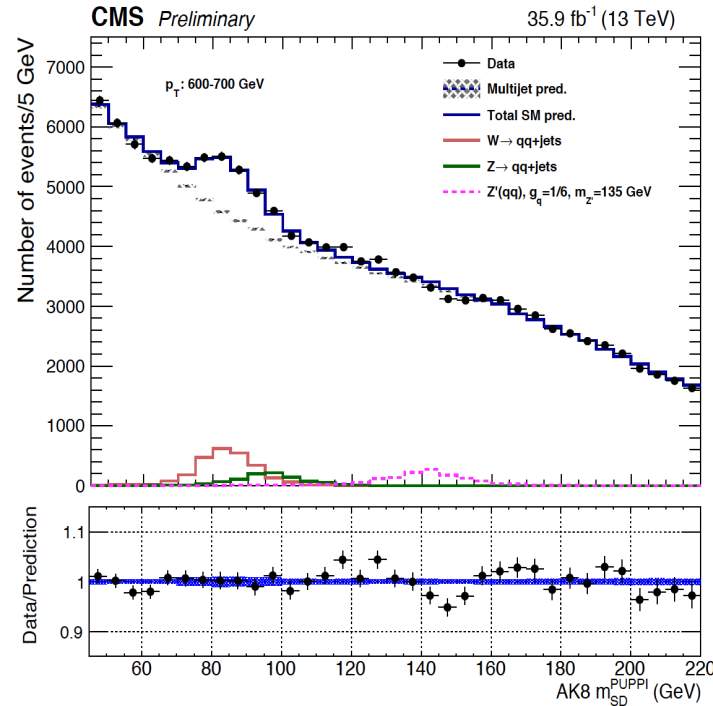
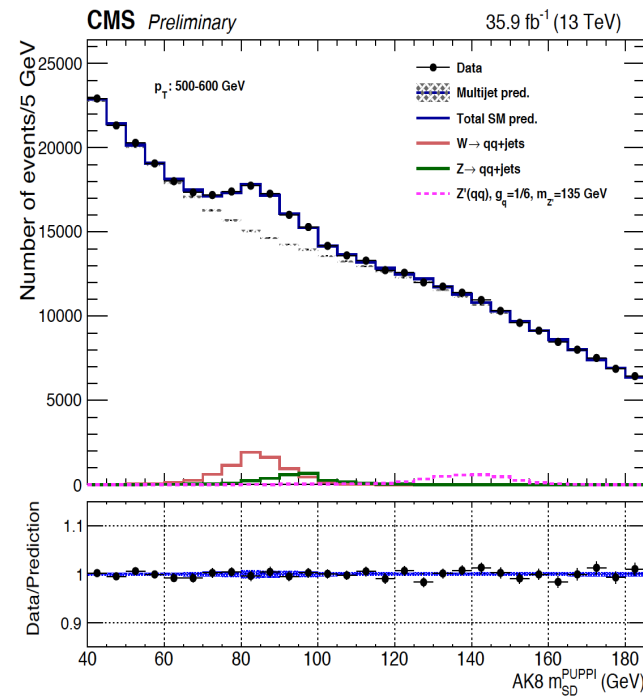
What do we conclude?

- What is driving the results is the coupling

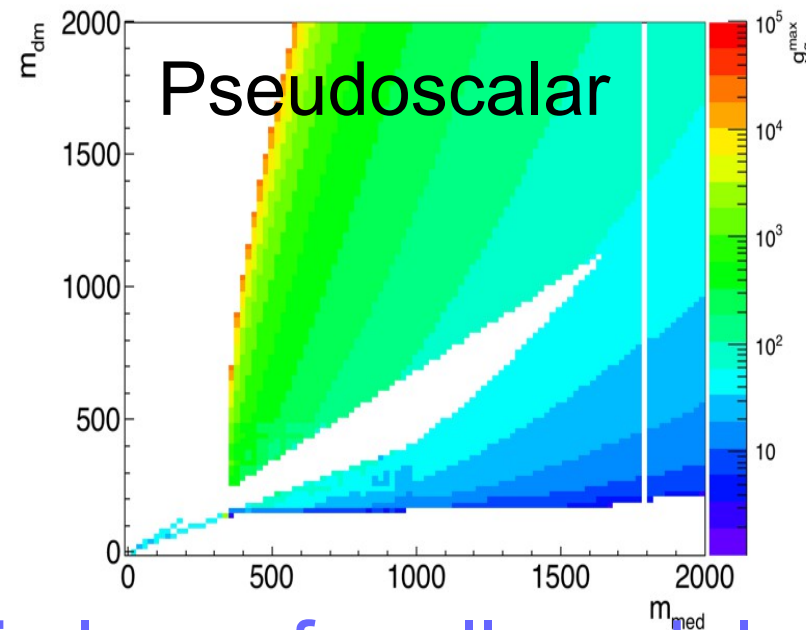
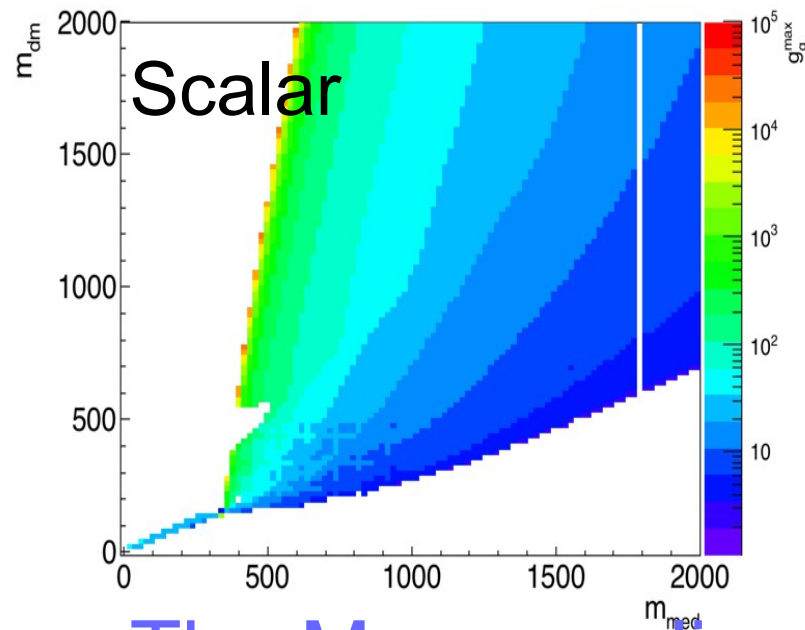
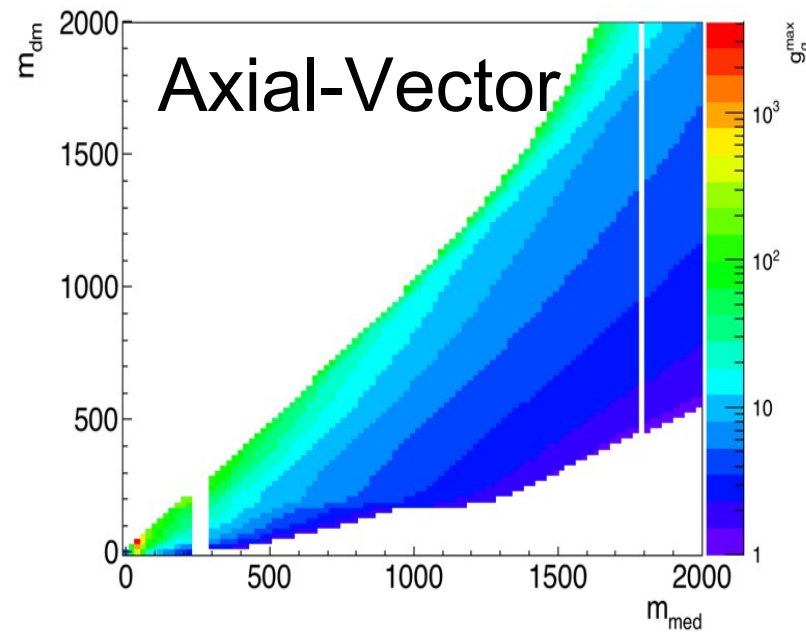
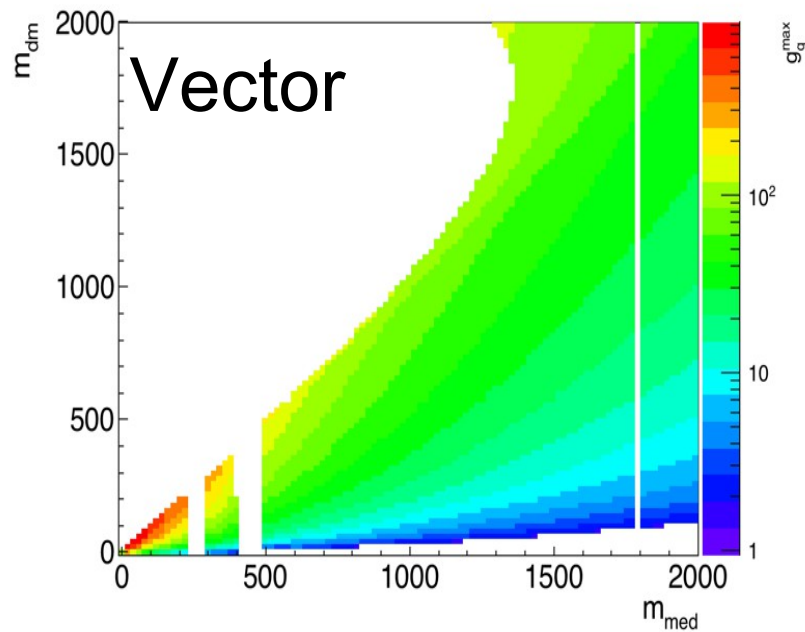


What we see

Excess present
near 115-120

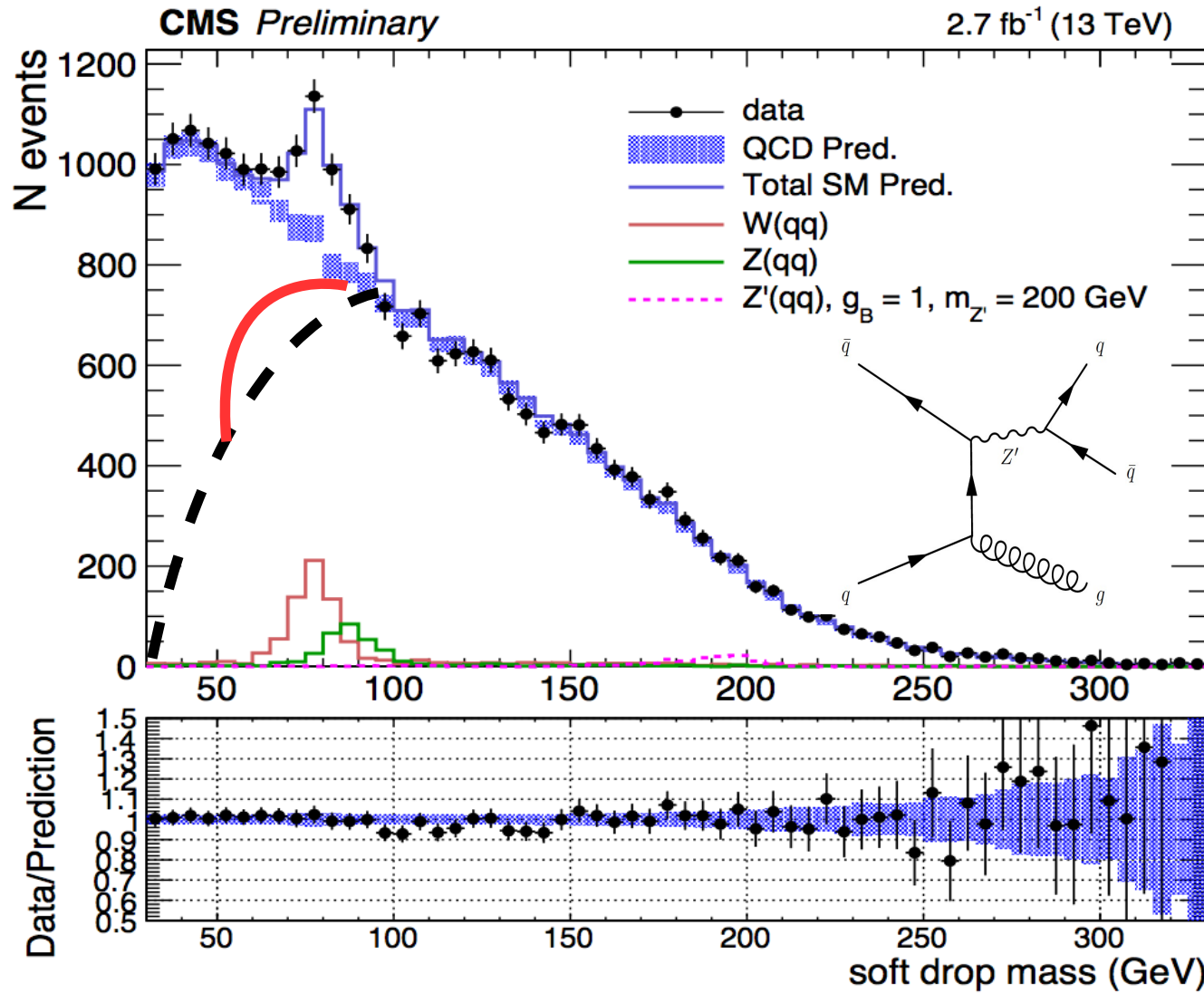


Max Couplings for all¹⁰⁸

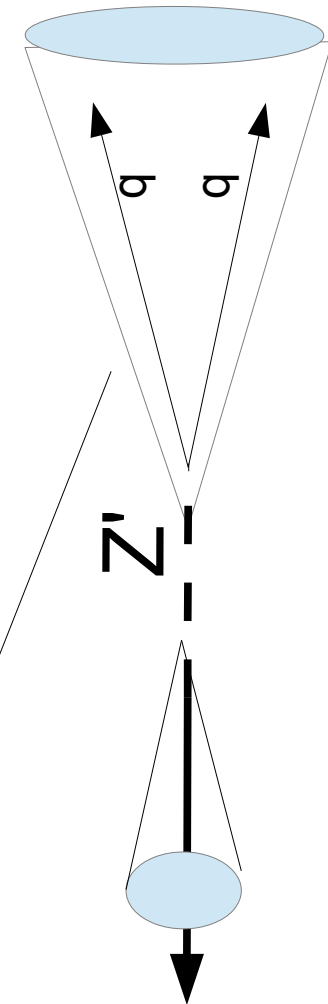


The Max coupling is large for all models

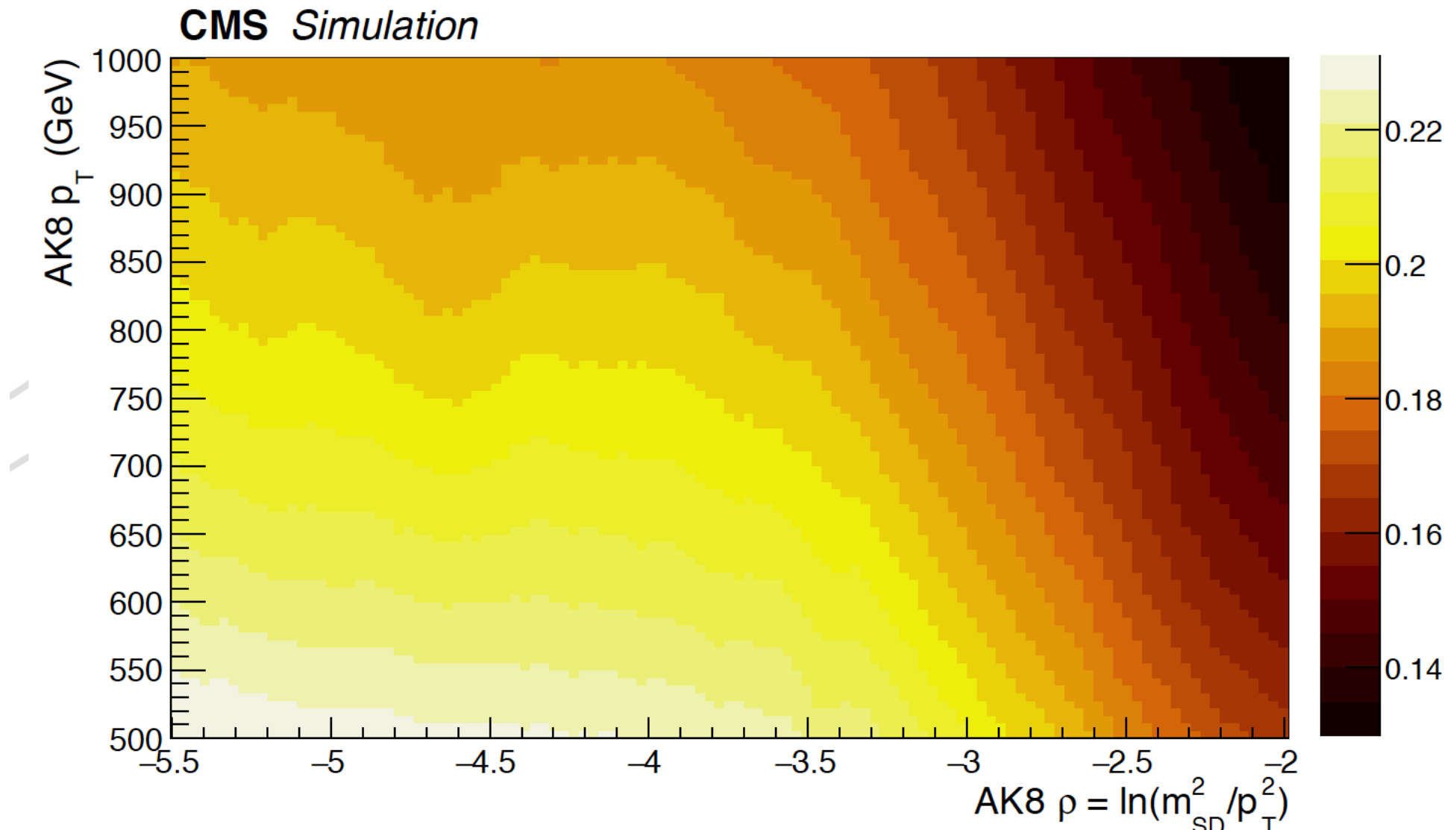
Going all the way down



Fat jet + ISR



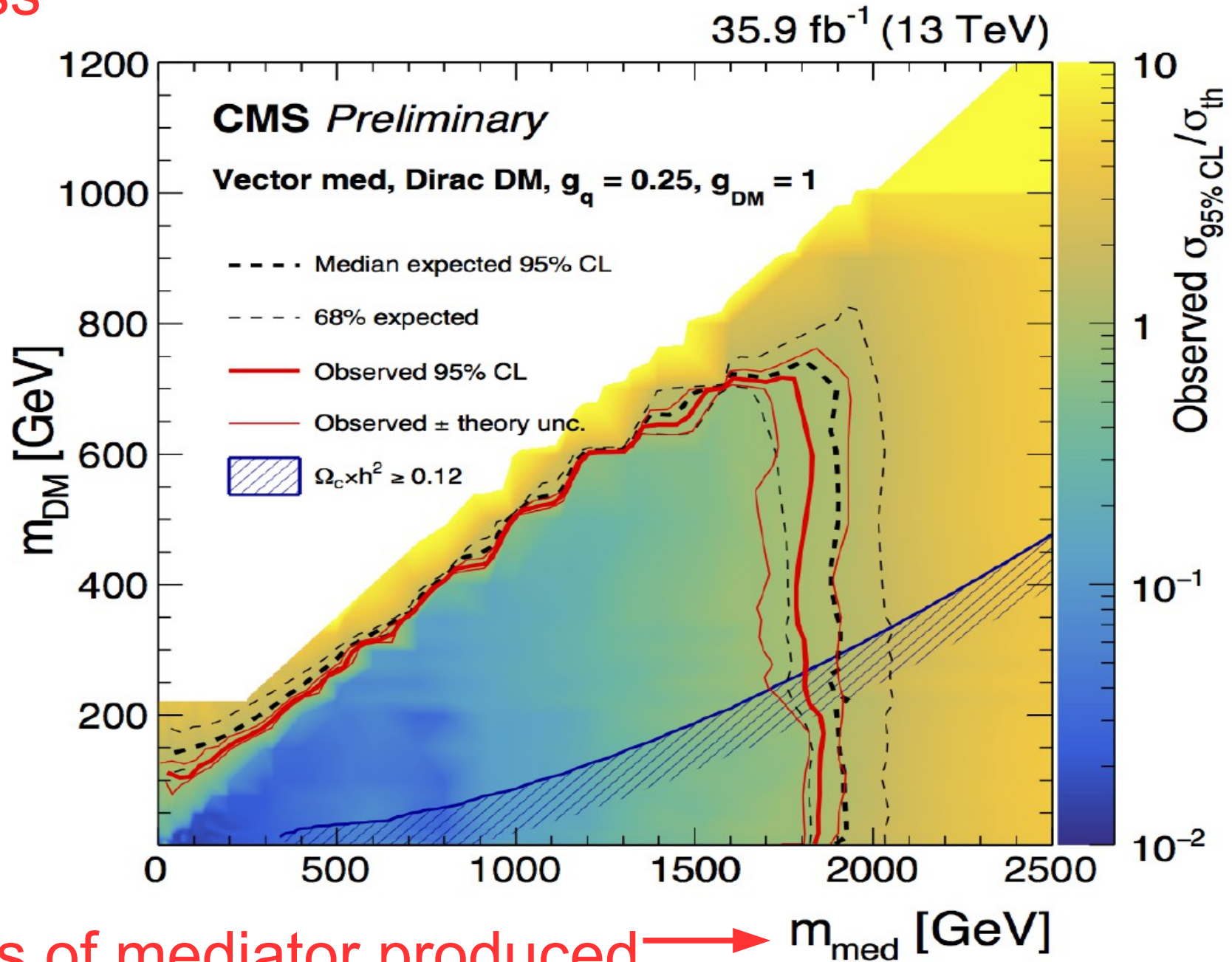
Full Decorrelation scheme



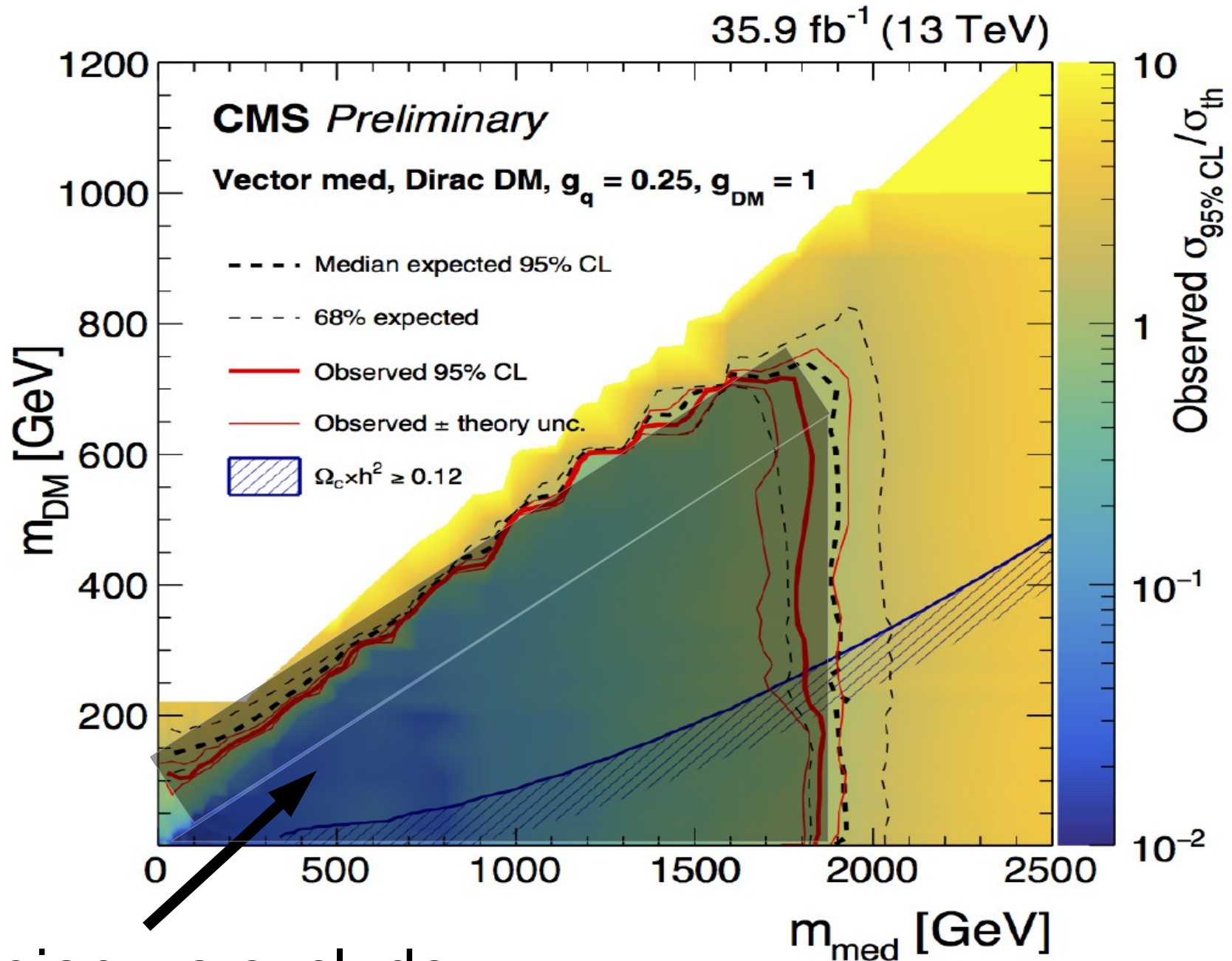
Use the k-Nearest Neighbor approach to determine N_2 cut

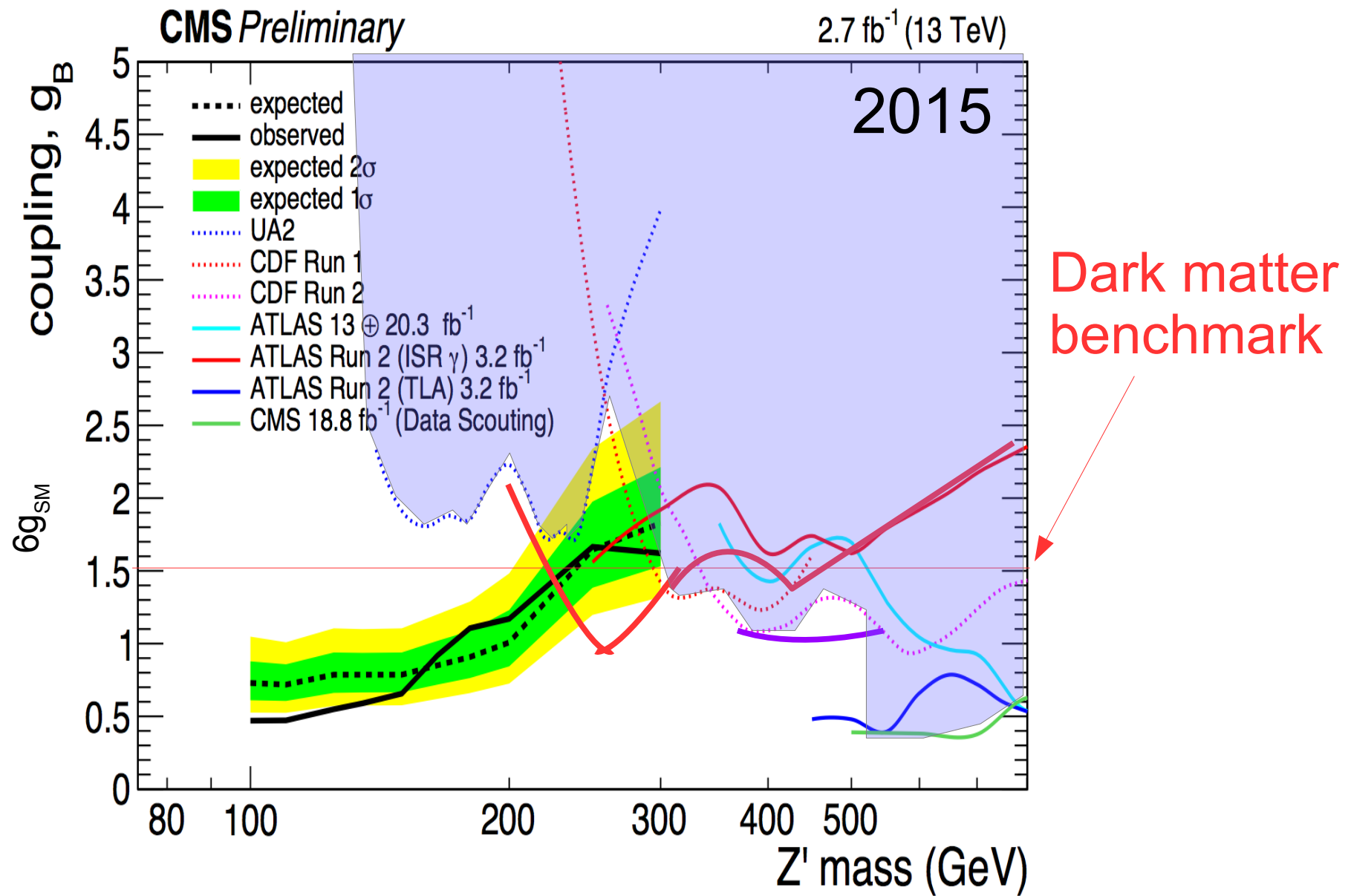
Dark Matter Mass

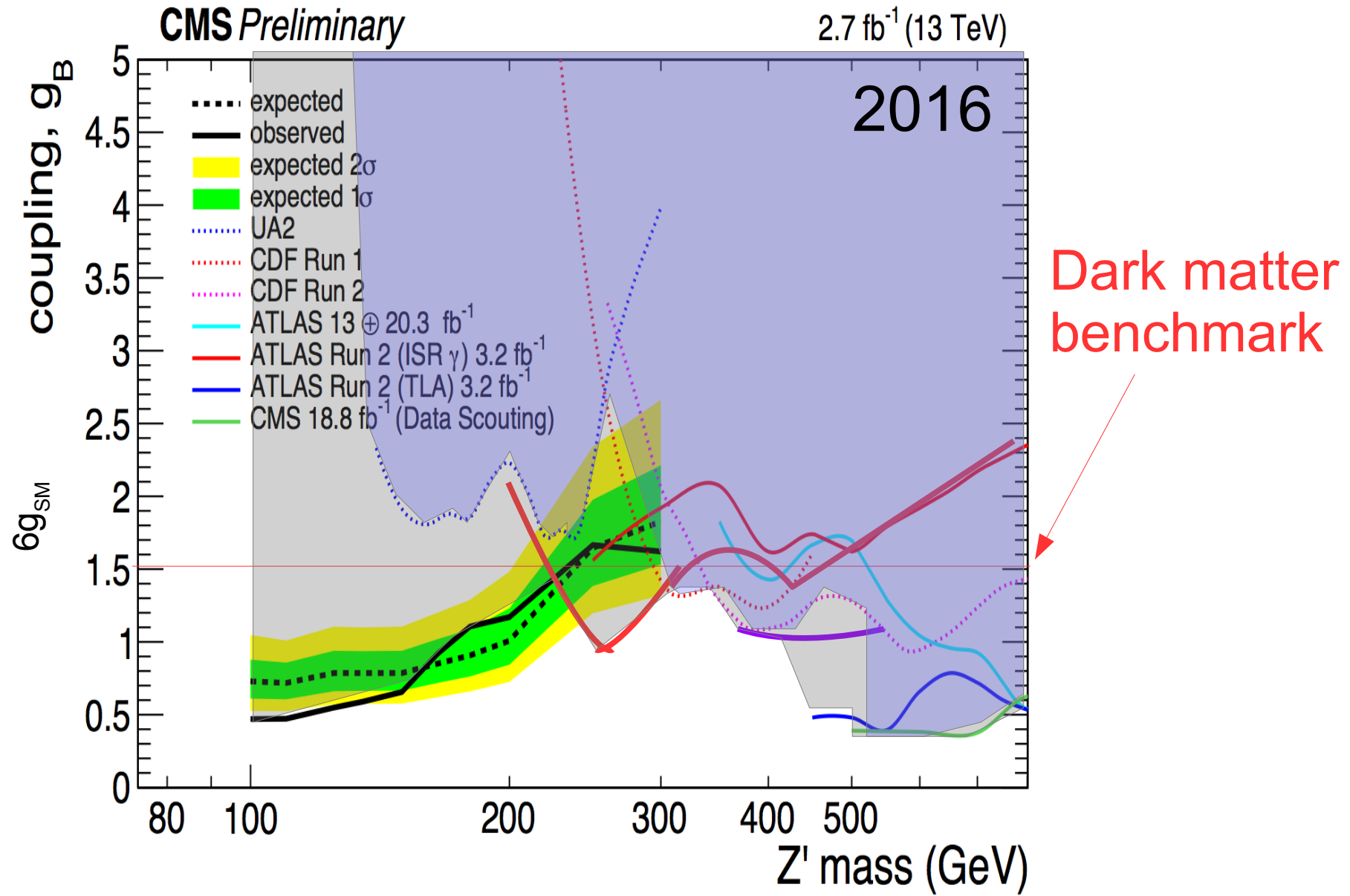
Pick a Model



Understanding sensitivity



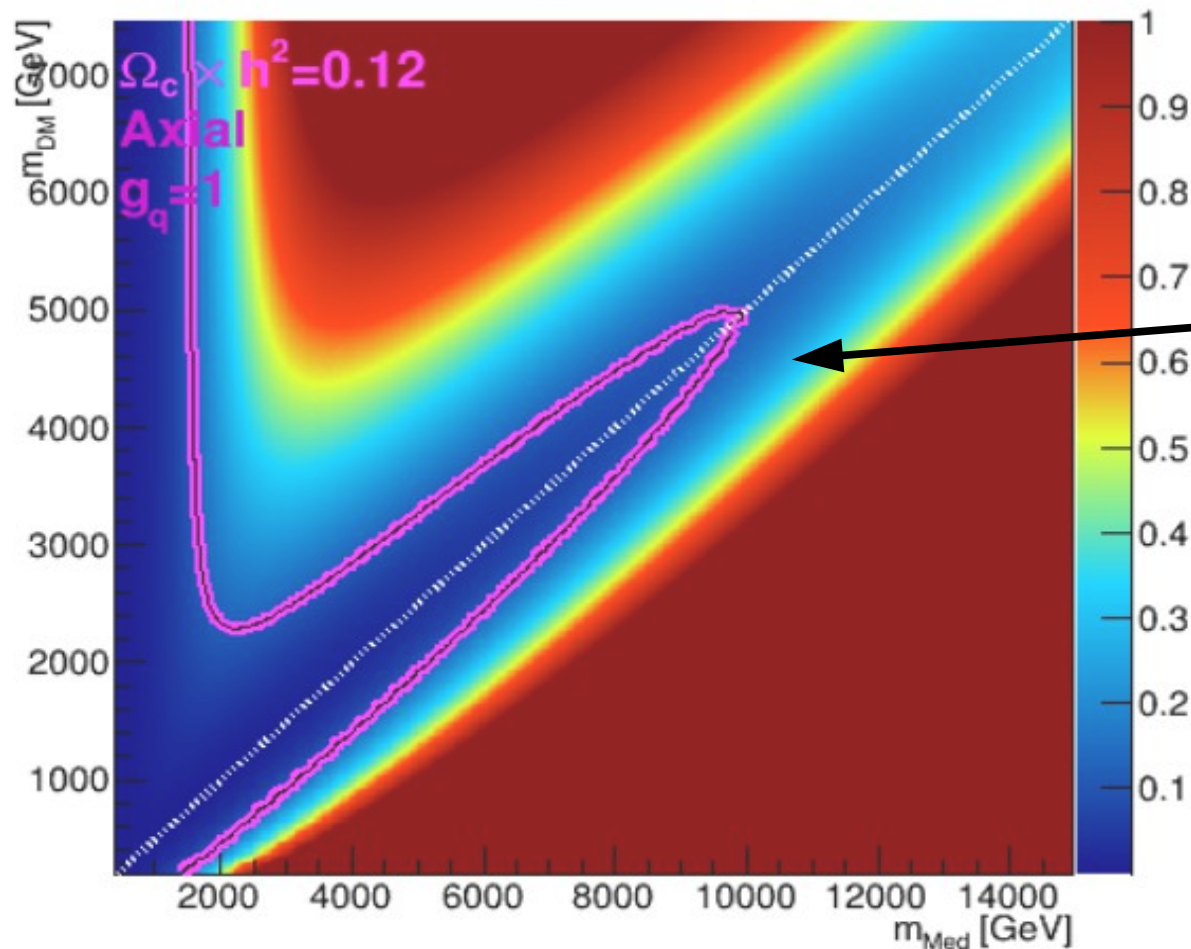




Outlook

Outlook

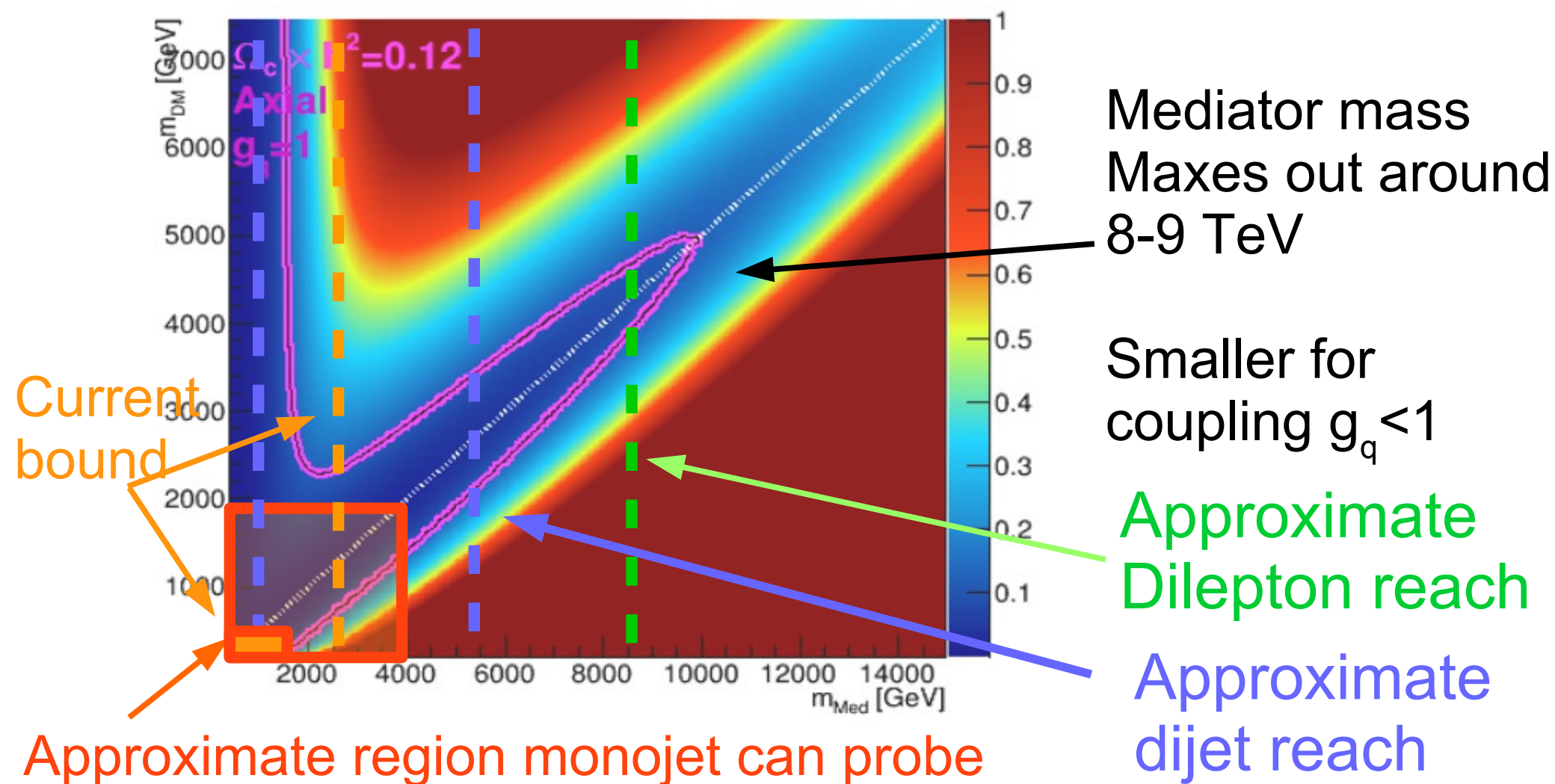
- Spin 1 :
 - Dijet and monojet will continue to push out the bounds
 - Have already crossed the neutrino wall



Outlook

- Spin 1 :

- Dijet and monojet will continue to push out the bounds
- Have already crossed the neutrino wall

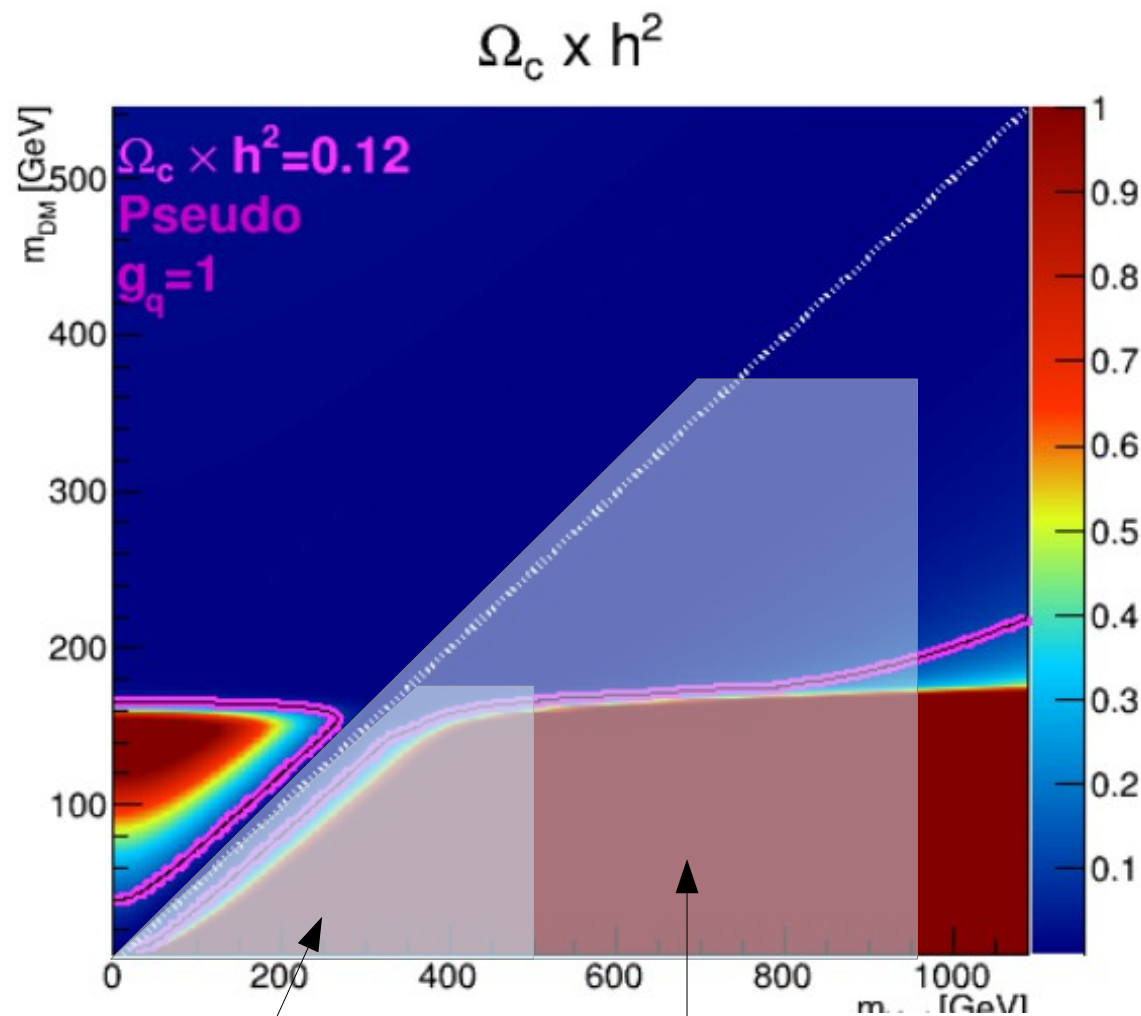


Hep-ph/1603.08525
 Hep-ph/1509.02904
 EXO-12-055
 HIG-16-012

Outlook

- Spin 0 :

- Just starting to probe interesting regions



Current monojet

Approx future

Current reach for
 Higgs-like Singlet
 Dominated by $H \rightarrow \text{Inv}$

BR < 25%

Will get to < 5%

Heavy Flavor :
 $bb/tt + \phi$ also excluding
 at low mass
 Expect mono-Z too

Mediator Arms Race

· CMS

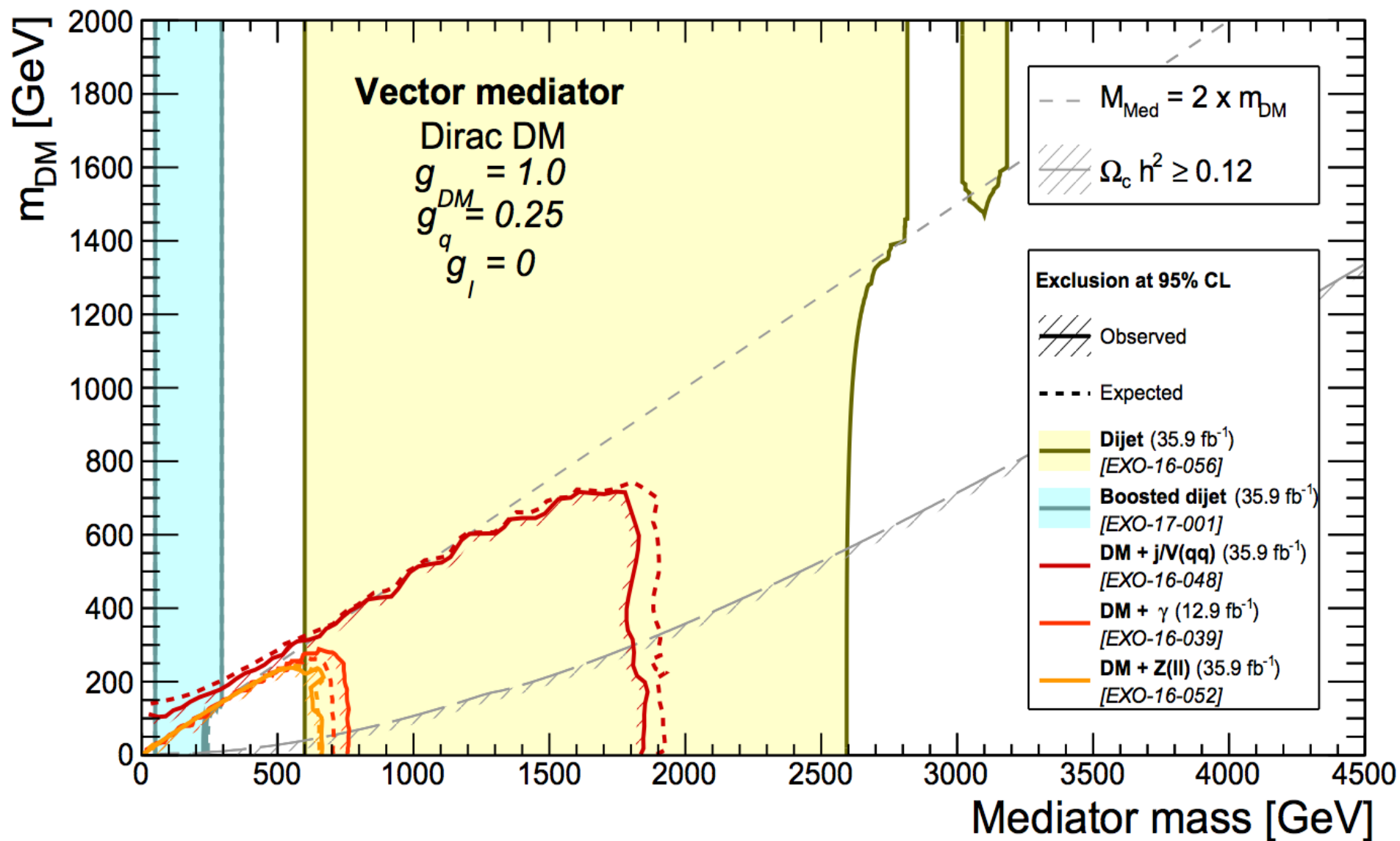
ATLAS



Who can cover the territory first?

CMS Preliminary

LHCP 2017



ATLAS & CMS

Dirac DM
 $g_q = 1.0$
 $g_q^{DM} = 0.25$
 $g_l = 0$

Med DM
 $\Omega_c h^2 \geq 0.12$

DM Mass [TeV]

DM Simplified Model Exclusions

ATLAS Preliminary March 2017

Exclusion at 95% CL

- Dijet
- Observed
- Expected
- Dijet 8 TeV
- Dijet (35.9 fb⁻¹)
- Dijet TLA
- Boosted dijet (35.9 fb⁻¹)
- Dijet + ISR
- DM + $\gamma/V(qq)$ (35.9 fb⁻¹)
- DM + γ/V (2.9 fb⁻¹)
- DM + Z(l $\bar{l}$) (35.9 fb⁻¹)
- DM + jet

Axial-vector mediator, Dirac DM
 $g_q = 0.25, g_l = 0, g_{DM} = 1$
 All limits at 95% CL

Mediator Mass [TeV]

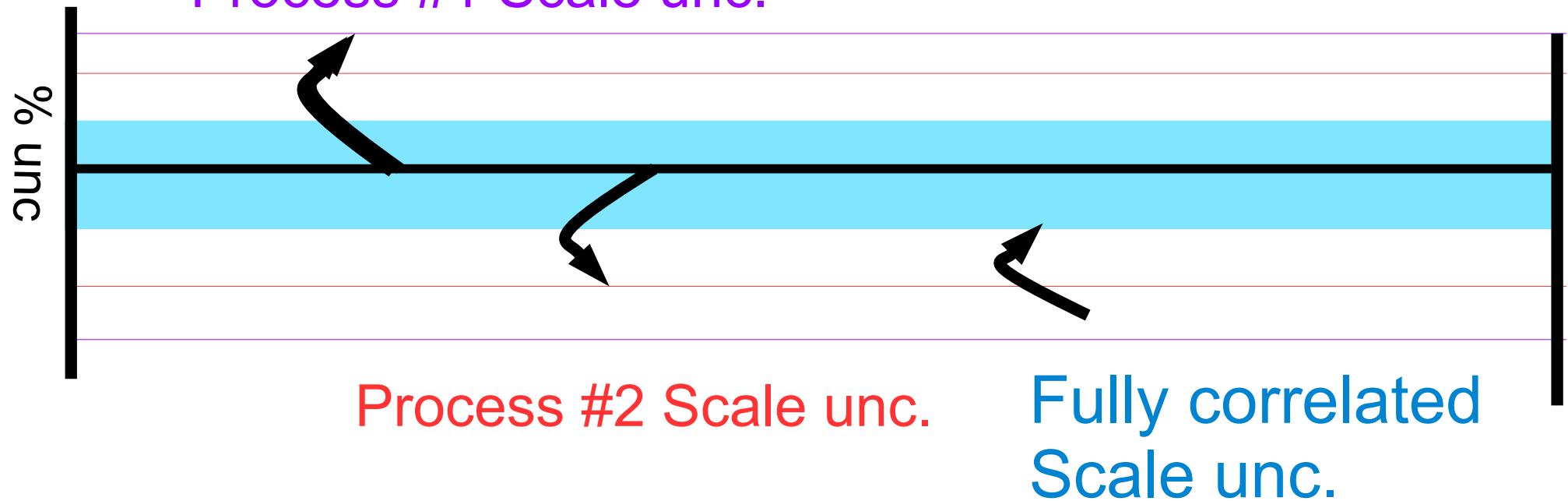
Mediator mass [GeV]

However we still have a problem!

Unc. $\longrightarrow \frac{d\sigma^y}{dp_T} / \frac{d\sigma^z}{dp_T} = d\sigma^y/d\sigma^z(\mu)$

We don't know how to do scale uncertainties on ratios!

Process #1 Scale unc.

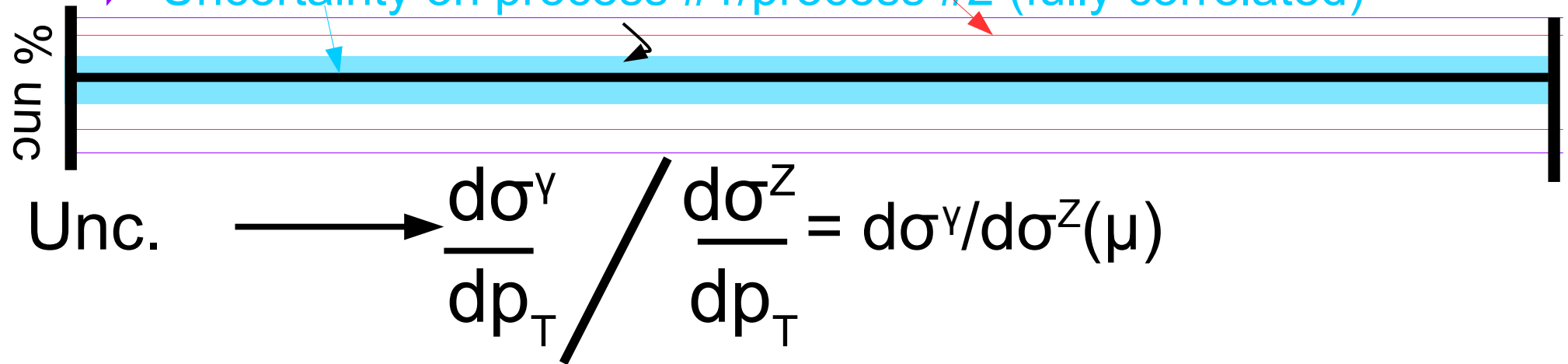


Uncertainty on ratio? How is it done?

Scale uncertainty on process #1

Scale uncertainty on process #2

Uncertainty on process #1/process #2 (fully correlated)



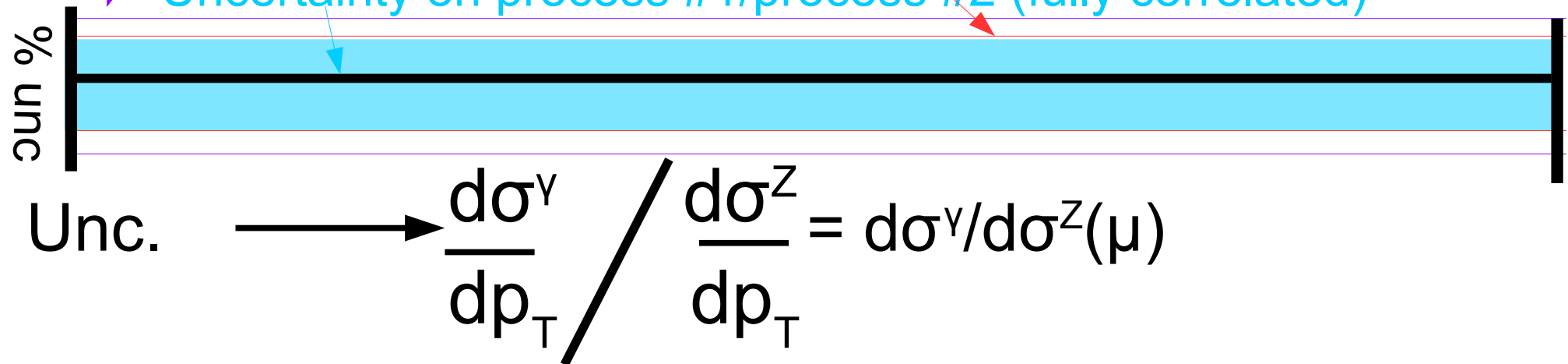
$$\begin{pmatrix} d\sigma^Y(+\sigma) \\ d\sigma^Z(+\sigma) \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} d\sigma^Y(\mu^{\text{up}})/d\sigma^i(\mu_0) \\ d\sigma^Z(\mu^{\text{up}})/d\sigma^i(\mu_0) \end{pmatrix}$$

Uncertainty on ratio? How is it done?

Scale uncertainty on process #1

Scale uncertainty on process #2

Uncertainty on process #1/process #2 (fully correlated)



$$\begin{pmatrix} d\sigma^Y(+\sigma) \\ d\sigma^Z(+\sigma) \end{pmatrix} = \begin{pmatrix} 1 & C \\ C & 1 \end{pmatrix} \begin{pmatrix} d\sigma^Y(\mu^{\text{up}})/d\sigma^i(\mu_0) \\ d\sigma^Z(\mu^{\text{up}})/d\sigma^i(\mu_0) \end{pmatrix}$$

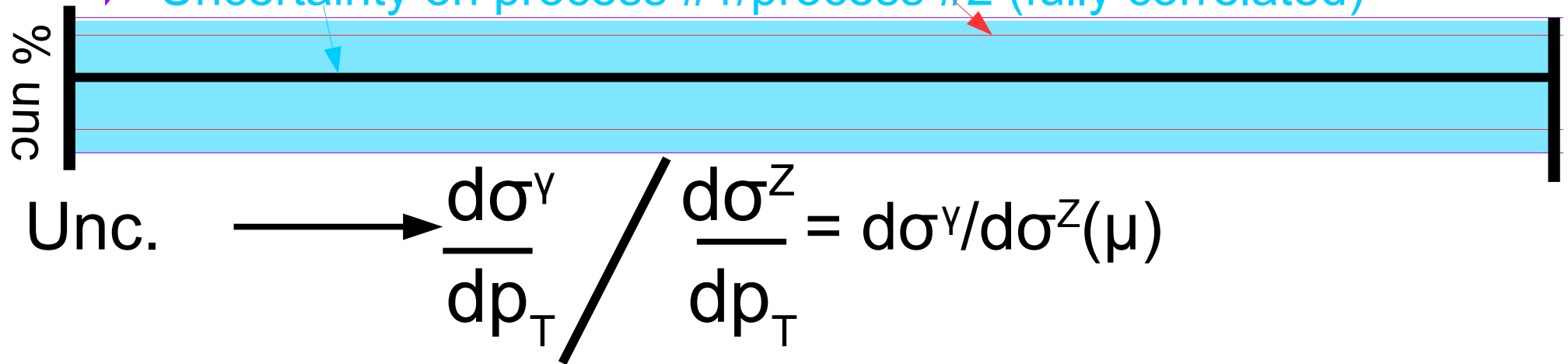
Adjust C until
uncertainty is

Uncertainty on ratio? How is it done?

Scale uncertainty on process #1

Scale uncertainty on process #2

Uncertainty on process #1/process #2 (fully correlated)



$$\begin{pmatrix} d\sigma^Y(+\sigma) \\ d\sigma^Z(+\sigma) \end{pmatrix} = \begin{pmatrix} 1 & C \\ C & 1 \end{pmatrix} \begin{pmatrix} d\sigma^Y(\mu^{\text{up}})/d\sigma^i(\mu_0) \\ d\sigma^Z(\mu^{\text{up}})/d\sigma^i(\mu_0) \end{pmatrix}$$

Decorrelate scale unc. until its max of either process

$$d\sigma^Y / d\sigma^Z(+\sigma) < \max_i (d\sigma^i(\mu^{\text{up}}) / d\sigma^i(\mu_0))$$

What is the previous unc?

Can we motivate this?

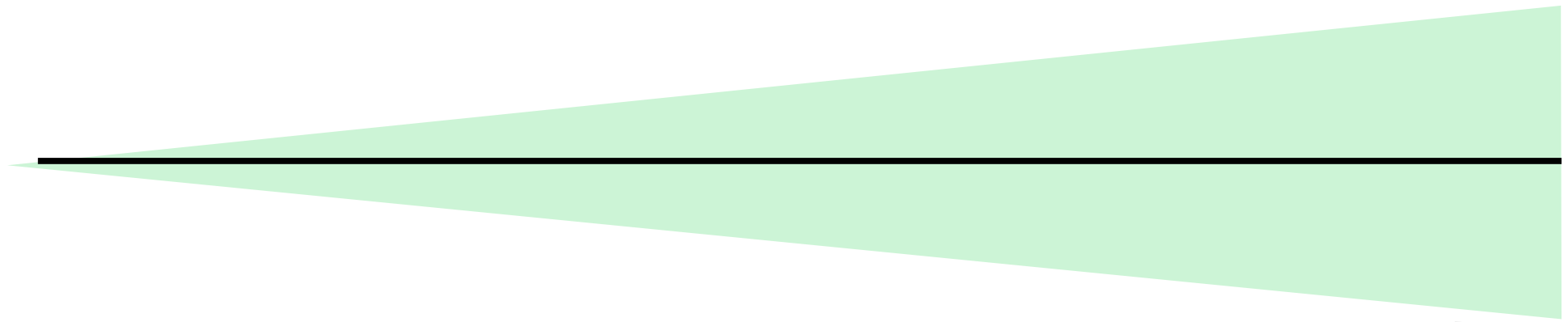


Unc. $\longrightarrow \frac{d\sigma^Y}{dp_T} / \frac{d\sigma^Z}{dp_T} = d\sigma^Y/d\sigma^Z(\mu)$

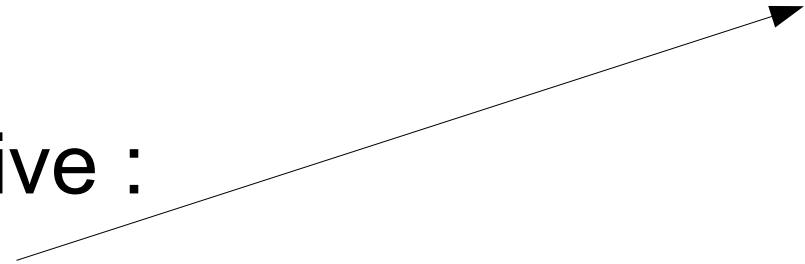
$$\begin{pmatrix} d\sigma^Y(+\sigma) \\ d\sigma^Z(+\sigma) \end{pmatrix} = \begin{pmatrix} 1 & C \\ C & 1 \end{pmatrix} \begin{pmatrix} d\sigma^Y(\mu^{\text{up}})/d\sigma^i(\mu_0) \\ d\sigma^Z(\mu^{\text{up}})/d\sigma^i(\mu_0) \end{pmatrix}$$

Makes Little Sense

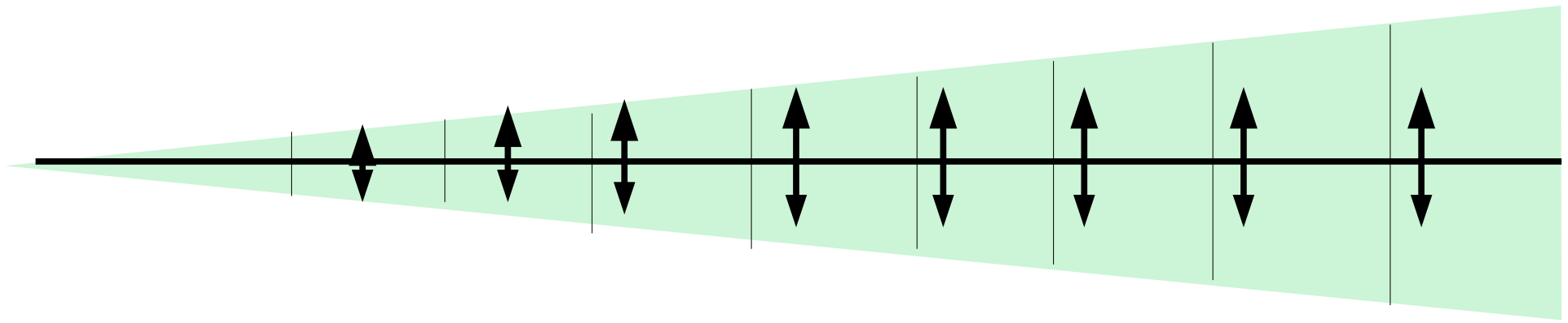
What about EWK corr. uncertainty?



In light of being conservative :
Take the full correction



What about EWK corr. uncertainty?



In light of being conservative :
Take the full correction

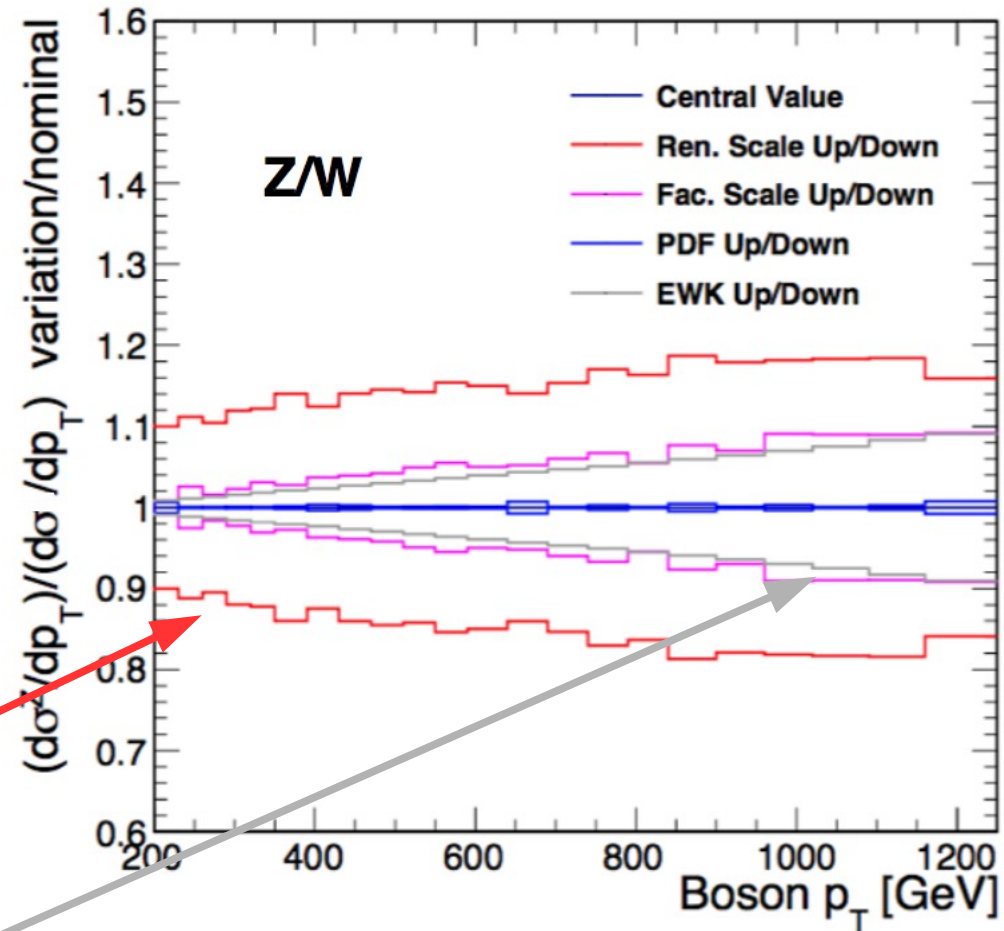
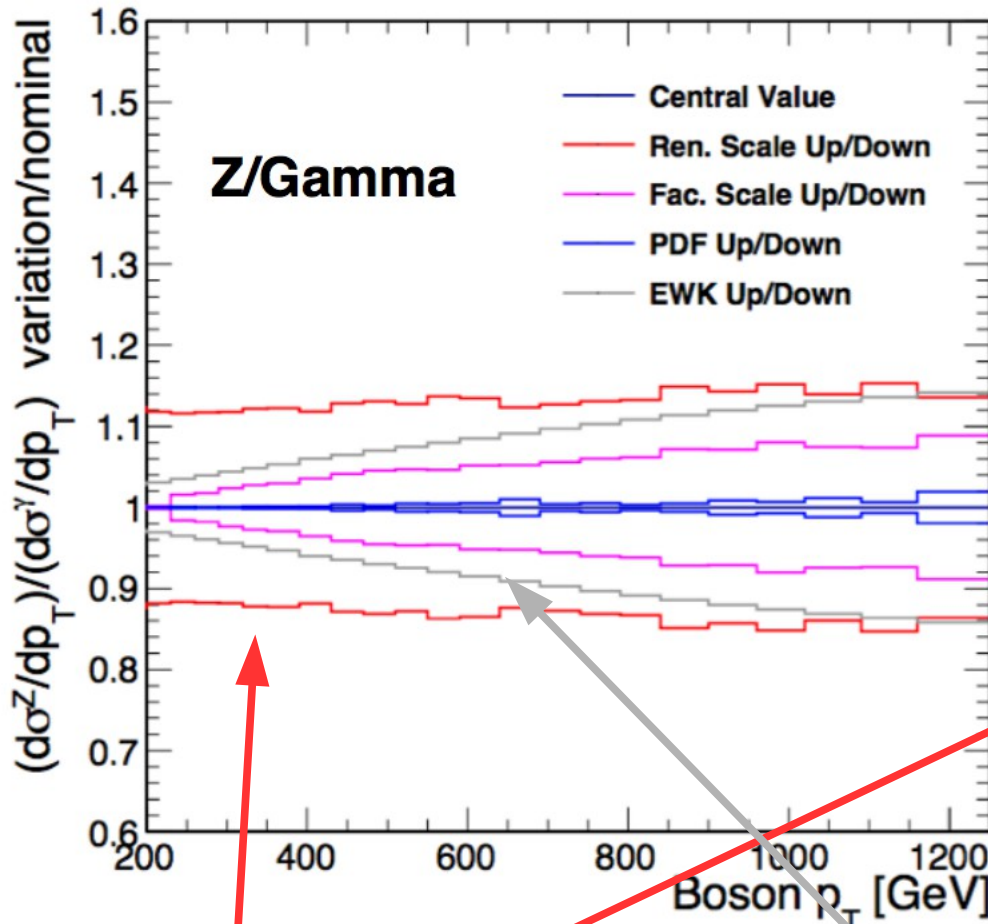
Additionally de-correlated this per bin
Avoids low MET to high MET constraints

We were forced to do this by management!
This makes us too conservative

2015...
We are stuck

How do we propagate this in?

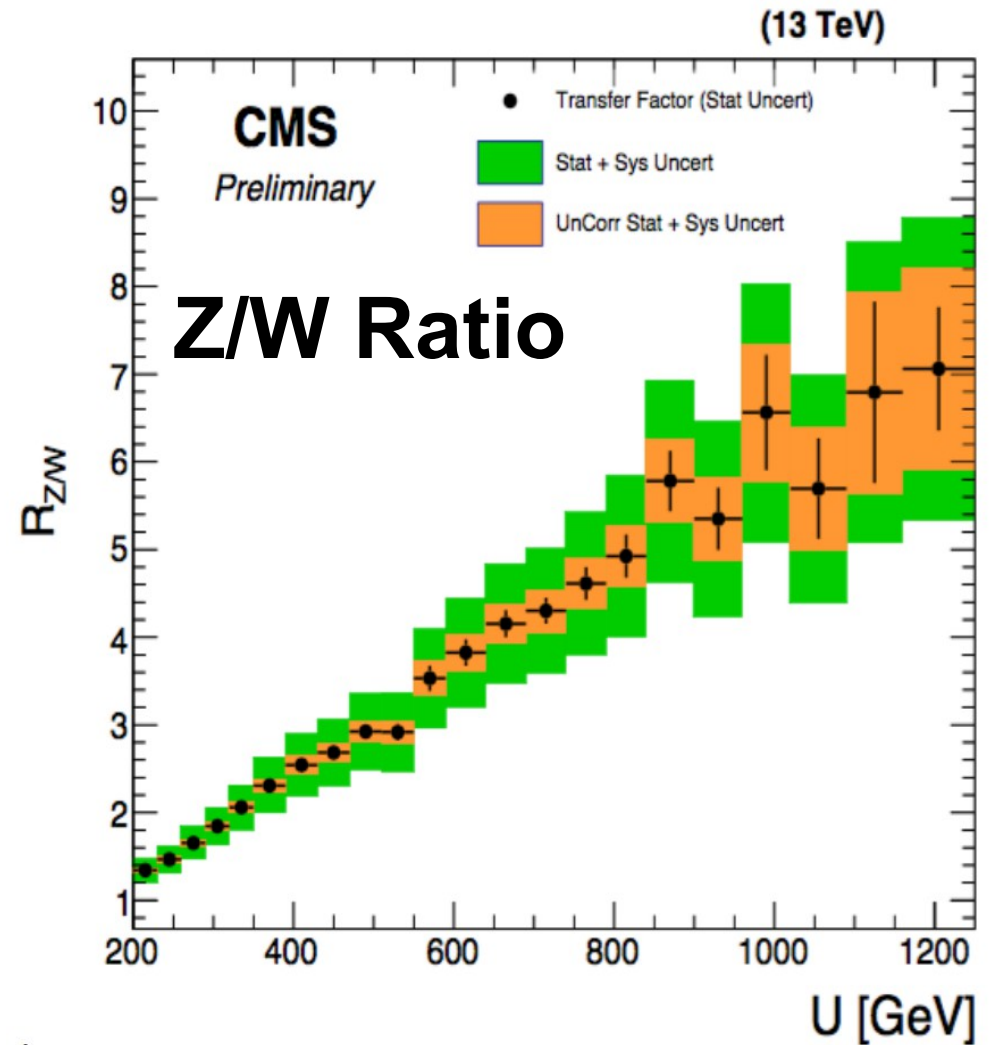
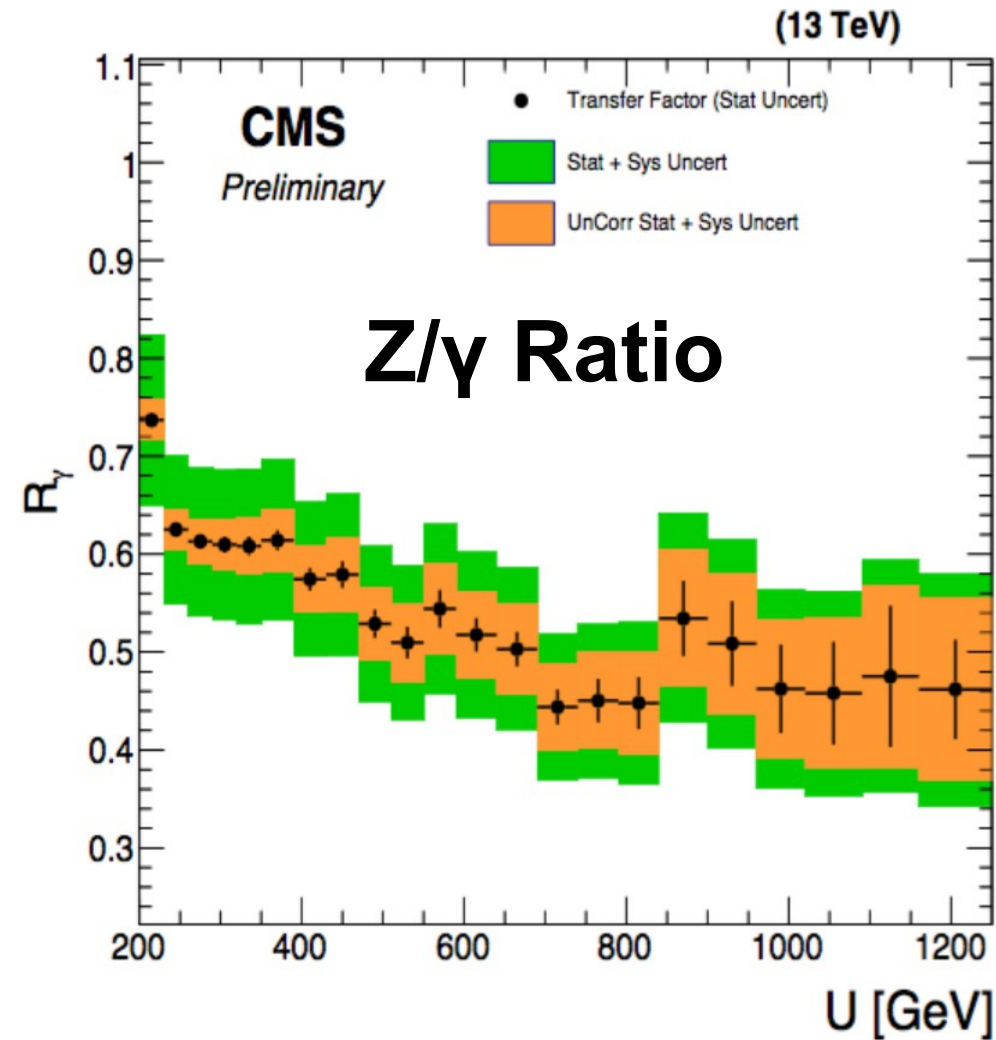
The actual uncertainties



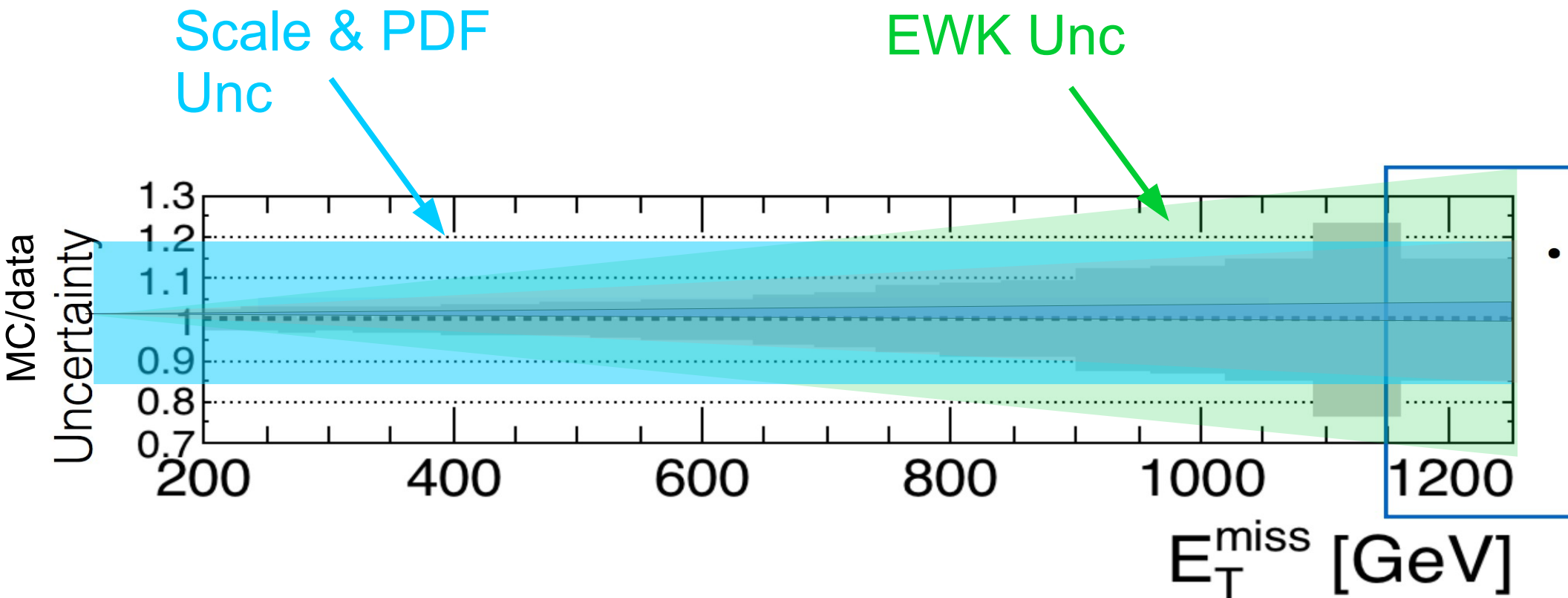
Correlated across
boson p_T

EWK: Since its not correlated
dominant in tails

How exactly do they look?



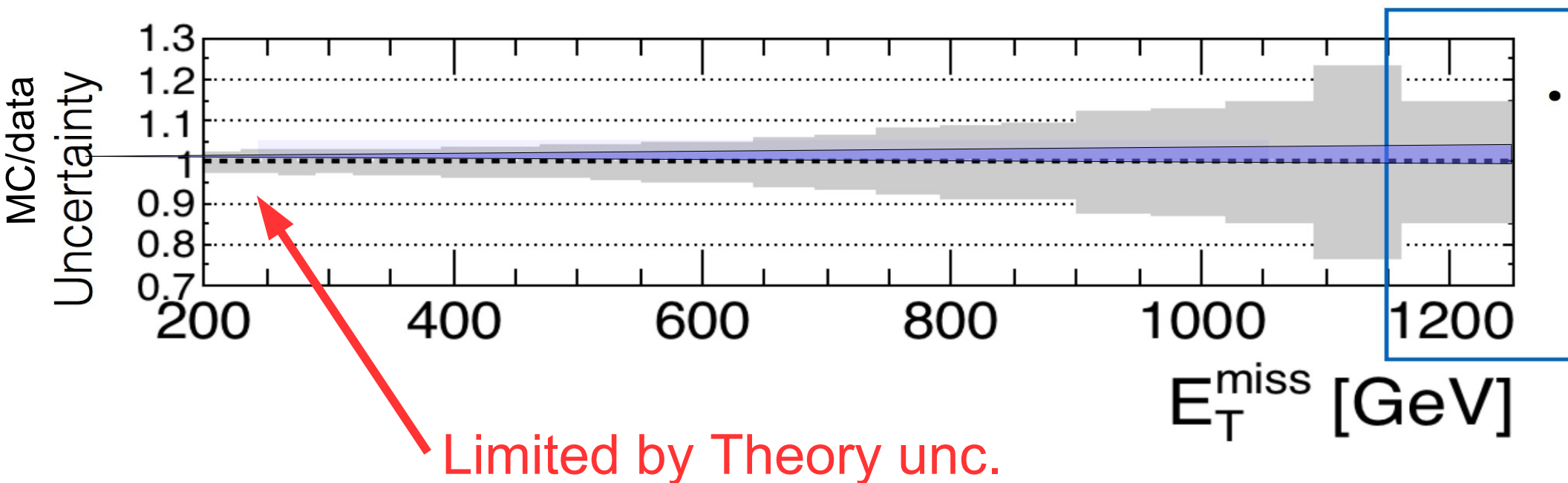
What do the uncertainties look like?



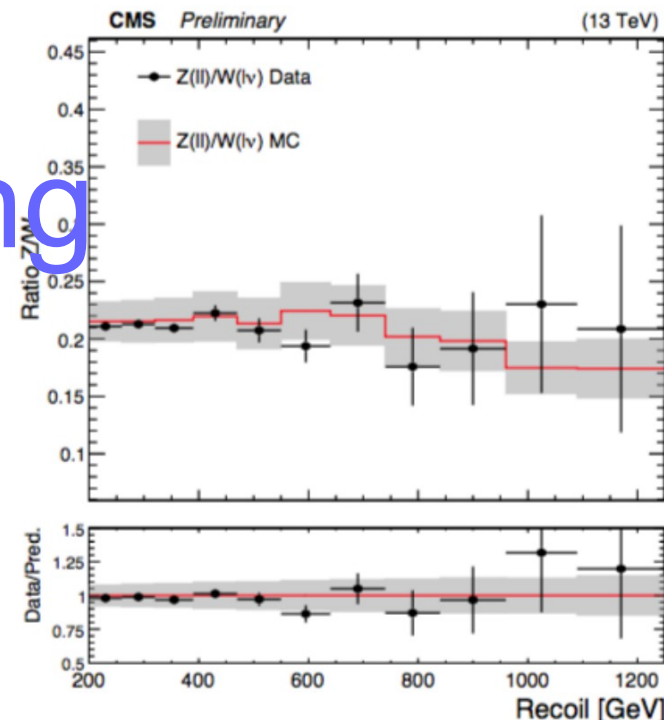
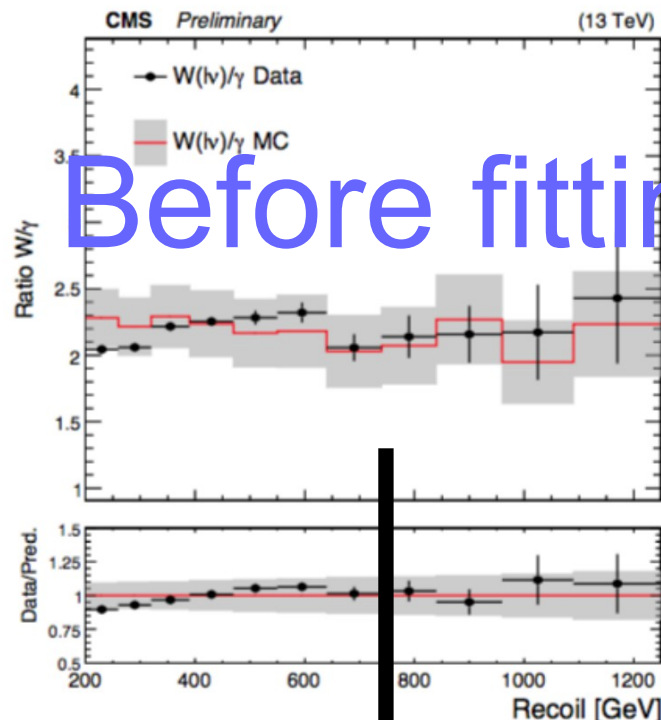
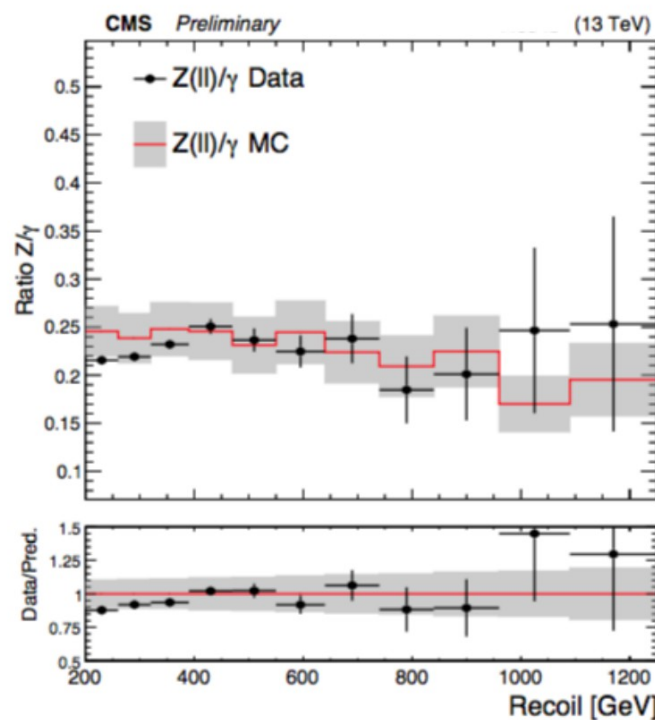
Updated unc still too large

Profiling them in the fit

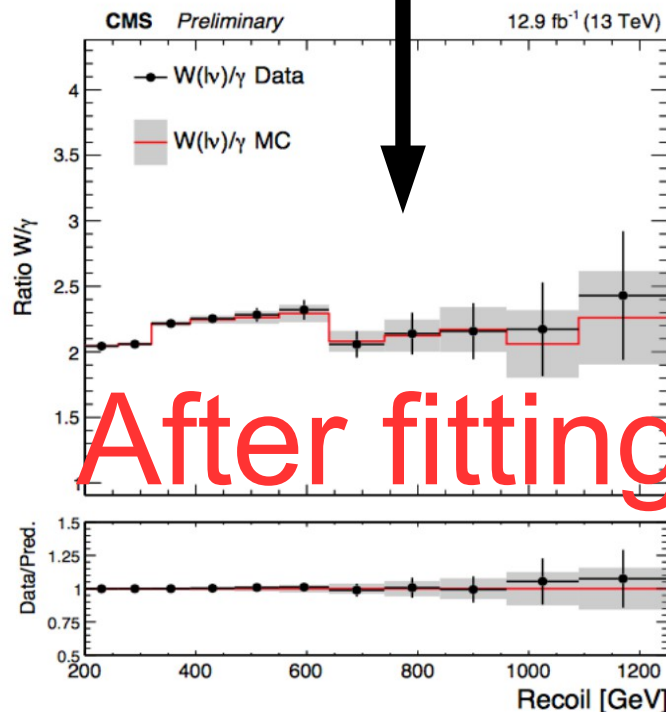
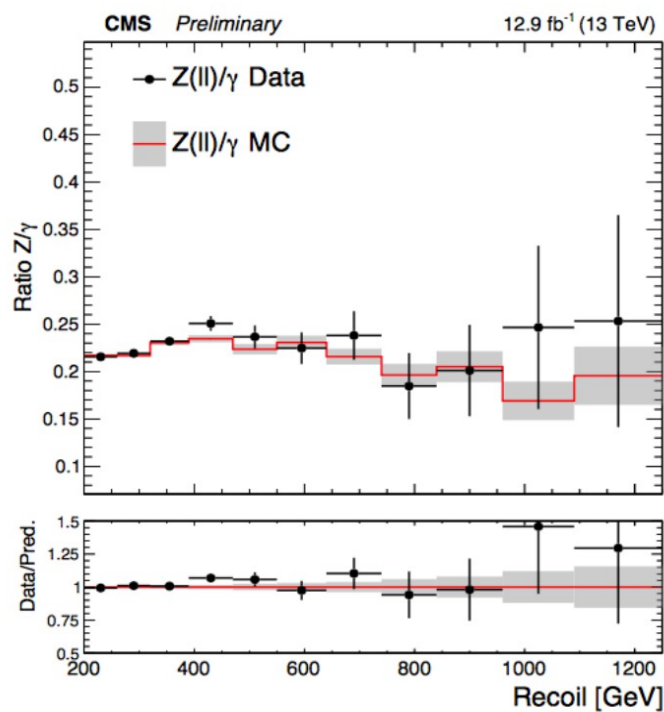
Constraints after the fit



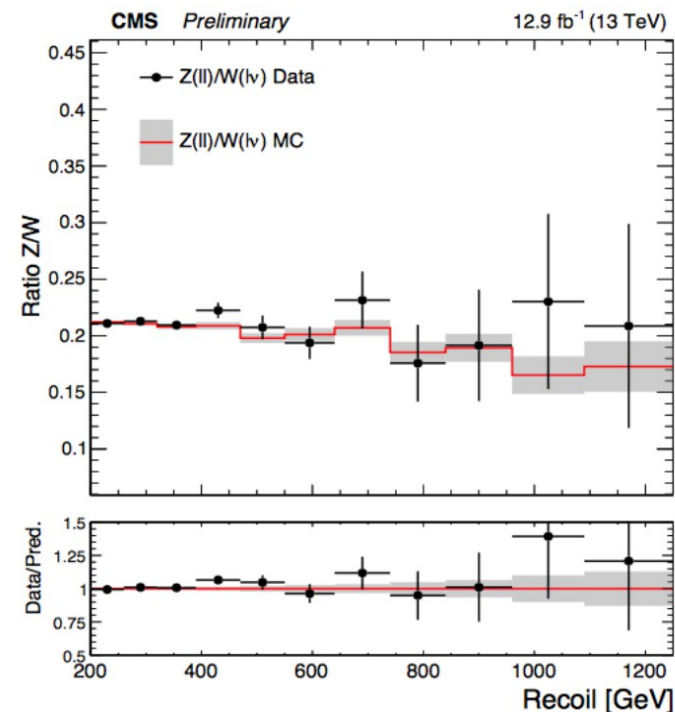
Still systematics limited @low *MET*
Not systematics limited @ high *MET*
→ Likely will *never* be



Before fitting



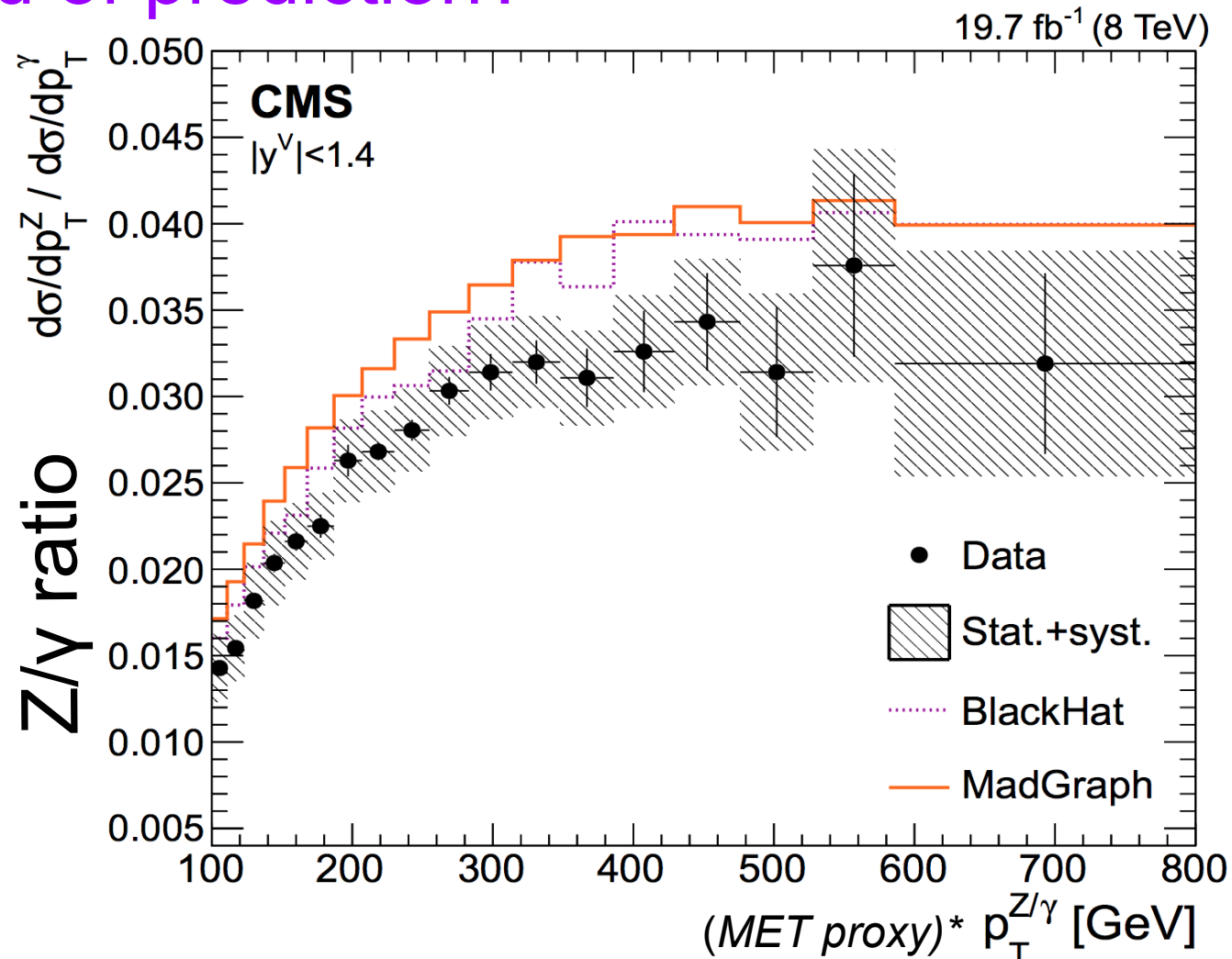
After fitting



2016: A new hope

A mystery? Understanding Z/γ p_T

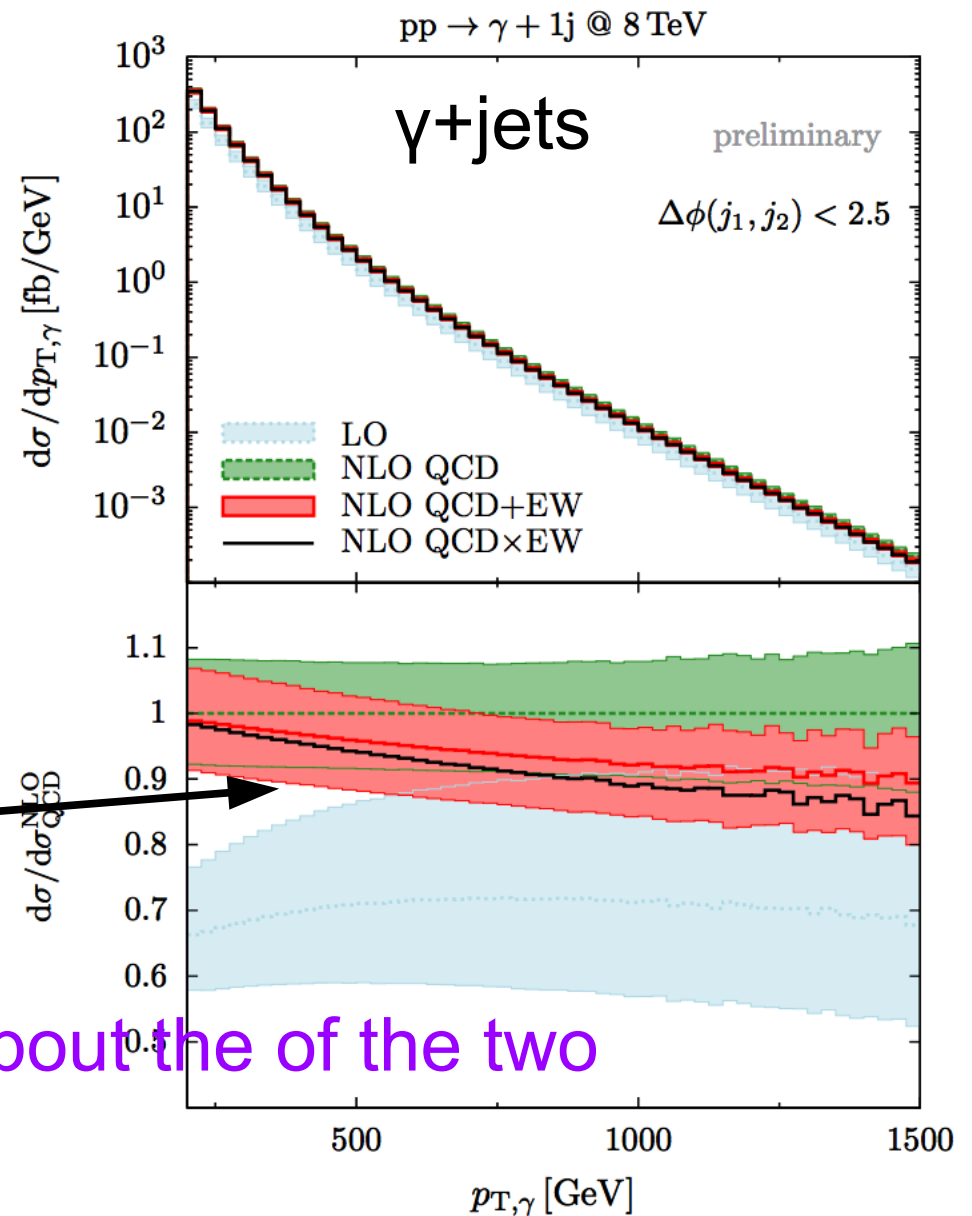
How are going to use photons for Z with this kind of prediction?



We need this to model the ratio perfectly to use it

2015

How do we fix this?

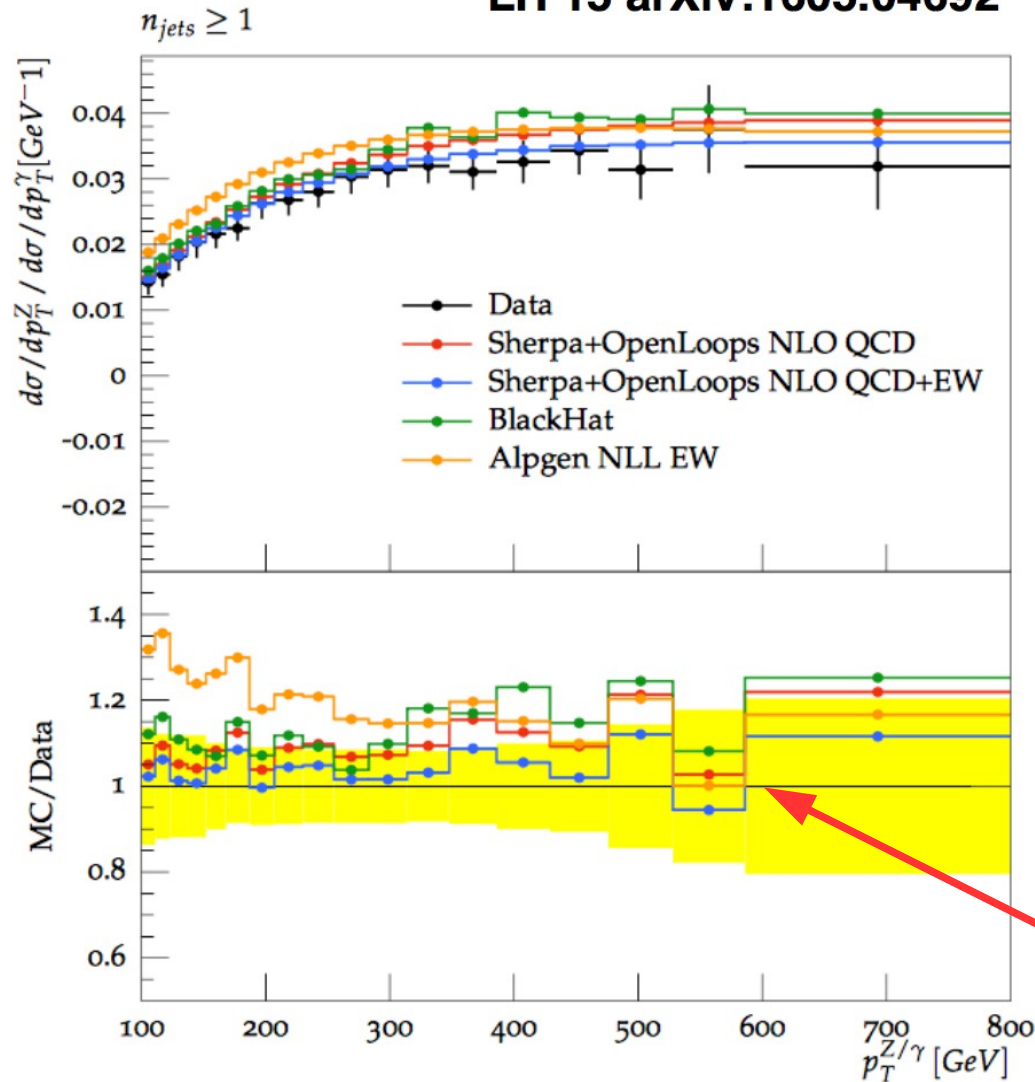


We care about the of the two

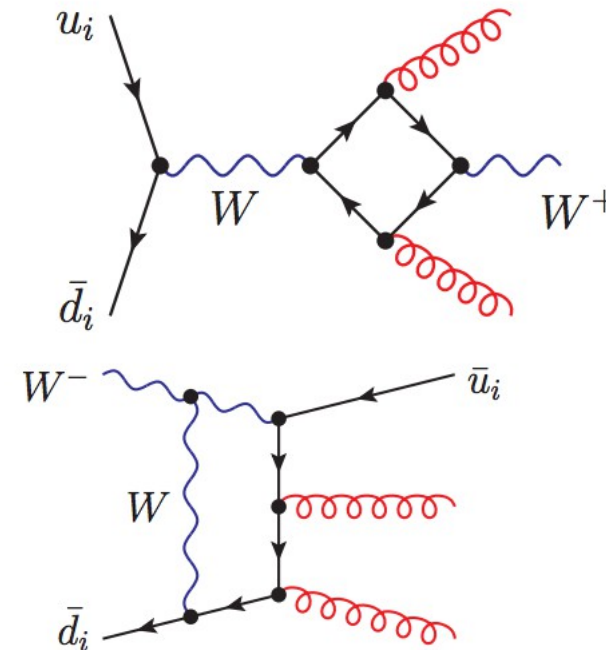
Solving the Mystery

LH'15 arXiv:1605.04692

Some of the more heinous diagrams



CMS Z/gamma ratio (8 TeV) measurement compared to different generators

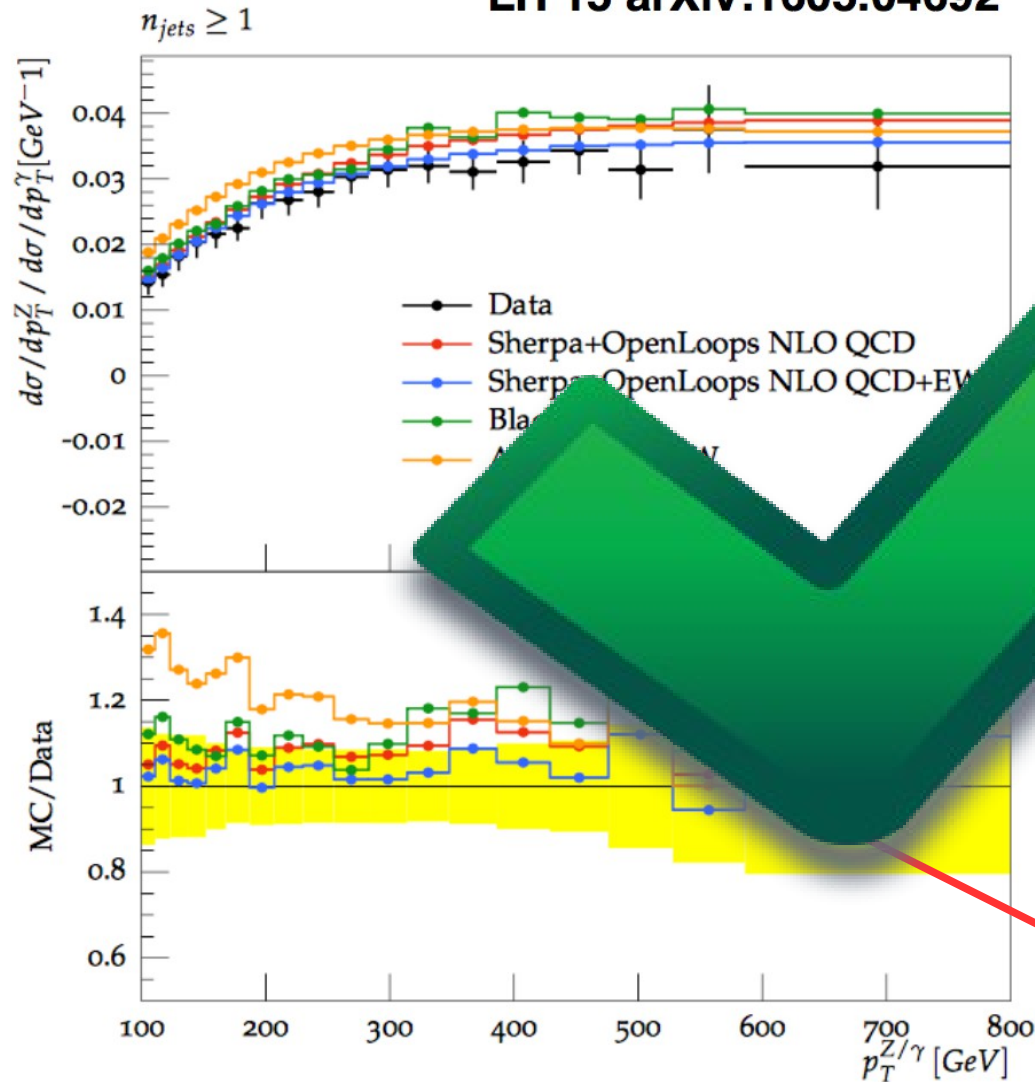


Adding the EWK corrections brings back agreement

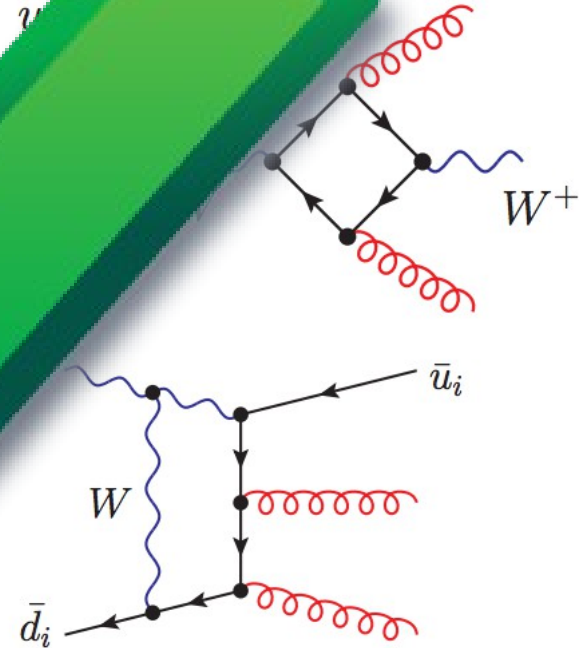
Solving the Mystery

LH'15 arXiv:1605.04692

Some of the more
heino diagrams



CMS Z/gamma ratio (8 TeV) measurement
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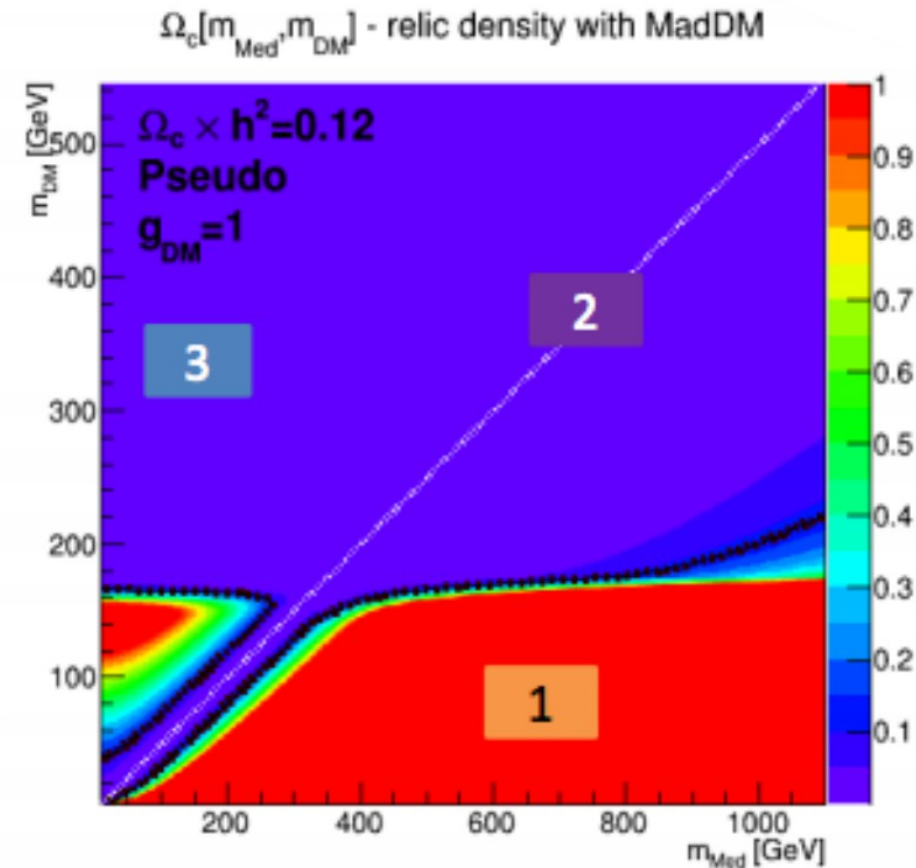


Adding the EWK
corrections brings
back agreement

2014...again!

Heuristic example to how it works

- [1] $m_\chi < m_t \rightarrow$ tight constraints
 - Suppression of the annihilation Process $\chi\chi \rightarrow Med \rightarrow SM$
- [2] $m_\Phi \sim 2 \cdot m_\chi \rightarrow$ weaker constraints
 - On-shell $\chi\chi \rightarrow Med$
- [3] $2 \cdot m_\chi > m_\Phi \rightarrow$ weaker constraints
 - Less suppressed annihilation $\chi\chi \rightarrow Med \rightarrow SM$



With Direct Detection

Vector(SI)



Spin independent

Extremely good

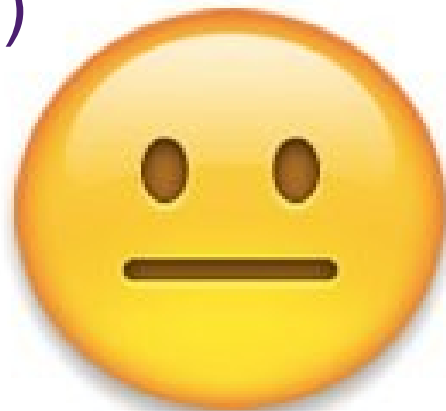
Axial(SD)



Spin dependent

Not so great

Scalar(SI)



So-so

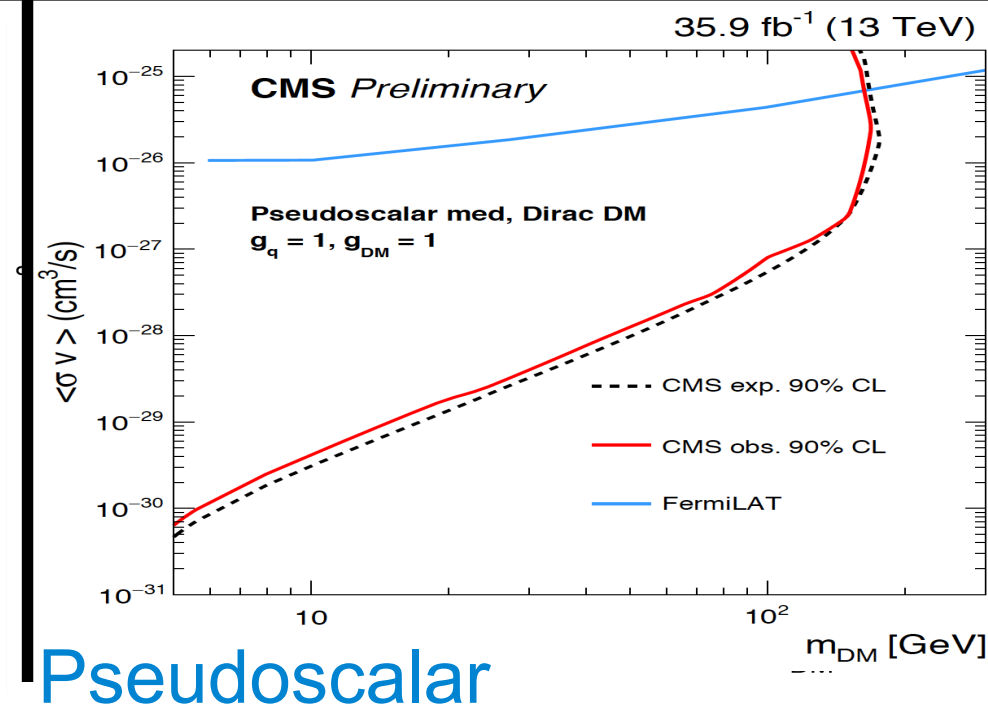
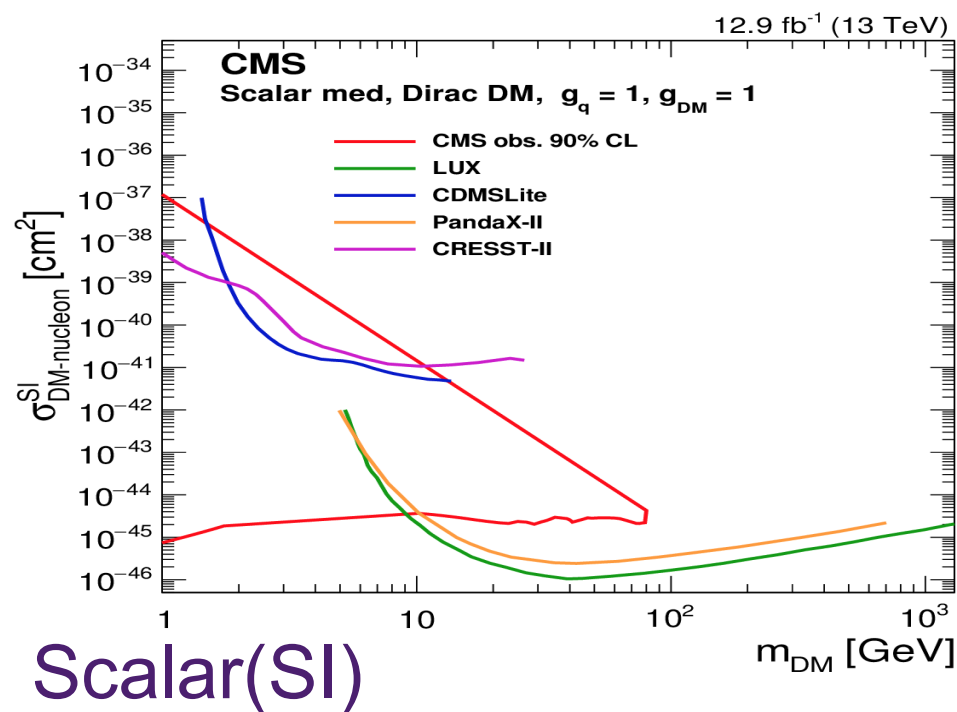
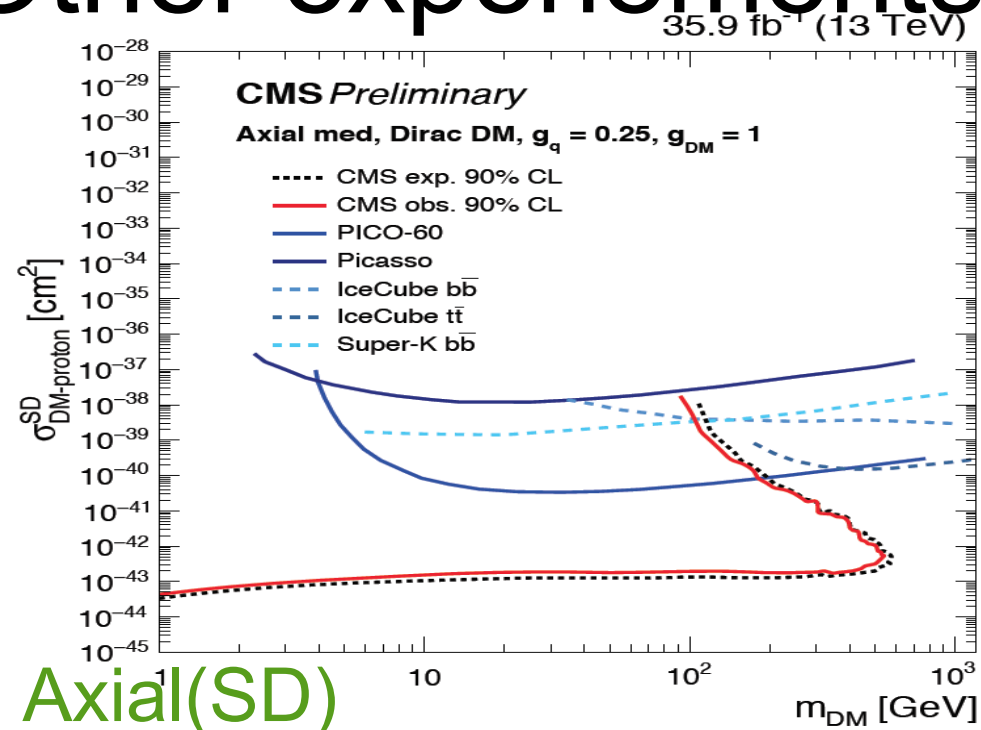
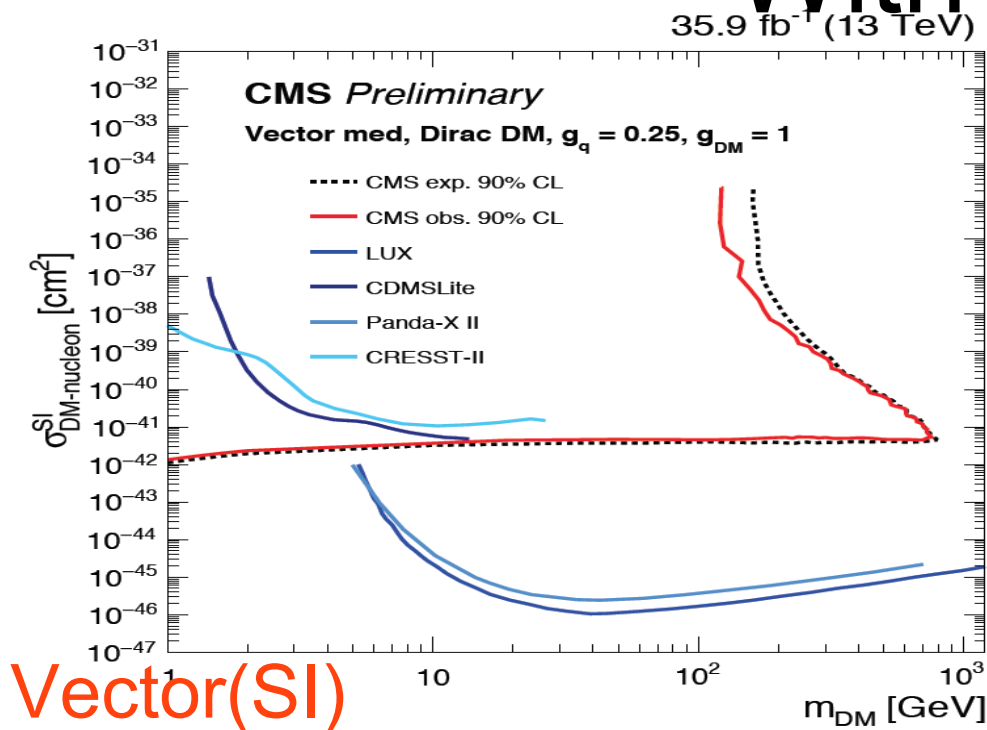
Spin independent

Pseudoscalar

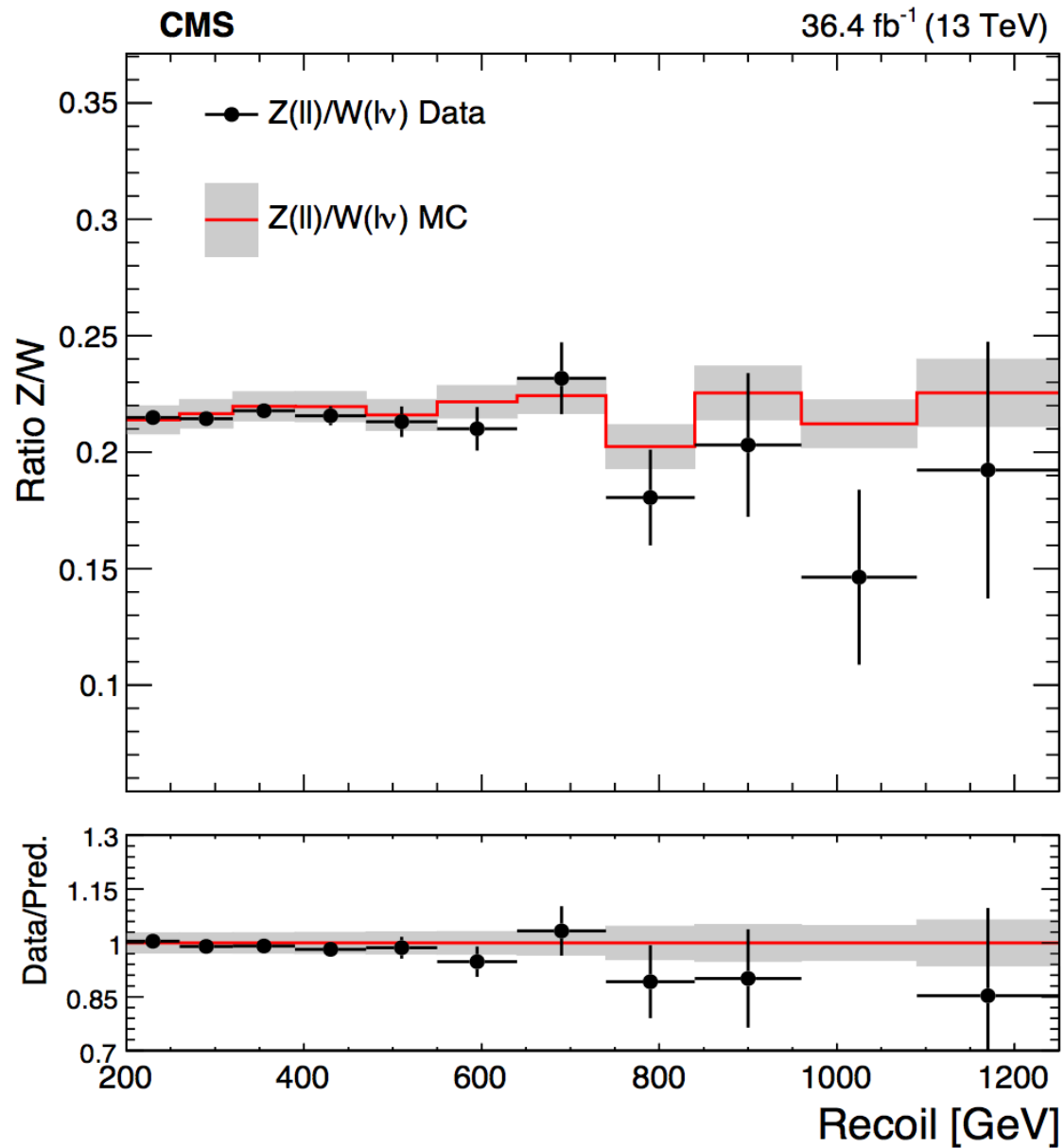


Use indirect detection

With Other experiments



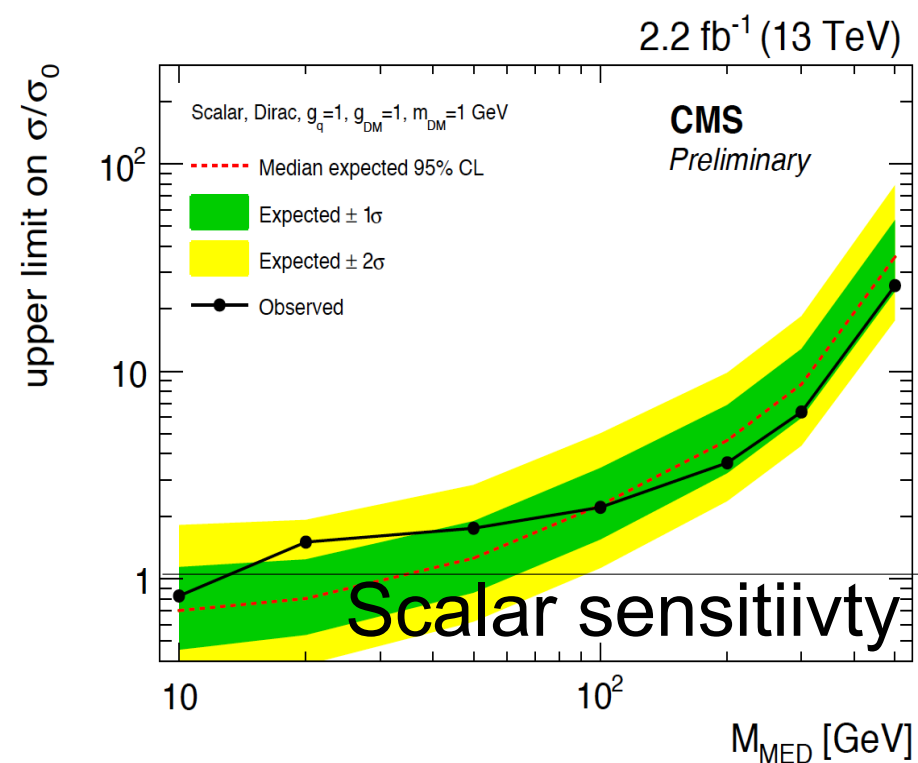
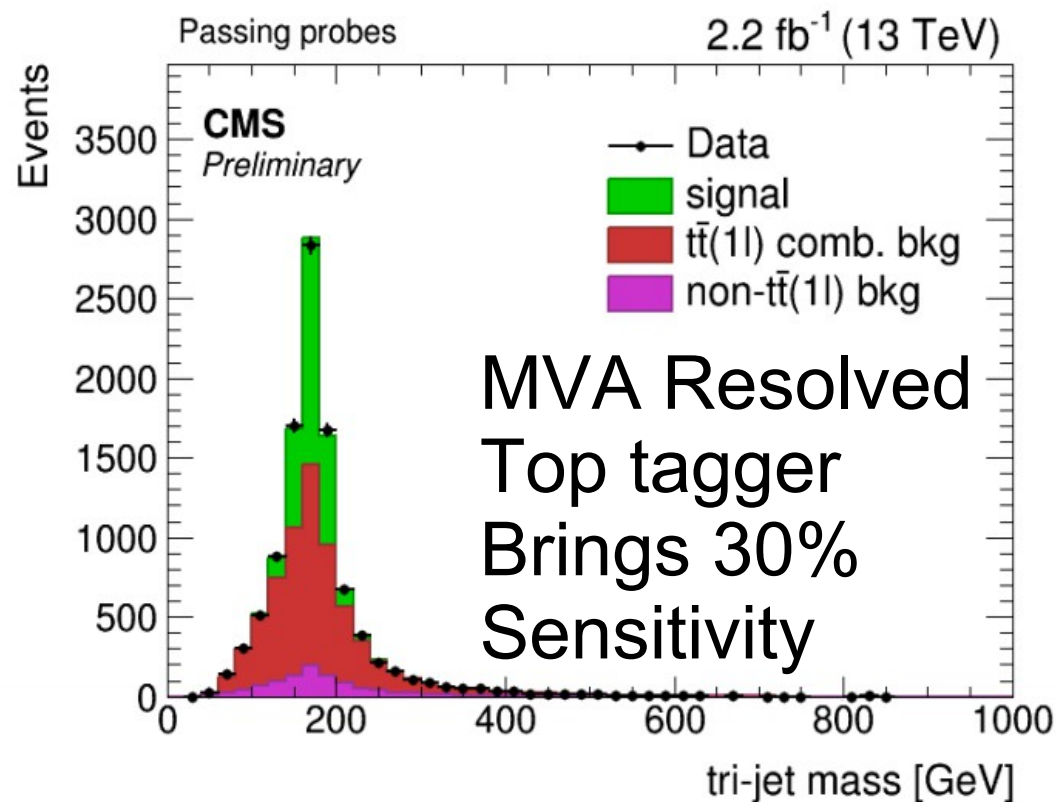
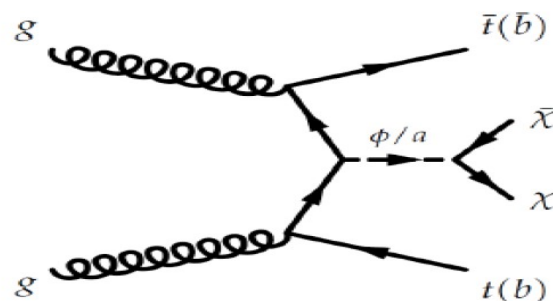
Can we really shrink unc?



Total Luminosity : 2.3 fb^{-1} (2015)

Final State : $t\bar{t} + MET$

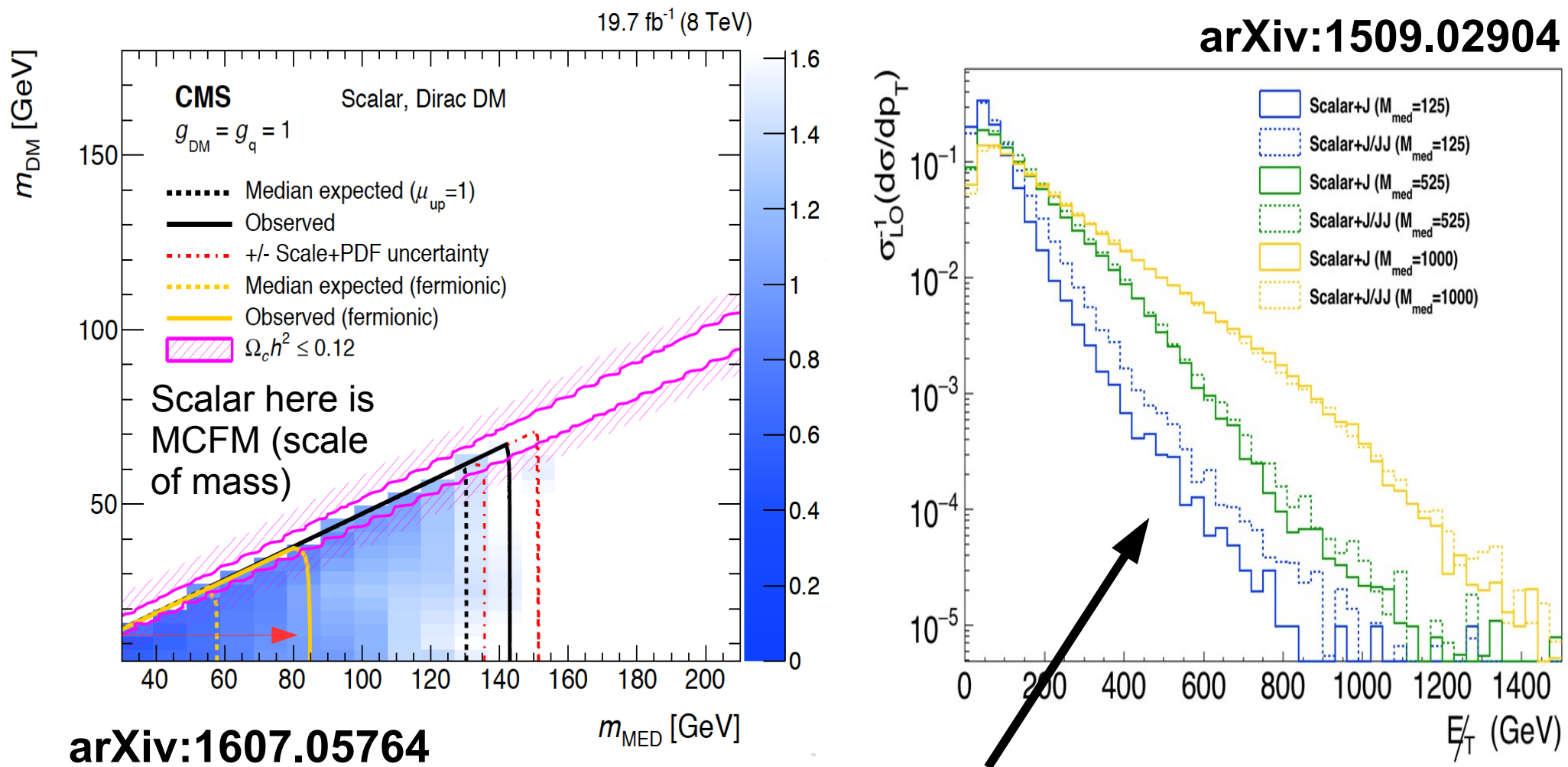
$t\bar{t} + \text{DM}$



Sensitivity pushed to exclude Scalar mediator models

Comparing with Old results

- Exclusion for scalar up to 80 GeV

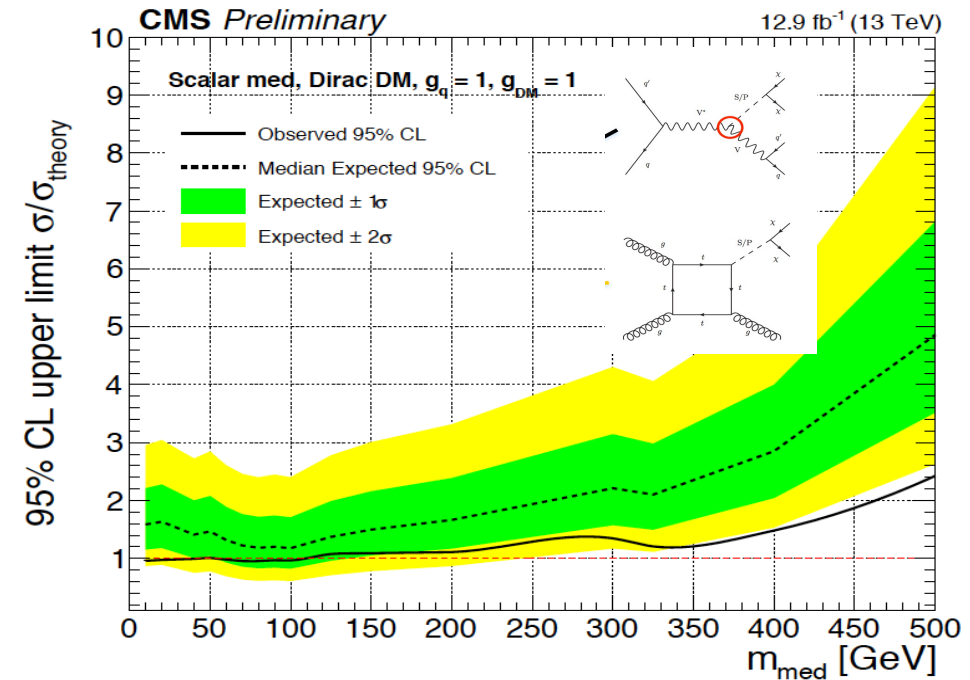
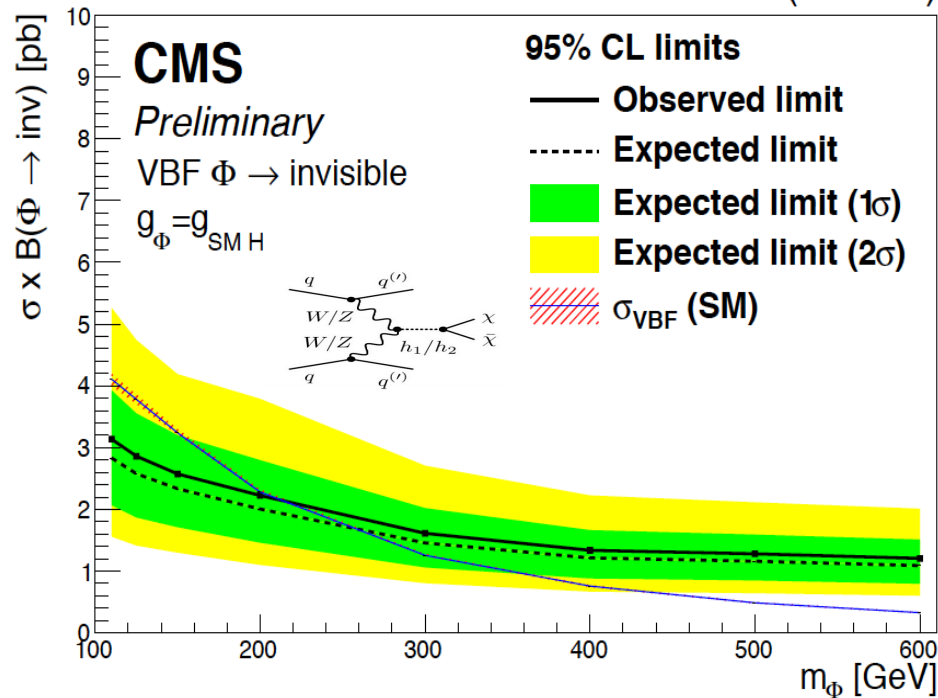


Expect the sensitivity to grow just by adding the 2 jet multiplicity back

Varying the mass

- When we vary the mass

2.3 fb⁻¹ (13 TeV)



Changing mass is a search for a different mediator

This state allows for combination of vector boson copuligns

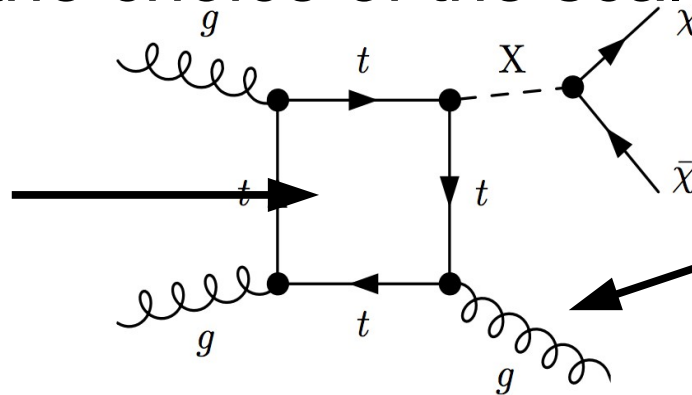
Adding more jets

- Can split the signal models into 2 different sources

- What is the choice of the scale?

Generate the two scales and compare

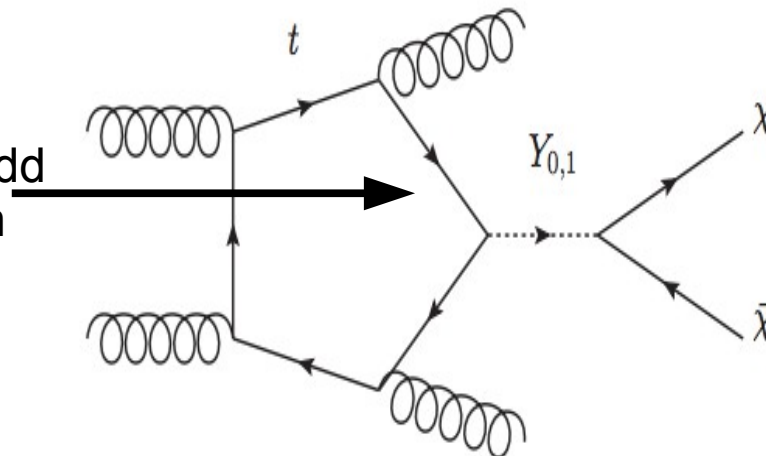
Inside the loop best motivated scale is Top mass



Outside the loop Best motivated scale is the jet

- What is the contribution from higher order effects?

Diagrams like these add significant contribution

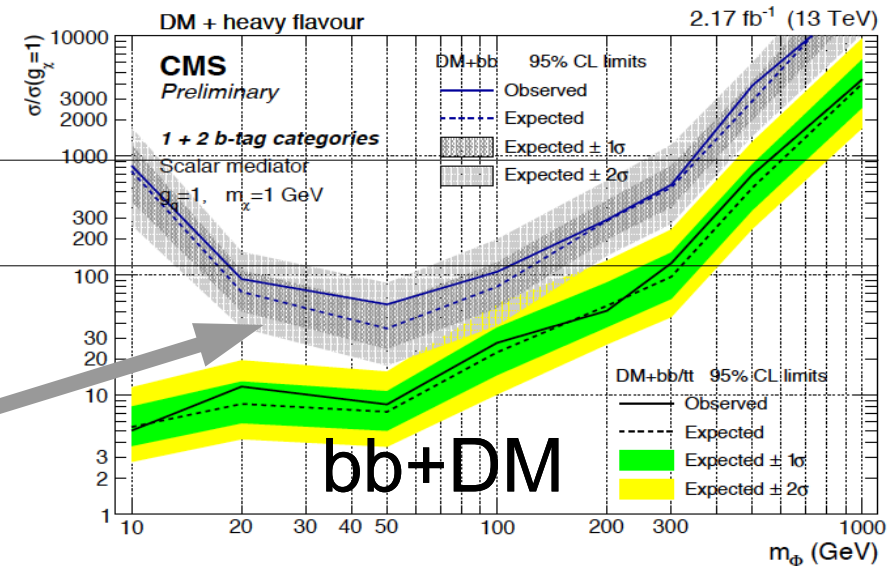
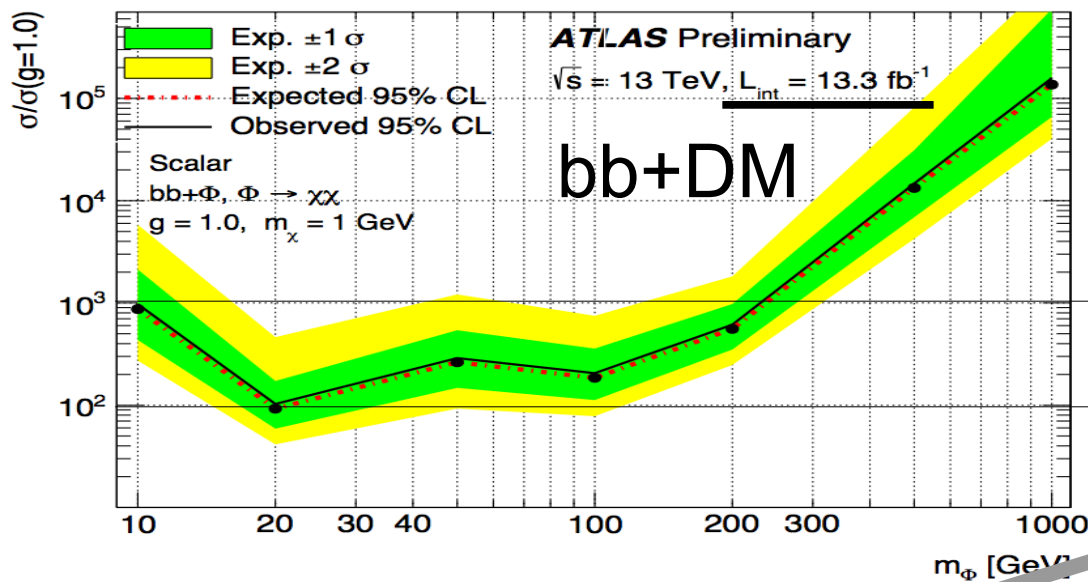


Generate the two jet diagrams and merge them with the one jet

We will soon take advantage of the 2nd jet

ATLAS vs CMS : Scalar HF results

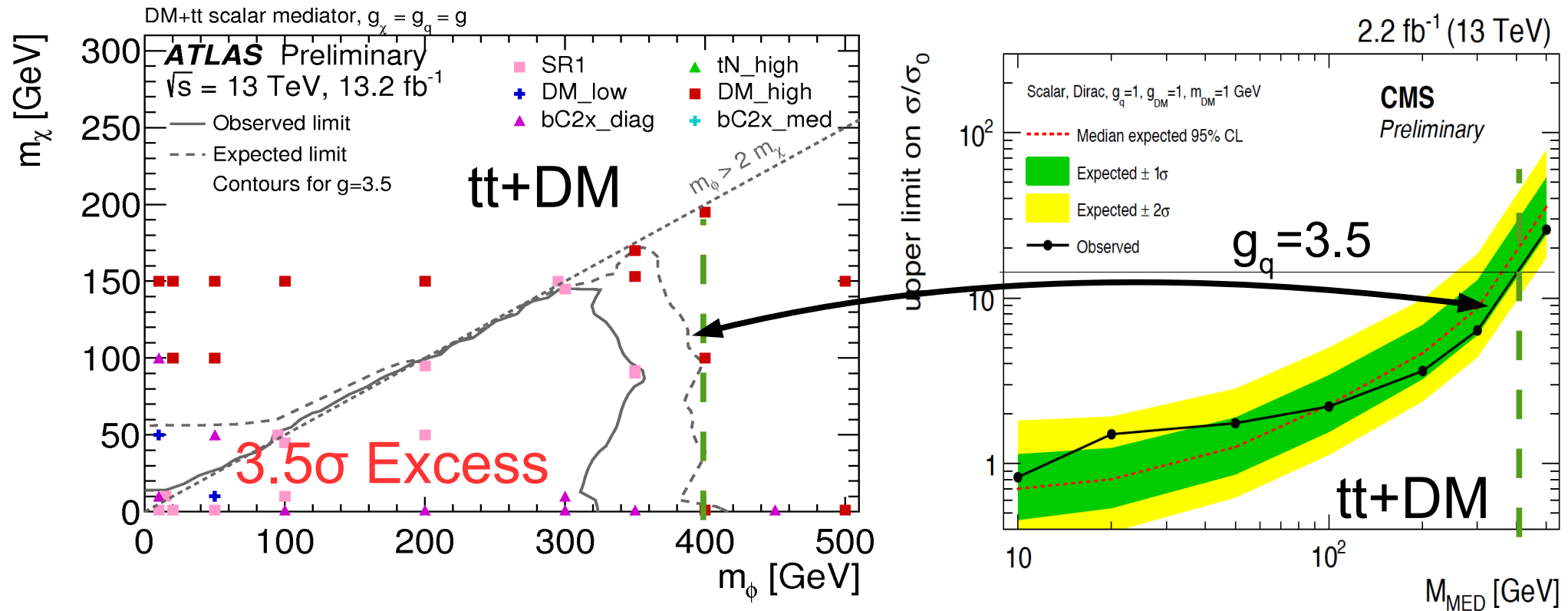
ATLAS result 5x data



CMS results roughly the same
 Even though 1/5 luminosity

ATLAS neglects tt +DM and bb +DM combined interpretation
 Neglects a potentially large contribution

ATLAS vs CMS : Scalar HF results



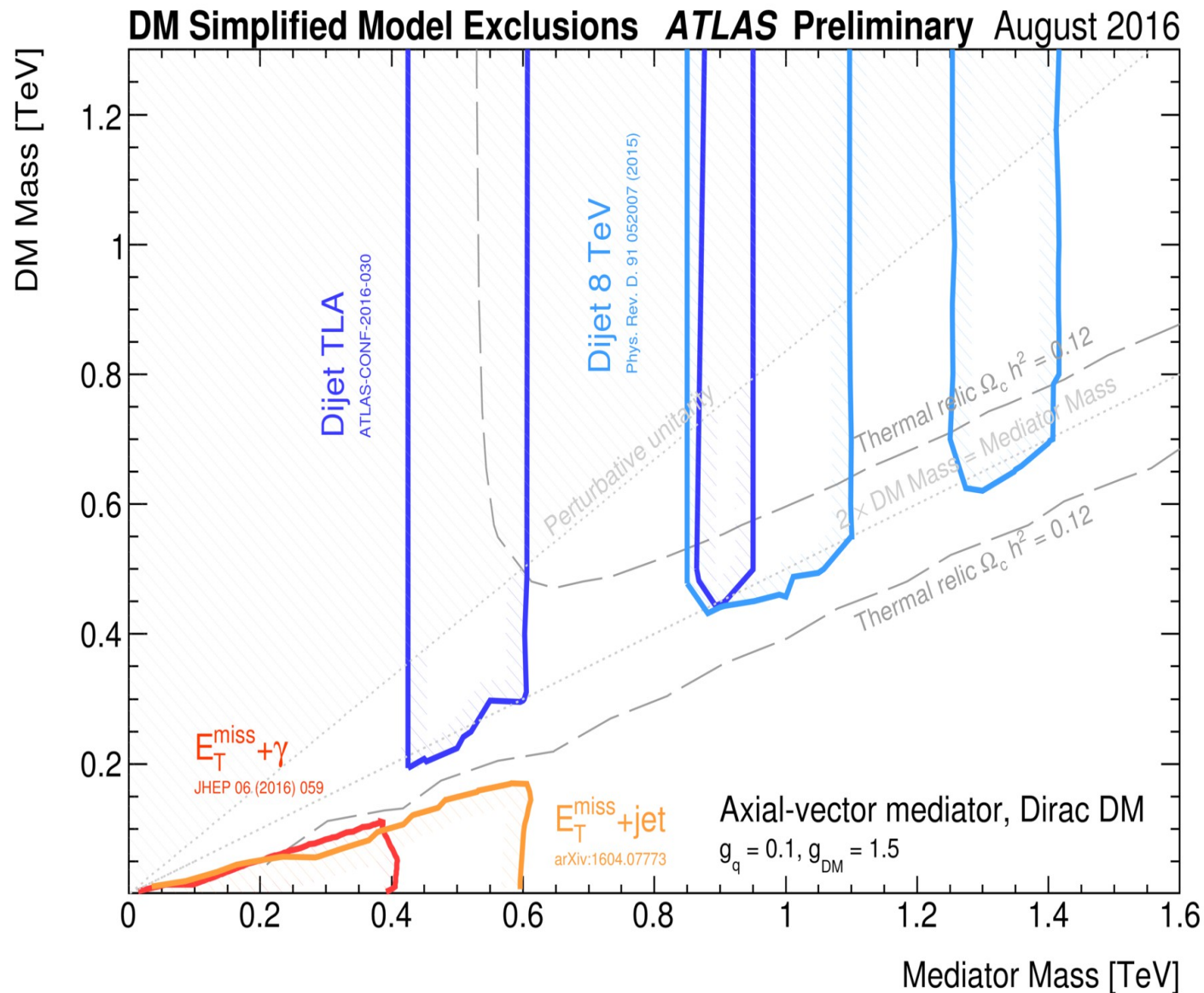
ATLAS result 5x data

CMS results have roughly the same sensitivity

CMS excludes the tt+DM excess in ATLAS

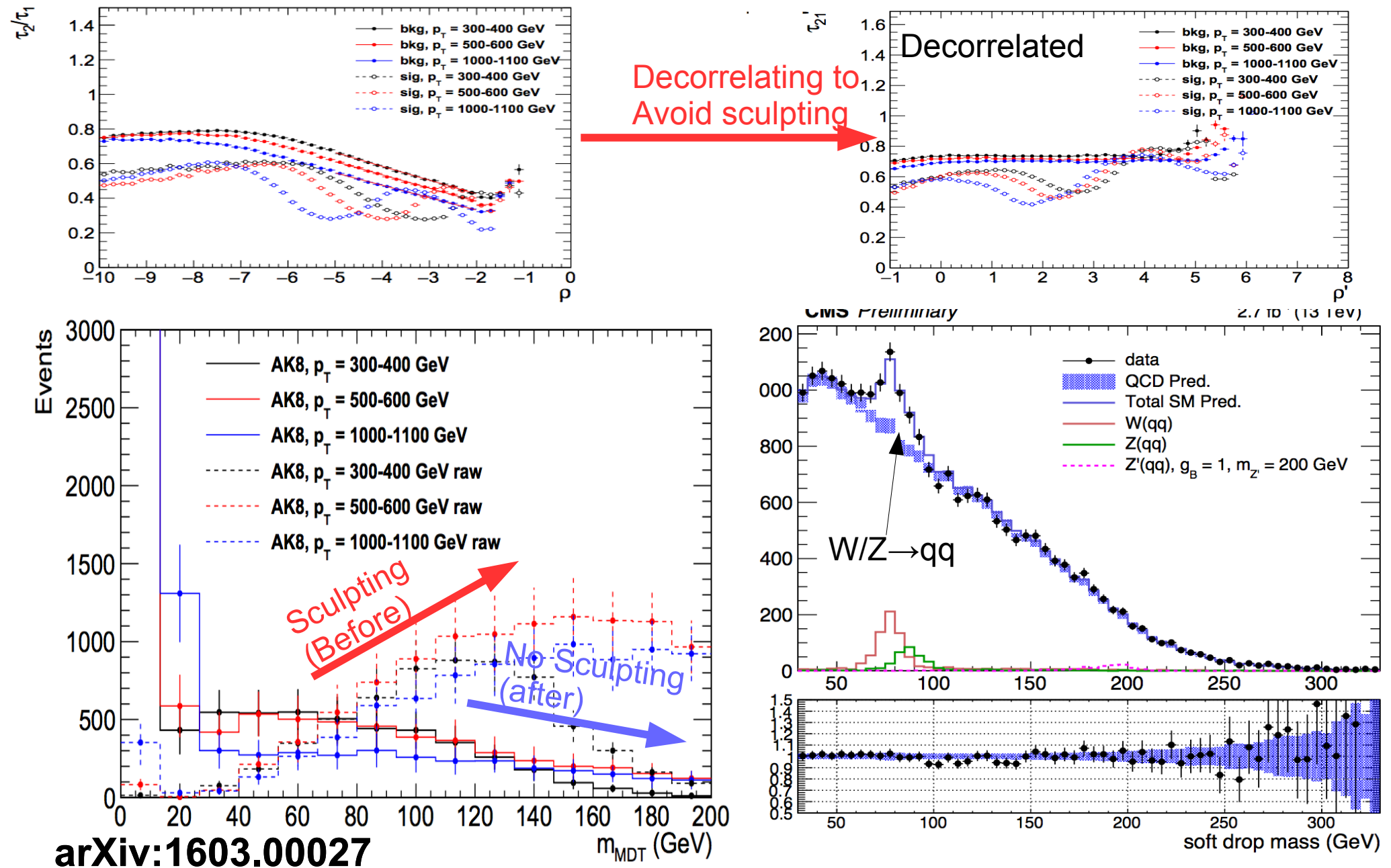
ATLAS modified couplings

<https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CombinedSummaryPlots/EXOTICS/index.html>



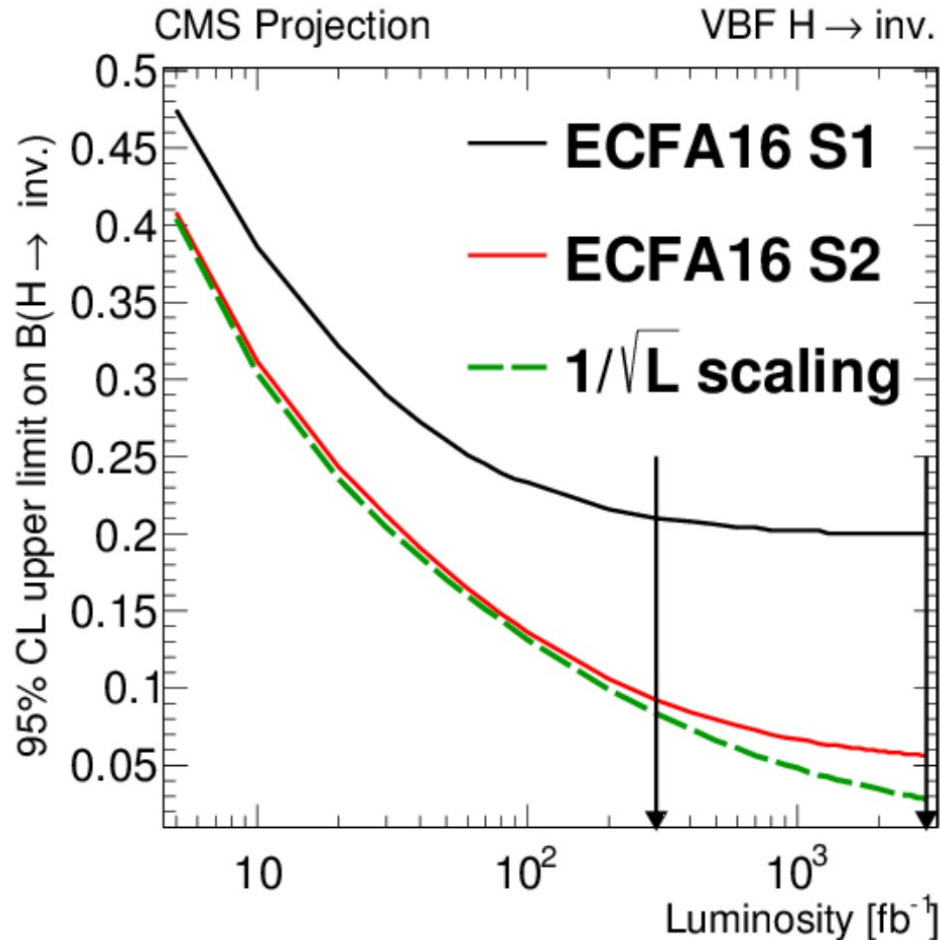
$$g_{\text{SM}} = 0.1$$
$$g_{\text{DM}} = 1.5$$

Design a transform to decorrelate against mass and p_T



How will things scale?

- Higgs to Invisible



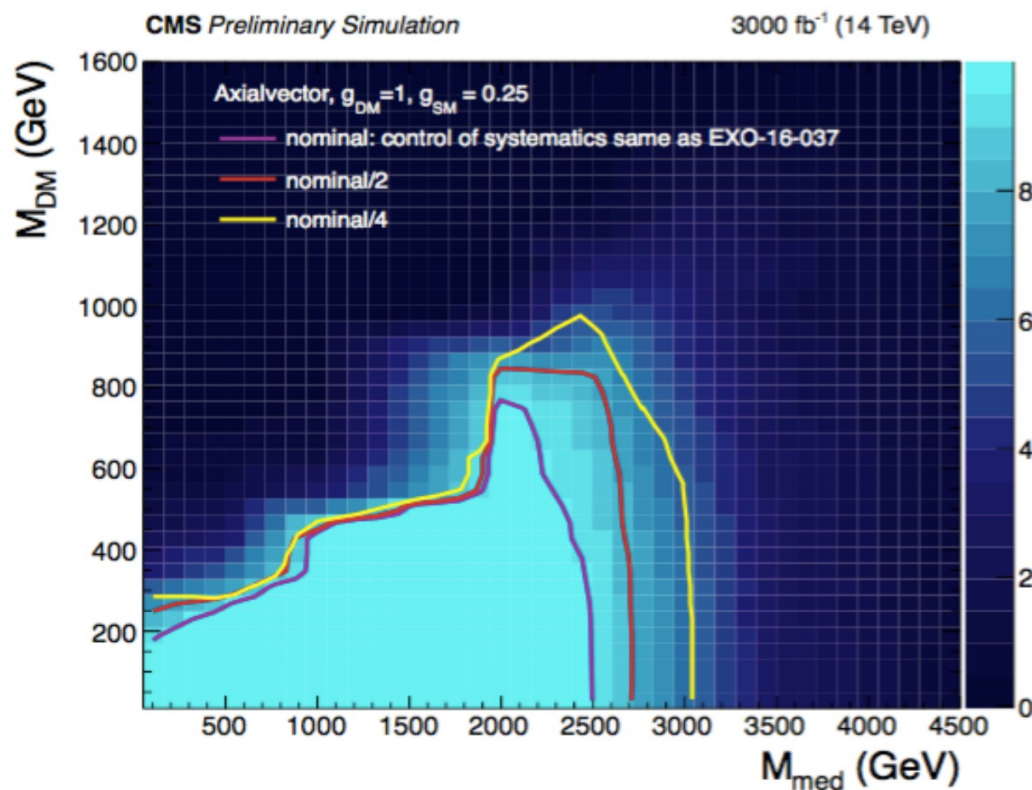
Can reach

$BR(H \rightarrow \text{Inv}) < 5.6\%$ in VBF

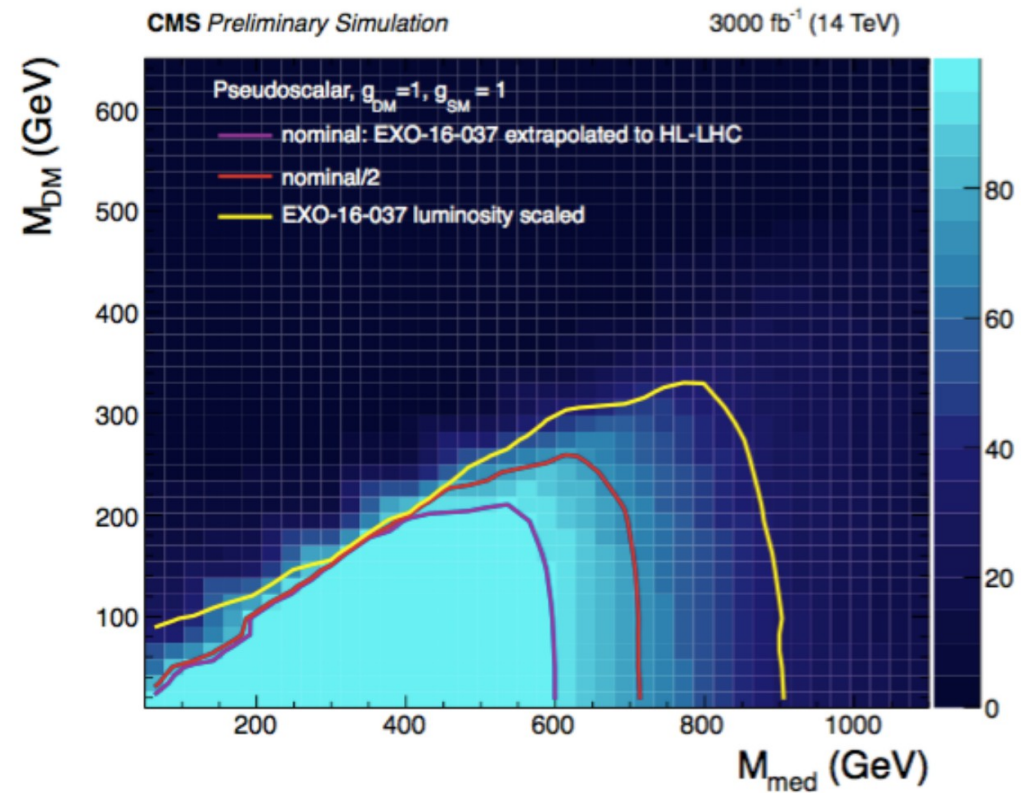
All channels $BR(H \rightarrow \text{Inv}) < 3\%$

	ECFA16 S1	ECFA16 S2	$1/\sqrt{L}$ scaling
300 fb^{-1}	0.210	0.092	0.084
3000 fb^{-1}	0.200	0.056	0.028

Monojet Projections



For Vector/Axial-Vector
reach is 3 TeV



For Pseudoscalar reach
is 800-900 GeV
Scalar reach is
600-700 GeV

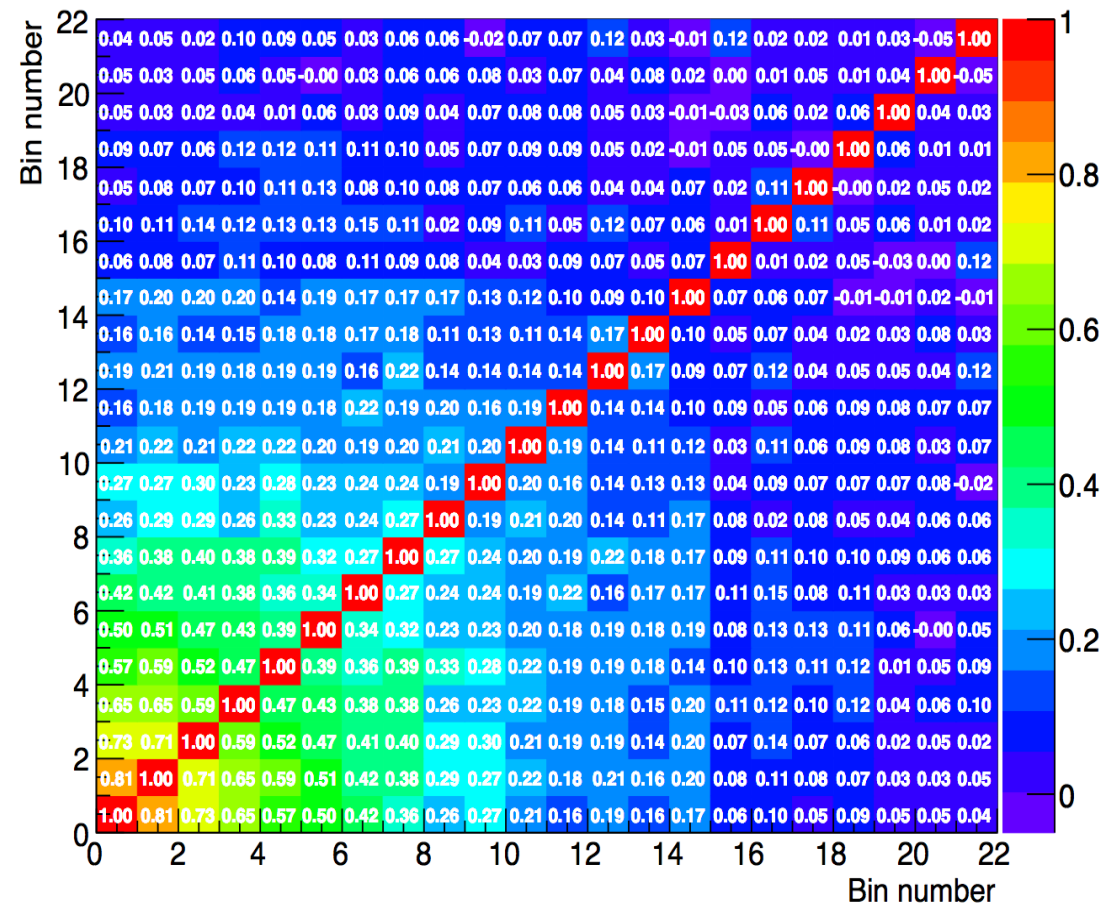
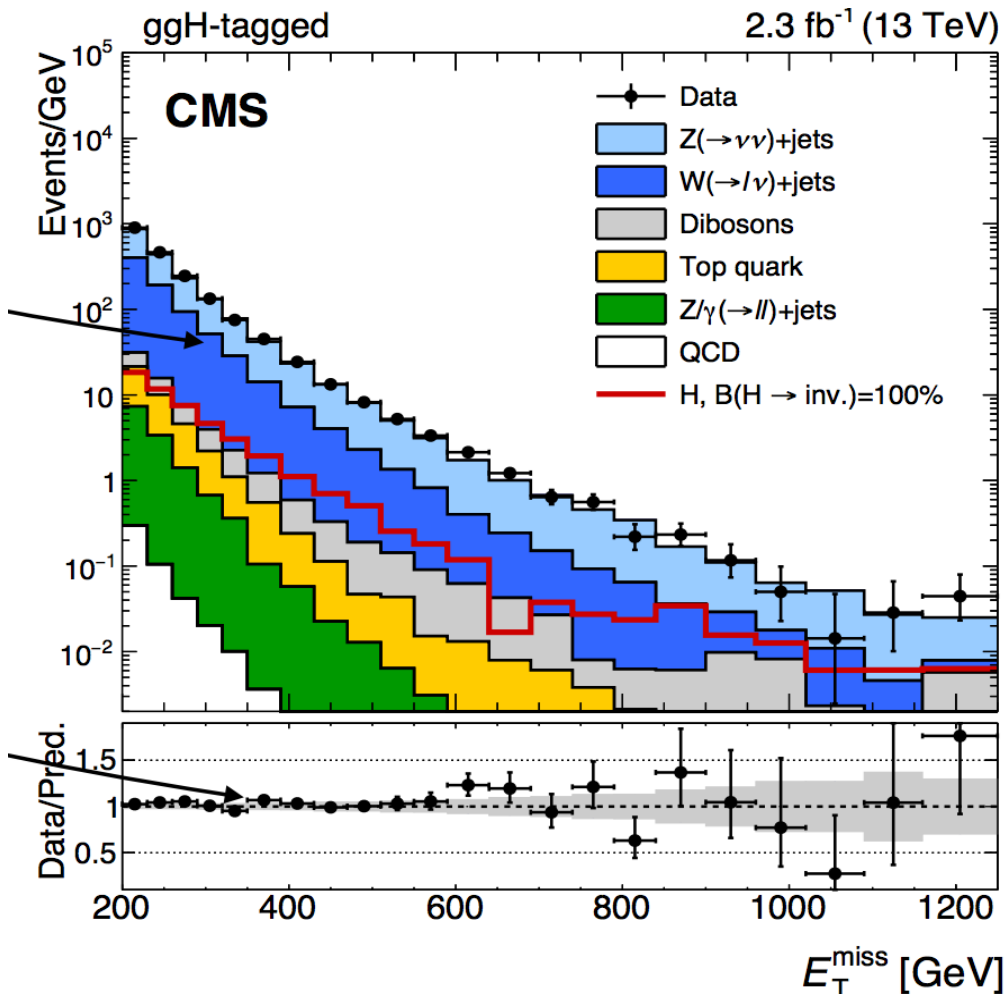
Re-interpreting the Analysis

- Future plans from LHC→simplified likelihoods

- What is it? => Reduced control fit to 2 objects

MET distribution

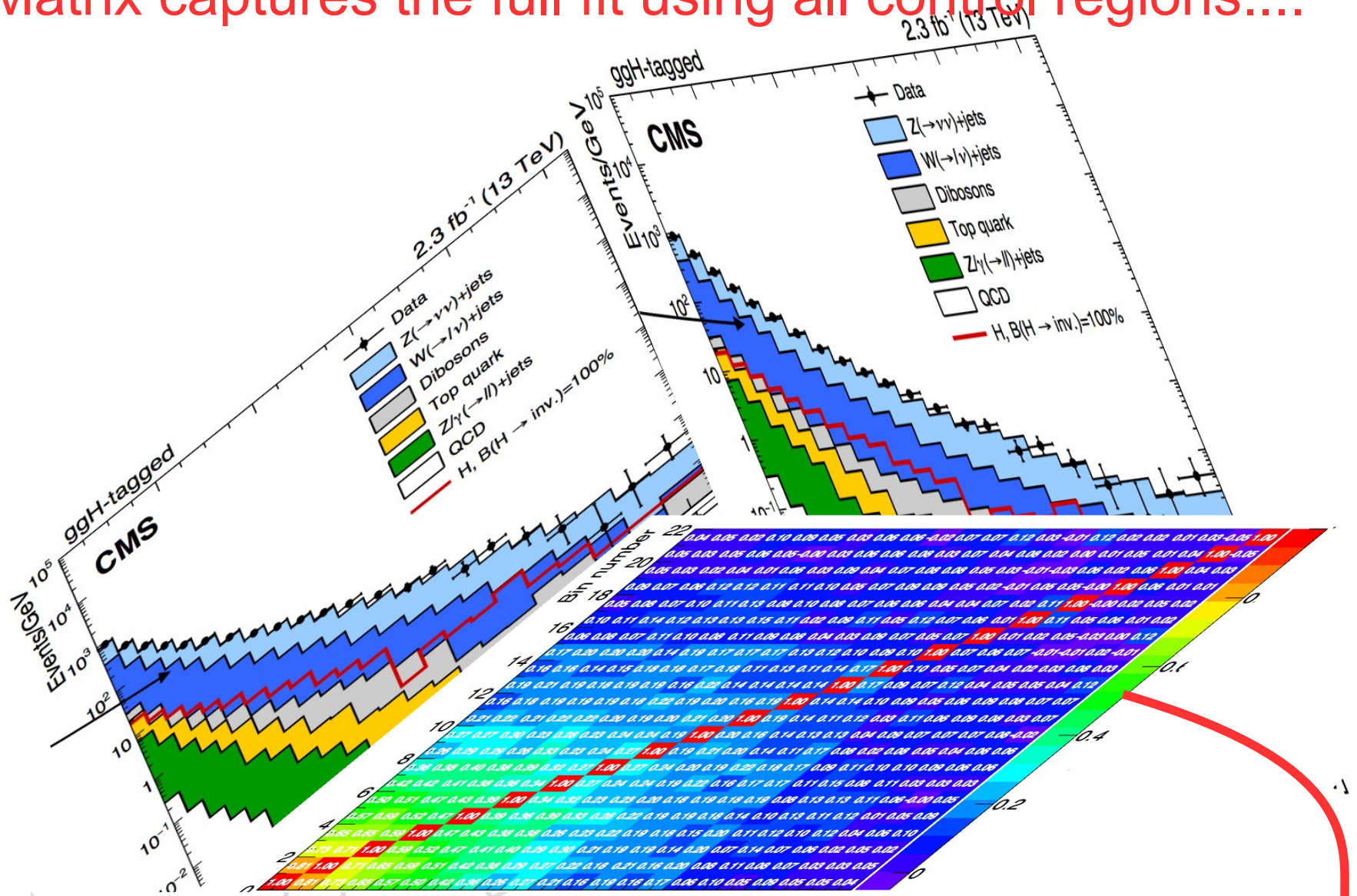
Toy Bin by bin covariance



From this: setup full CL_s get both expected and observed

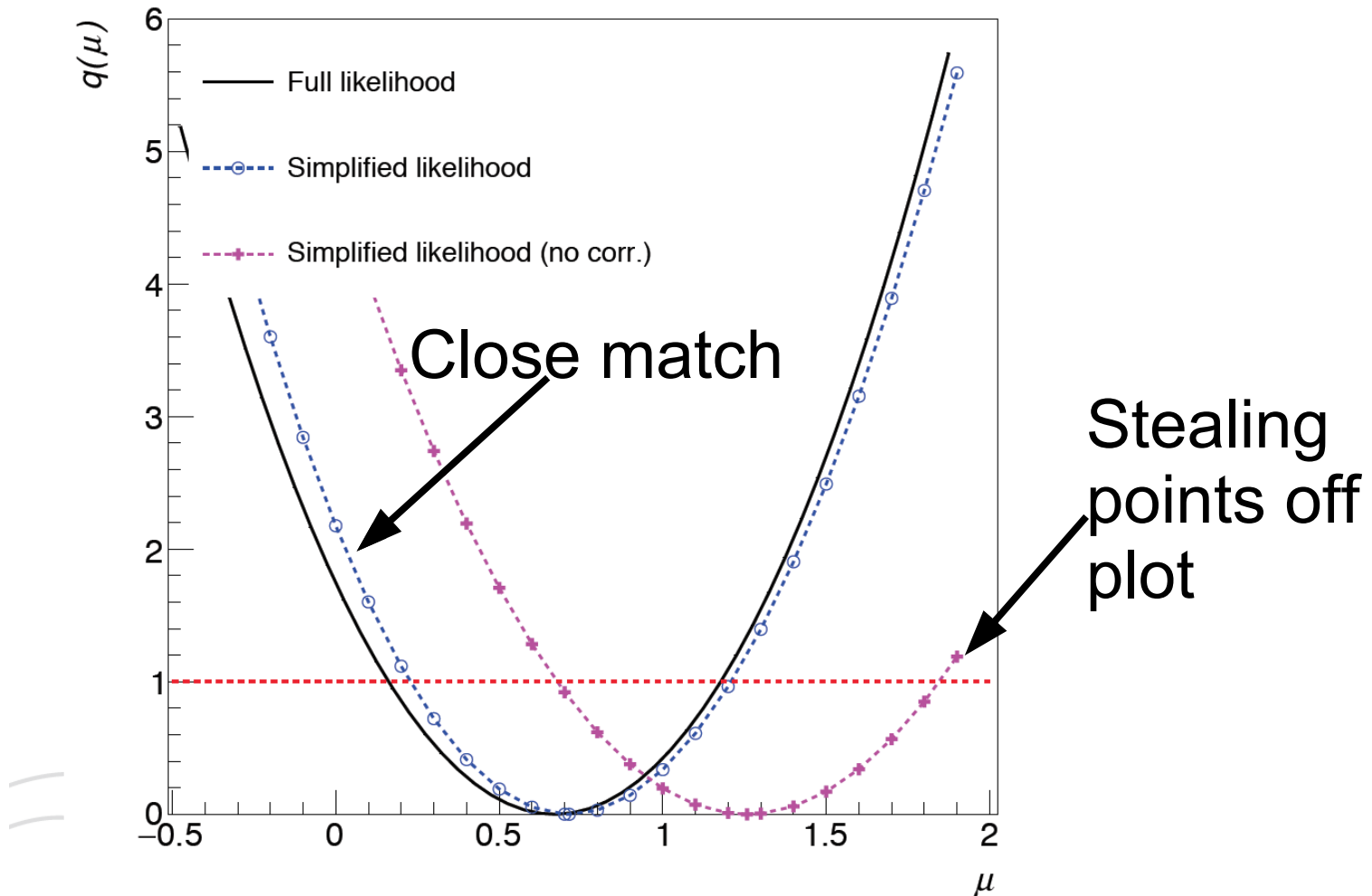
Simplified likelihood¹⁵⁷

Matrix captures the full fit using all control regions....



$$\mathcal{L}_S(\mu, \theta) = \prod_{i=1}^N \frac{(\mu \cdot \underset{\text{Signal}}{s_i} + \underset{\text{background}}{b_i} + \theta_i)^{n_i} e^{-(\mu \cdot s_i + b_i + \theta_i)}}{n_i!} \cdot \exp \left(-\frac{1}{2} \theta^T \mathbf{V}^{-1} \theta \right)$$

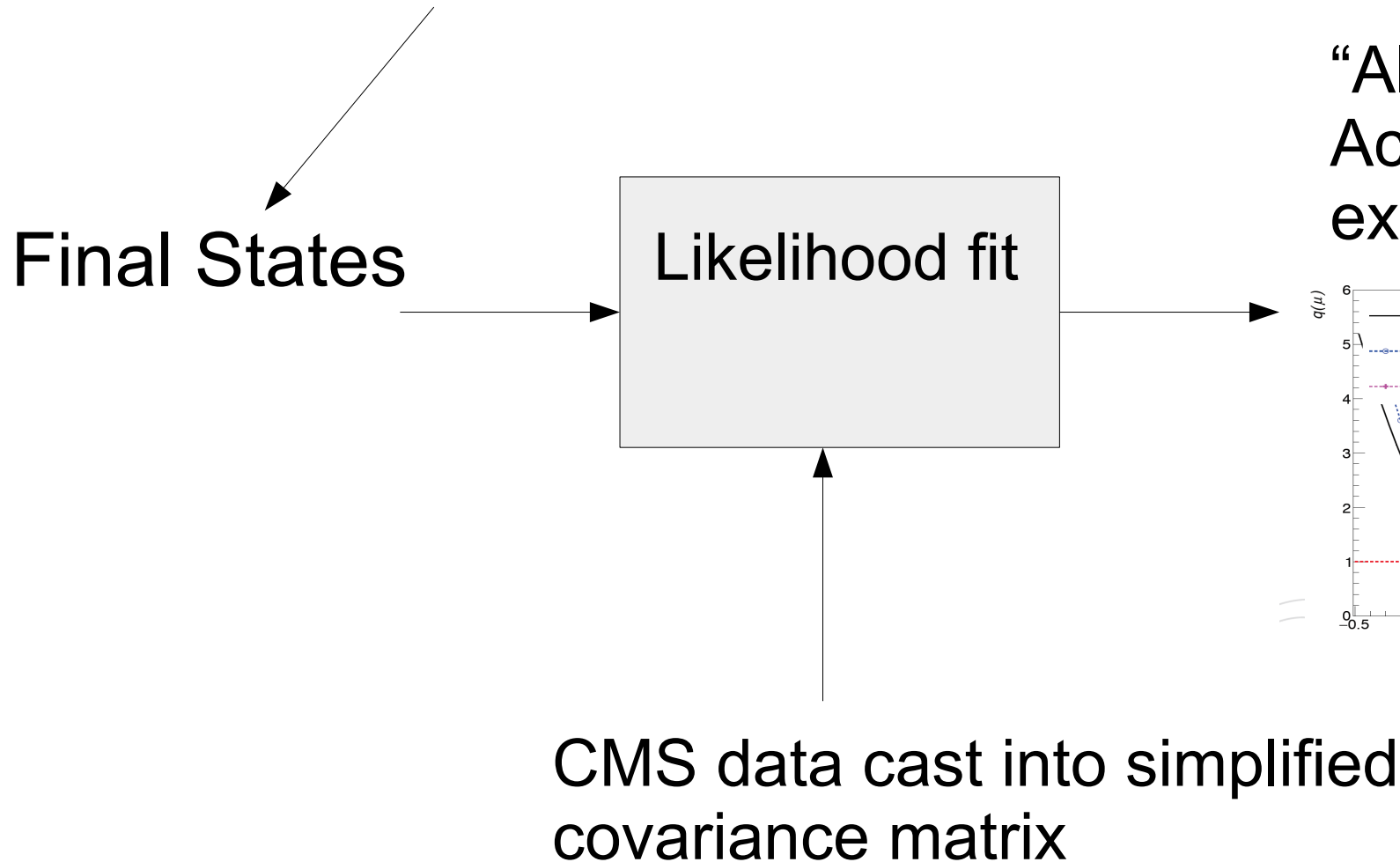
Whats the accuracy?



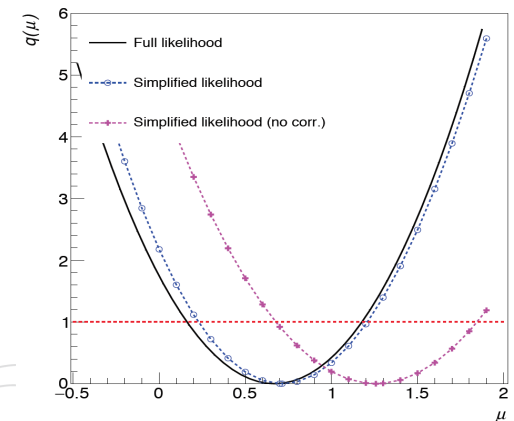
CMS is aiming to release simplified likelihoods for
all future SUSY and Dark matter searches

What does this mean?

Lets say you have a new model



“Almost”
Actual LHC
exclusion

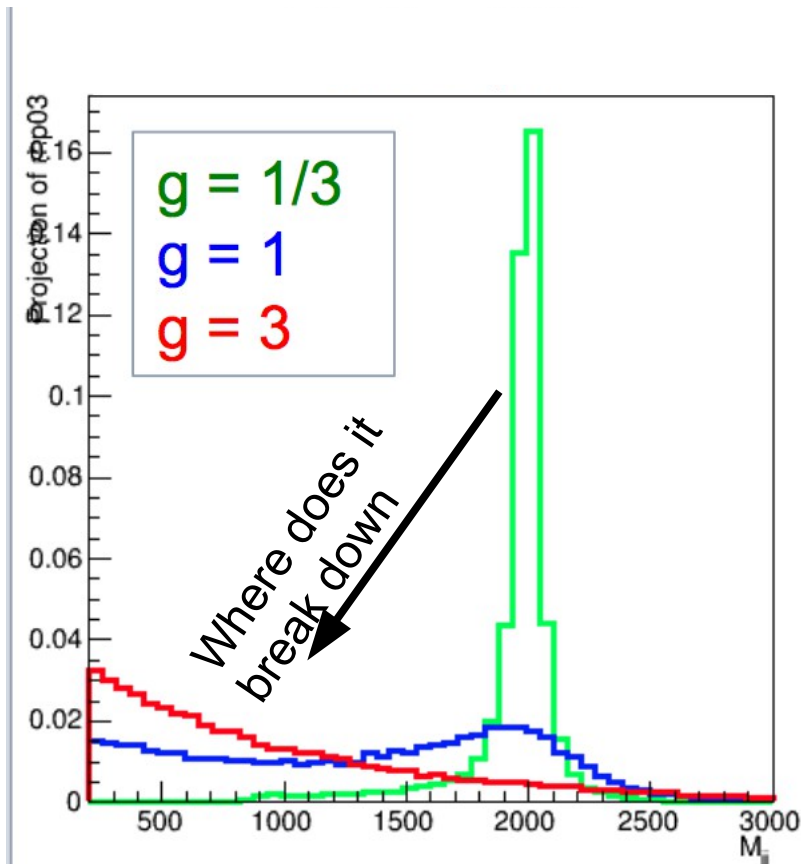


Real reason for this :

We are tired of running every theorists' model

Follow up

- The resonant searches break down
 - Larger the coupling the wider the resonance
 - Bump hunts only work with a relatively narrow bump
 - With large couplings cross section is larger



- To probe wide resonance
 - Alternative search strategy
 - Typically with angular analysis
 - Searches tend to be weaker
- Good to know if there is missing phase space

Adding EWSB in a better way

Khoze *et al.*, 1505.03019

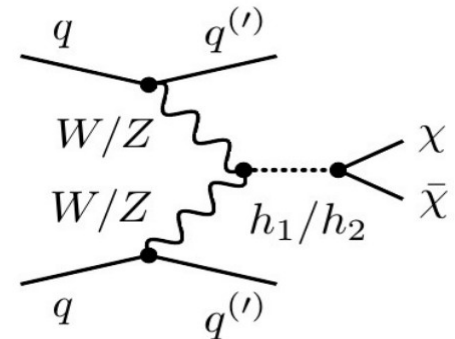
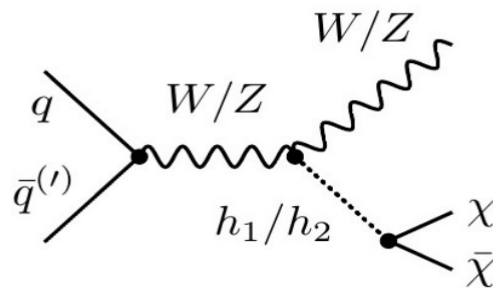
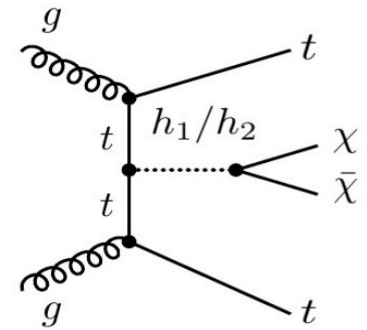
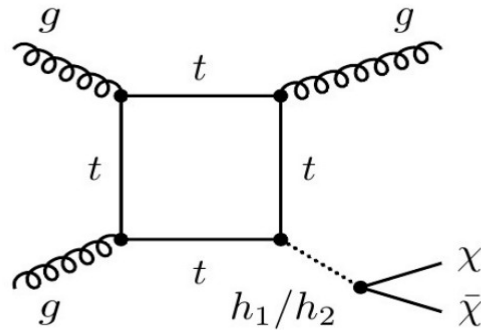
$$\begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} h \\ s \end{pmatrix} \quad \mathcal{L} \supset -y_{\text{DM}} (\sin \theta h_1 + \cos \theta h_2) \bar{\chi} \chi$$

$$+ (\cos \theta h_1 - \sin \theta h_2) \left(\frac{2M_W^2}{v} W_\mu^+ W^{-\mu} + \frac{M_Z^2}{v} Z_\mu Z^\mu - \sum_f \frac{m_f}{v} \bar{f} f \right)$$

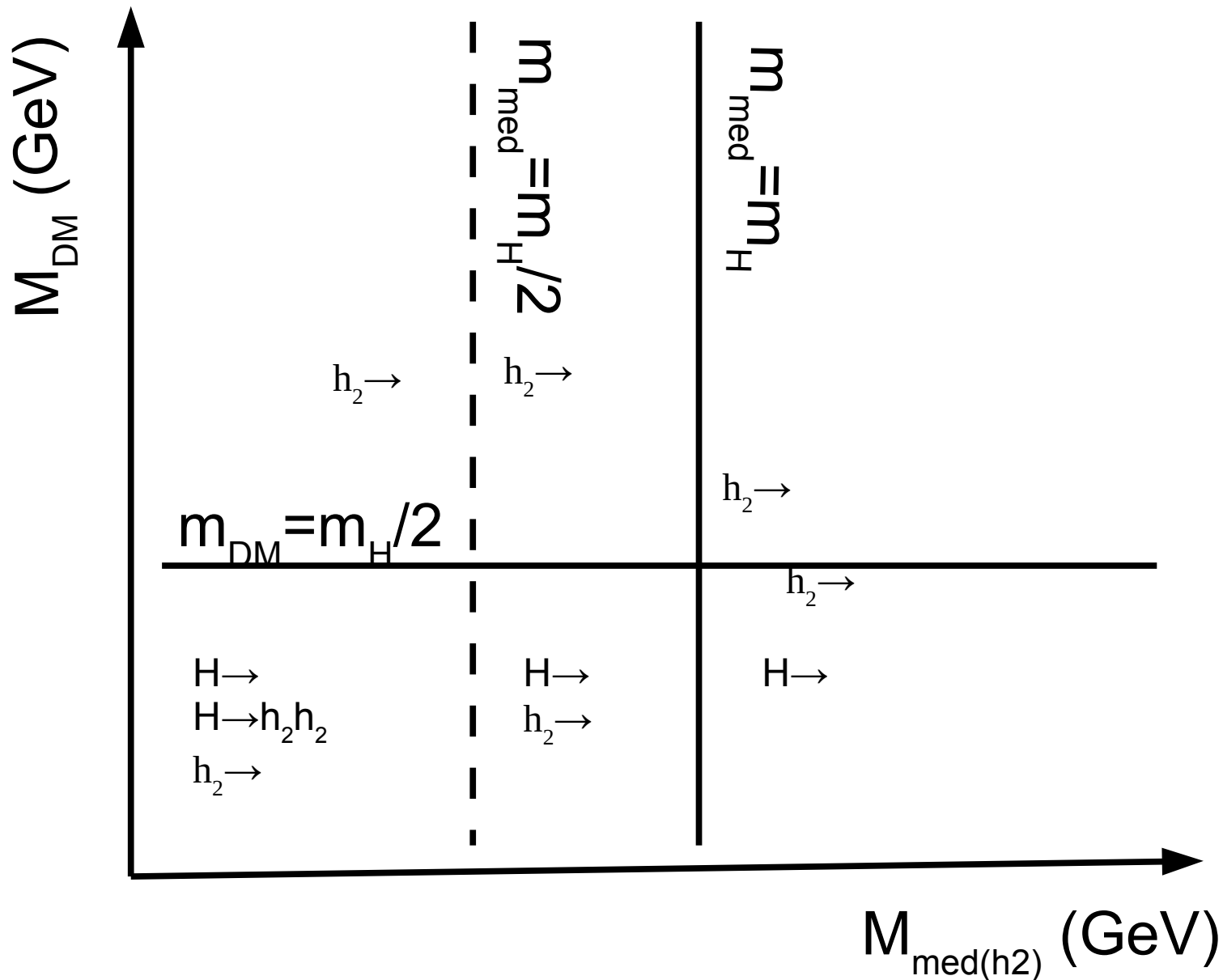
$$g_{\text{DM}} \rightarrow y_{\text{DM}} \cos \theta$$

$$g_{\text{SM}} \rightarrow -\sin \theta$$

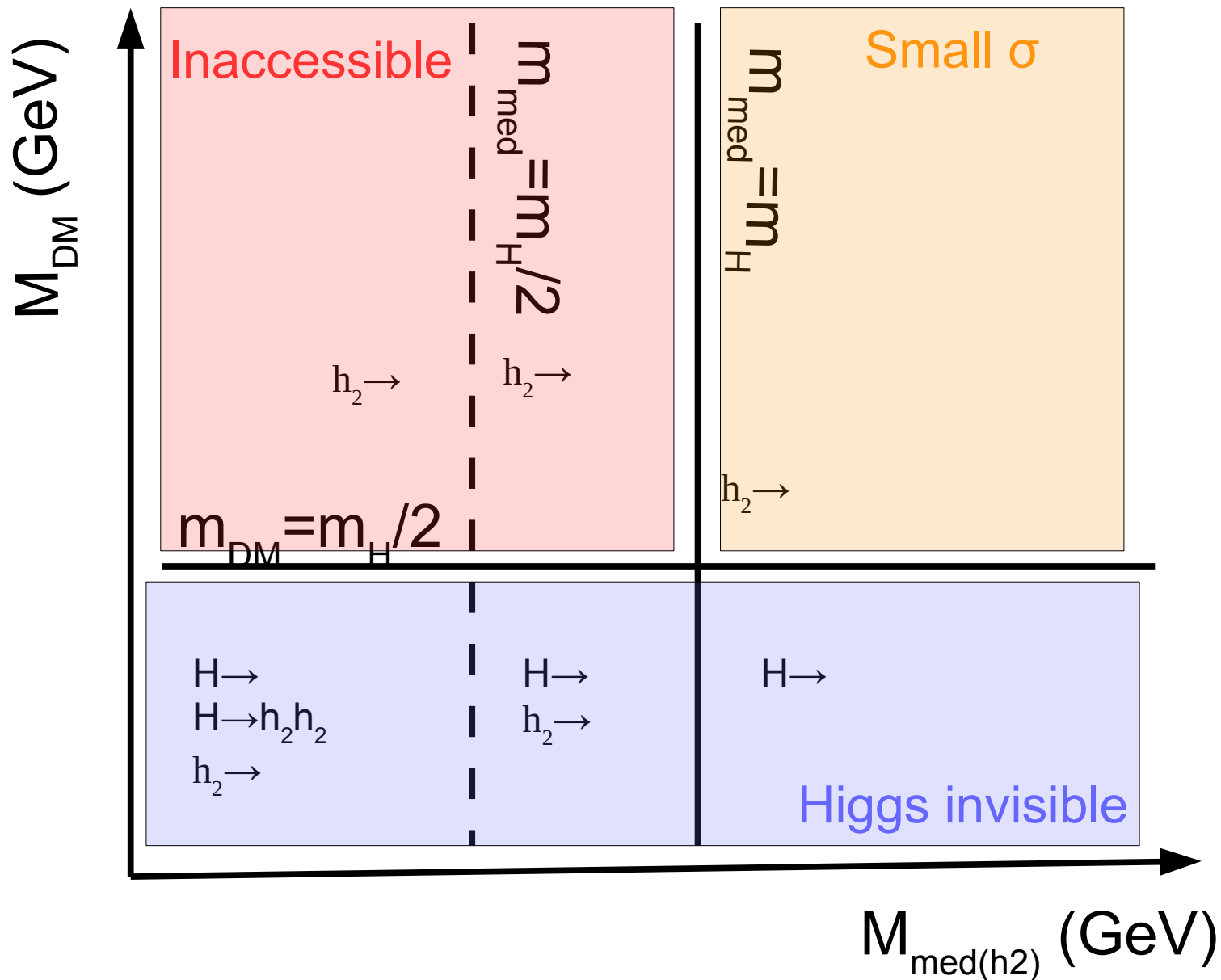
$$\Gamma(h_1 \rightarrow \chi \bar{\chi}) = \frac{y_{\text{DM}}^2 \sin^2 \theta m_{h_1}}{8\pi} \left(1 - \frac{4m_\chi^2}{m_{h_1}^2} \right)^{3/2}$$



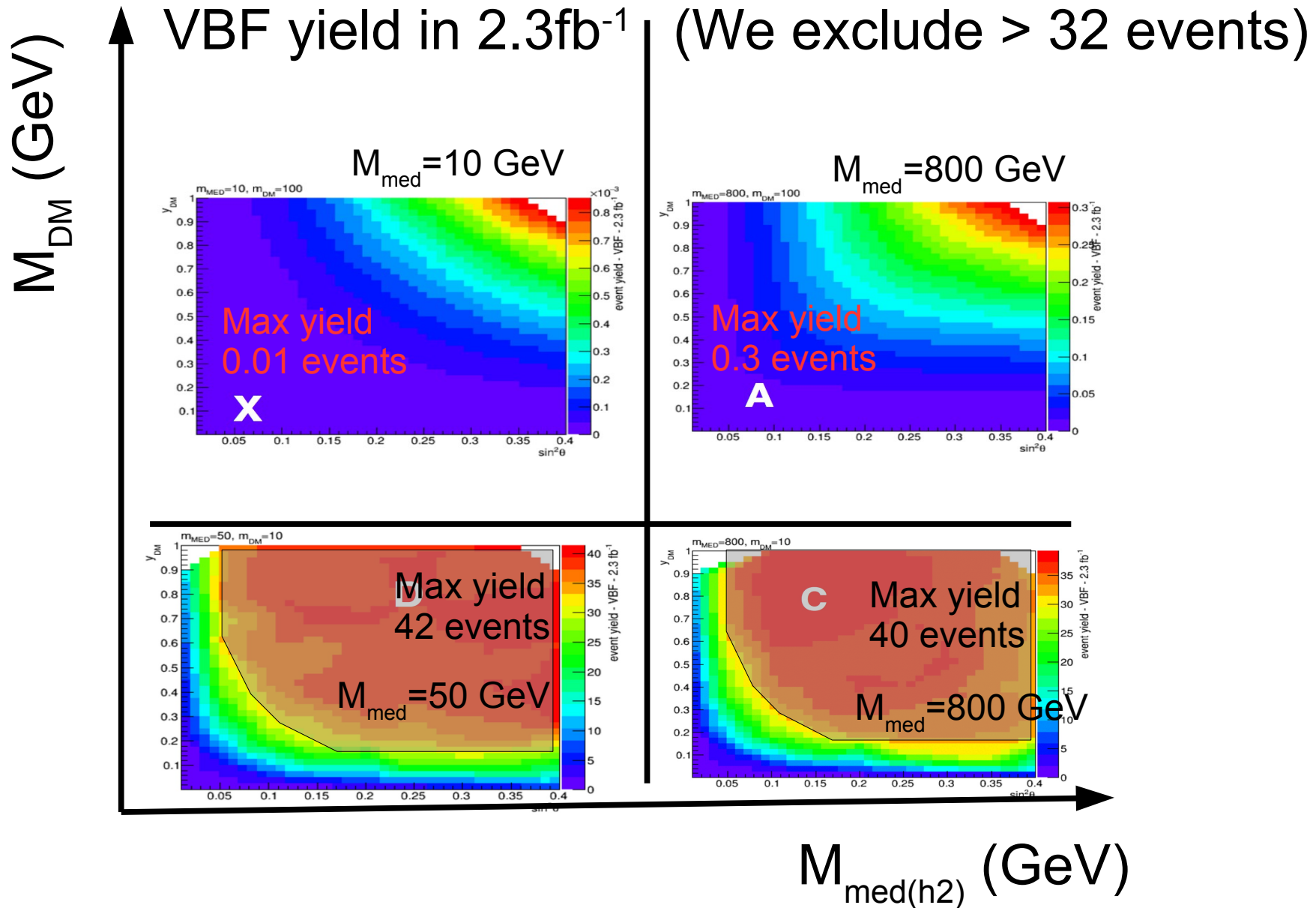
Whats the scope of phase space?



Whats the scope of phase space?



Whats the scope of phase space?



Adding EWSB in an even better way

Looking to embed dark matter in 2HDMs

There has been a lot of work already

Are the current simplified models sufficient?

Dark matter working group actively investigating

Previous slides are a simple example

Composite Models

*** Discussion title: Review of EXO-16-030

Dear Authors

I noticed that in your PAS the introduction refers to UA1 and UA2 results at $\sqrt{s}=300$ GeV. These results came out when I was on UA2 in fact, and the SppS ran at $\sqrt{s}=630$ GeV.

Best regards Joe[Incandela]

Broke a record almost 30 years old



Incidentally this did not happen at the olympics

How does this translate to DM?

With the standard cross section formula

$$\sigma = \frac{g_{SM}^4}{\Gamma} \frac{C}{M_{med}^4}$$

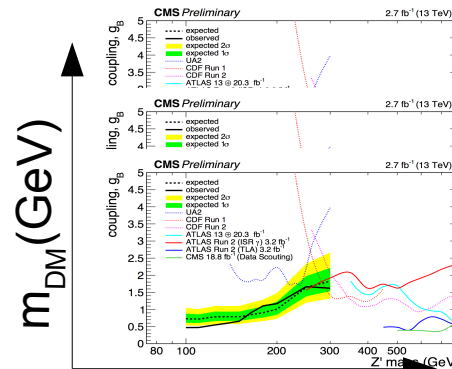
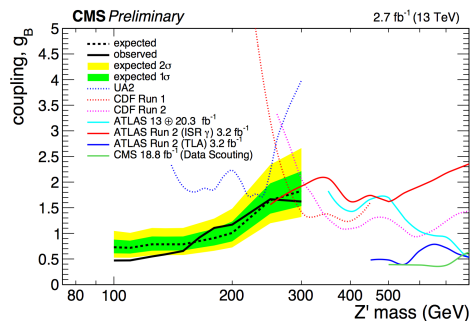
Dijet cross
section@LO

$$\sigma_0(g_B) = \sigma_{DM}(g'_B, g_{DM} = 1, m_{DM}): \text{Add DM contribution}$$

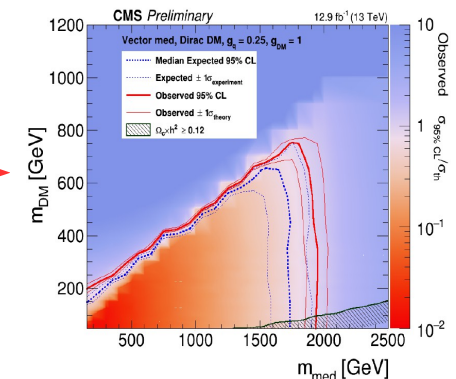
Solve for coupling
bound with DM

$$(g'_B)^2 = \frac{g_B^2}{2} \left(1 + \sqrt{1 + 4 \frac{\Gamma_{DM}}{\Gamma(g_B)}} \right)$$

Coupling
bound



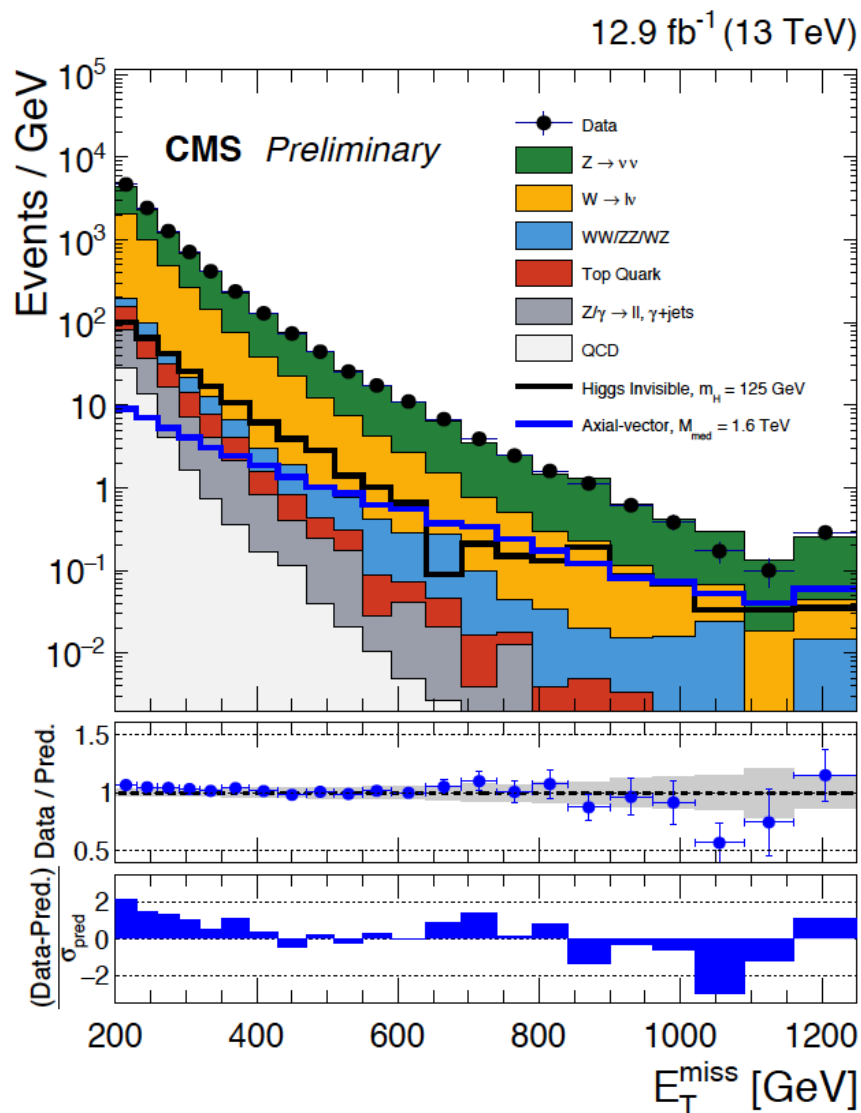
m_{med} (GeV)



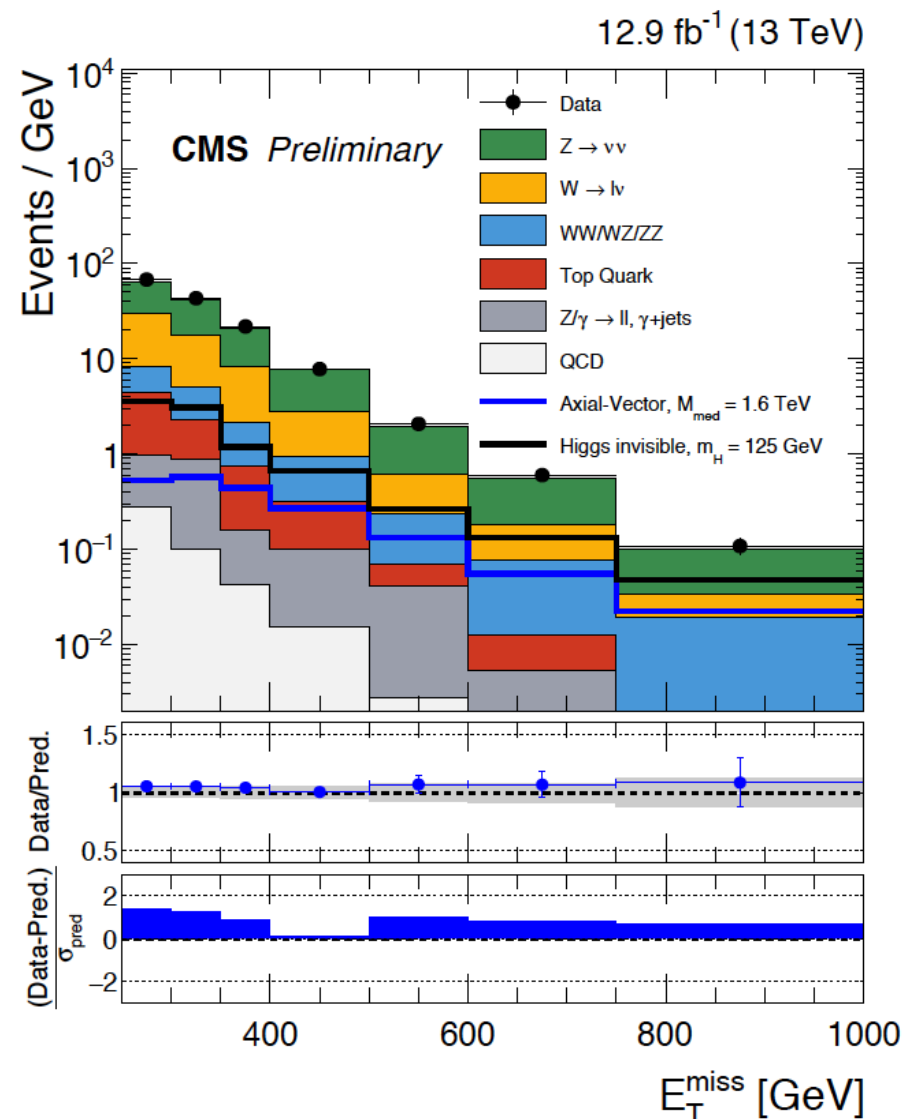
Total Luminosity : 12.9 fb⁻¹ (2016)
 Final State : J/V → qq + MET

Monojet/V

Monojet Category



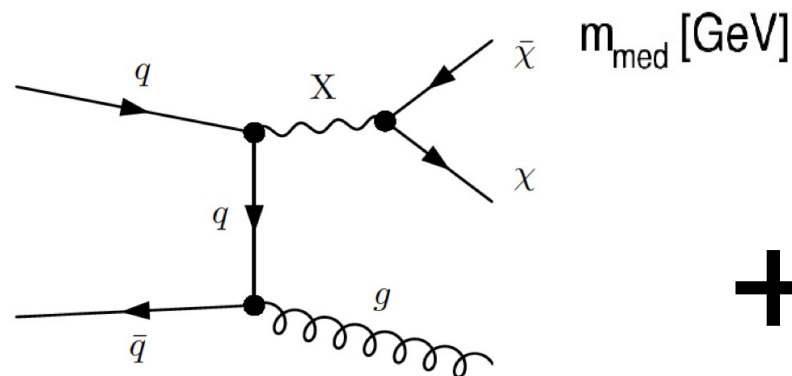
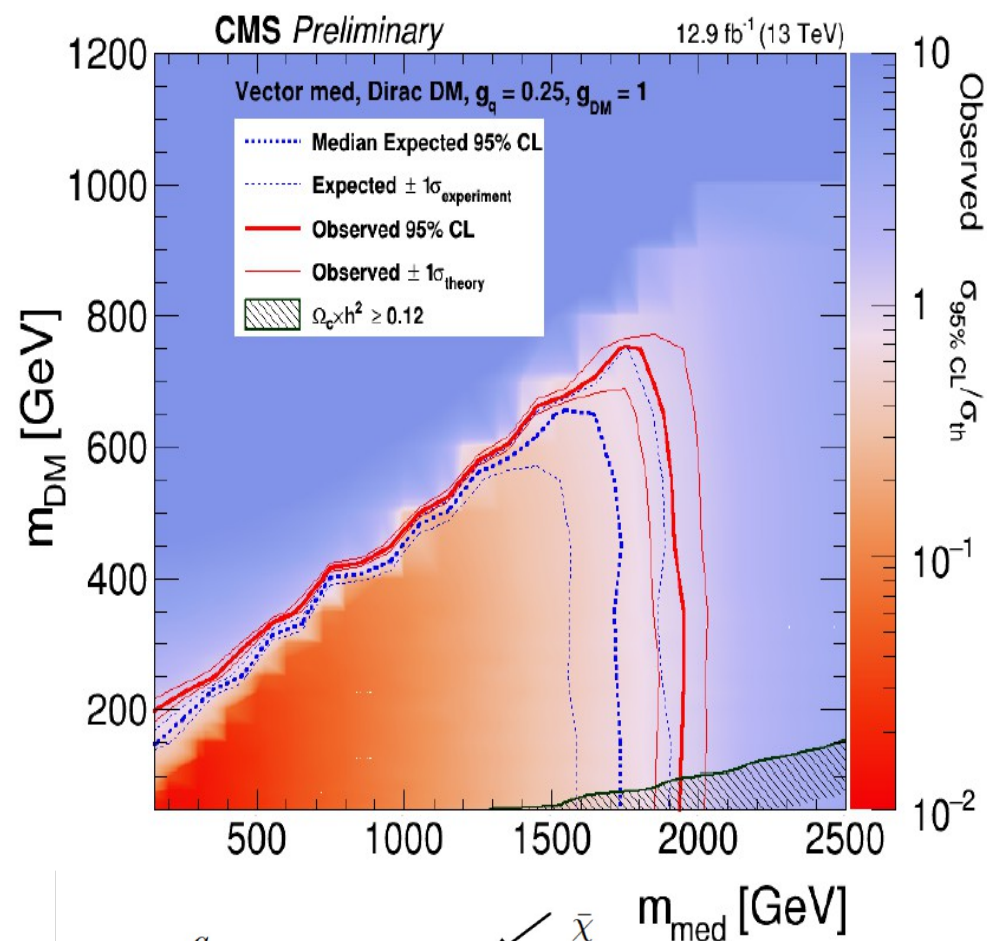
V-tagged Category



Total Luminosity : 12.9 fb⁻¹ (2016)

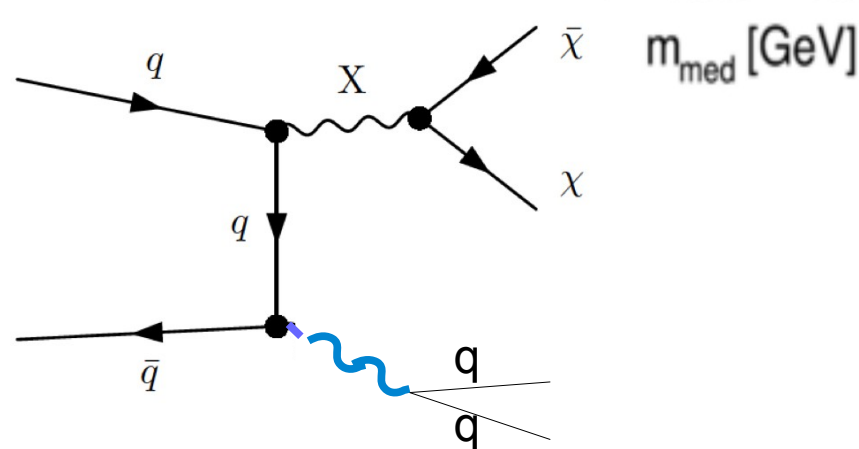
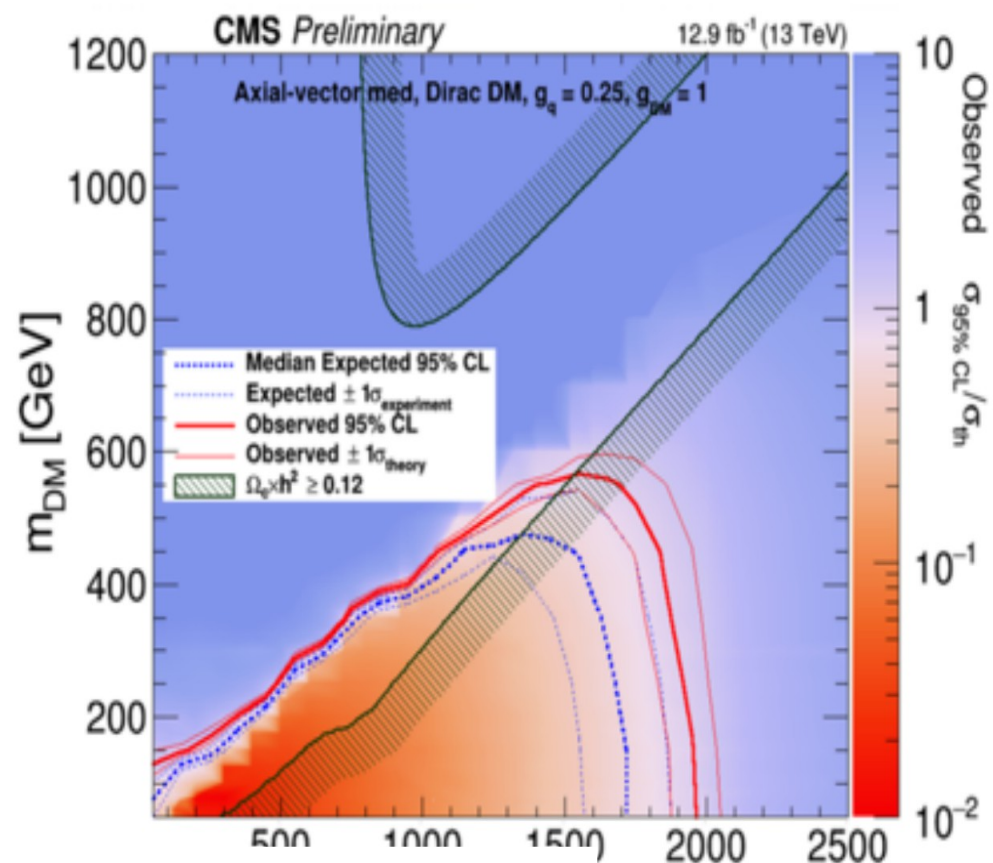
Final State : J/V → qq + MET

$u\gamma^\mu u$



+

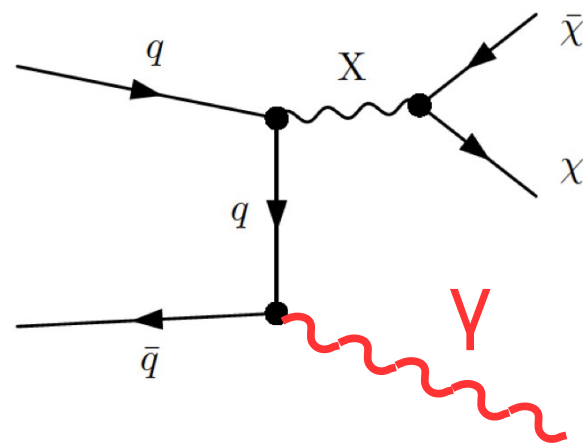
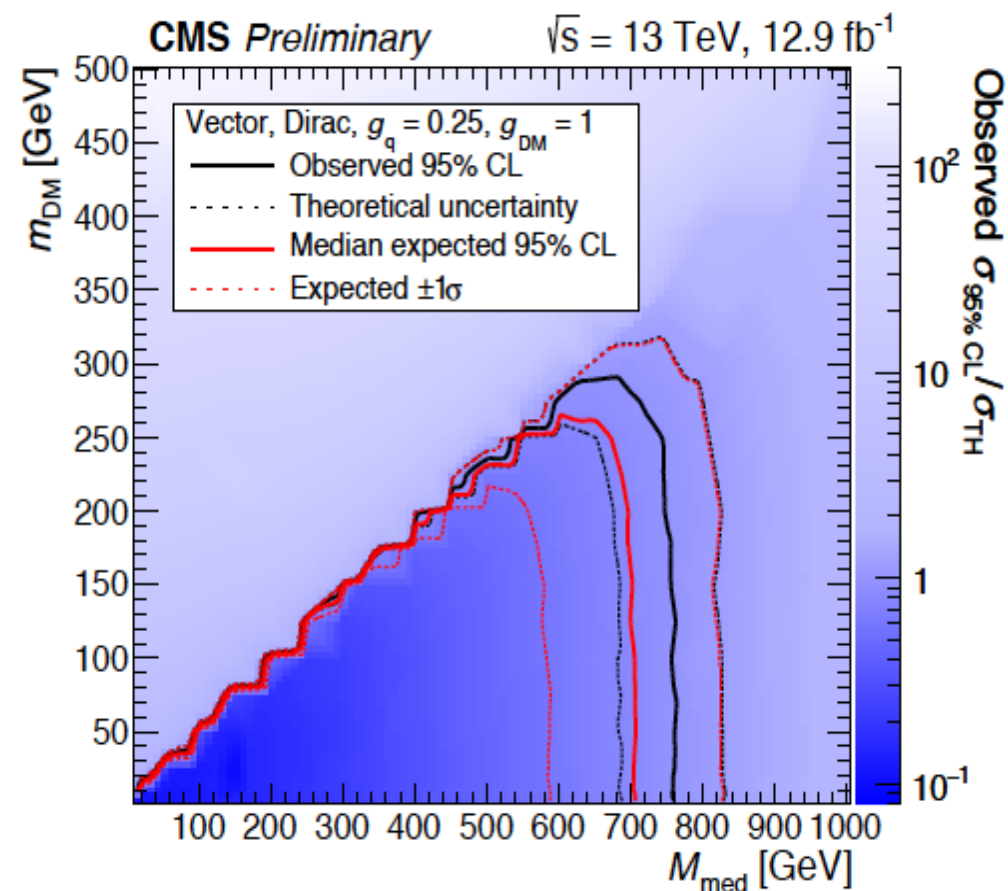
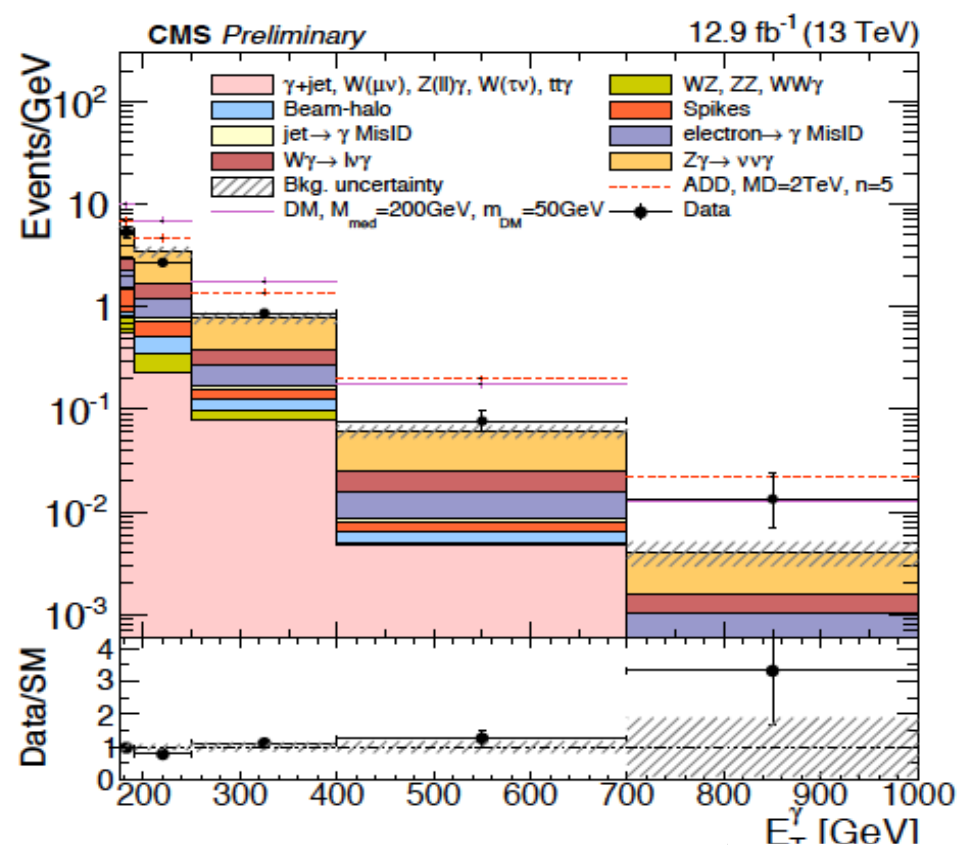
$u\gamma^5 u$ Monojet/V



Total Luminosity : 12.9 fb^{-1} (2016)

Final State : $\gamma + MET$

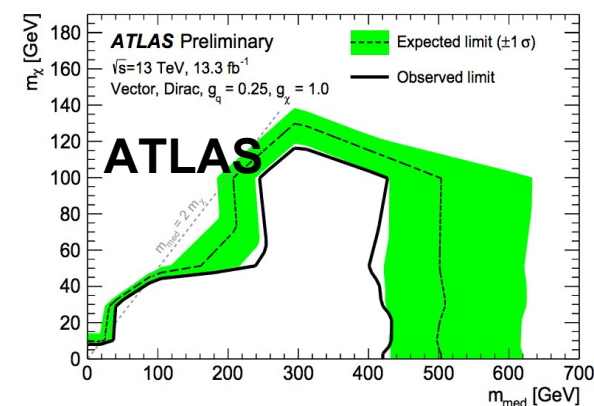
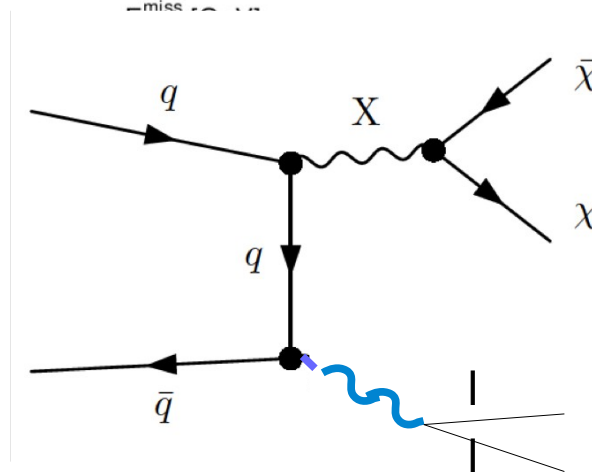
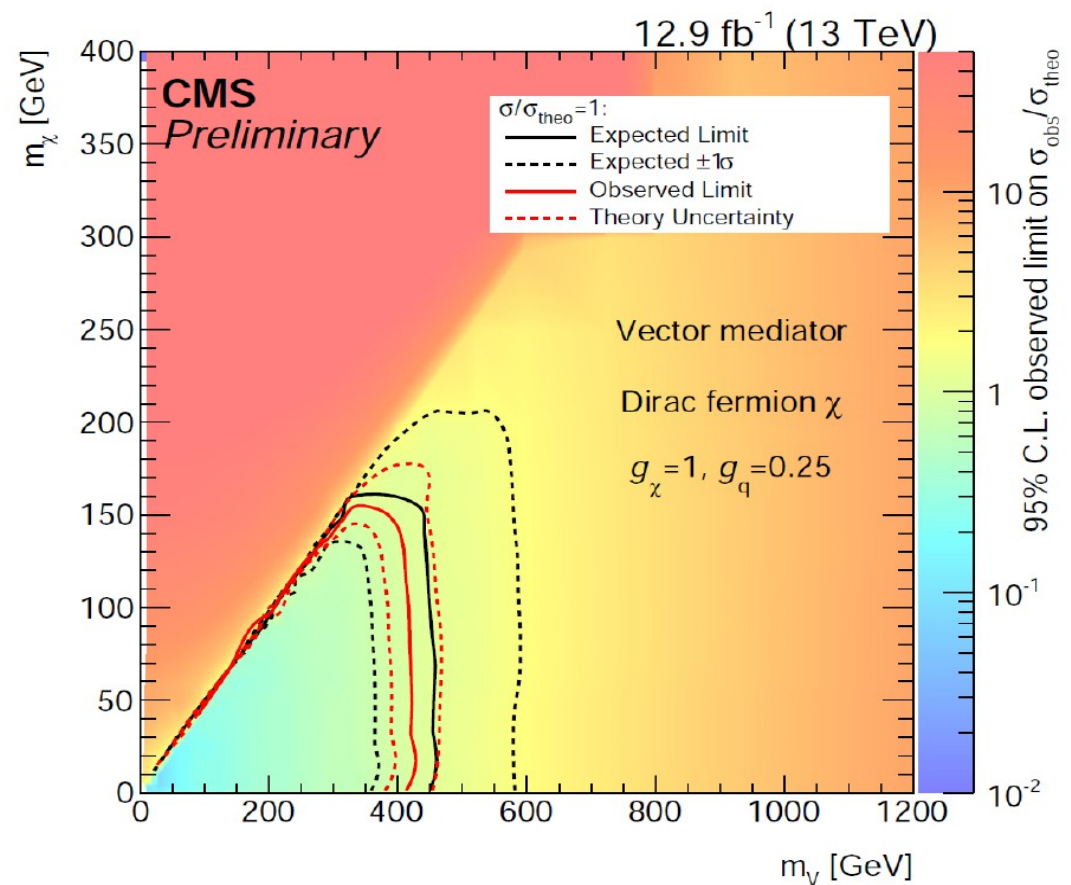
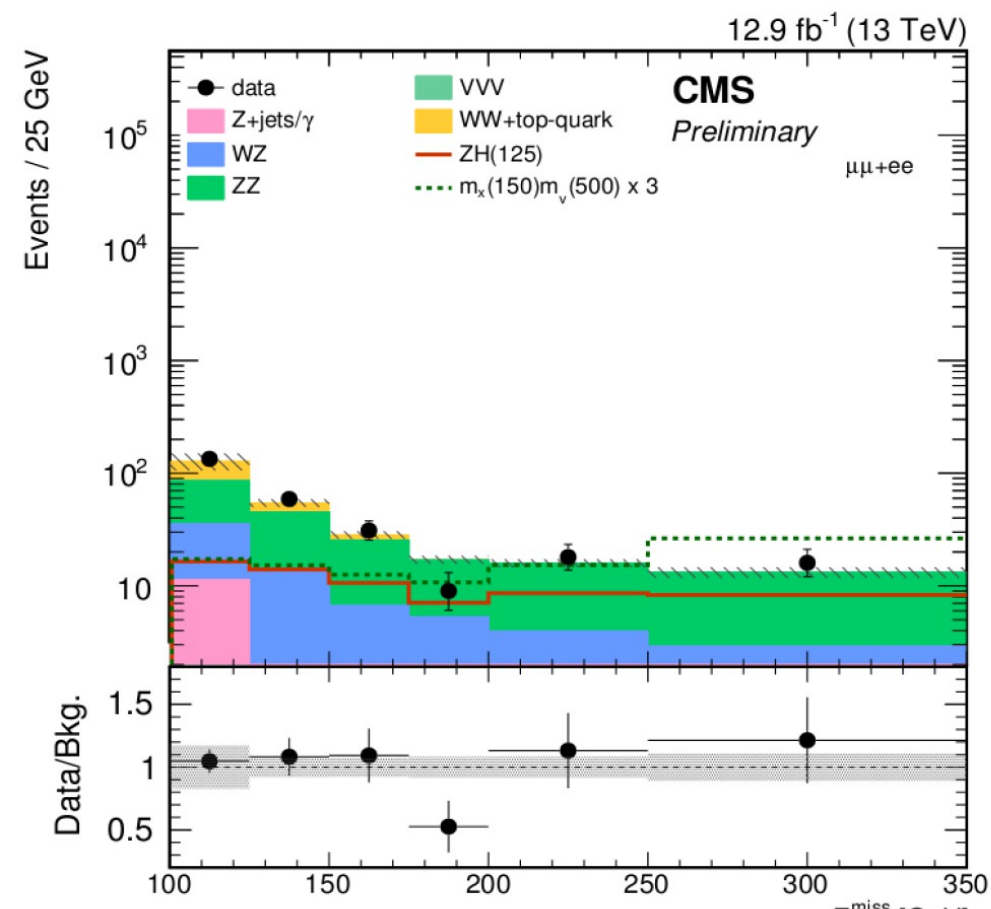
Mono- γ



Total Luminosity : 12.9 fb⁻¹ (2016)

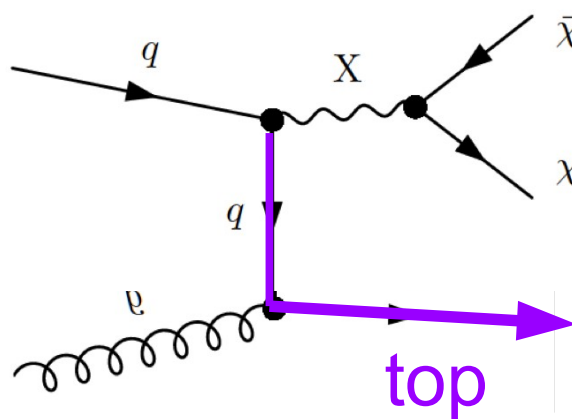
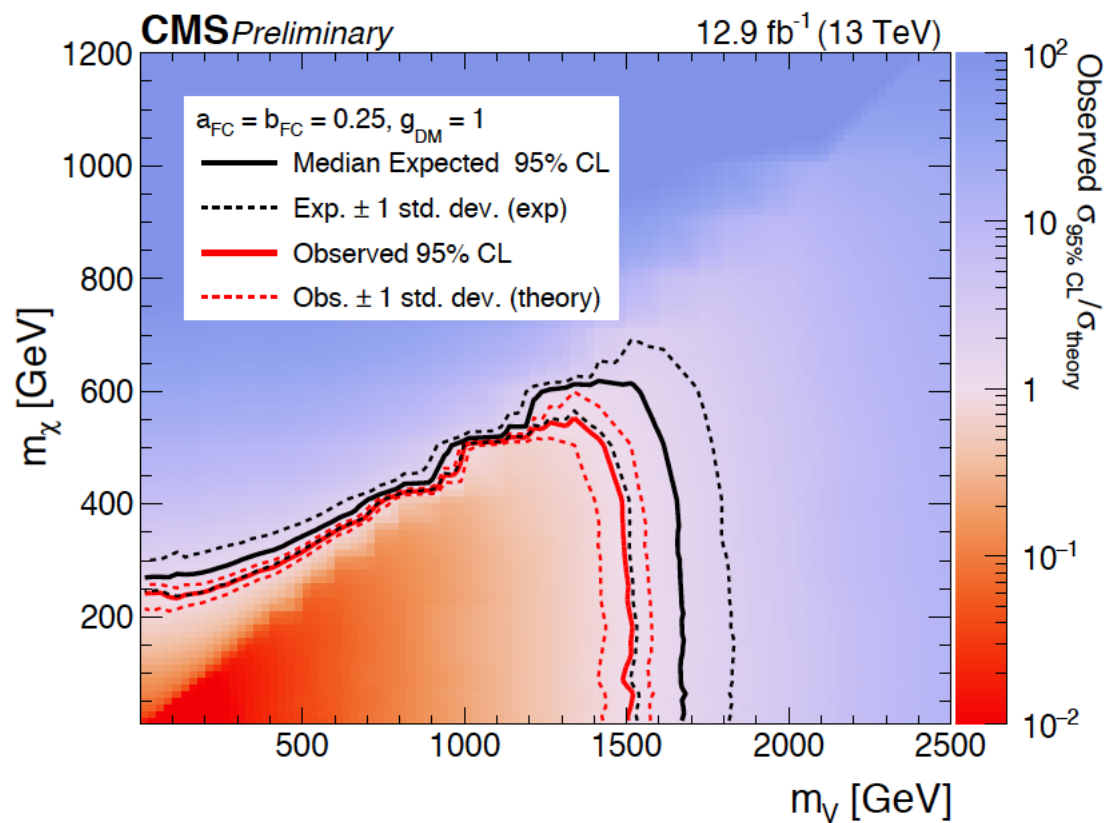
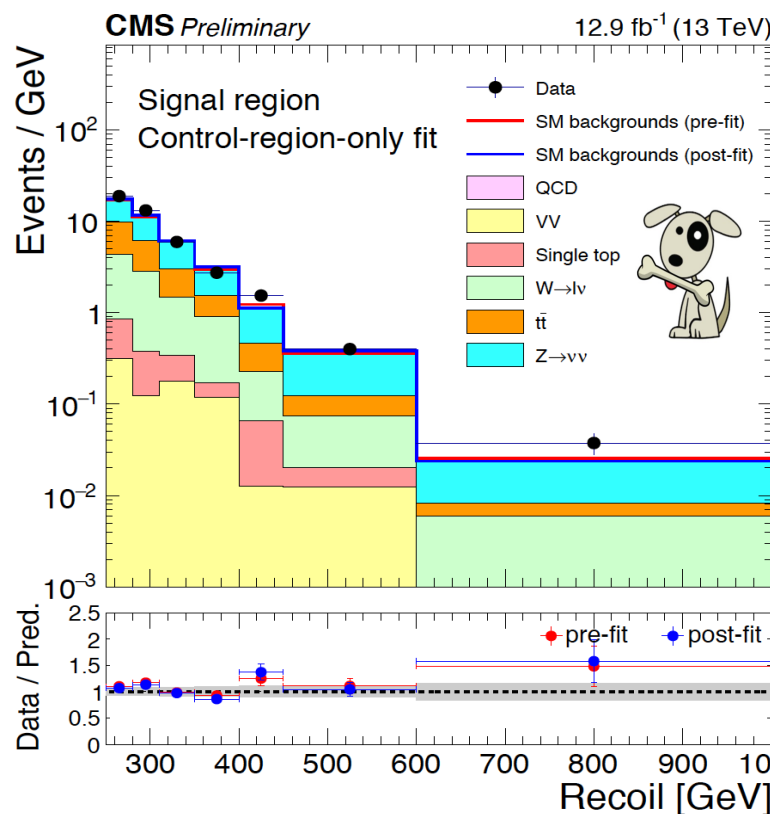
Final State : Z → ll + MET

Mono-Z



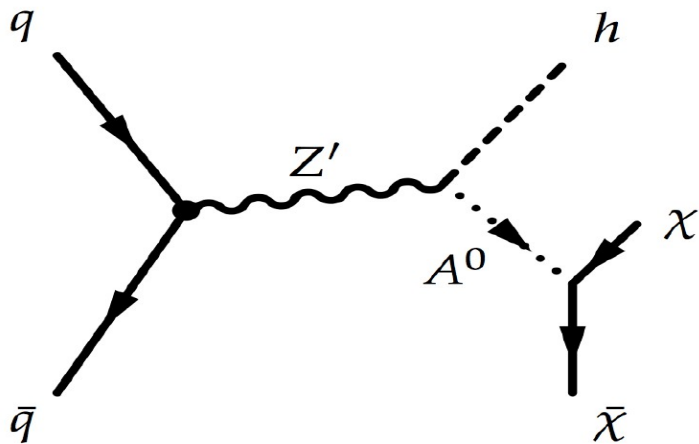
Total Luminosity : 12.9 fb⁻¹ (2016)
 Final State : top+MET

Mono-Top



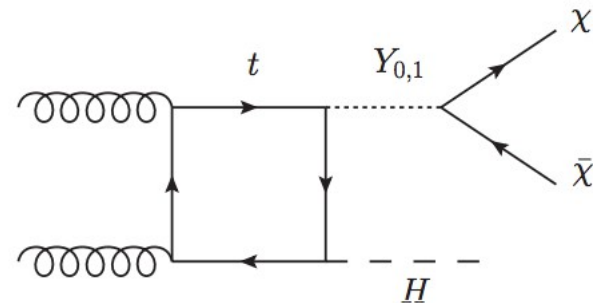
Going beyond one mediator

- Currently only being used for mono-Higgs



Searching for Z' and
and Pseudoscalar

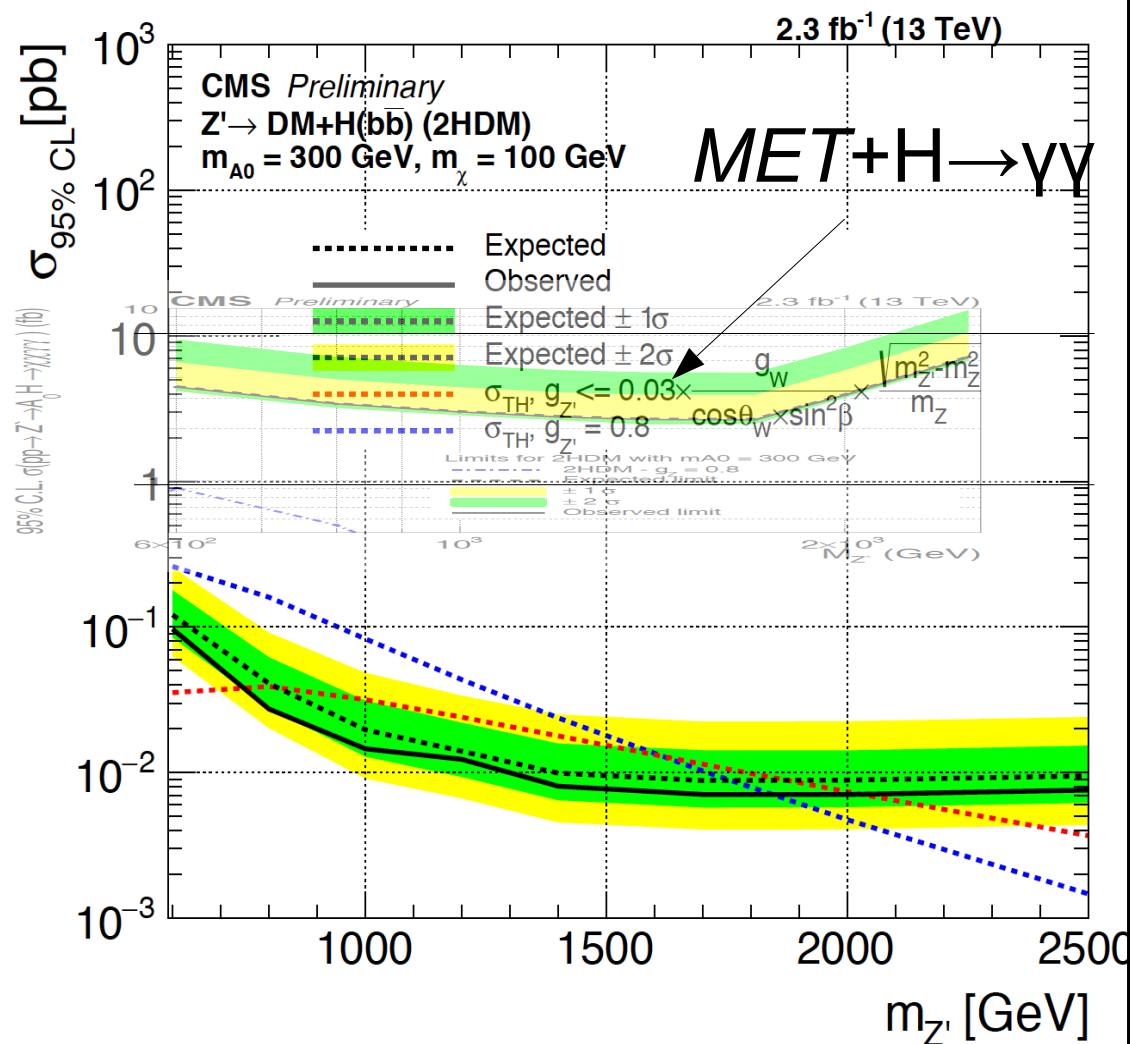
- Signature highlights the use of mono-H analysis
 - Mono-H can be added to the simplified models
 - However cross section is small



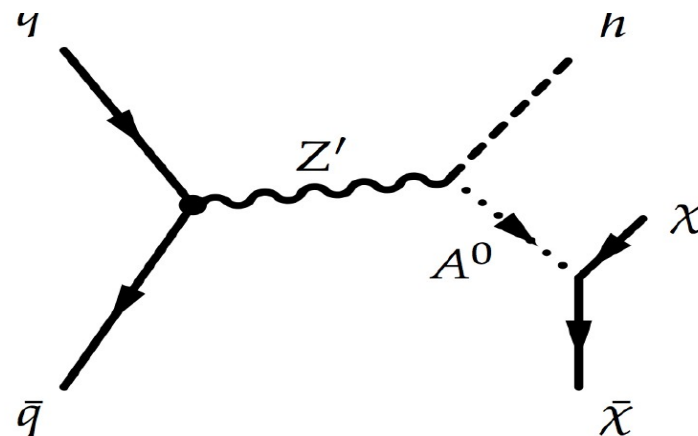
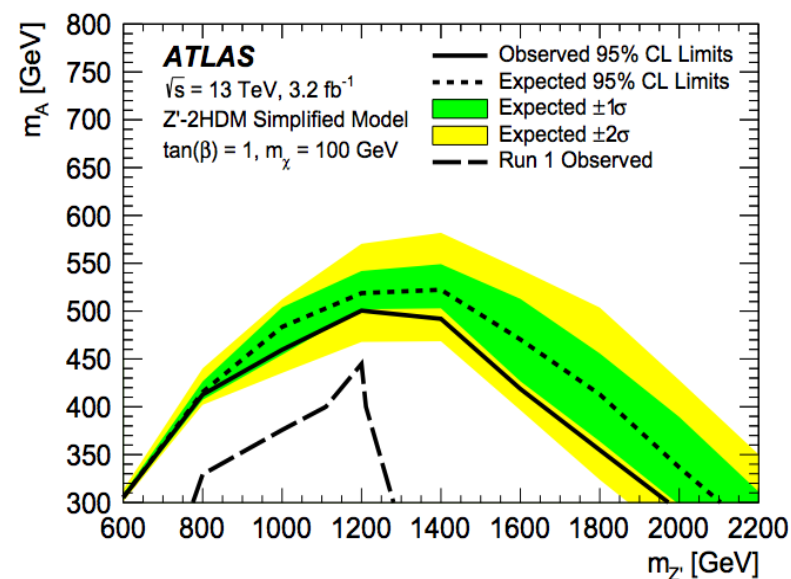
Total Luminosity : 2.3 fb^{-1} (2015)
 Final State : Higgs(bb/ $\gamma\gamma$)+MET

Mono-Higgs

$MET+H \rightarrow bb$



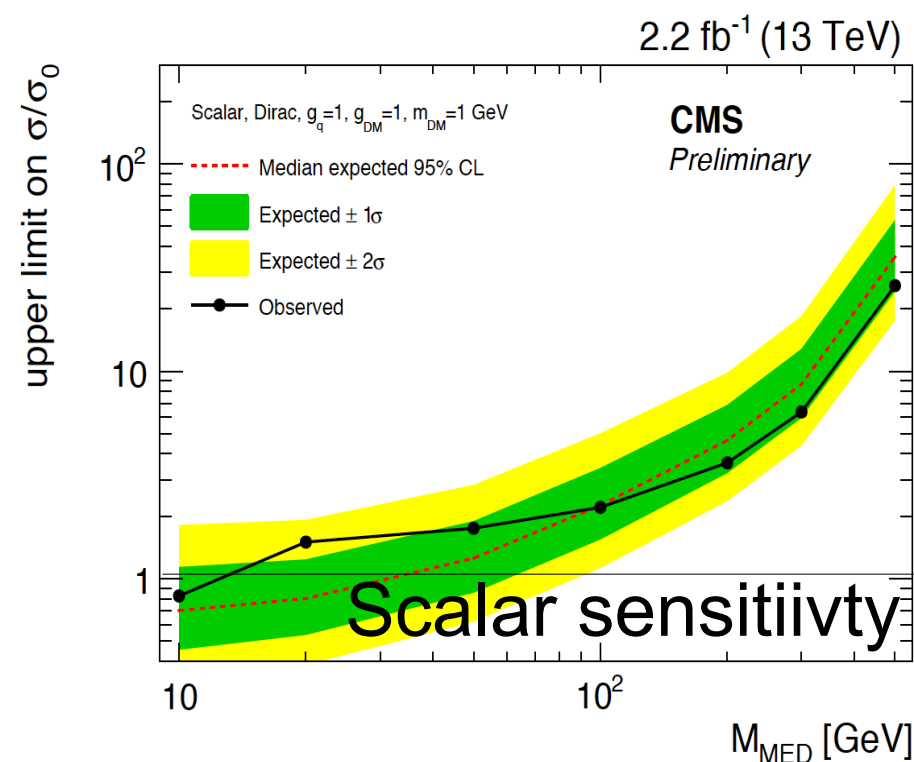
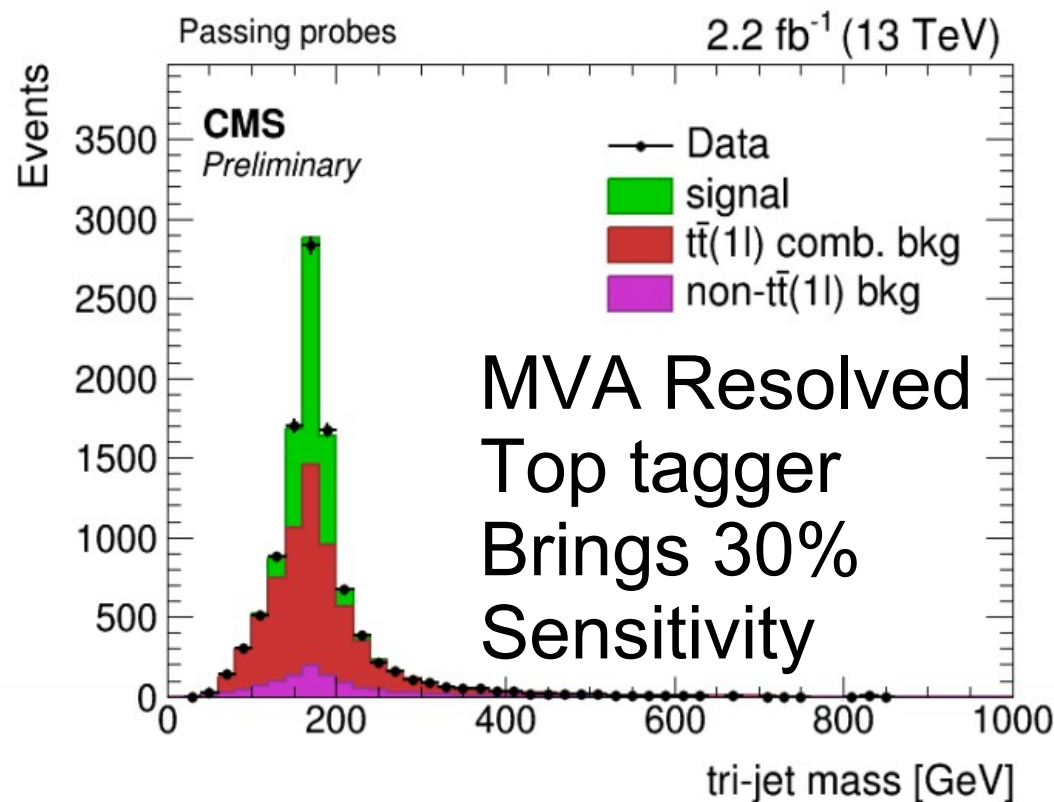
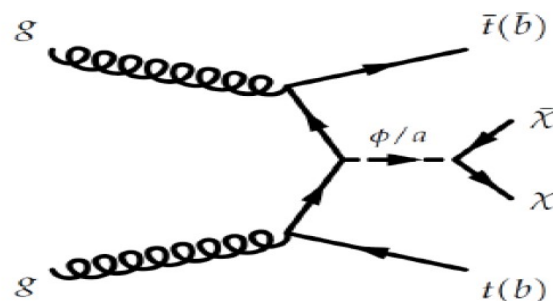
ATLAS results a little better



Total Luminosity : 2.3 fb^{-1} (2015)

Final State : $t\bar{t} + MET$

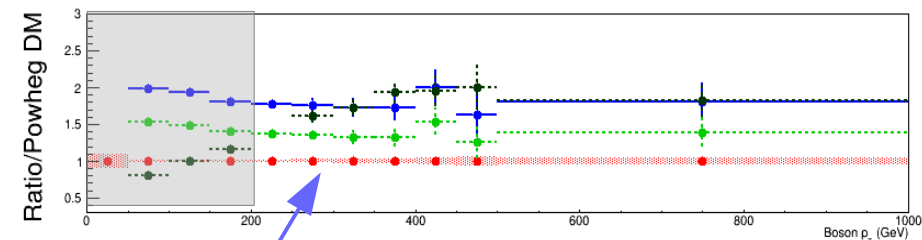
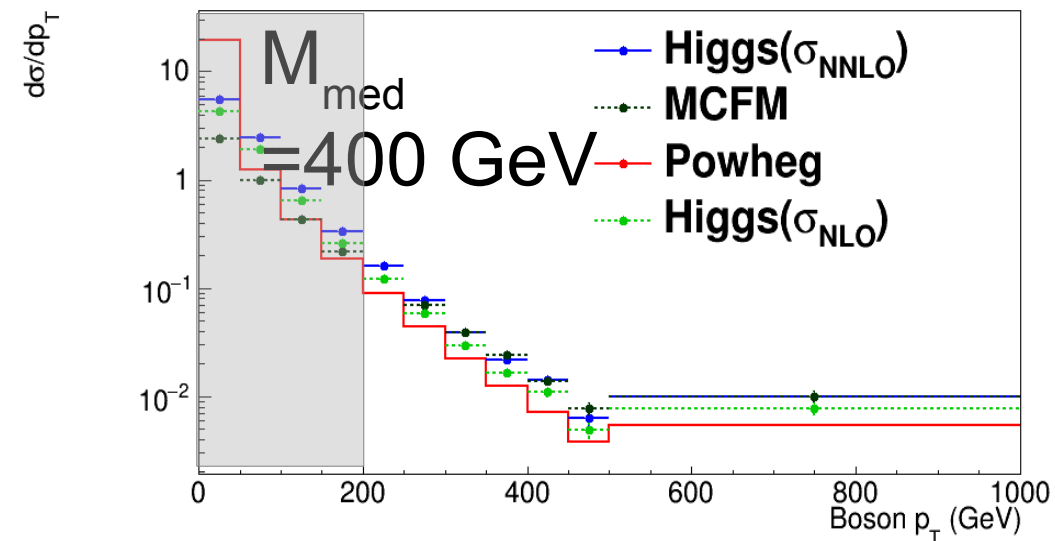
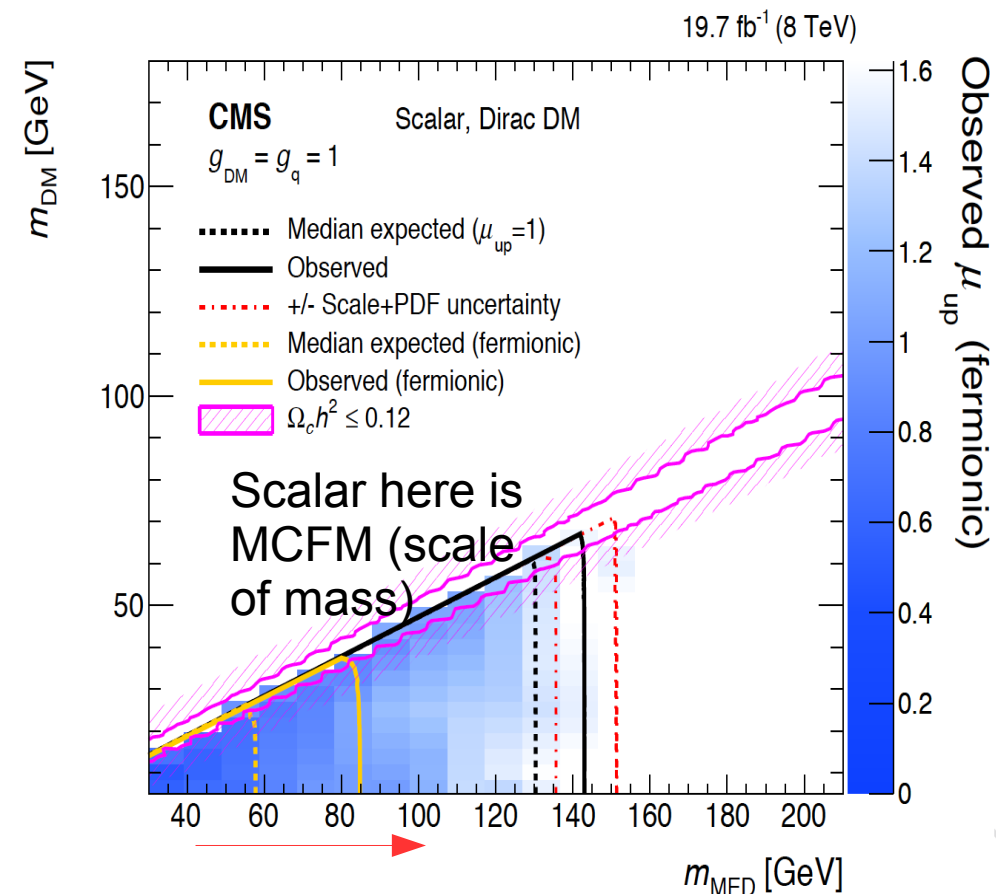
$t\bar{t} + \text{DM}$



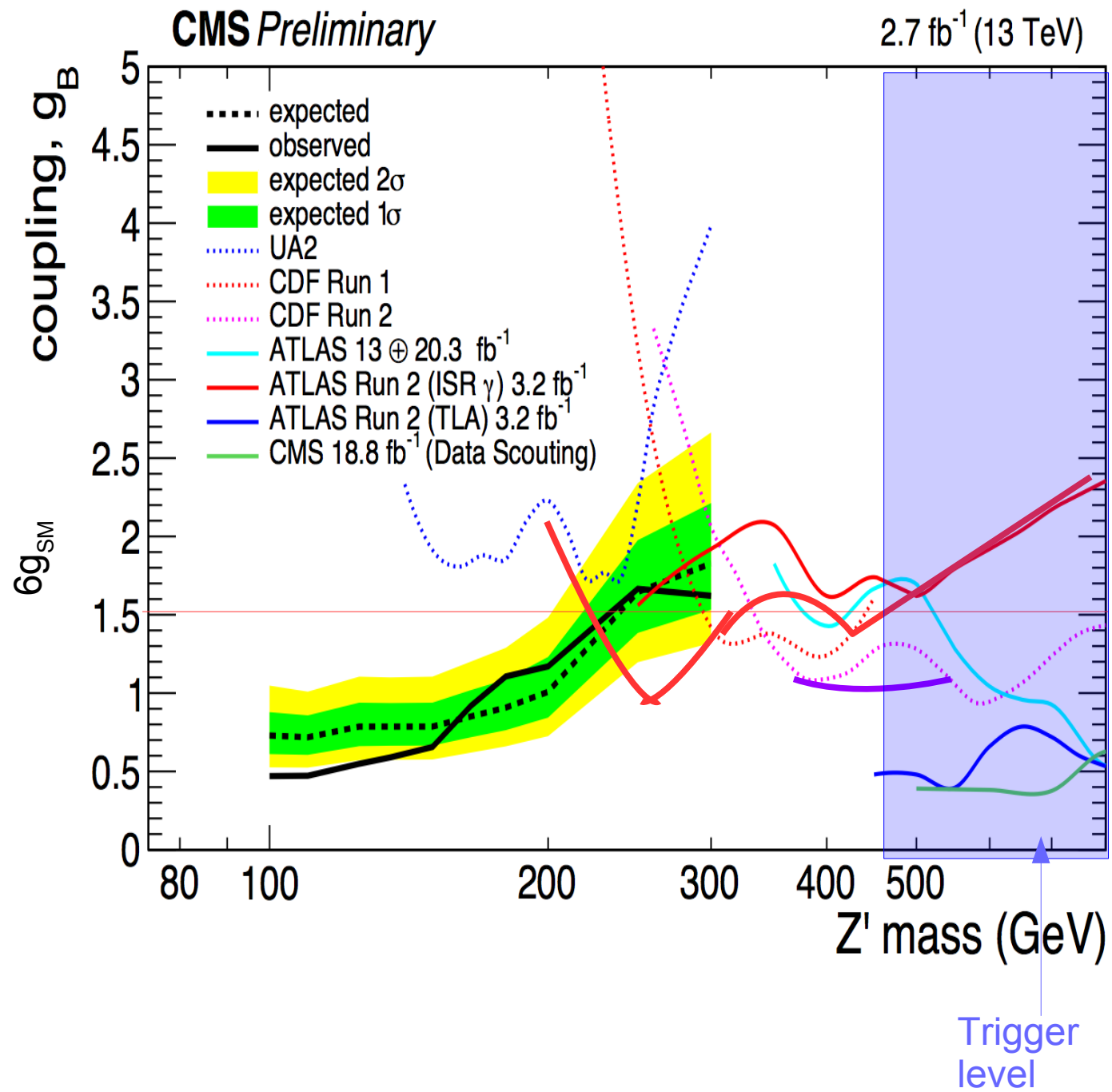
Sensitivity pushed to exclude Scalar mediator models

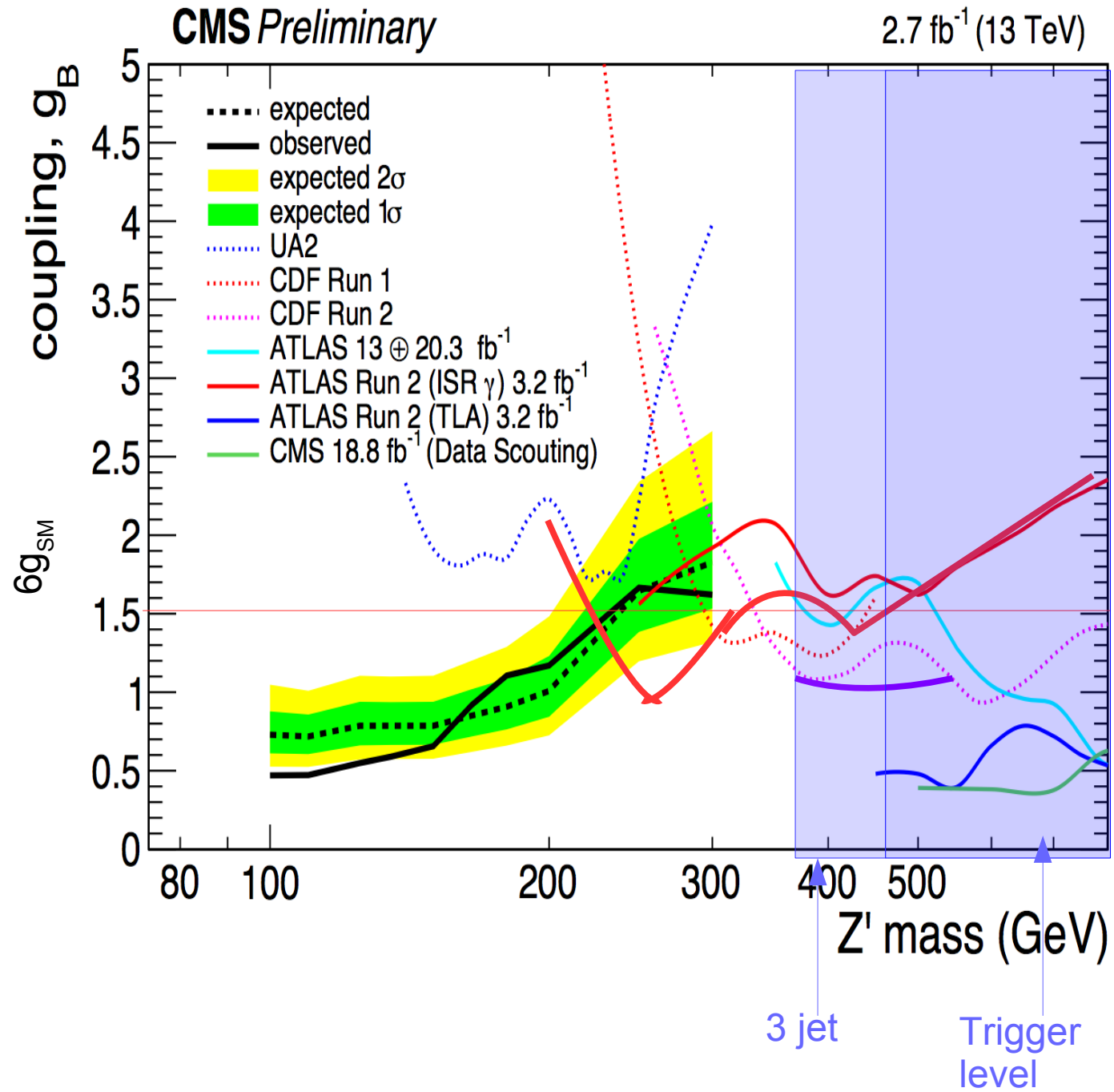
Scalar : Comparing with Old results

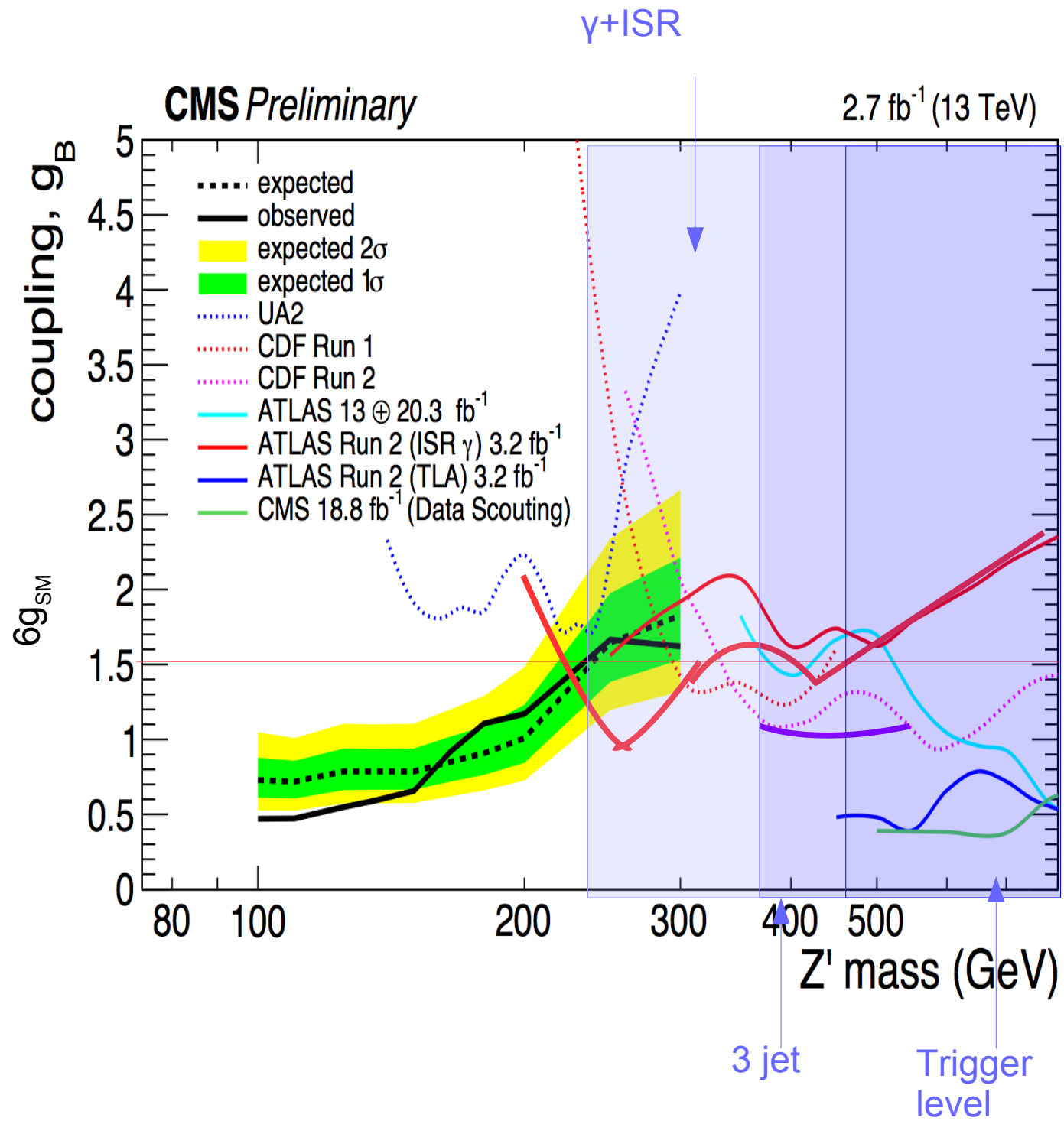
- Exclusion up to 80 GeV but not 13 TeV
 - Two reasons : 13 TeV is systematics limited (for now)

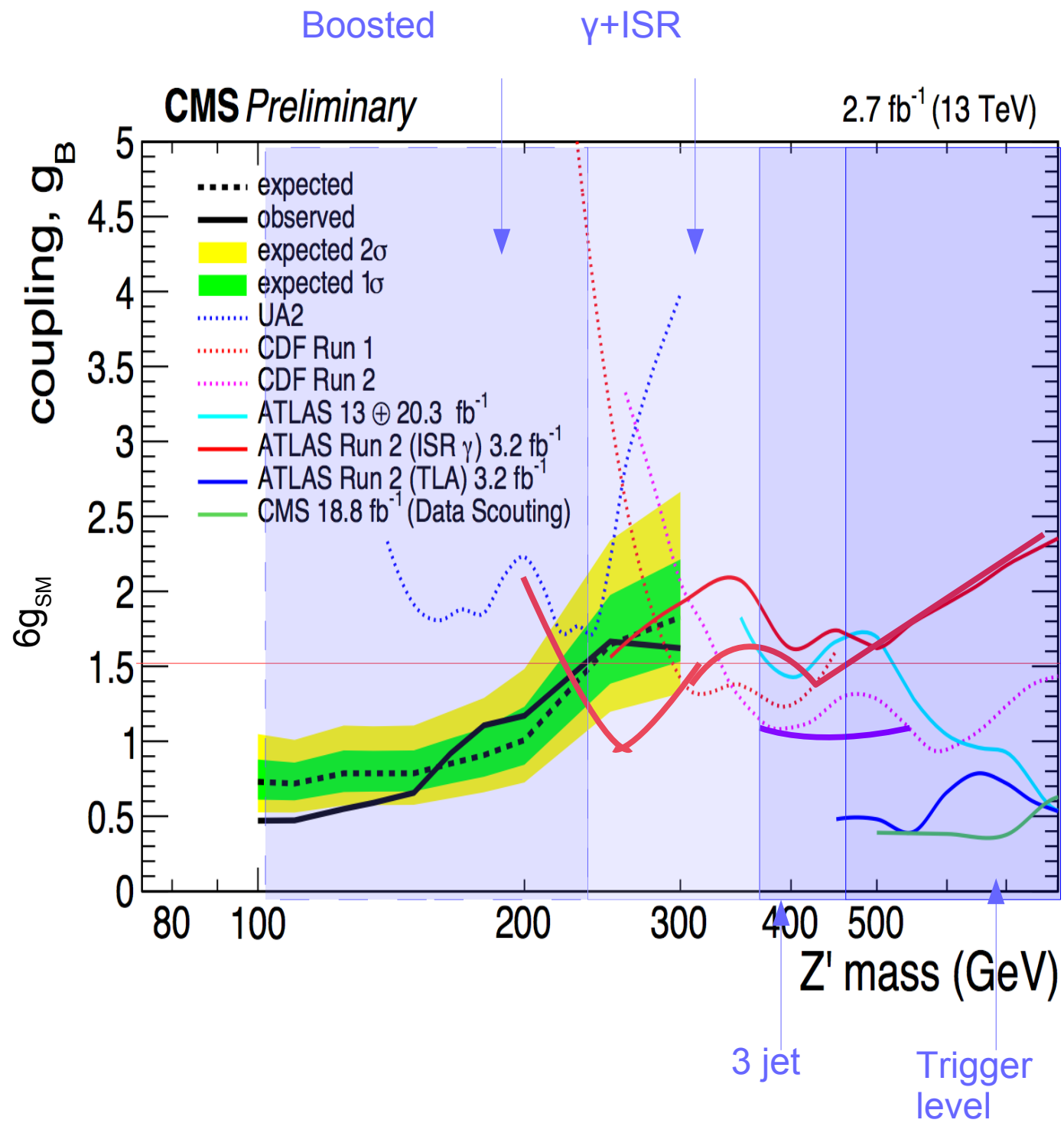


8 TeV is using MCFM (larger scale matches Higgs)







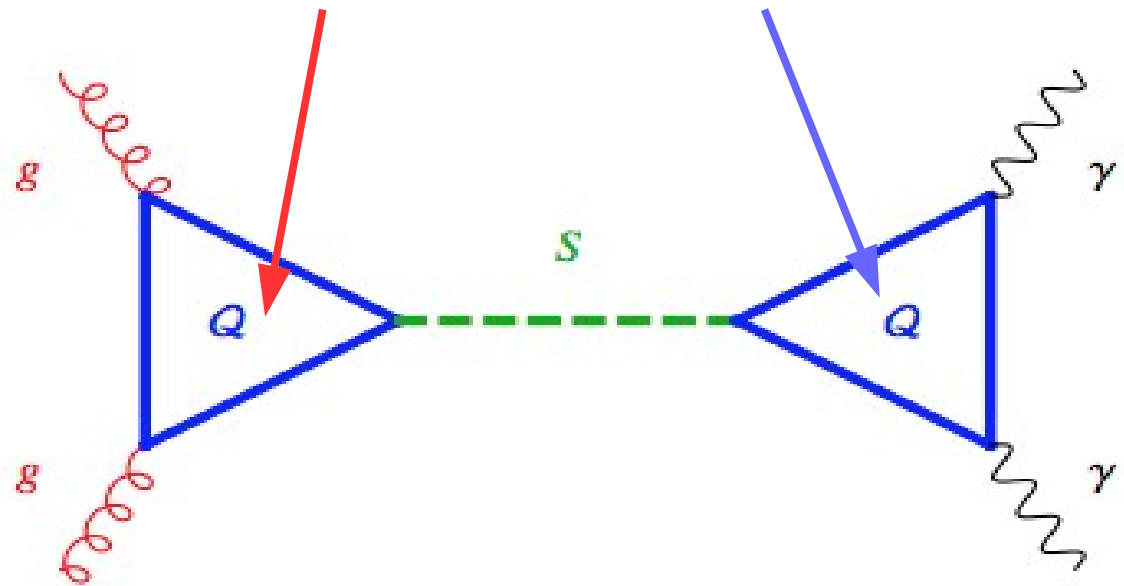


The simplest DM model

- Lets try to make something super basic
 - Basic model

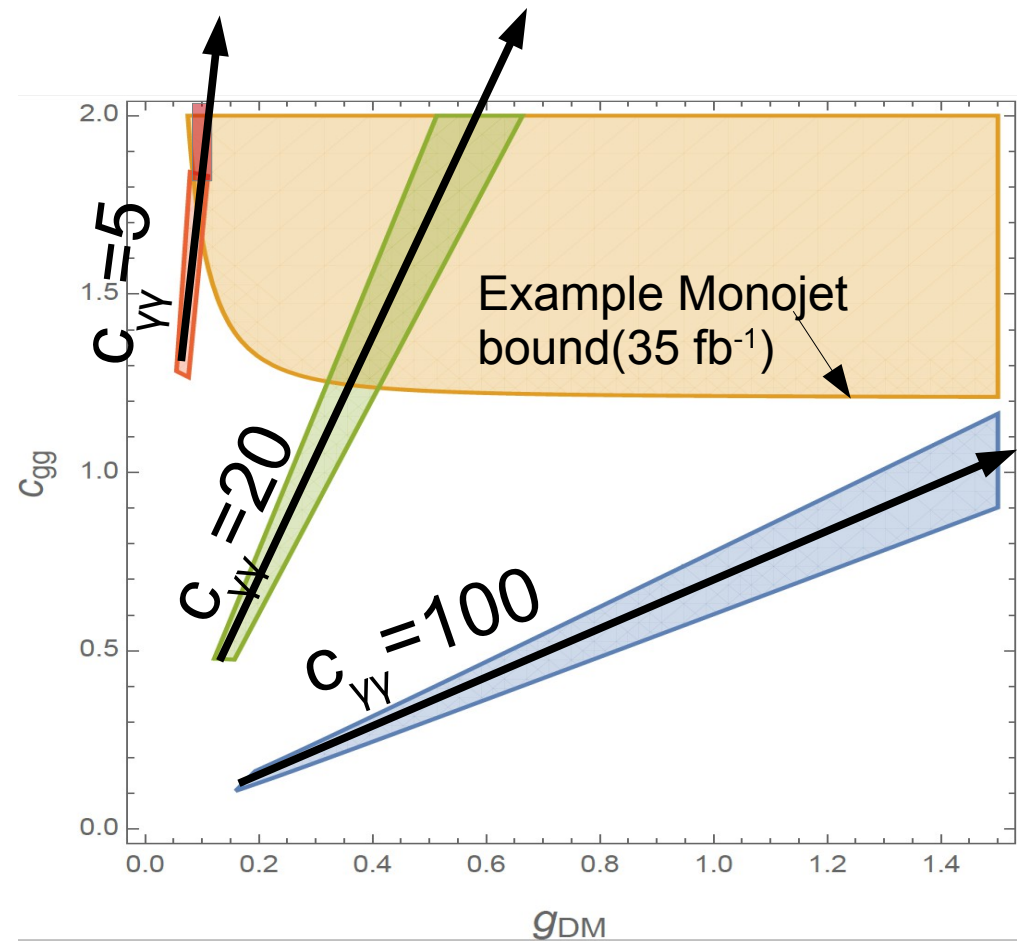
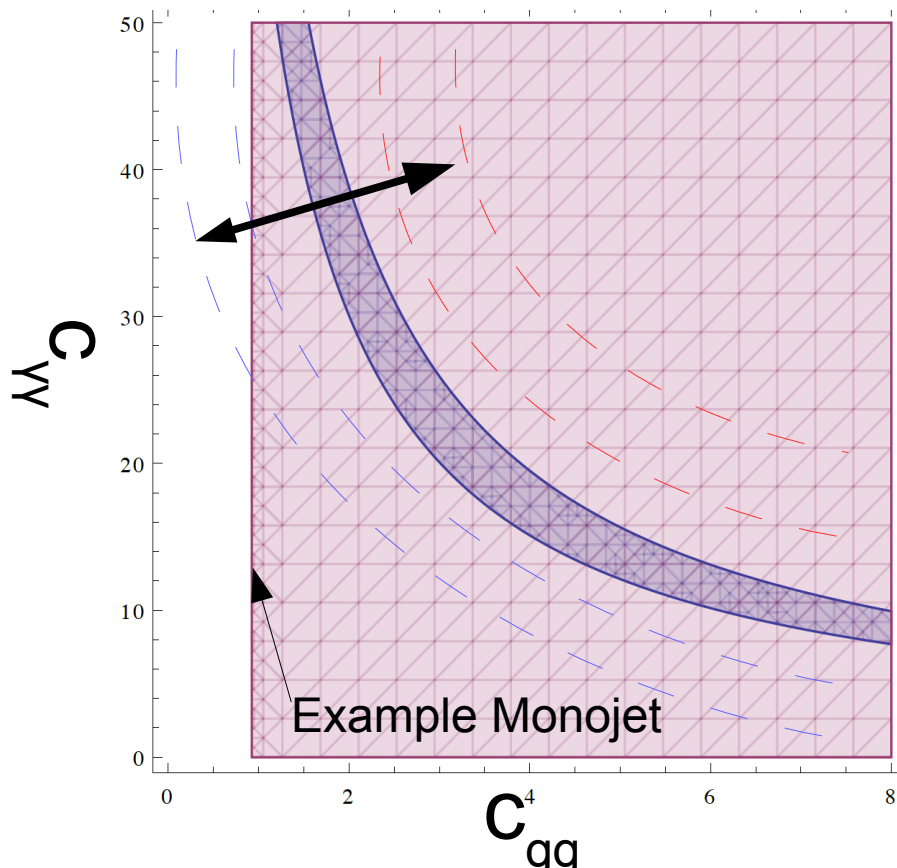
$$\mathcal{L}_S = \underbrace{g_\chi S \bar{\chi} \chi}_{\text{black}} + \underbrace{\frac{\alpha_s}{4\pi} \frac{c_g}{\Lambda} S G_{\mu\nu}^a G^{a\mu\nu}}_{\text{red}} + \underbrace{\frac{\alpha}{4\pi} \frac{c_\gamma}{\Lambda} S F_{\mu\nu} F^{\mu\nu}}_{\text{blue}}$$

??



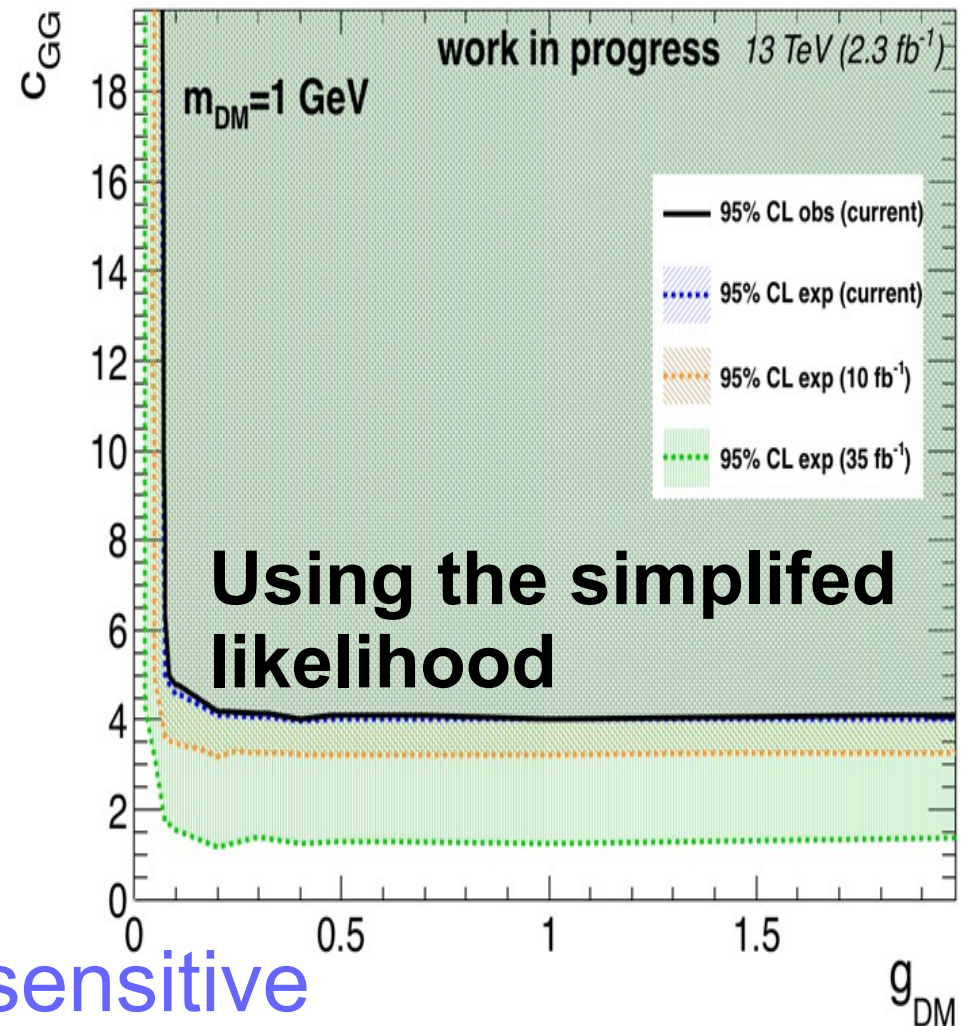
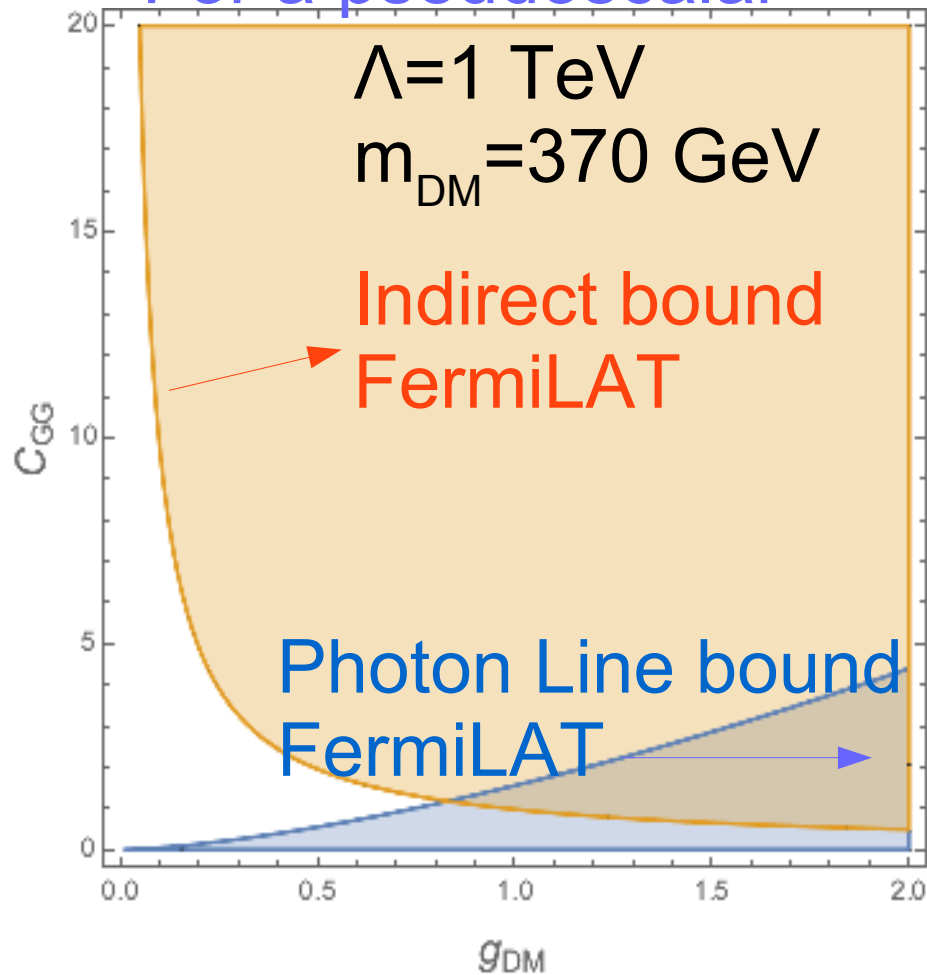
Using the best fit cross section

- We have 3 free couplings :
 - $g_{\text{DM}}, c_{\text{GG}}, c_{\text{YY}}$
- Taking the photon best fit can constrain one



Considering All Constraints

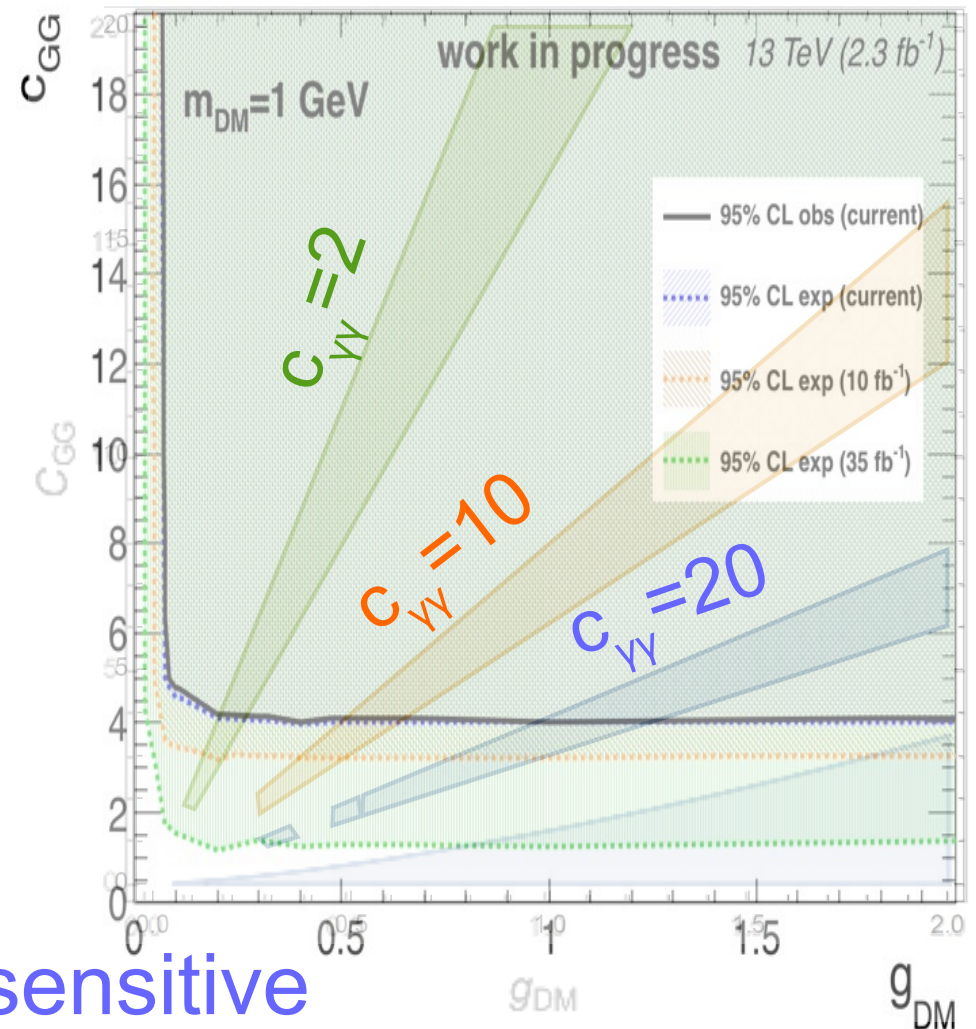
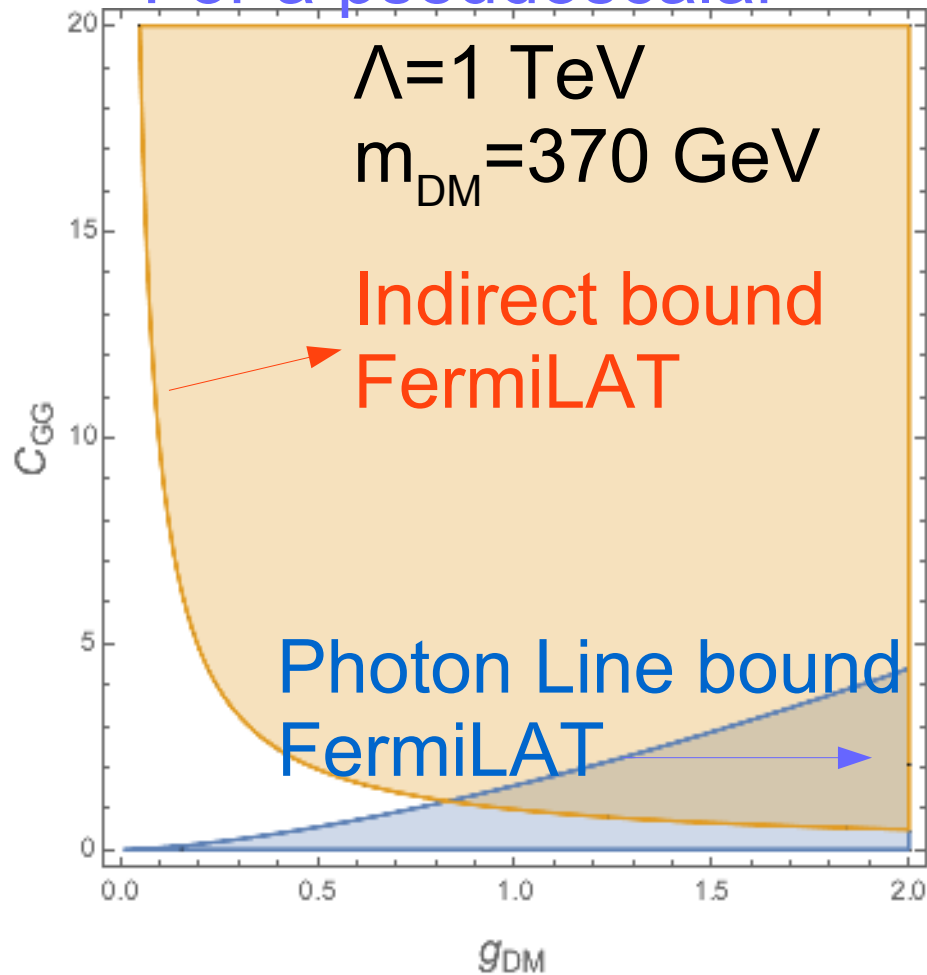
For a pseudoscalar



- Direct detection not yet sensitive
- Indirect detection limits on-shell production
- Photon Line bounds limit photon coupling < 100

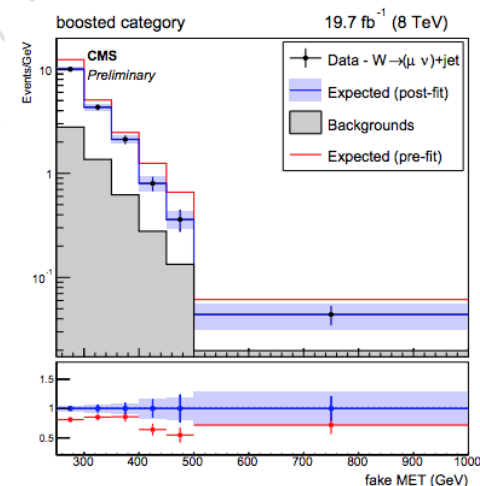
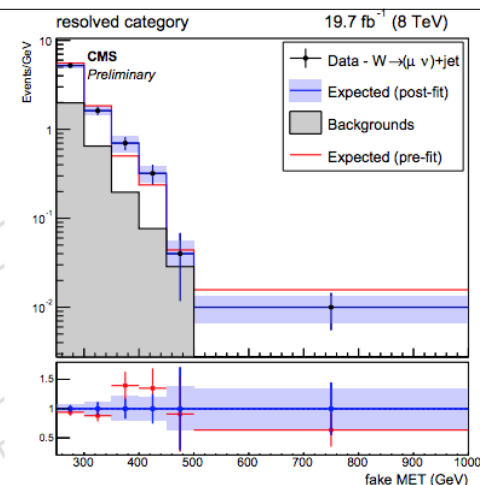
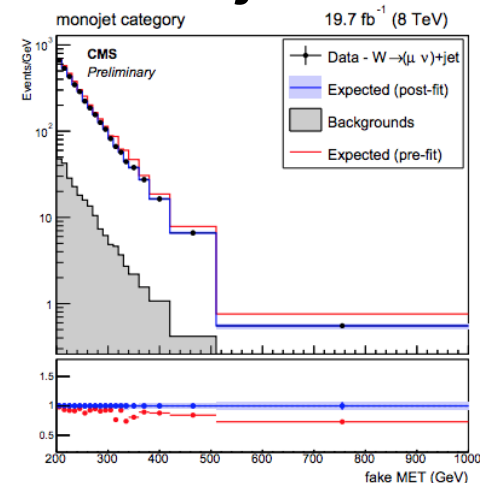
Considering All Constraints

For a pseudoscalar



- Direct detection not yet sensitive
- Indirect detection limits on-shell production
- Photon Line bounds limit photon coupling < 100

CMS-EXO-12-055

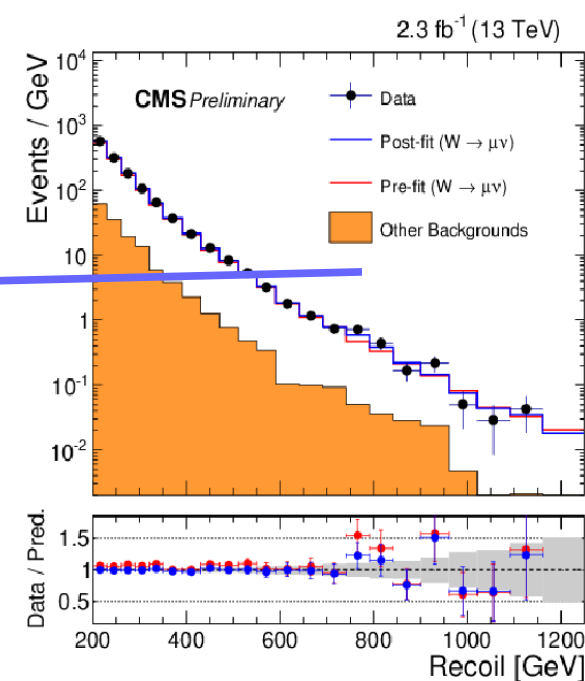
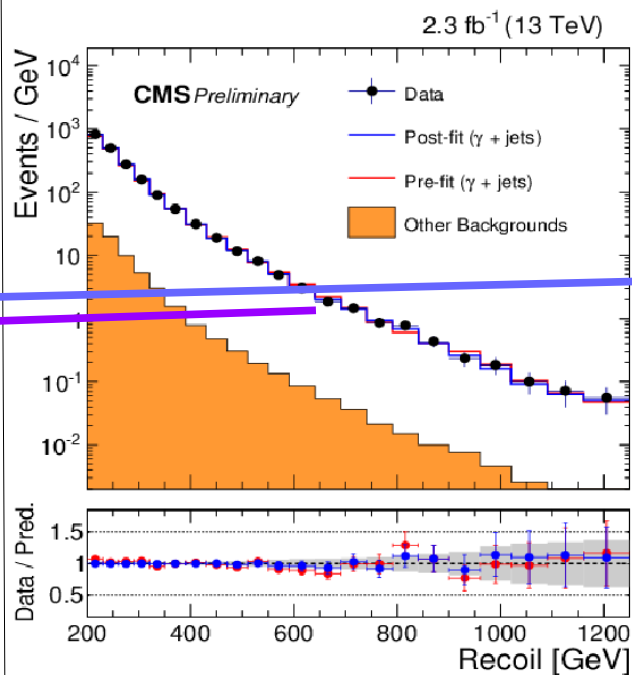
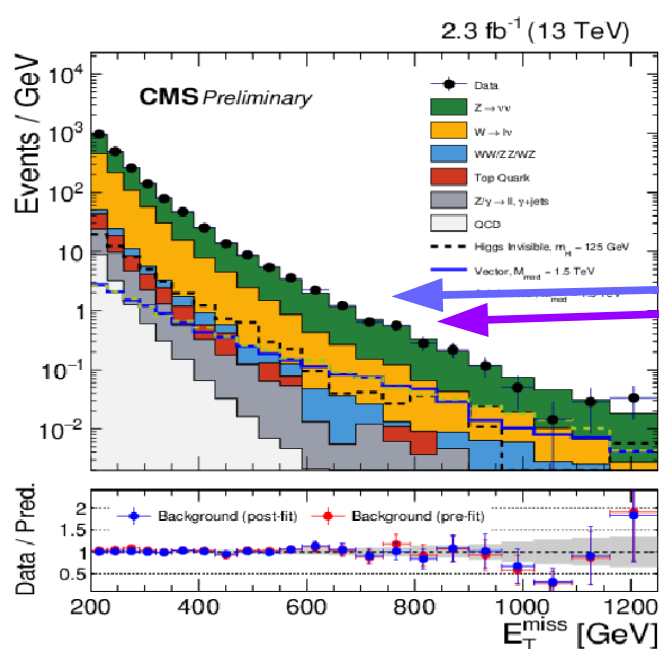


Signal

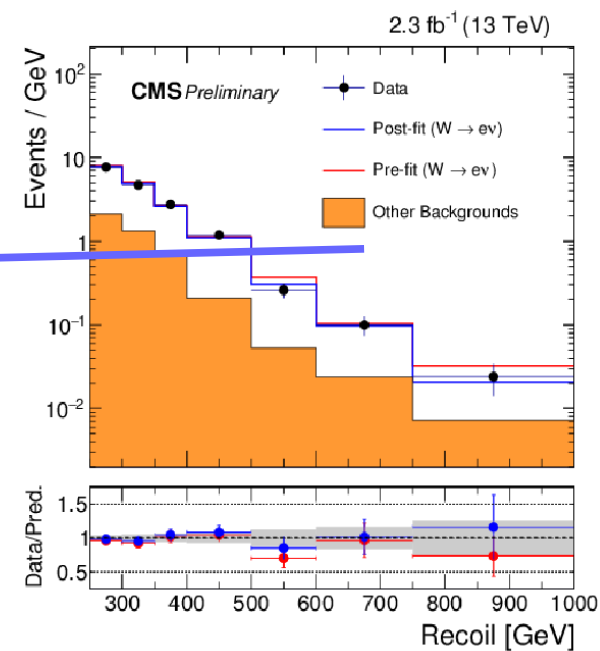
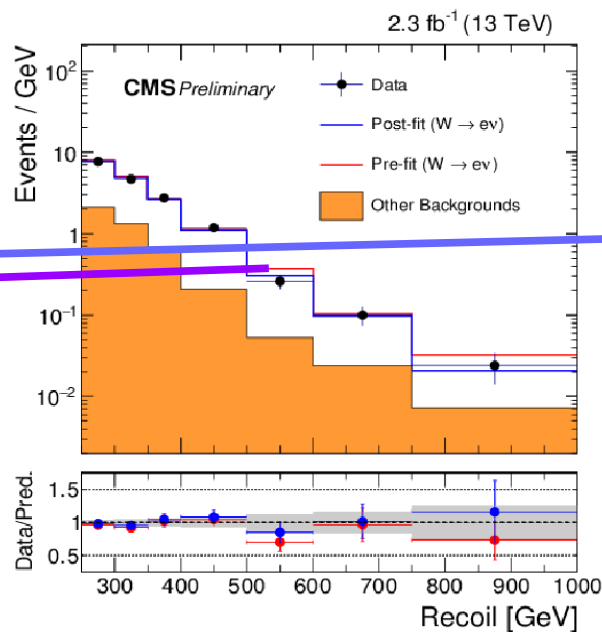
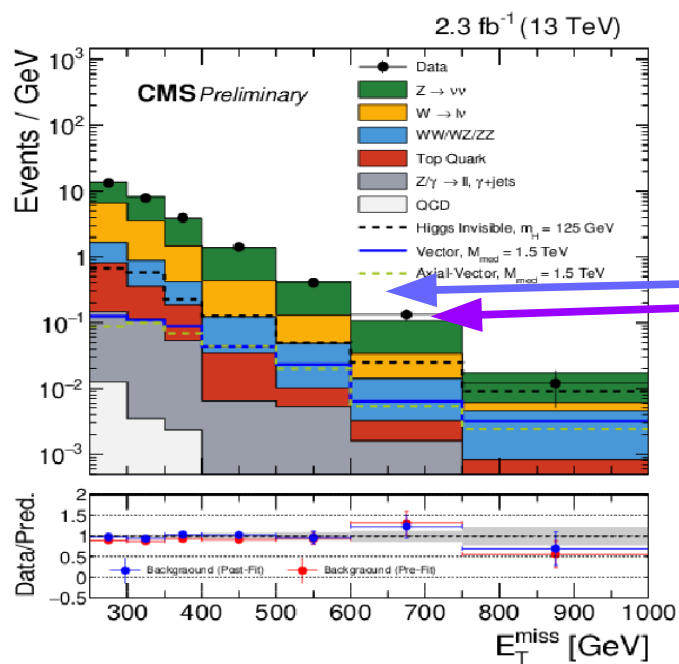
CRs: γ +jets

+ W +

Monojet

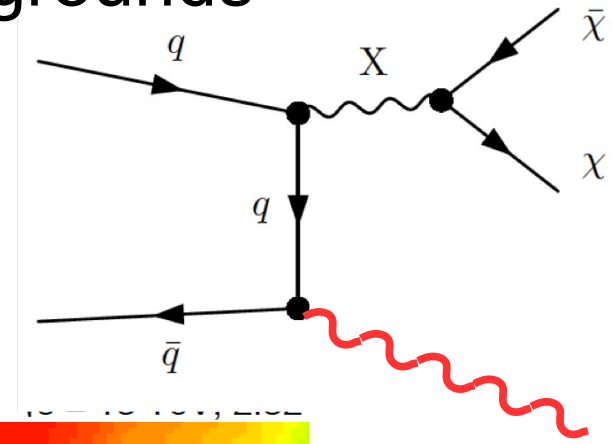
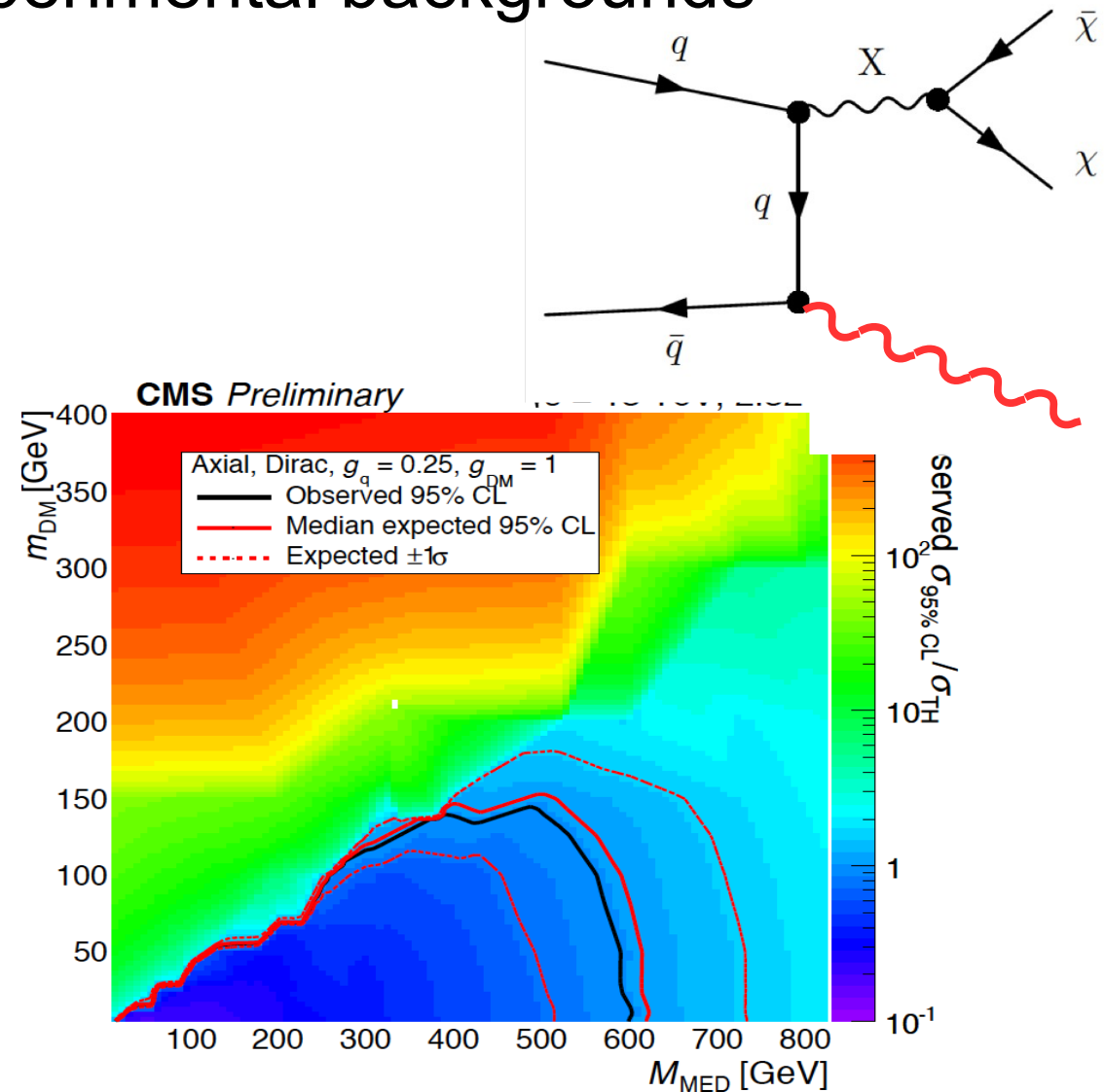
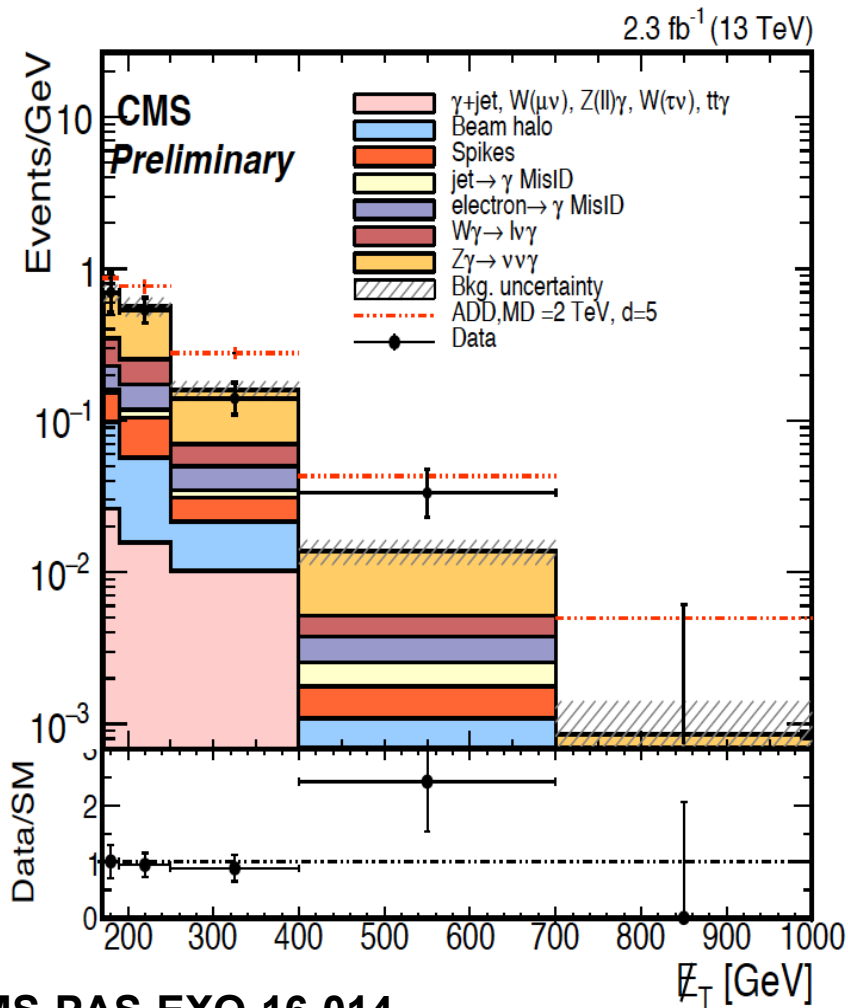


Mono-V



Monophoton

- Tag a photon and look for MET
 - Many challenging experimental backgrounds



Monotop

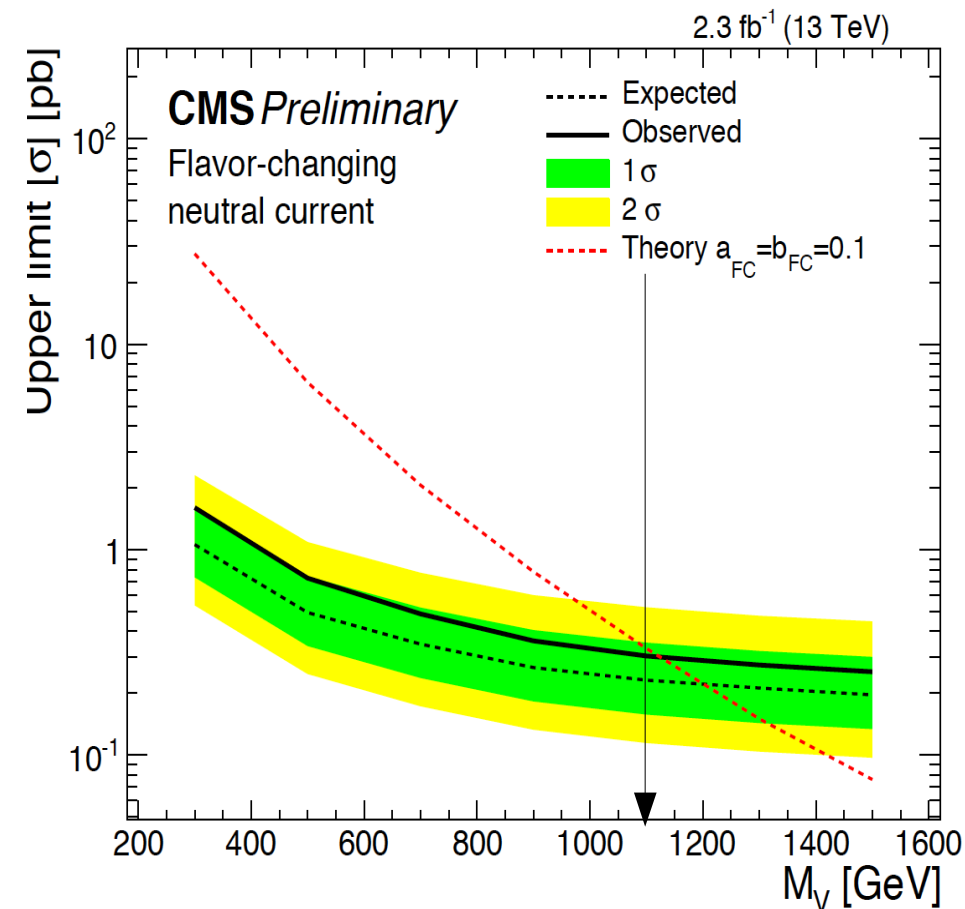
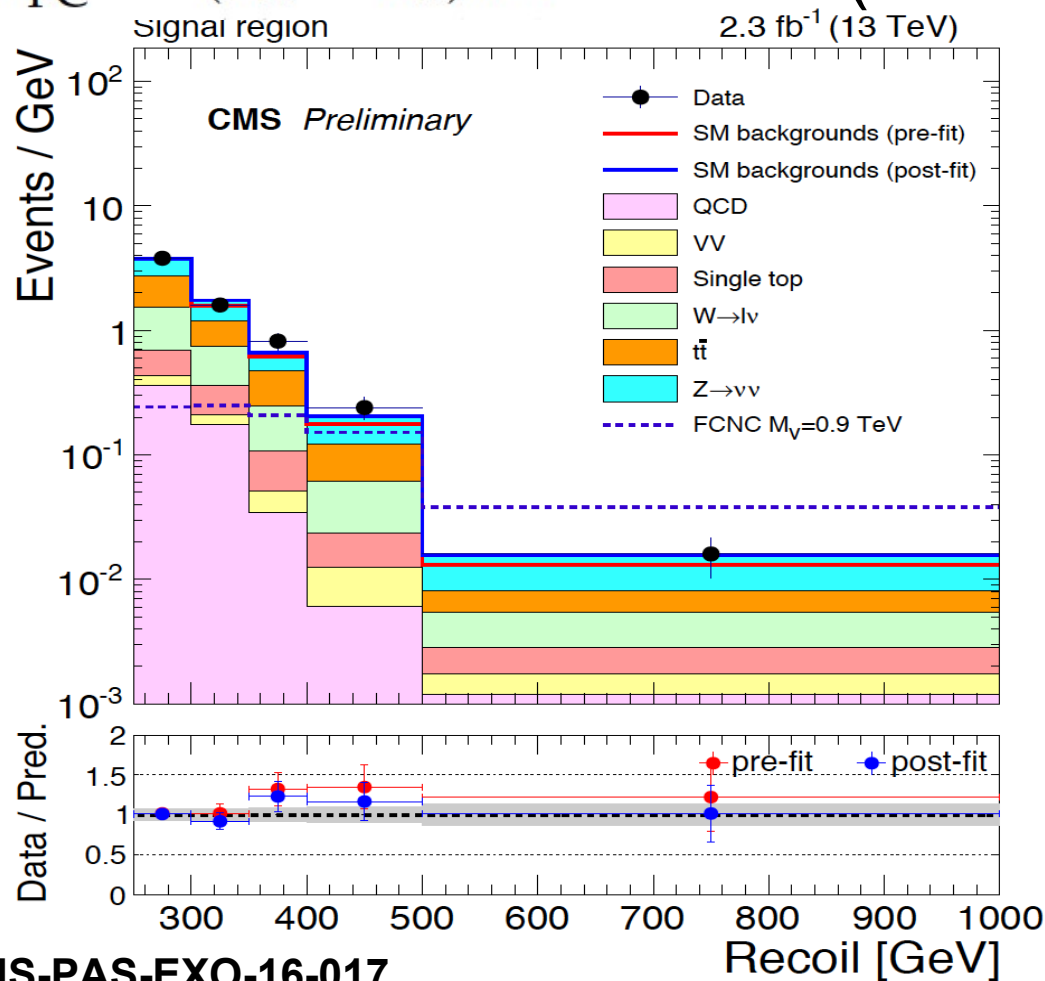
$$V_\mu \bar{u}_i [(a_{FC}^1)^{ij} \gamma^\mu + (b_{FC}^1)^{ij} \gamma^\mu \gamma^5] u_j$$

$$a_{FC}^1 = (a_R + a_L)/2$$

$$b_{FC}^1 = (a_R - a_L)/2$$

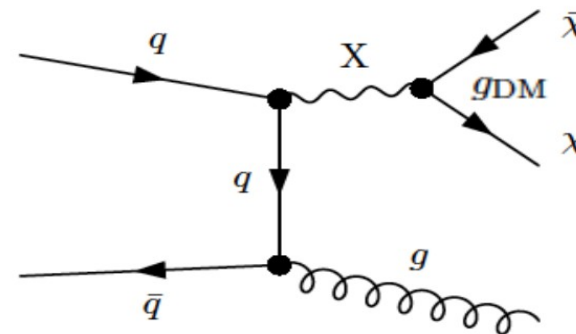
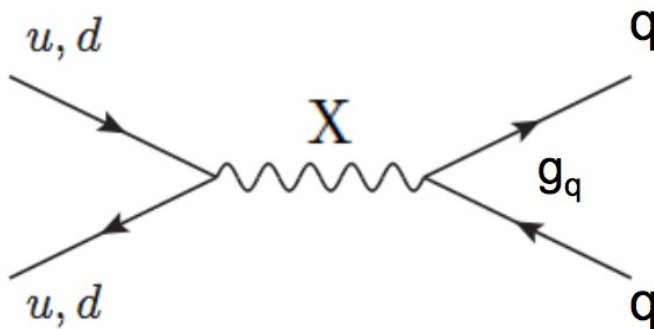
Taking $a_i=0.1$

(DM forum choice $0.25 \rightarrow \times 10$)

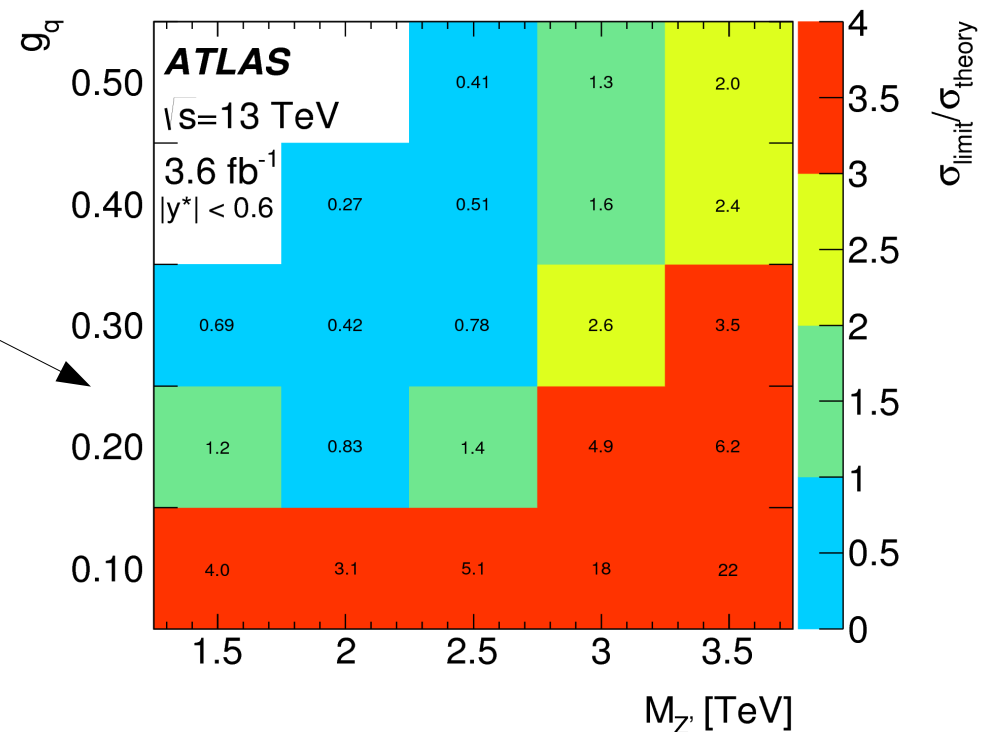
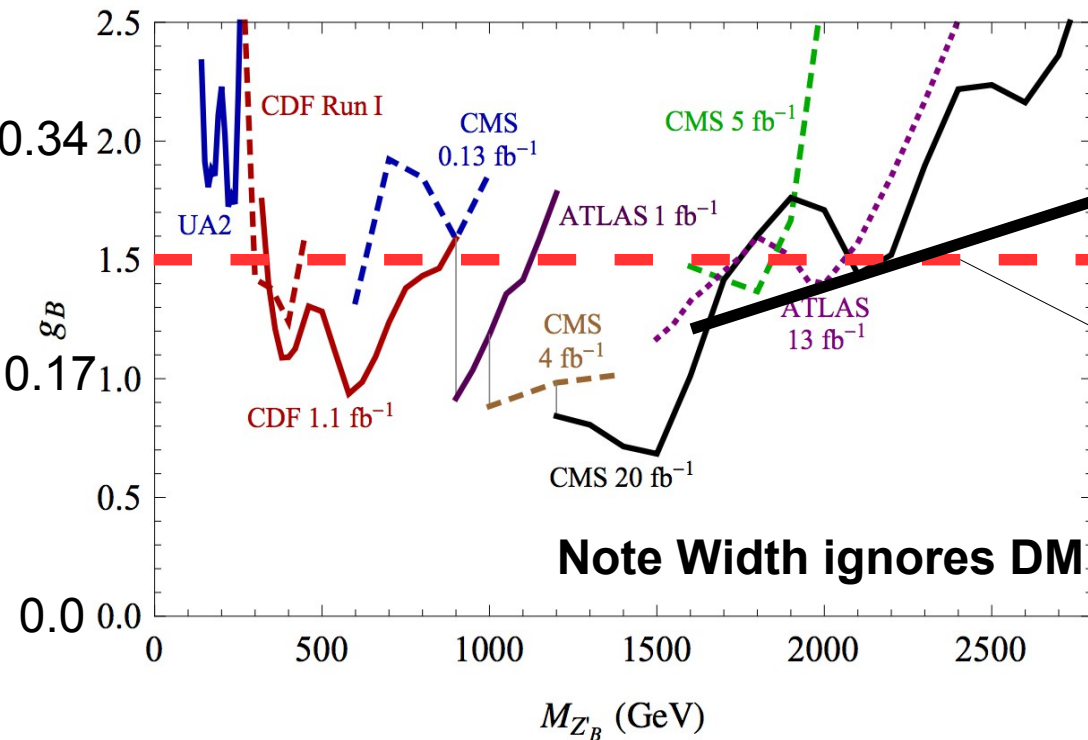


Mediator Search

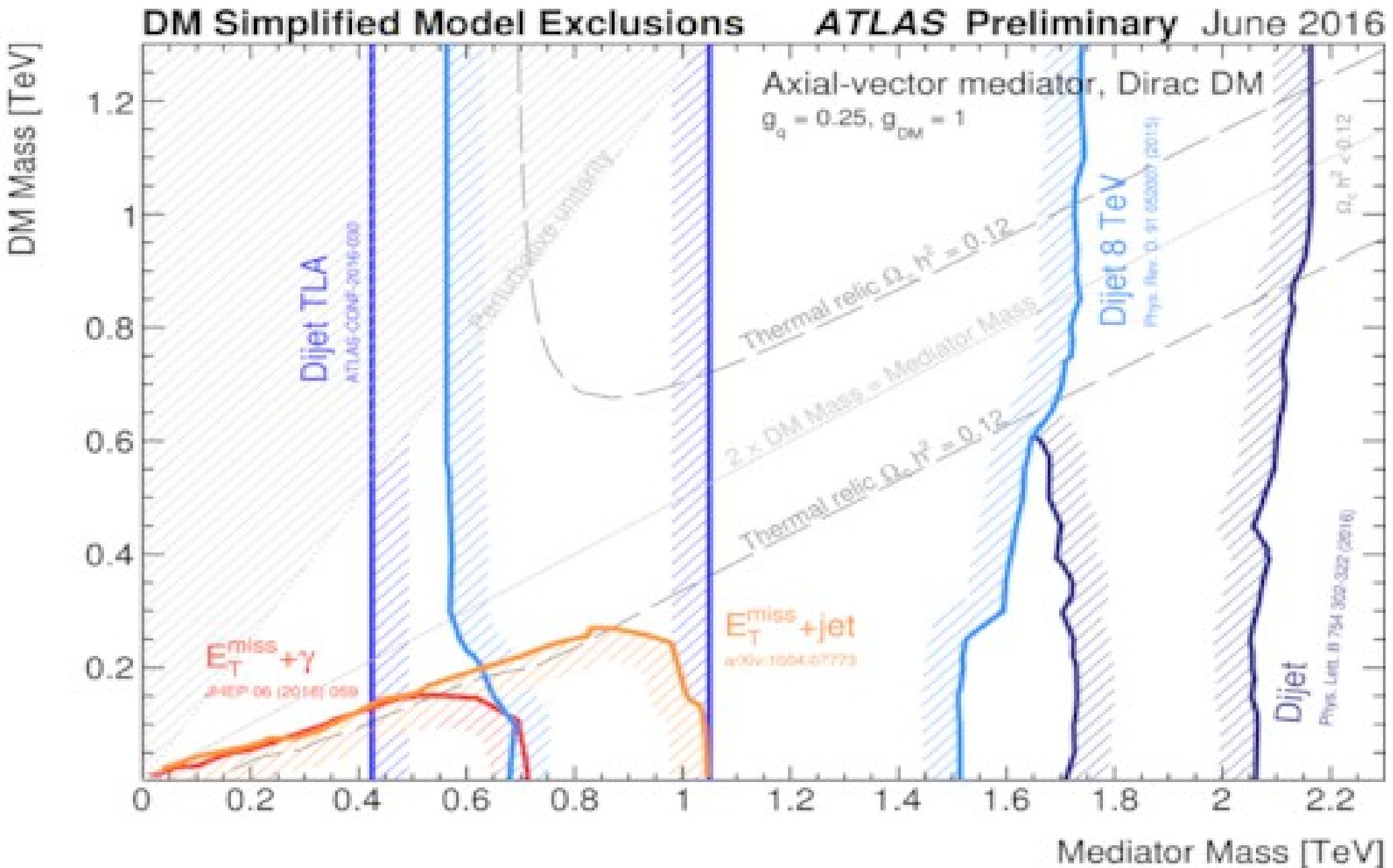
- In addition we can just look for the mediator



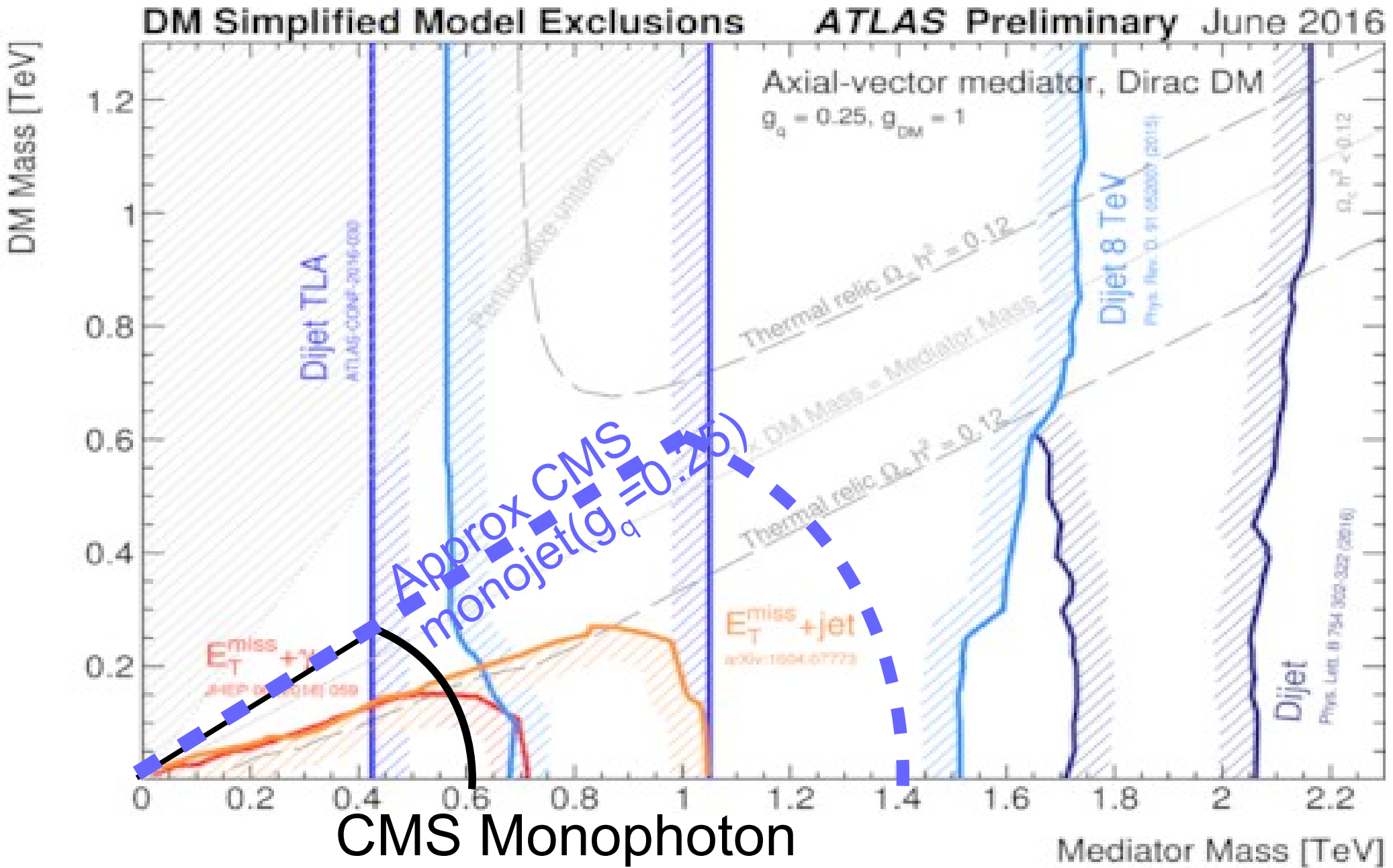
See Krisztian Peters' talk for more



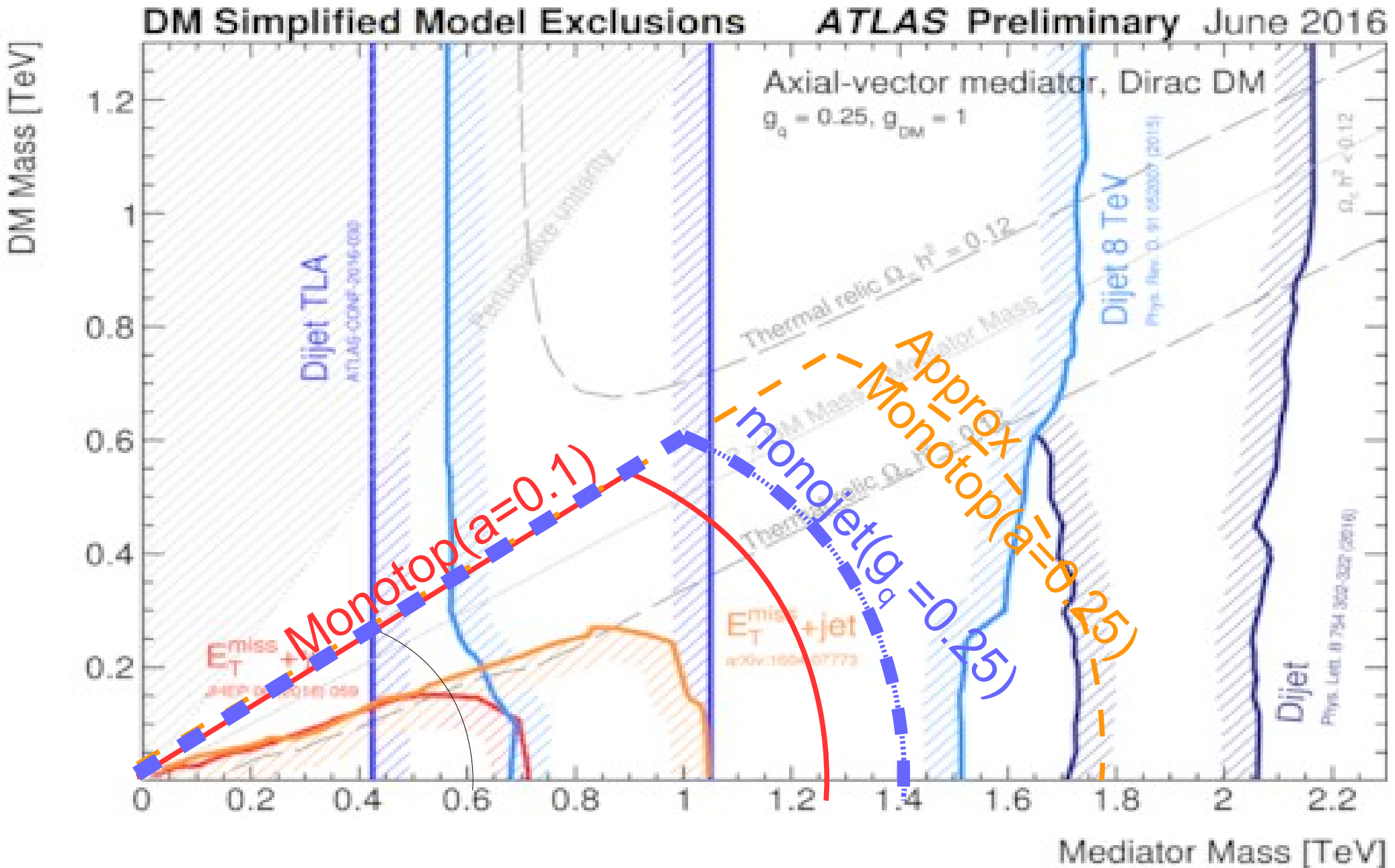
Putting it all together



Putting it all together



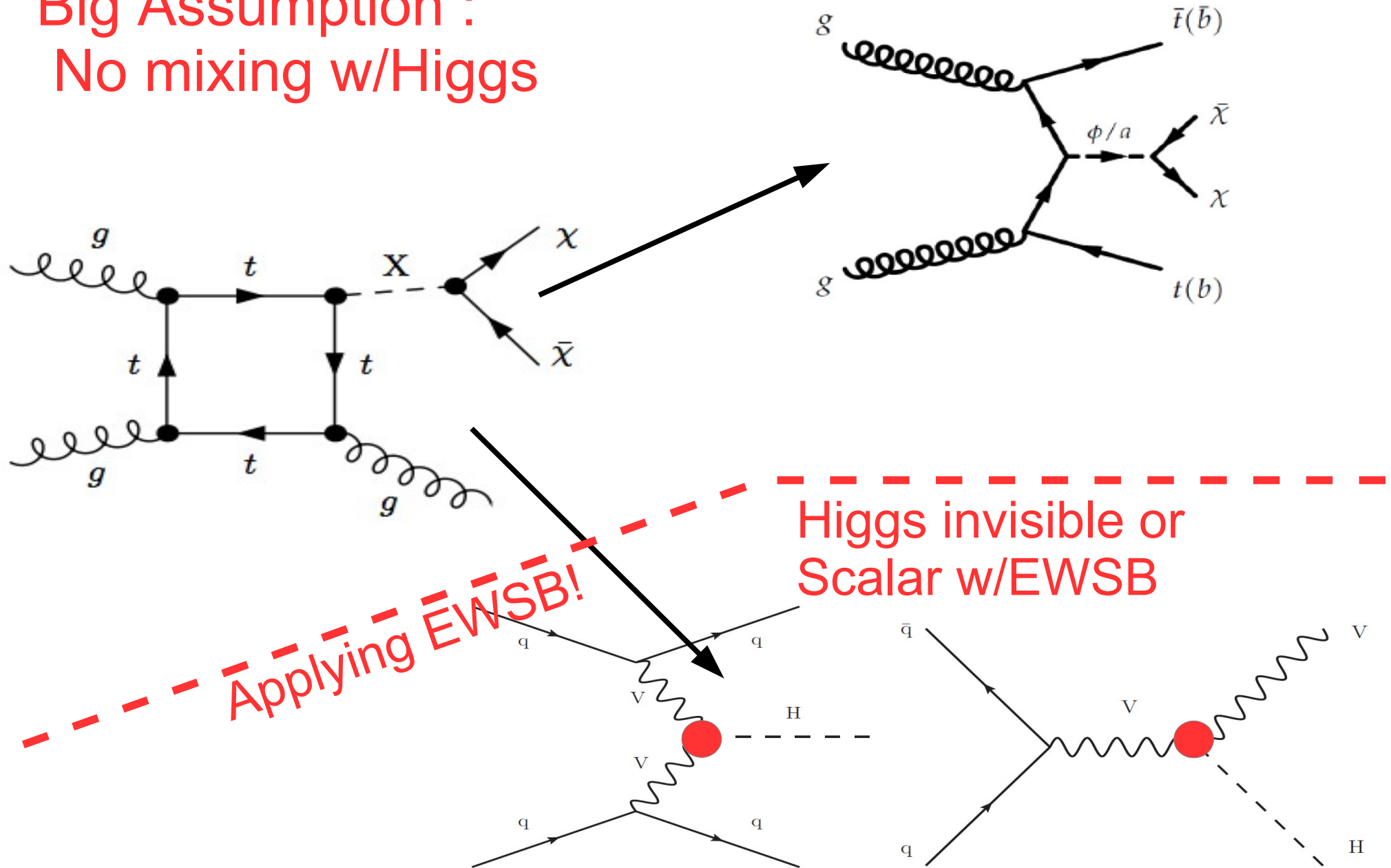
Putting it all together



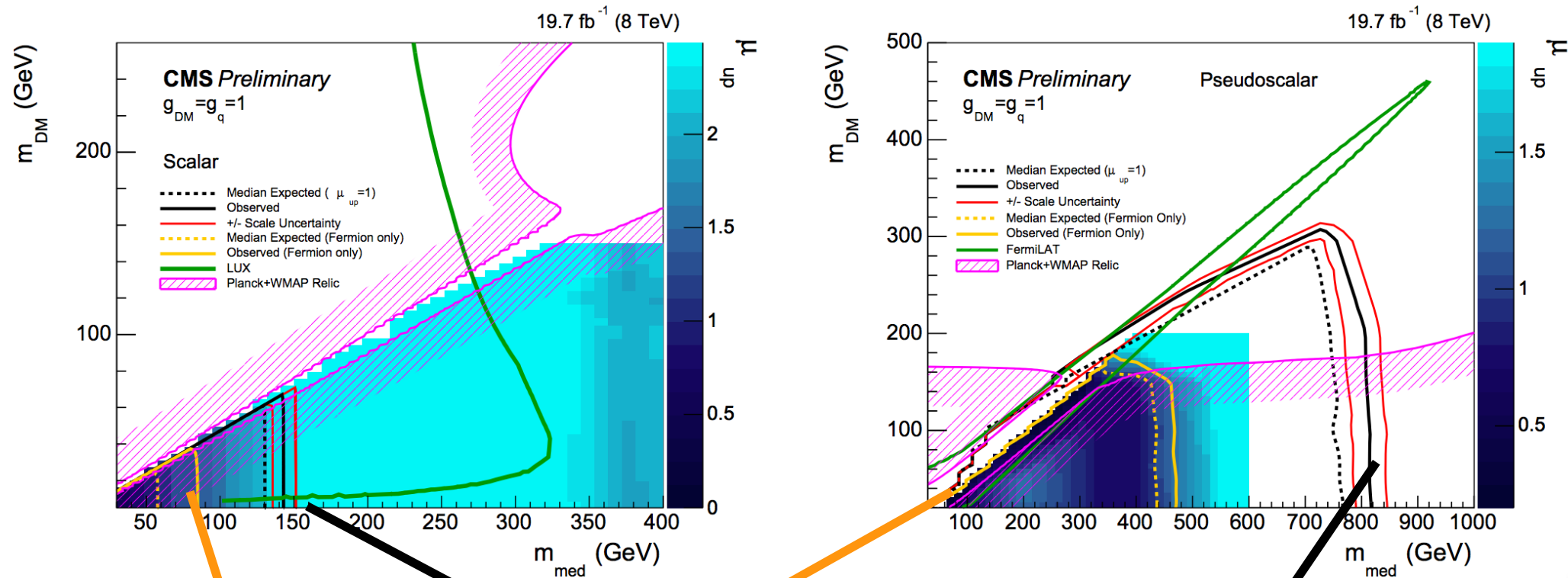
Spin 0

What can you do with Spin 0?

Big Assumption :
No mixing w/Higgs

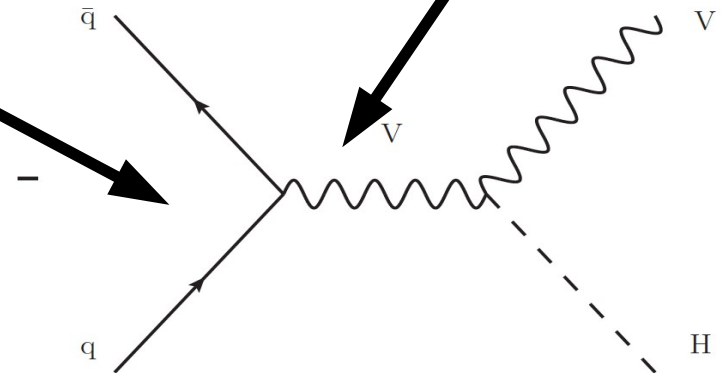
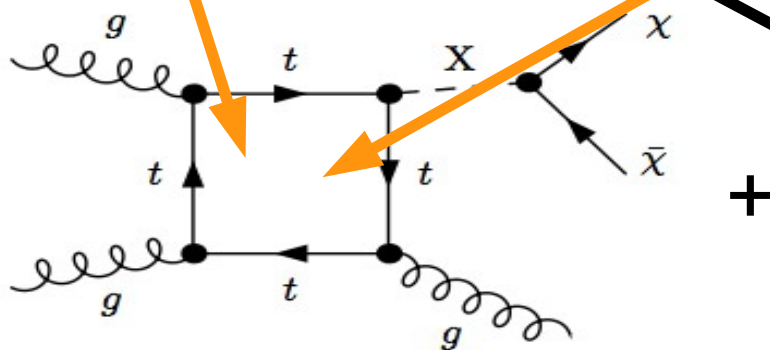


Our Current Scalar & Psuedo results

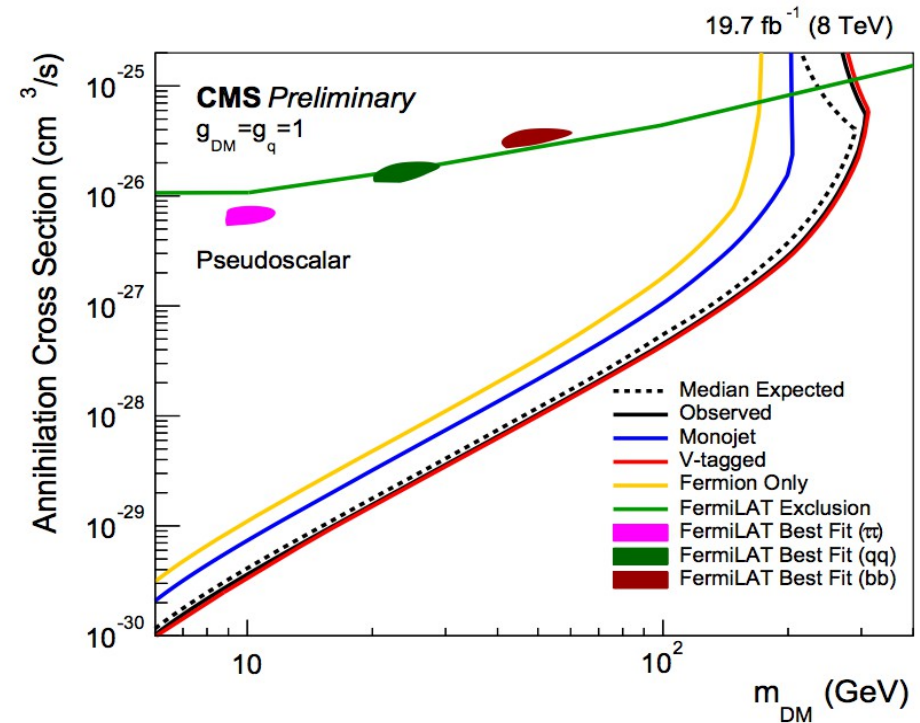
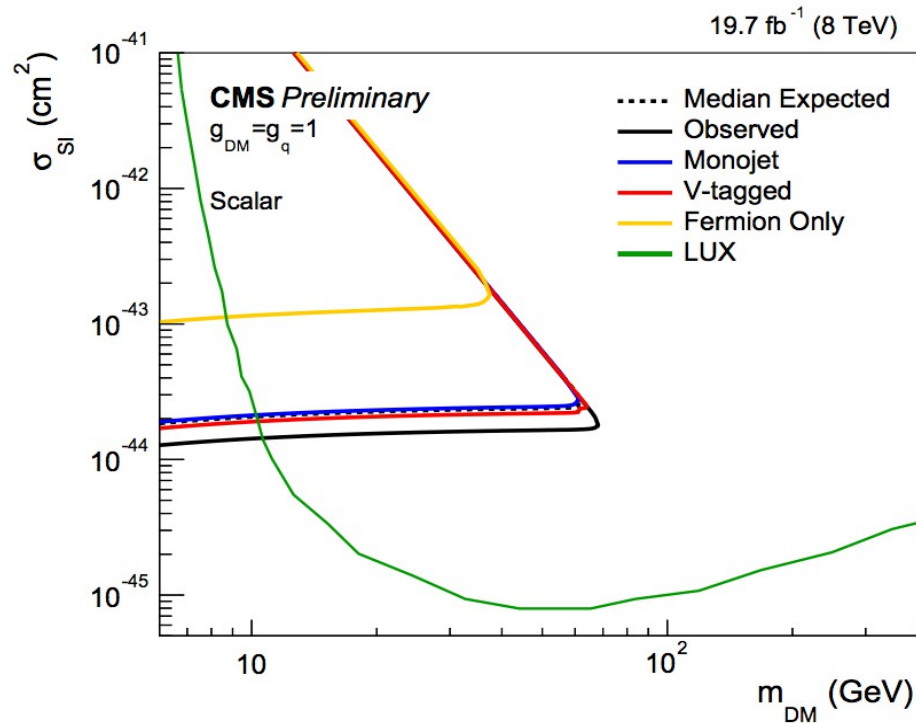


- Currently driven by 8TeV exclusion

Clear that there are issues



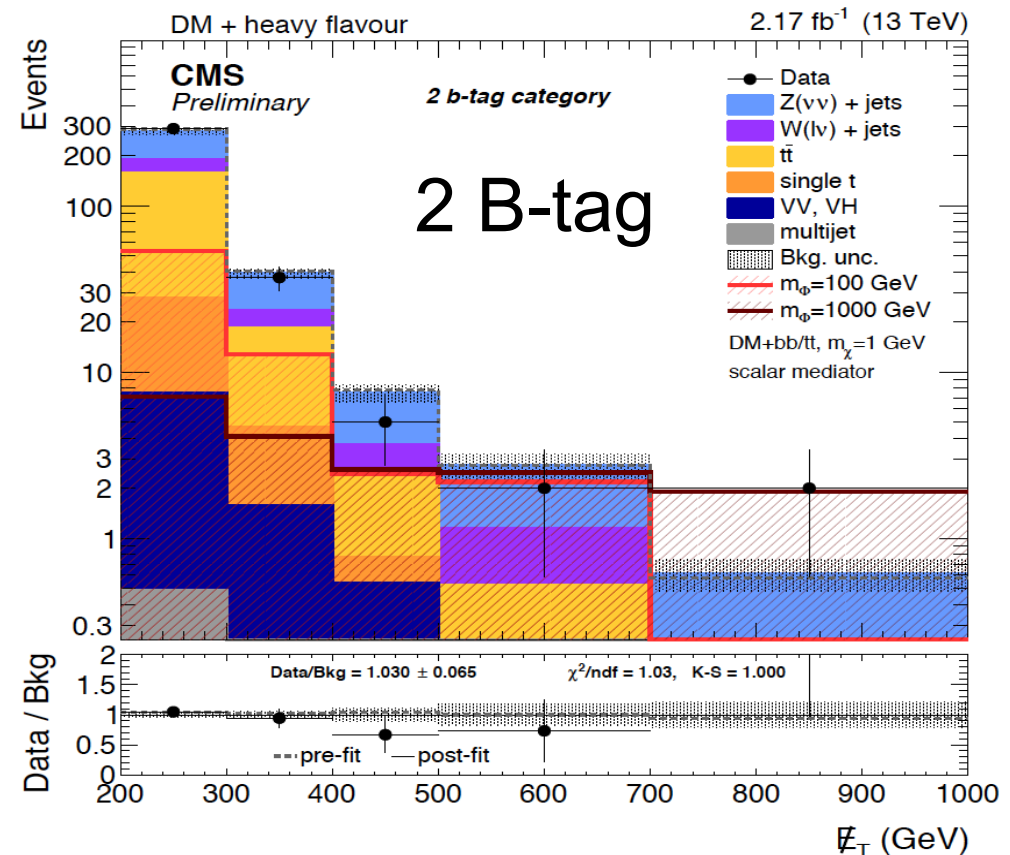
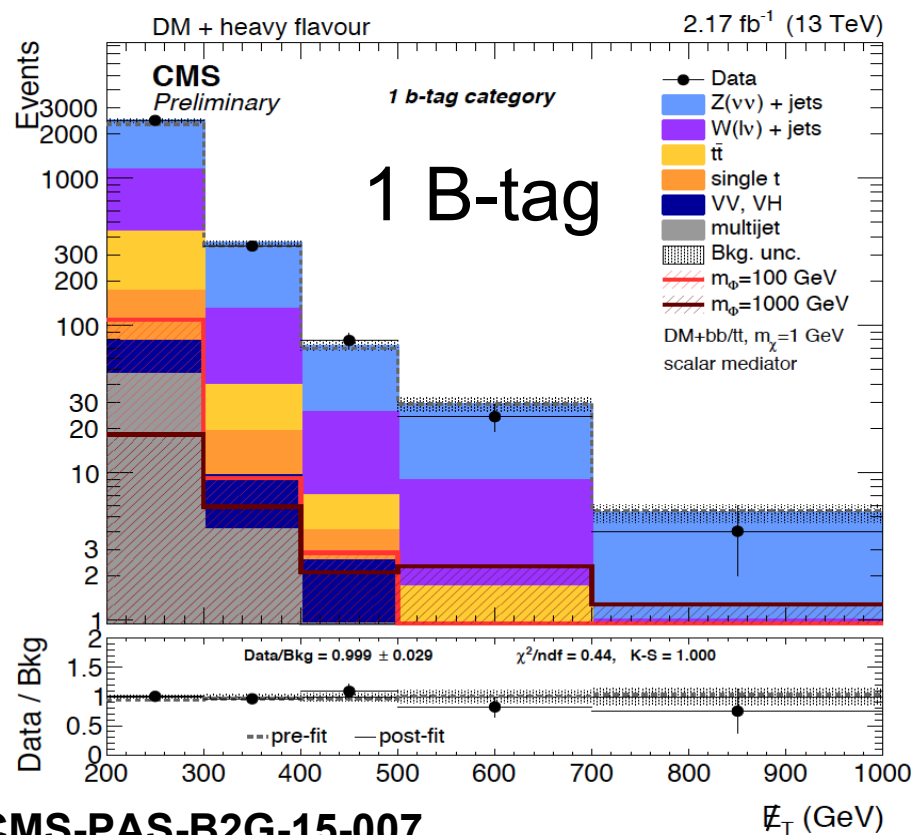
Our Current Scalar & Psuedo results



- When the dark matter is not onshell
 - Strong exclusion of pseudoscalar interpretation of LAT
 - Scalar and Direct detection are in close competition
 - Expect LHC to pass LUX this summer

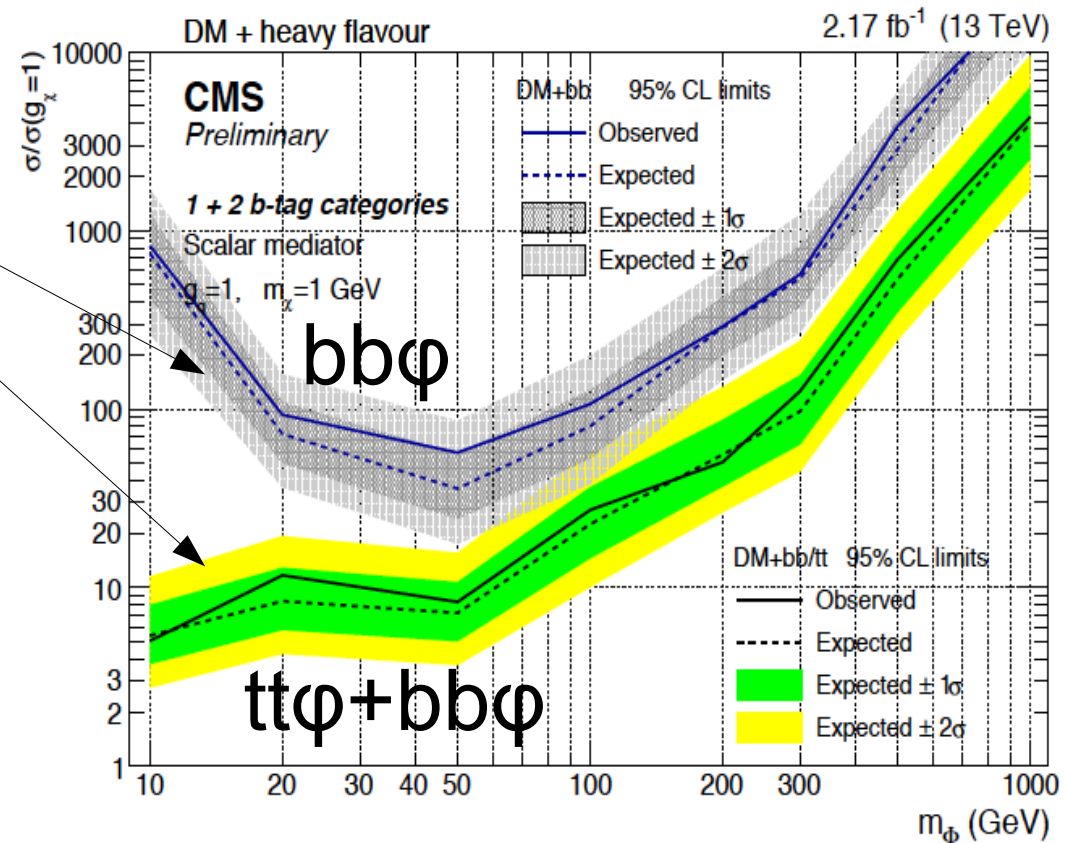
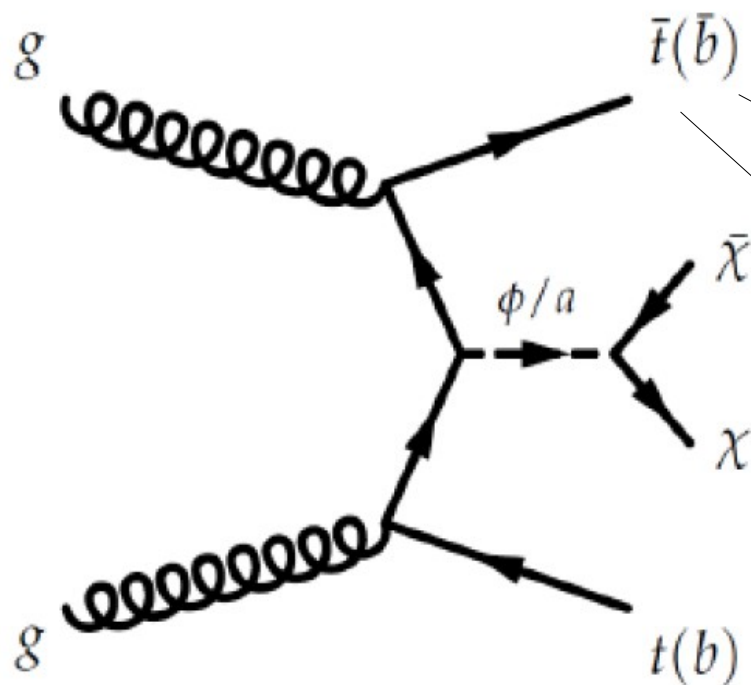
Heavy Flavor

- Mono-B or B(s)
 - Require less than 4 jets
 - Basically the monojet analysis with either 1 or 2 bs
 - Inject both $tt+DM$ and $bb+DM$ into the analysis



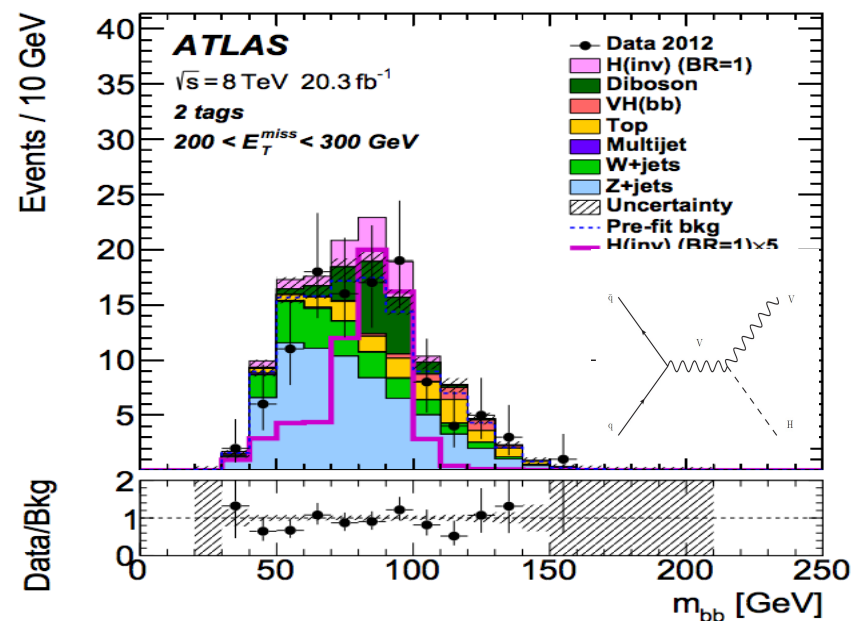
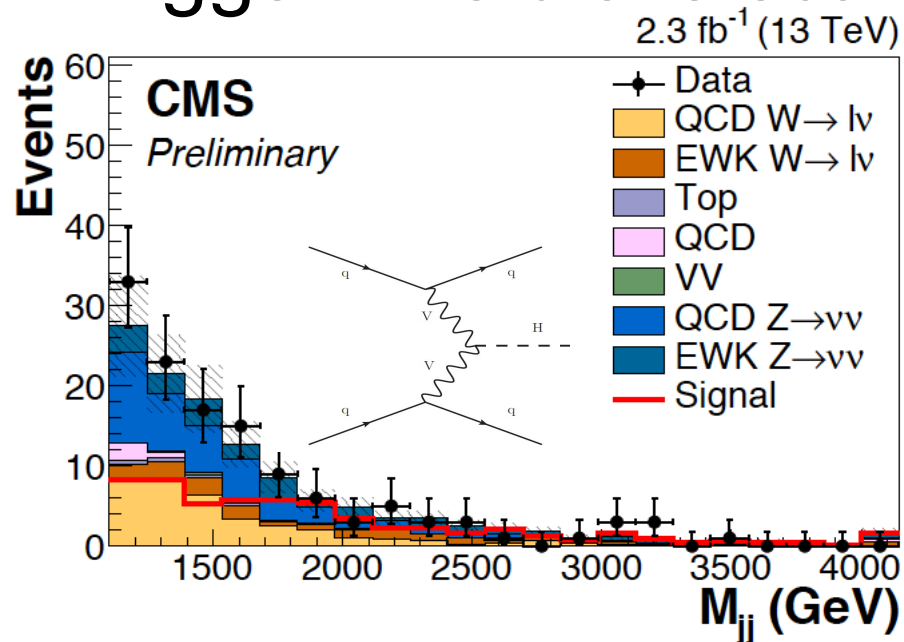
Heavy Flavor Results

- Mono-B or B(s)
 - Note that this is only < 4 jets
 - Inject both $tt+DM$ and $bb+DM$ into the analysis



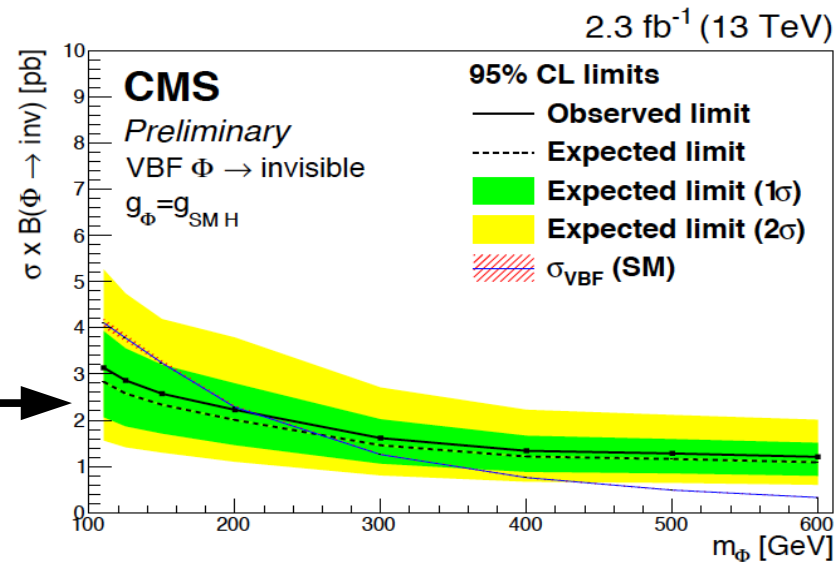
Higgs Invisible Interpretation

- Higgs Invisible is scalar model with EWSB

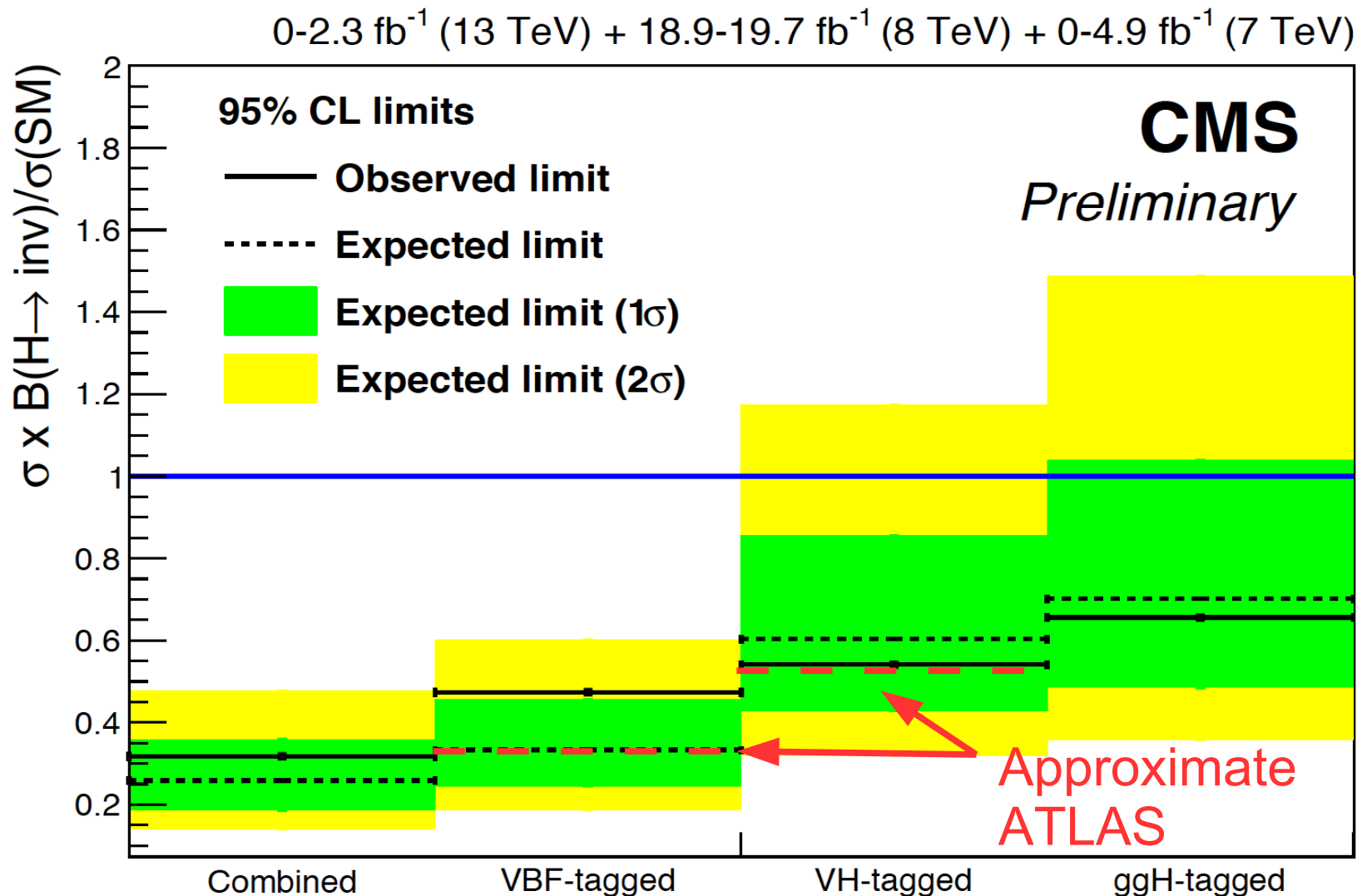


30% improvement in VBF from tying the $W \rightarrow Z$ constraints together in control regions

Starting to scan mass



Results

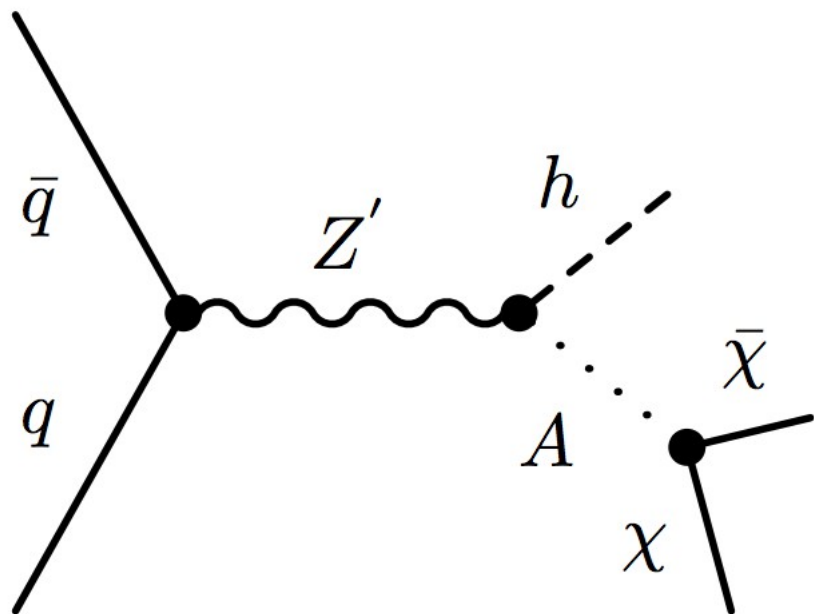


CMS : $\text{BR}(H \rightarrow \text{Inv}) < 0.32$ (0.26 expected)

ATLAS : $\text{BR}(H \rightarrow \text{Inv}) < 0.25$ (0.27 expected)

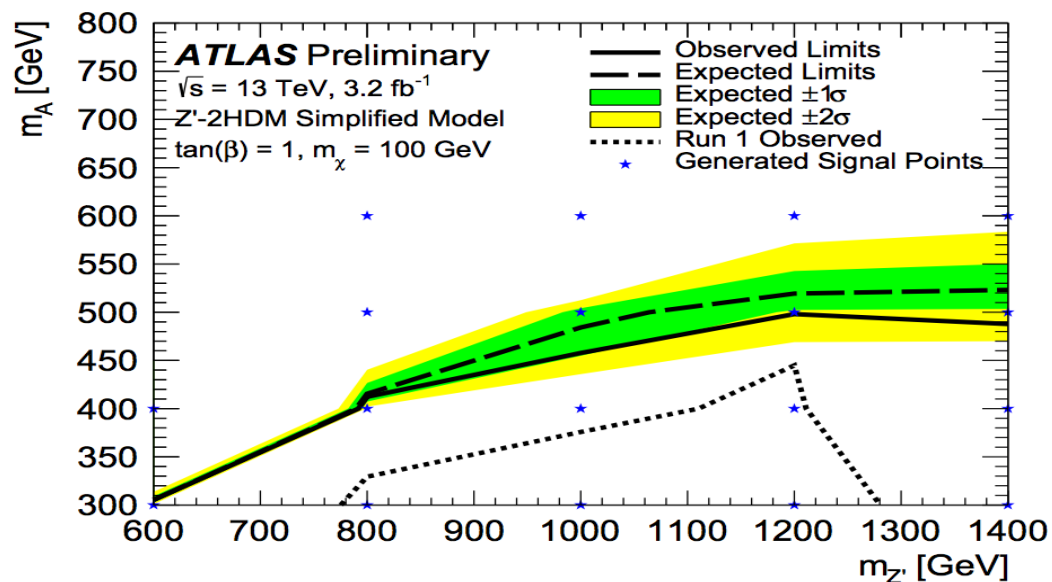
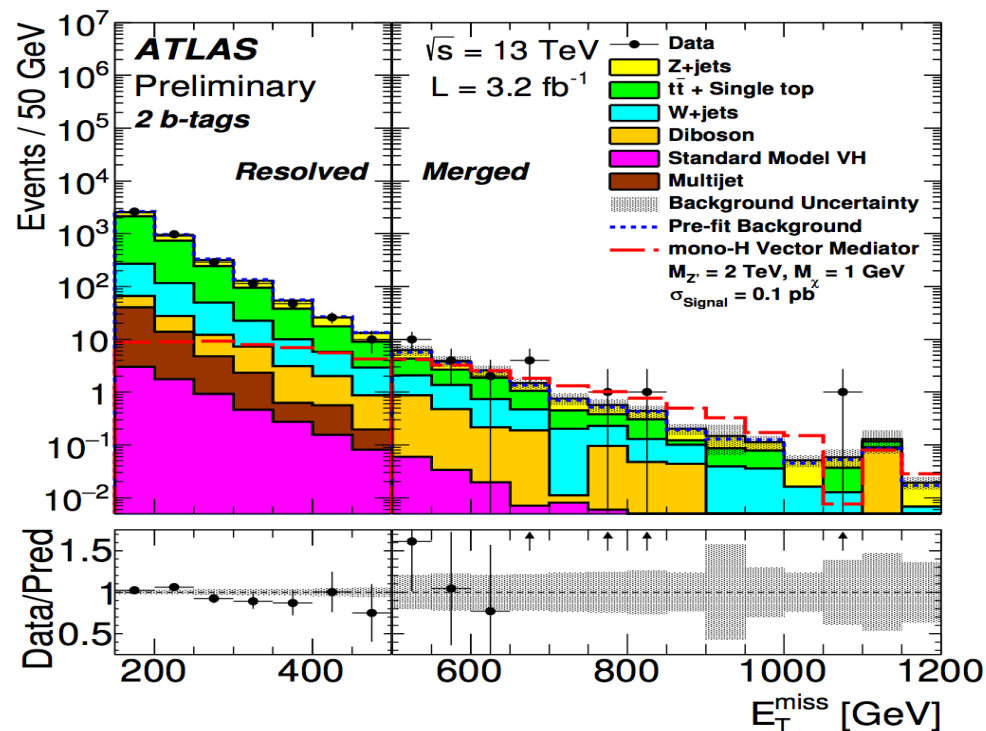
Mono-Higgs

Adding Spin 1 and Spin 0 mediators

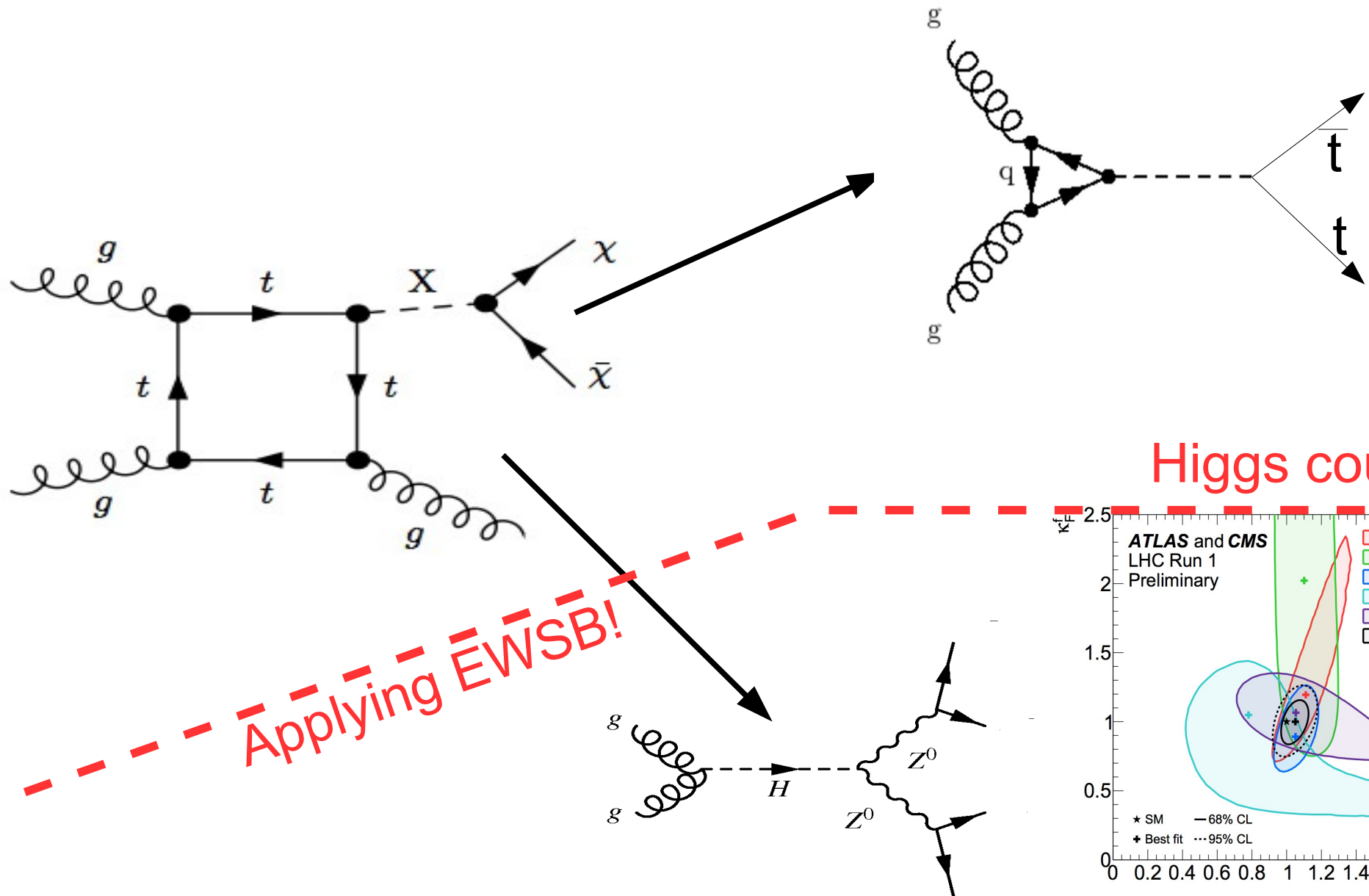


$h \rightarrow b\bar{b}$ bounds drive
mono-Higgs

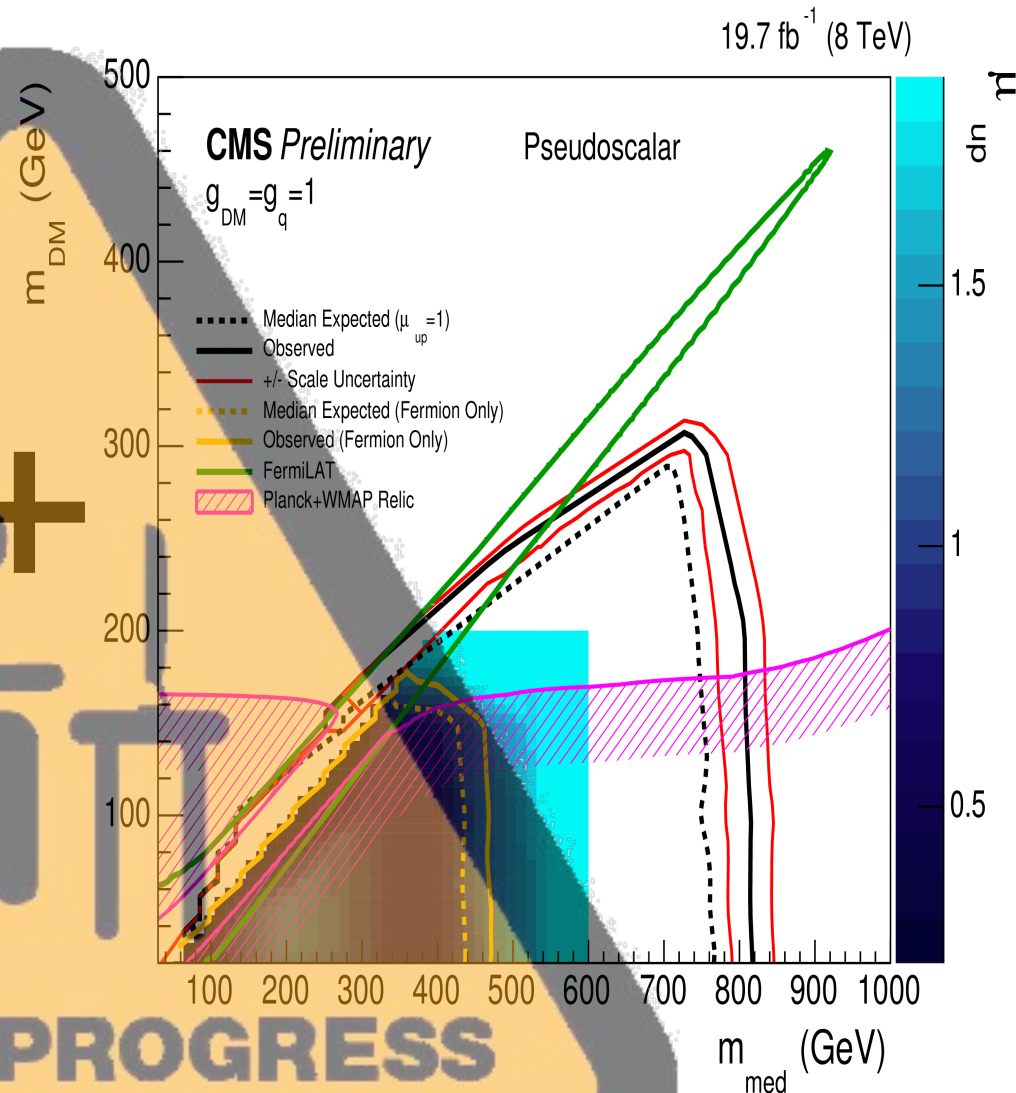
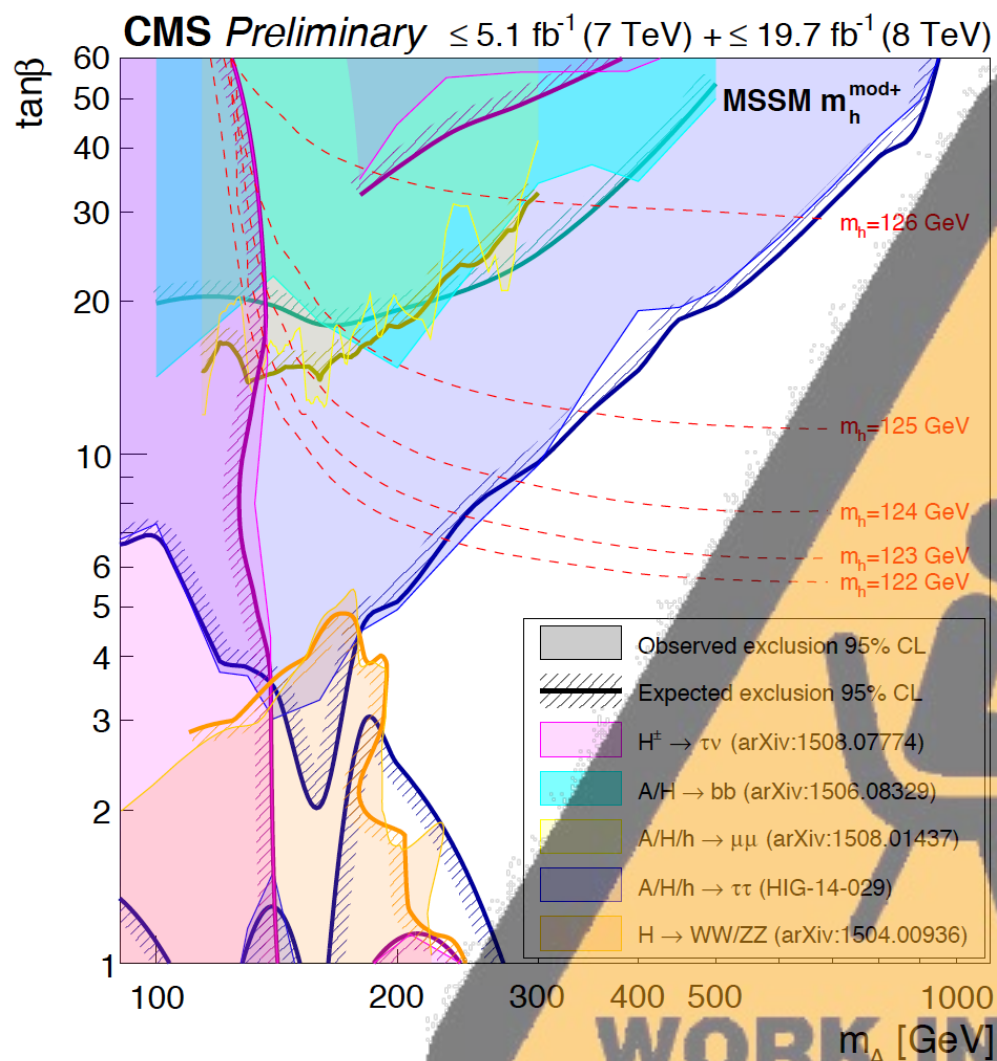
B-tagging forces $t\bar{t}b\bar{b}$
background to drive
analysis



What about the visible?



Not yet available



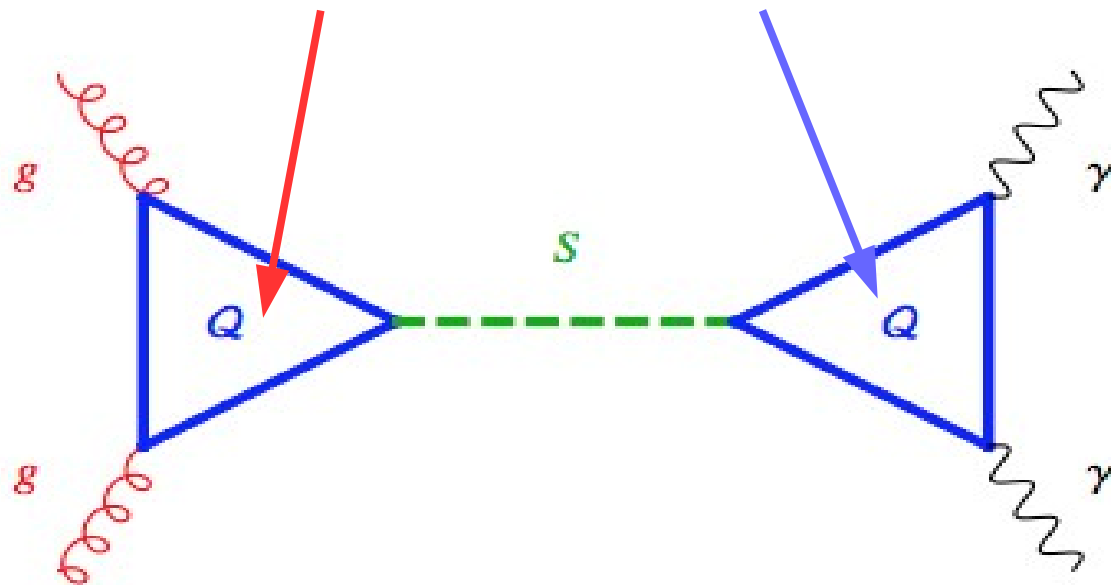
Spin 0 Di-photon (미안 해요)

The simplest DM model

- Lets try to make something super basic
 - Basic model

$$\mathcal{L}_S = \underbrace{g_\chi S \bar{\chi} \chi}_{\text{black}} + \underbrace{\frac{\alpha_s}{4\pi} \frac{c_g}{\Lambda} S G_{\mu\nu}^a G^{a\mu\nu}}_{\text{red}} + \underbrace{\frac{\alpha}{4\pi} \frac{c_\gamma}{\Lambda} S F_{\mu\nu} F^{\mu\nu}}_{\text{blue}}$$

??



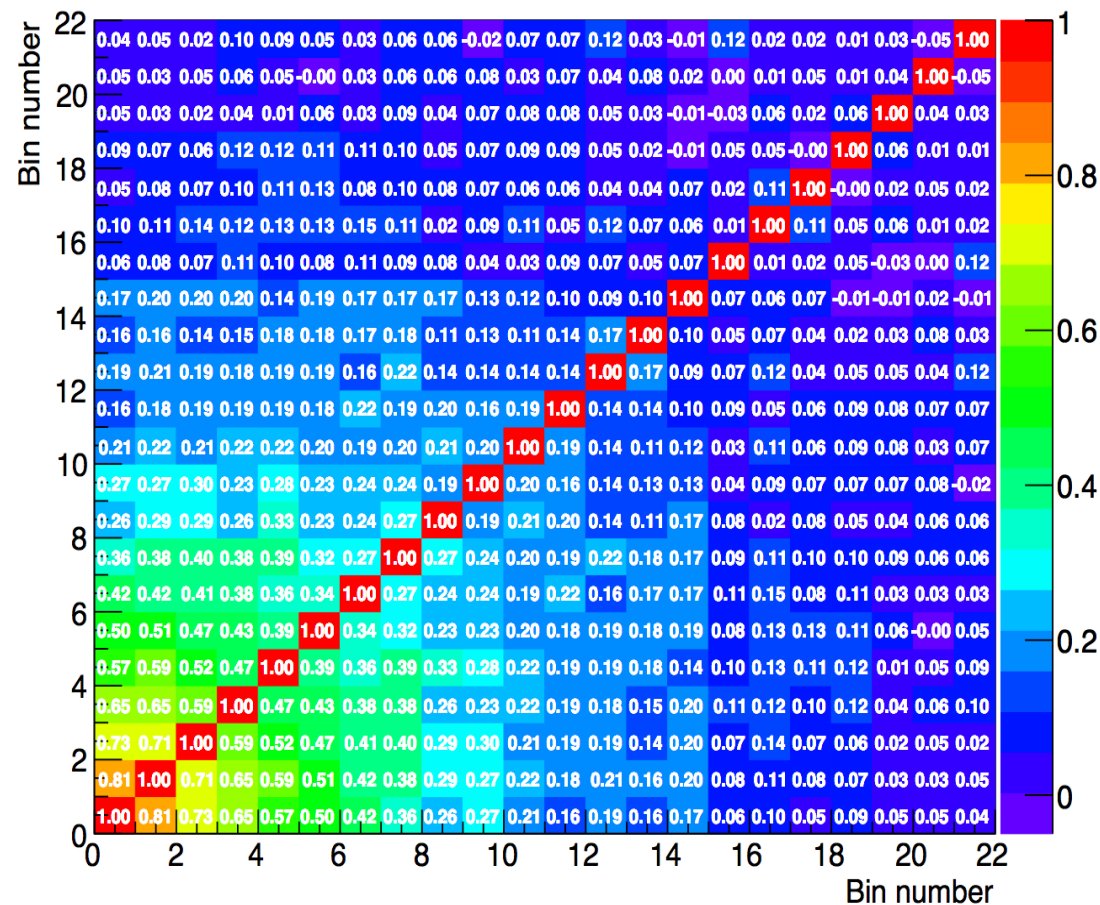
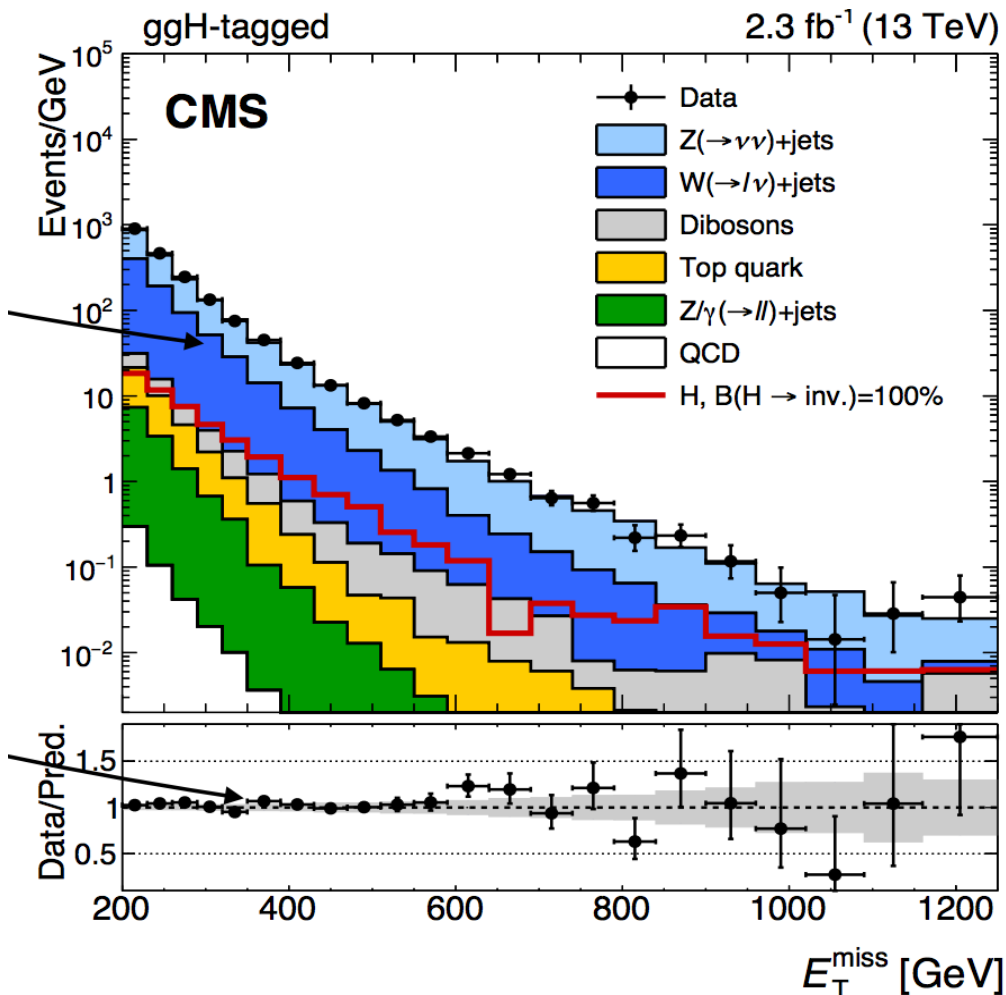
Re-interpreting the Analysis

- Future plans from LHC→simplified likelihoods

- What is it? => Reduced control fit to 2 objects

MET distribution

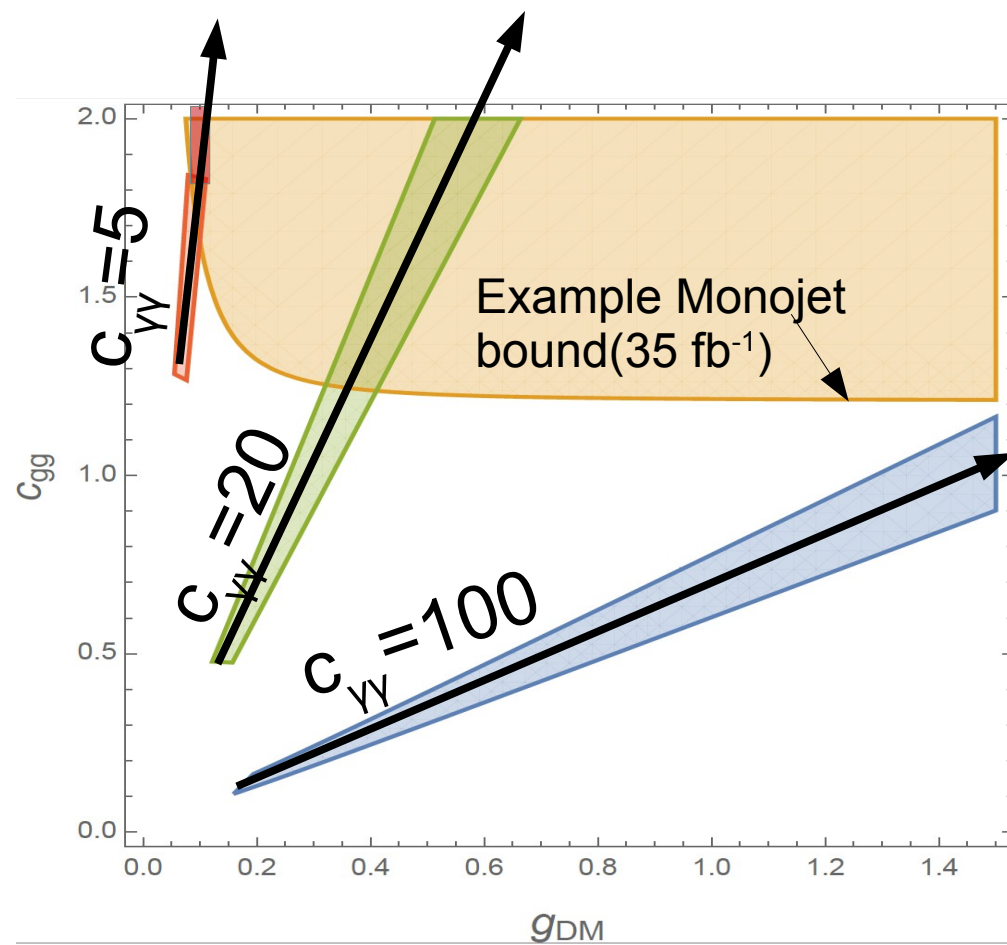
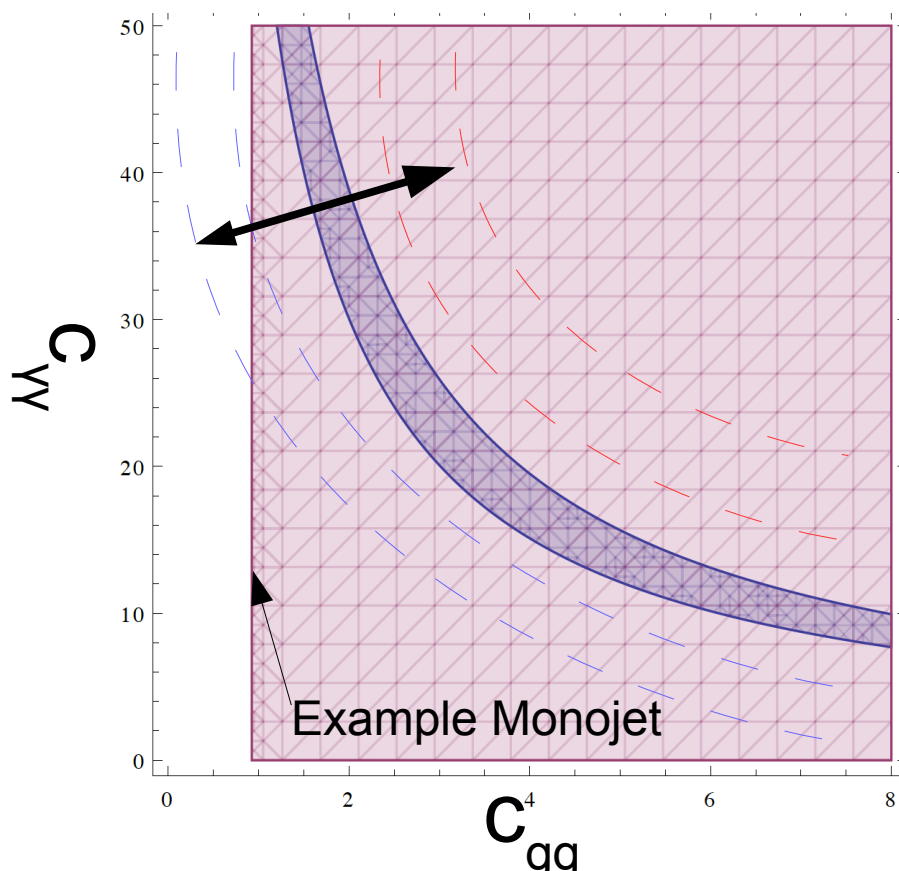
Toy Bin by bin covariance



From this: setup full CL_s get both expected and observed

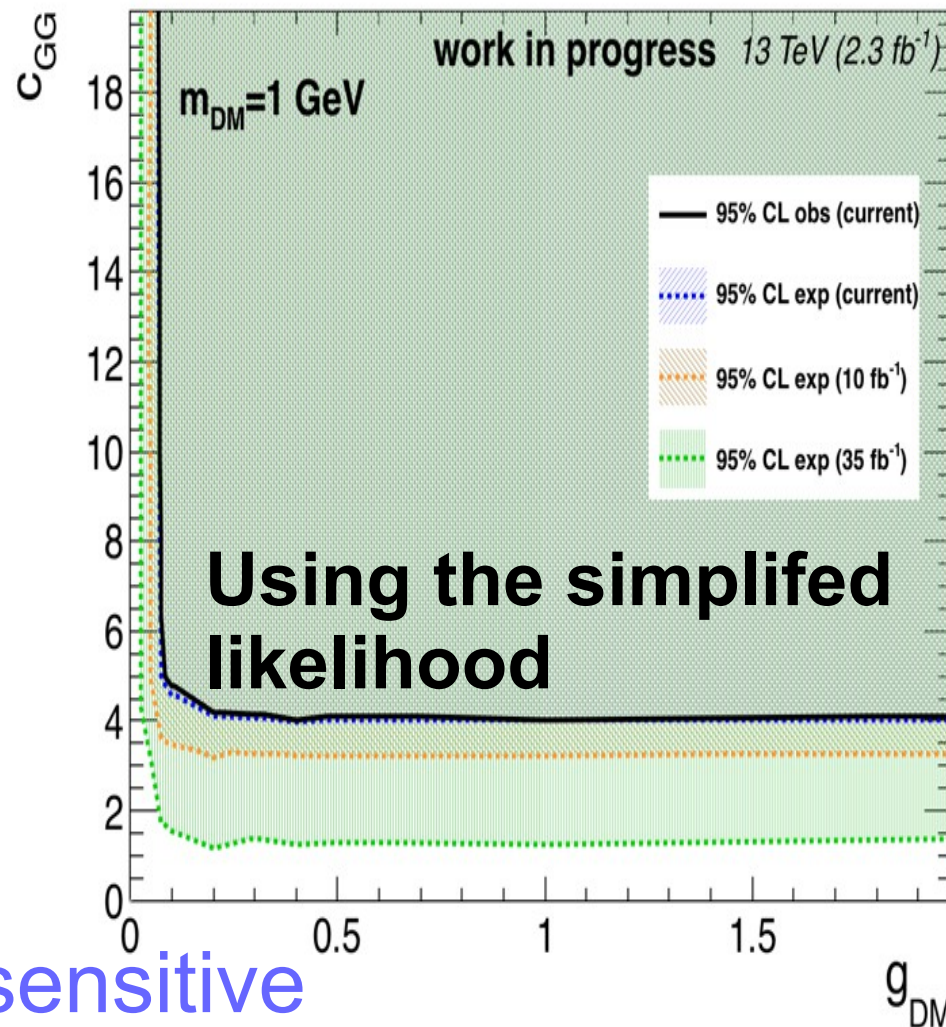
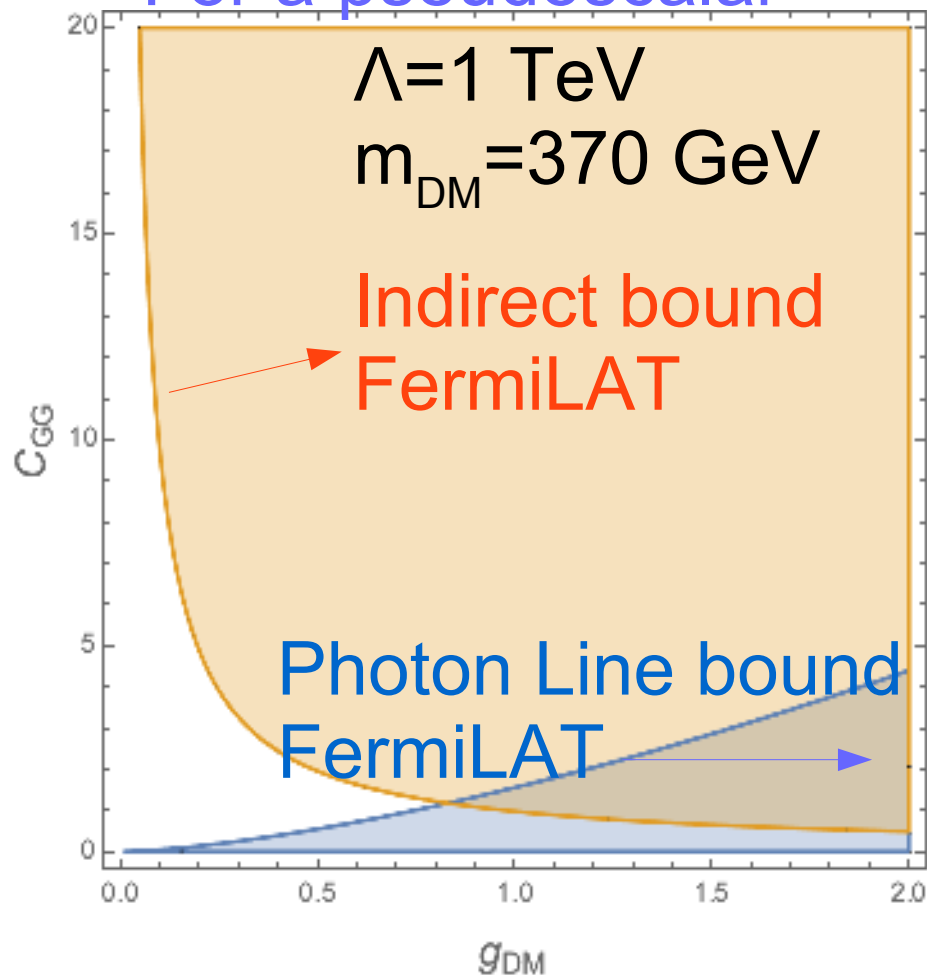
Using the best fit cross section

- We have 3 free couplings :
 - $g_{\text{DM}}, c_{\text{GG}}, c_{\text{YY}}$
- Taking the photon best fit can constrain one



Considering All Constraints

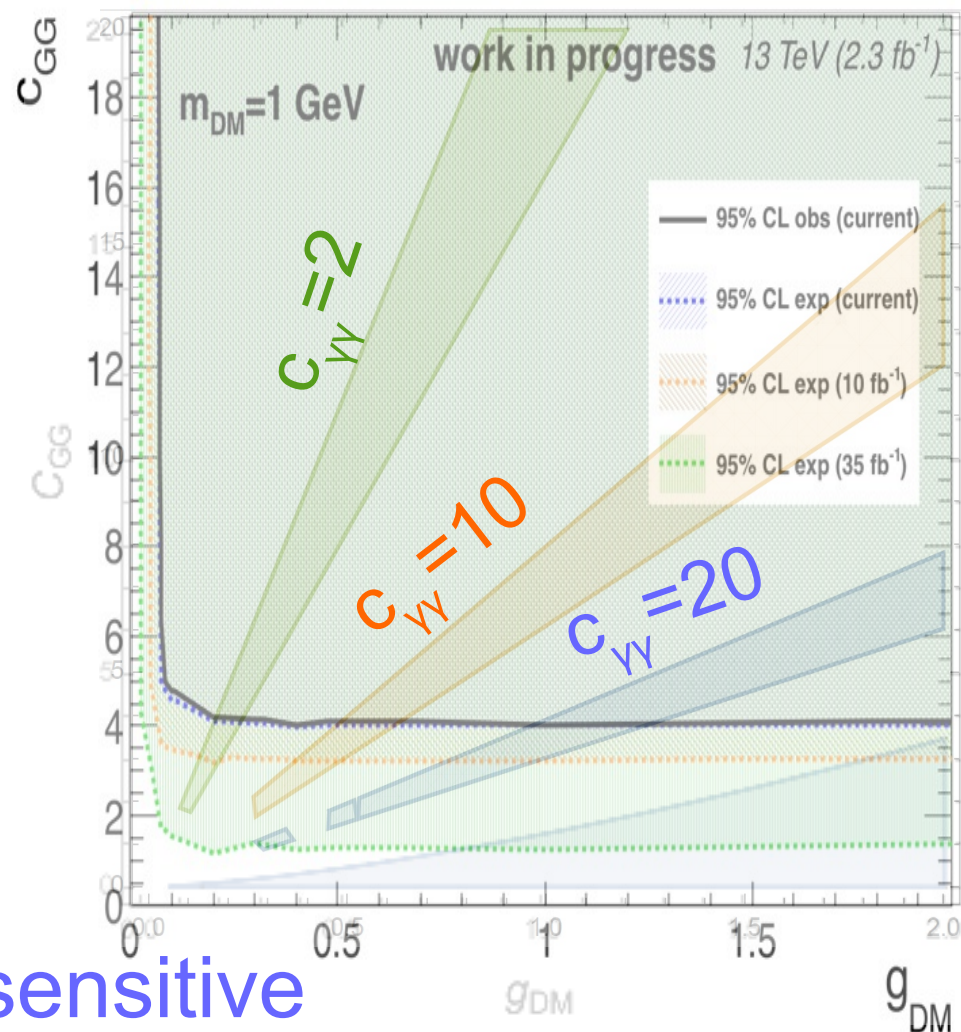
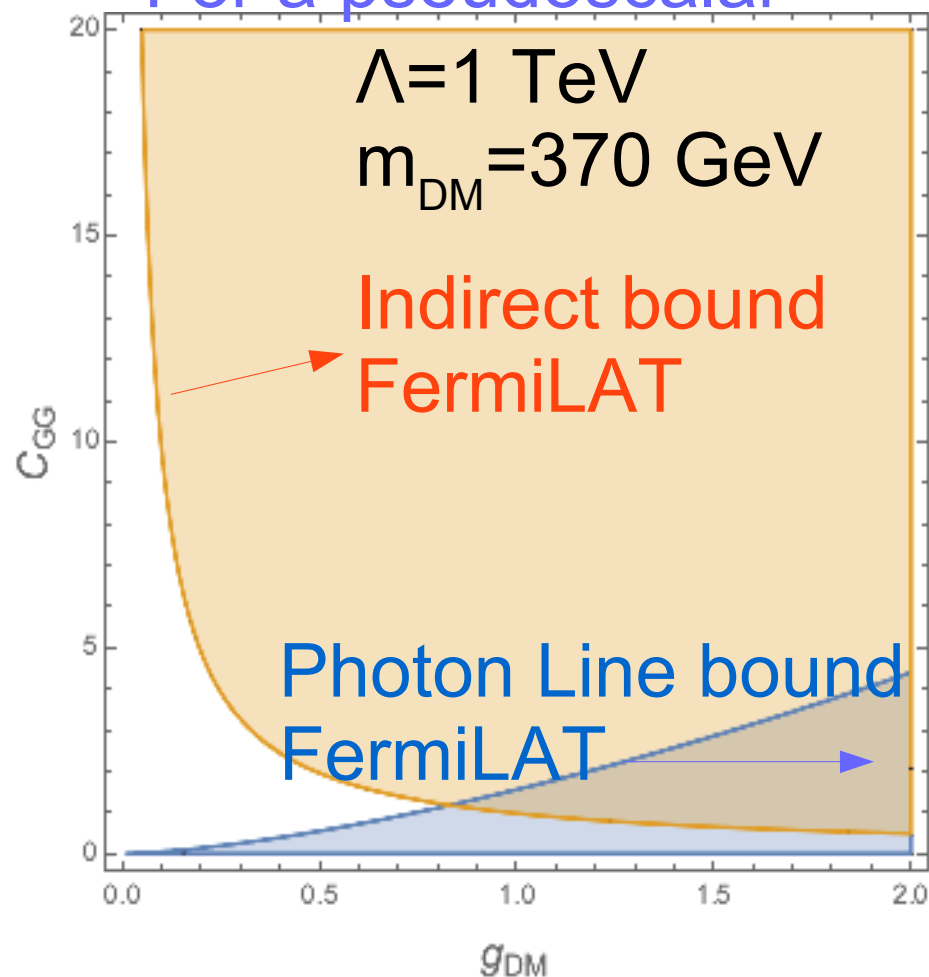
For a pseudoscalar



- Direct detection not yet sensitive
- Indirect detection limits on-shell production
- Photon Line bounds limit photon coupling < 100

Considering All Constraints

For a pseudoscalar



- Direct detection not yet sensitive
- Indirect detection limits on-shell production
- Photon Line bounds limit photon coupling < 100

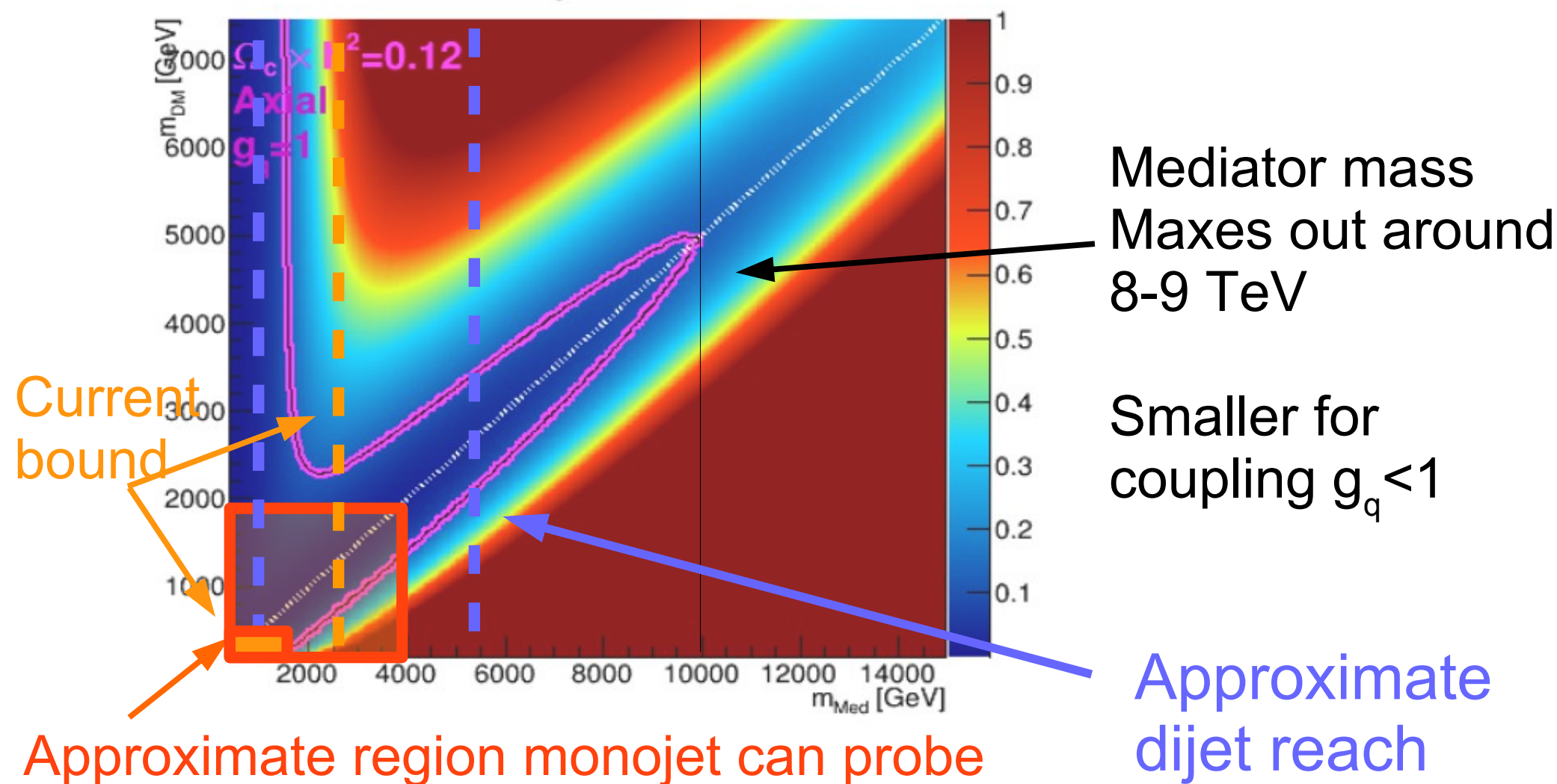
Outlook

Outlook

- Spin 1 :

- Dijet and monojet will continue to push out the bounds

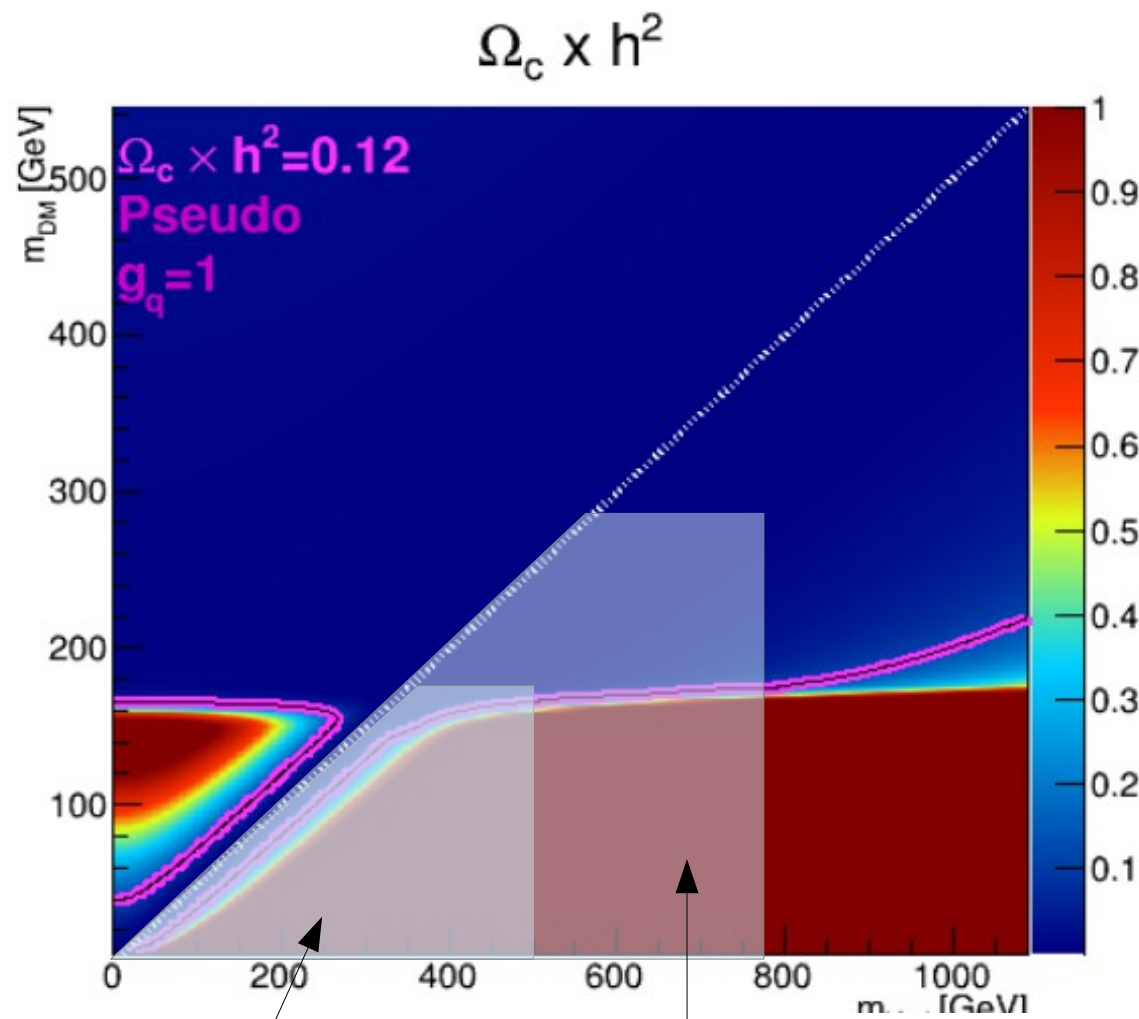
$$\Omega_c \times h^2$$



Hep-ph/1603.08525
 Hep-ph/1509.02904
 EXO-12-055
 HIG-16-012

Outlook

- Spin 0 :
 - Yet to truly coalesce in 13 TeV



Current monojet Approx future

Current reach for
 Higgs-like Scalar
 w/EWSB + no mixing

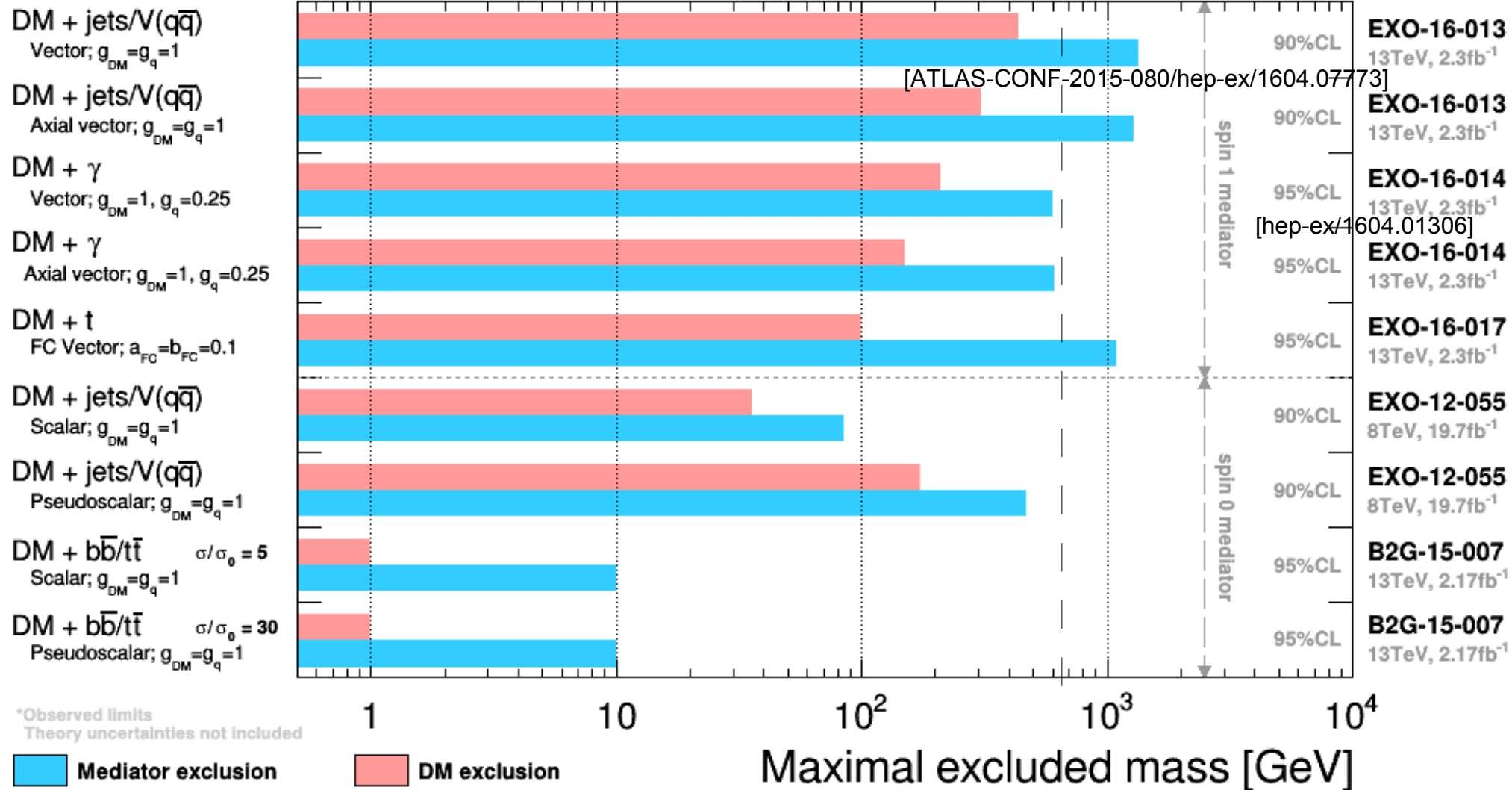
Mono-Z : 130 GeV
 VBF : 250 GeV
 Mono-V/j : 150 GeV

Heavy Flavor :
 1st bb/tt+ ϕ result
 No exclusion yet

Summary

CMS Preliminary

Dark Matter Summary* - June 2016



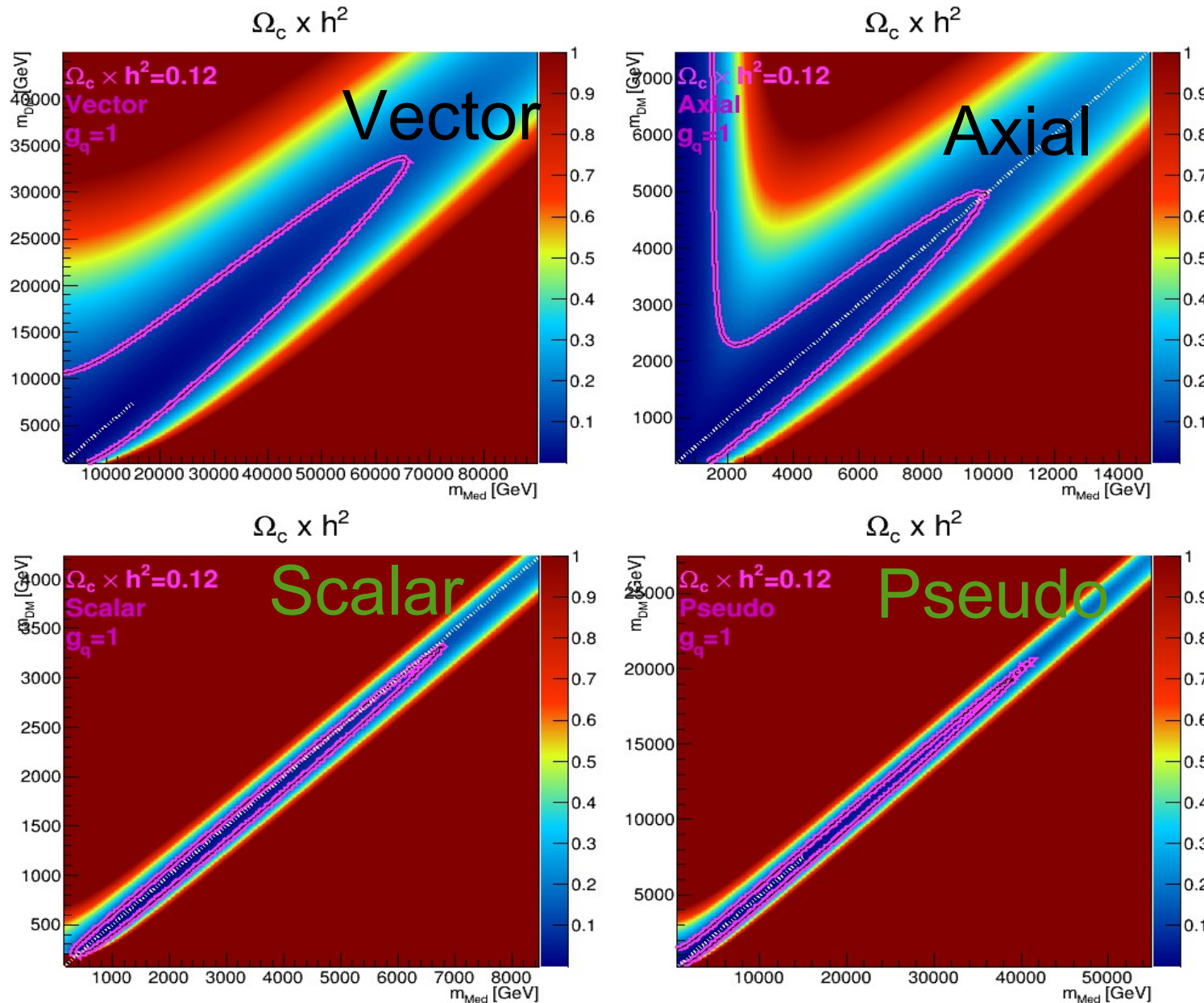
ATLAS-CONF-2016-019
ATLAS-CONF-2016-011

Thanks

감사합니다

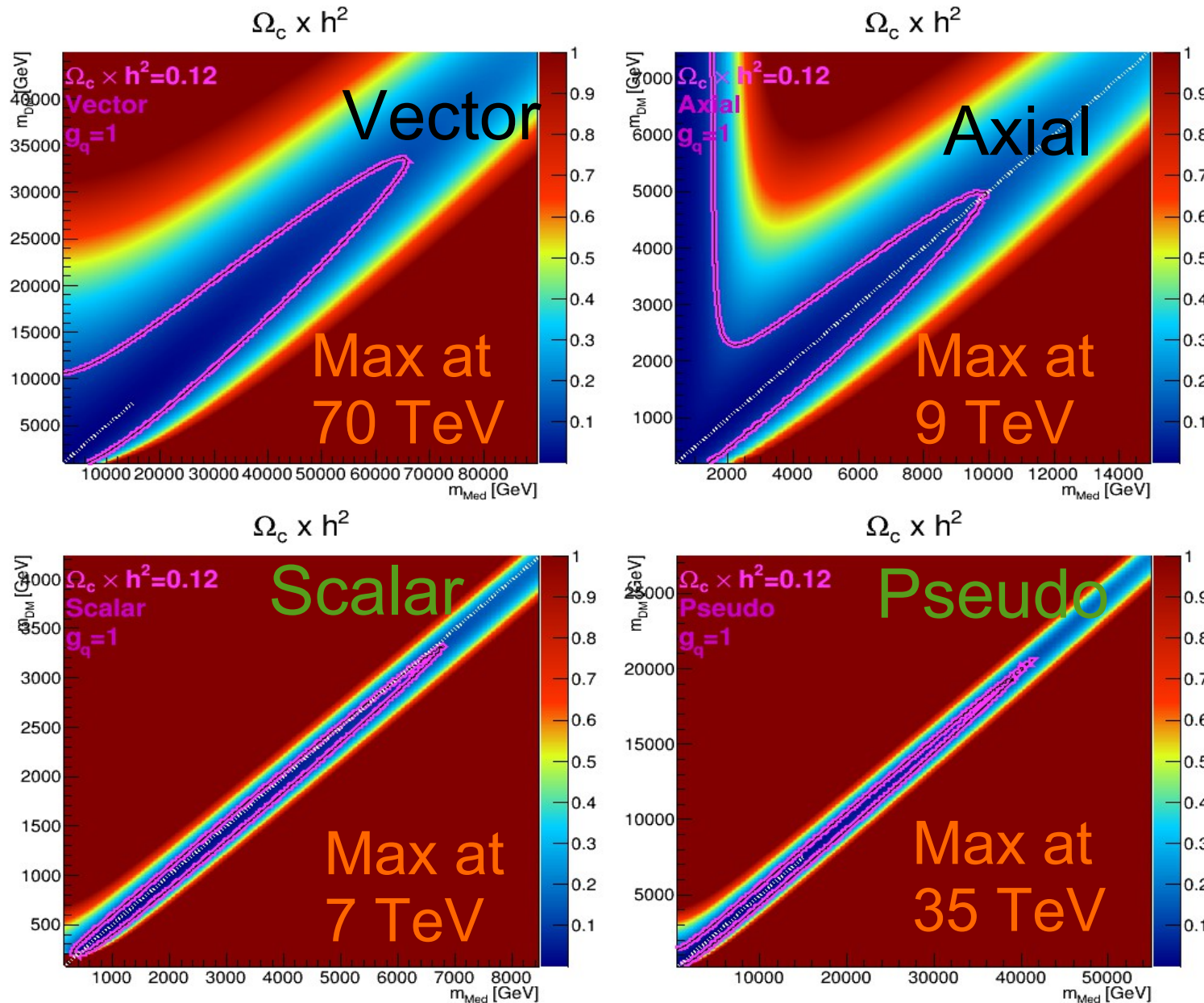
How does it look for our 4 friends?

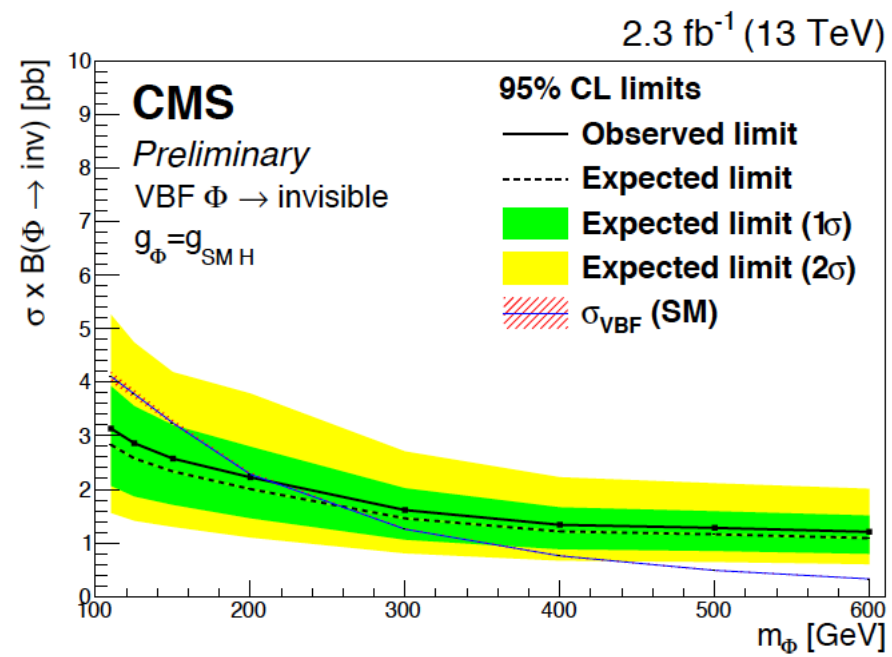
Coupling is set to 1 for all plots



How does it look for our 4 friends?

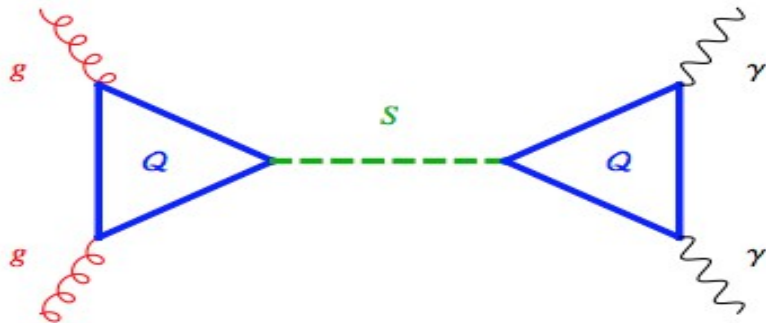
Coupling is set to 1 for all plots



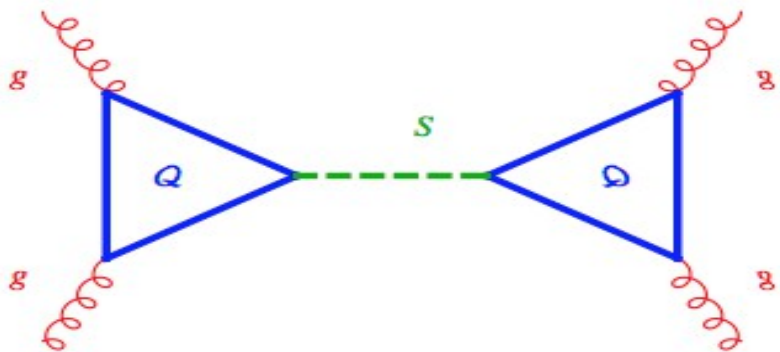


What are the decays

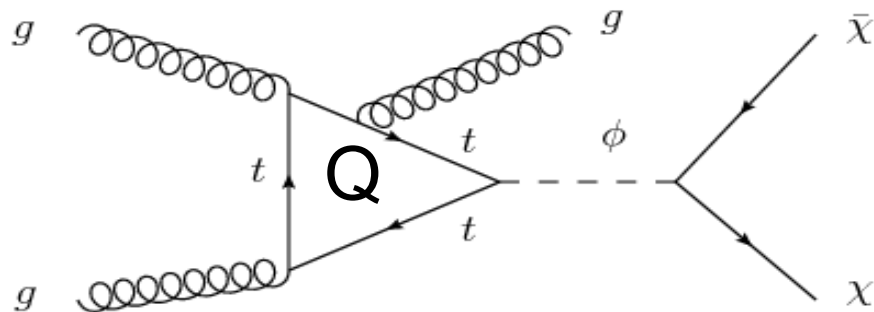
- We only really have a few decays:



Diphoton decays



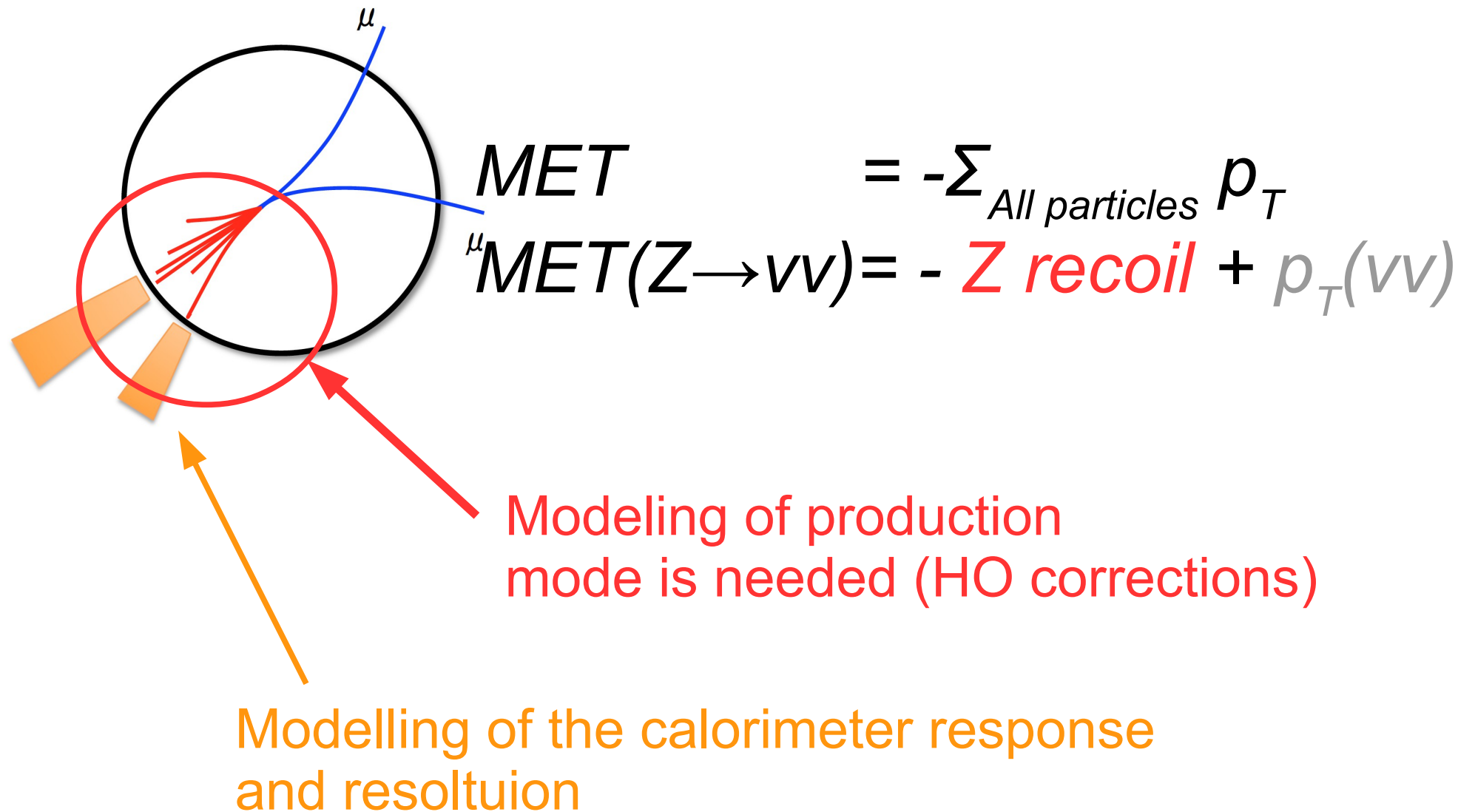
Dijet decays



Monojet decays

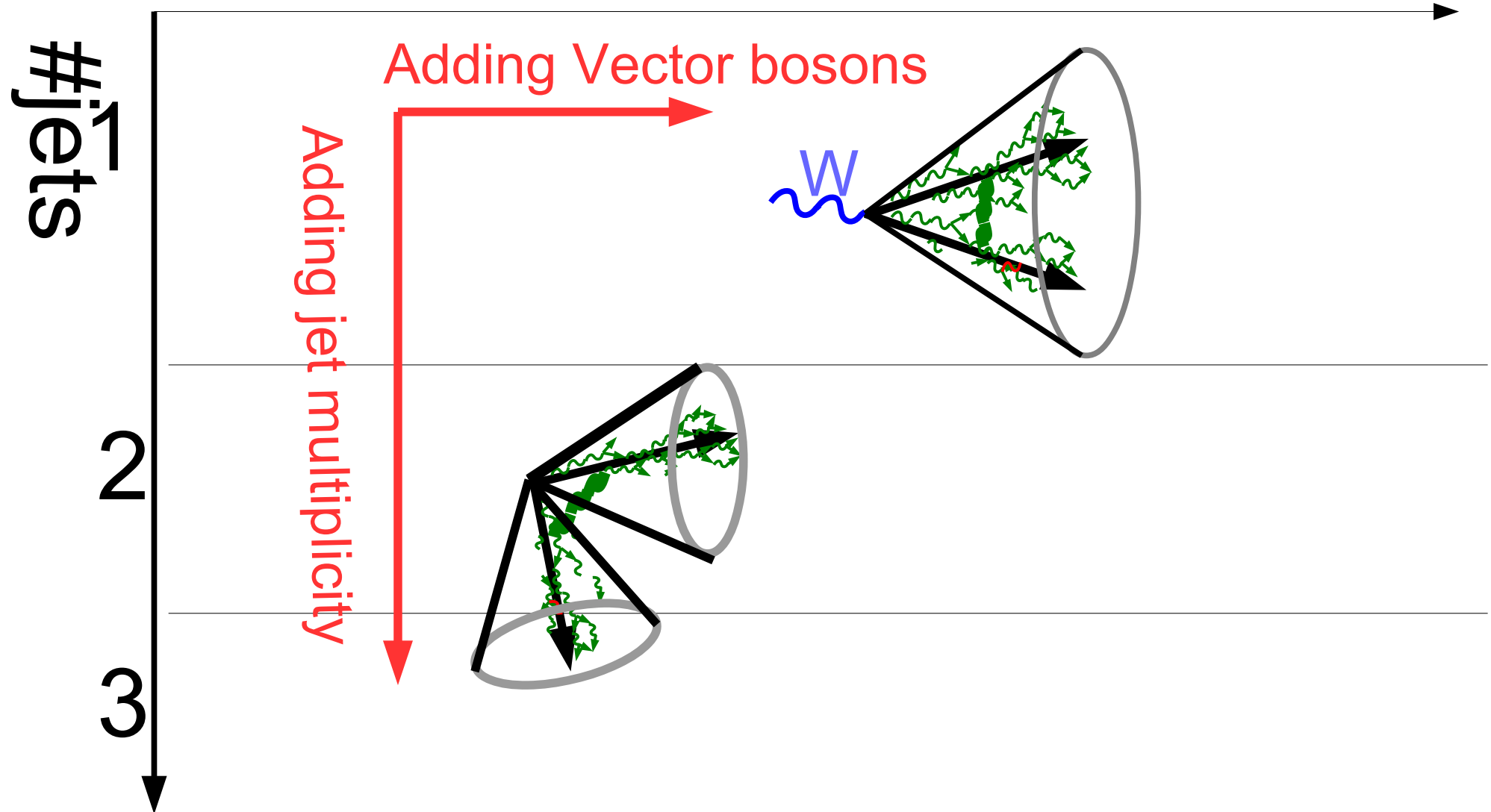
What goes into this?

- To find a signal we look for high MET :



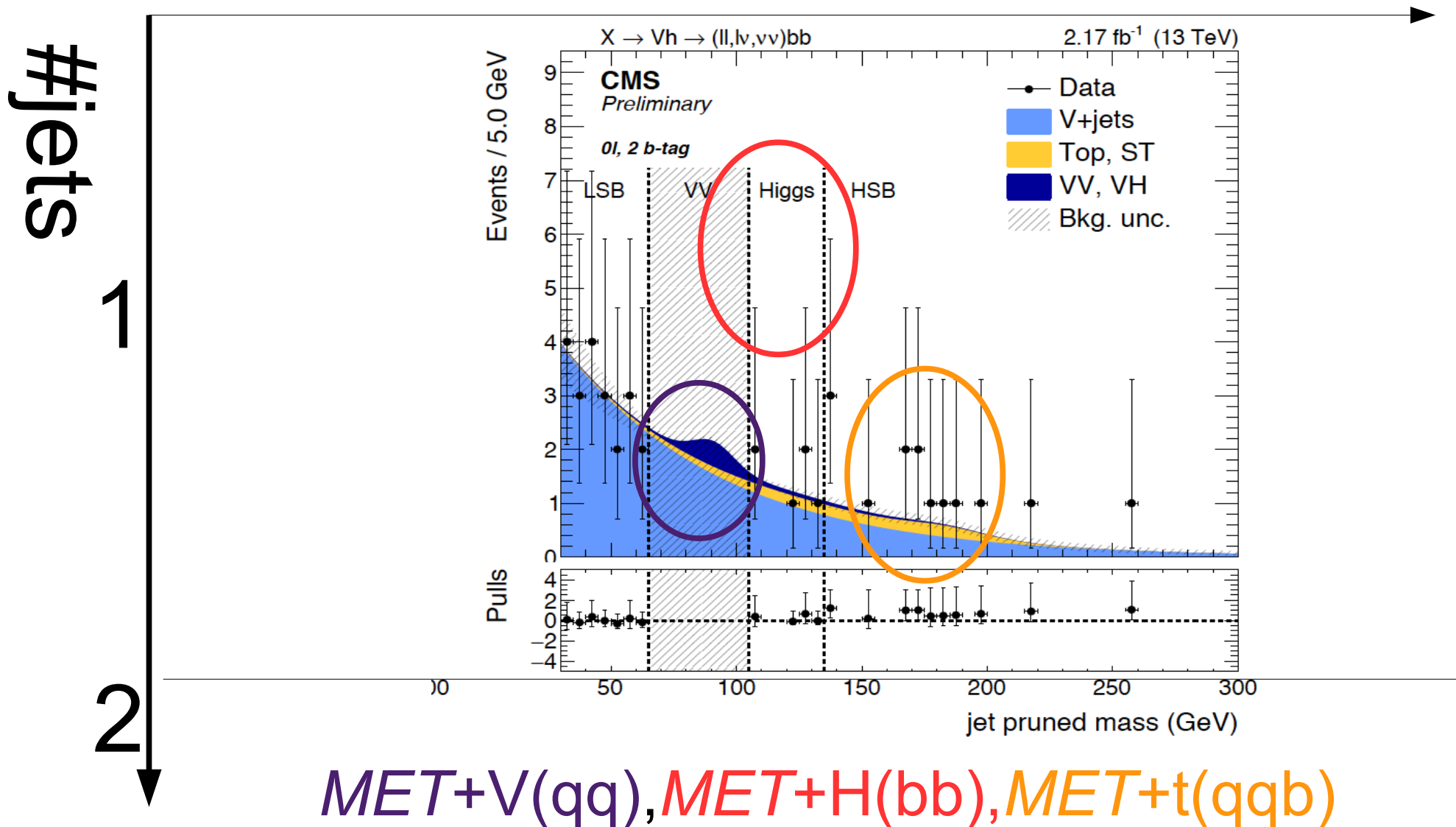
Monojet Extension Plane

Jet Mass



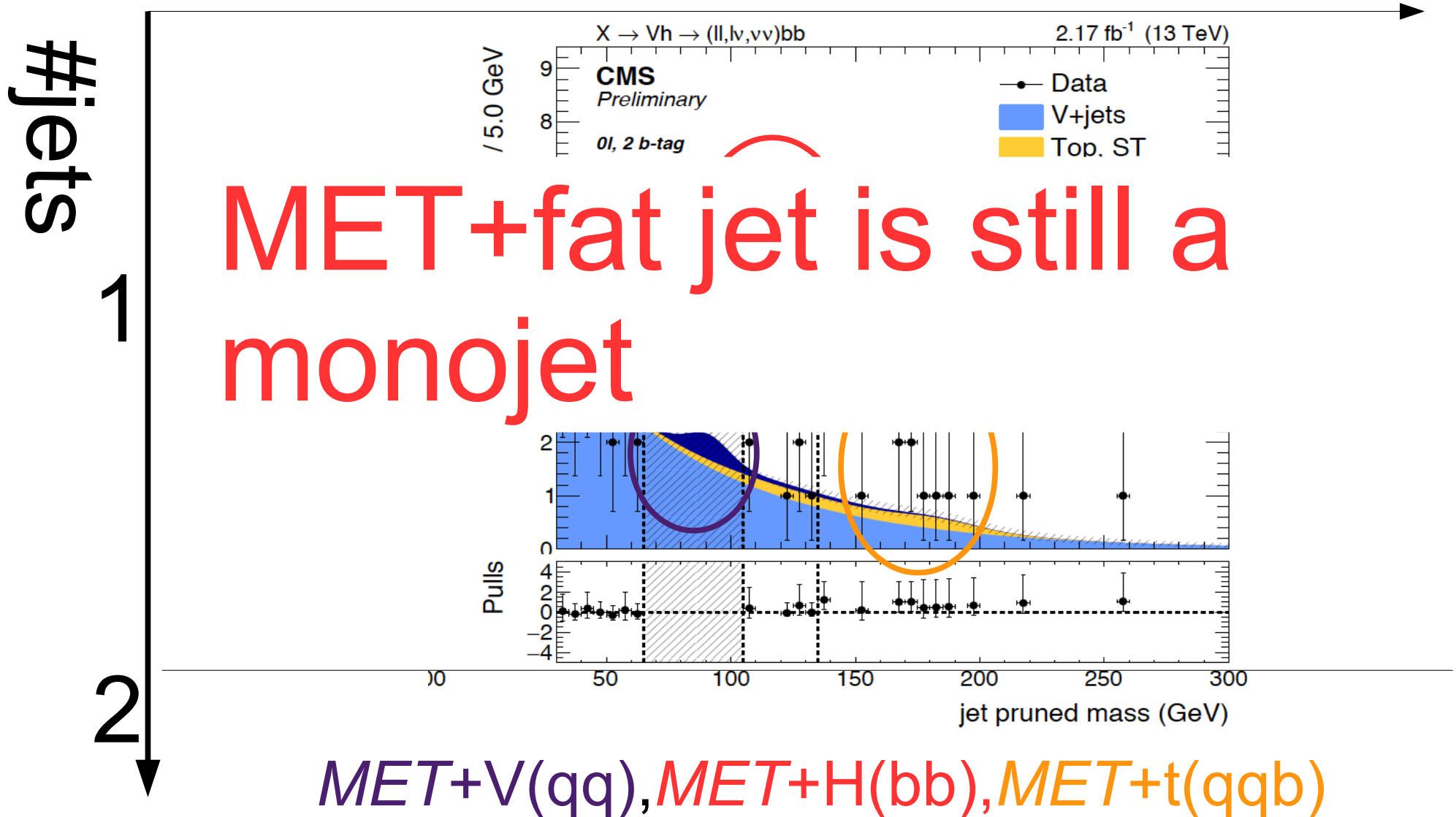
Monojet Extension #1 ($V \rightarrow qq$)

Jet Mass



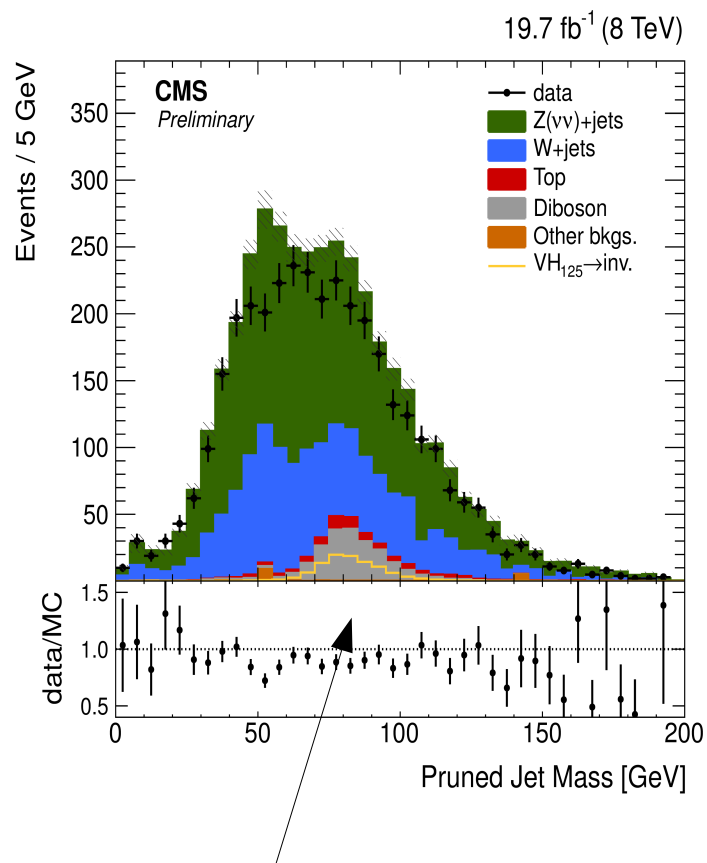
Beyond Monojet

Jet Mass



MET + fat jet

- There is no clean way to separate fat jets from jets



Is there room for improvement?

Yes

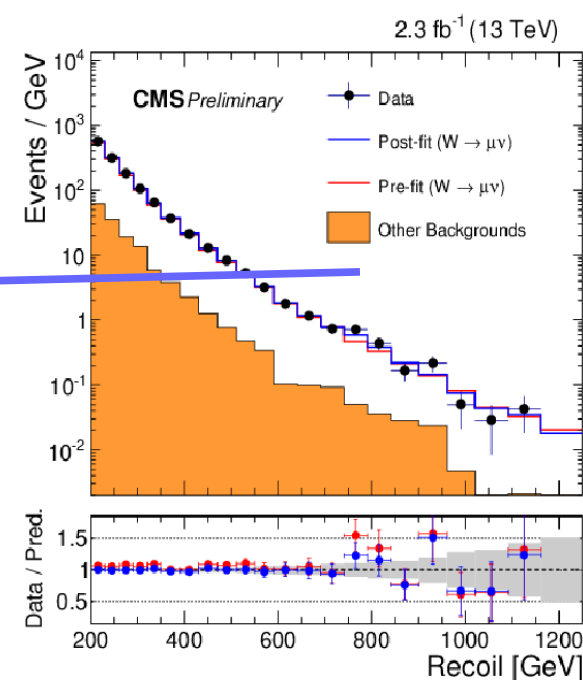
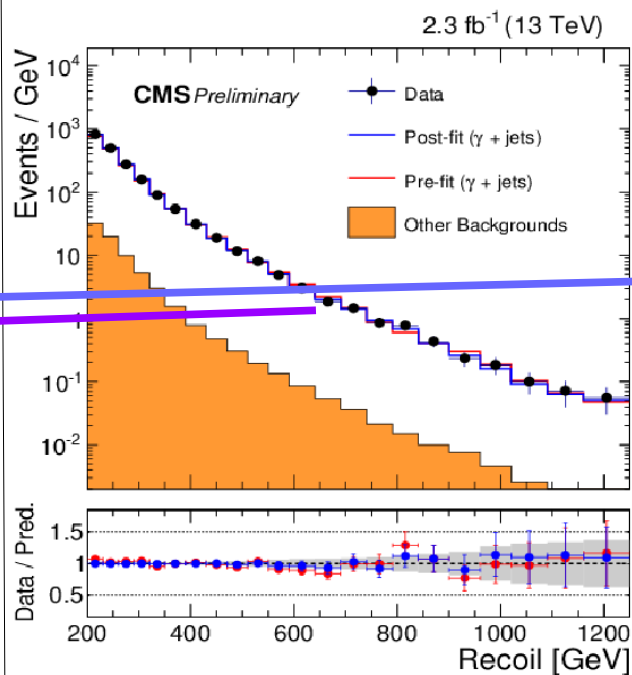
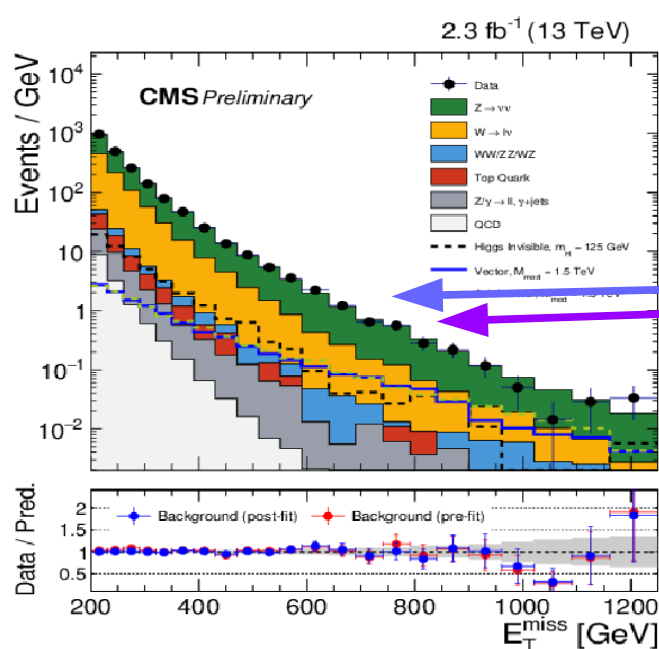
Currently require a simple :
jet mass cut + τ_2/τ_1

Signal

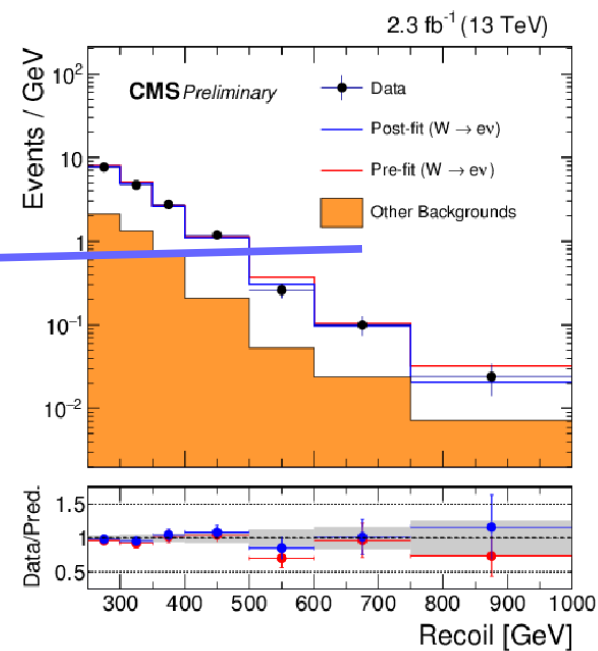
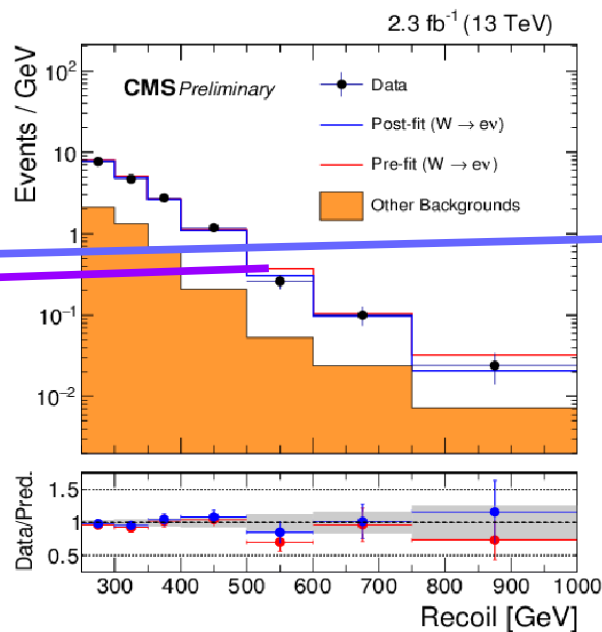
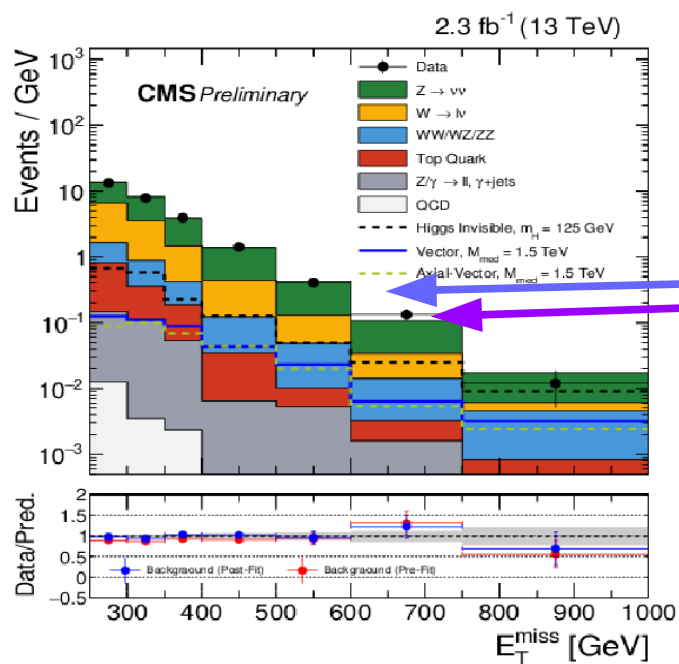
CRs: γ +jets

+ W +

Monojet

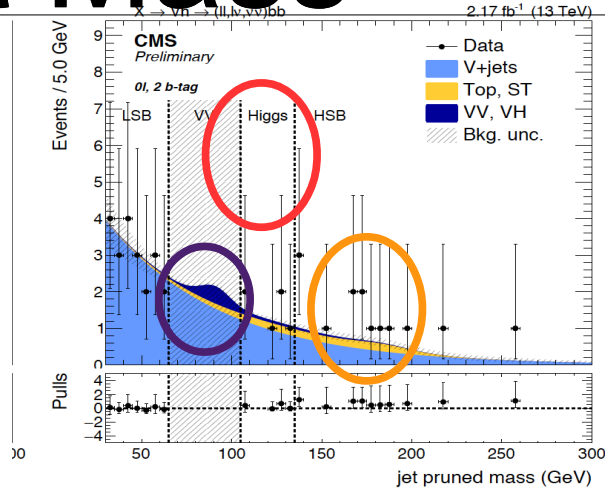


Mono-V



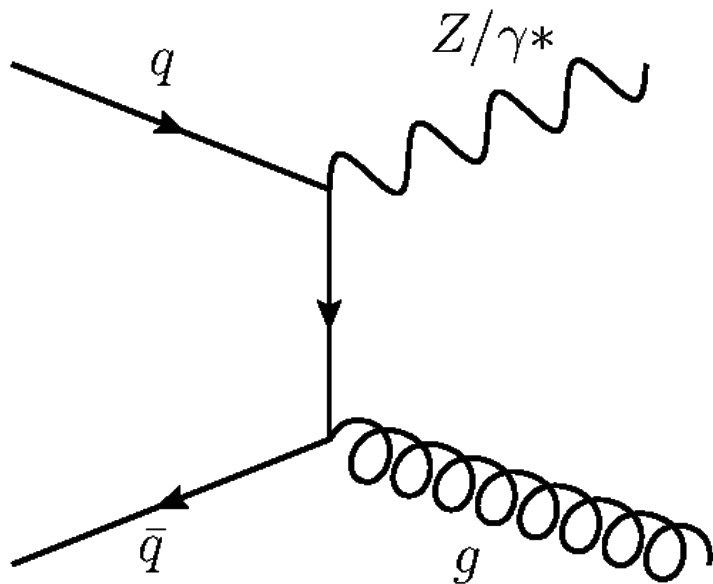
Beyond Monojet

Jet Mass



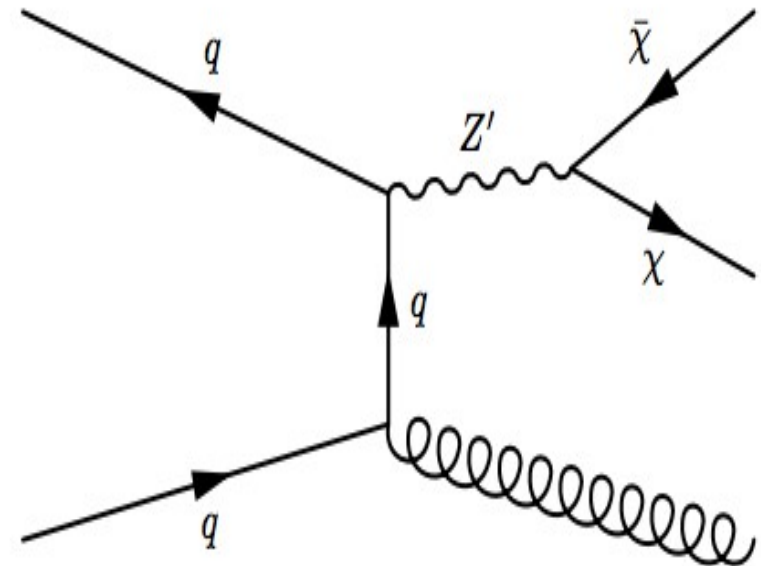
Using the 2nd jet or more
can add to discrimination

Where do we gain from 2nd Jet?



Background

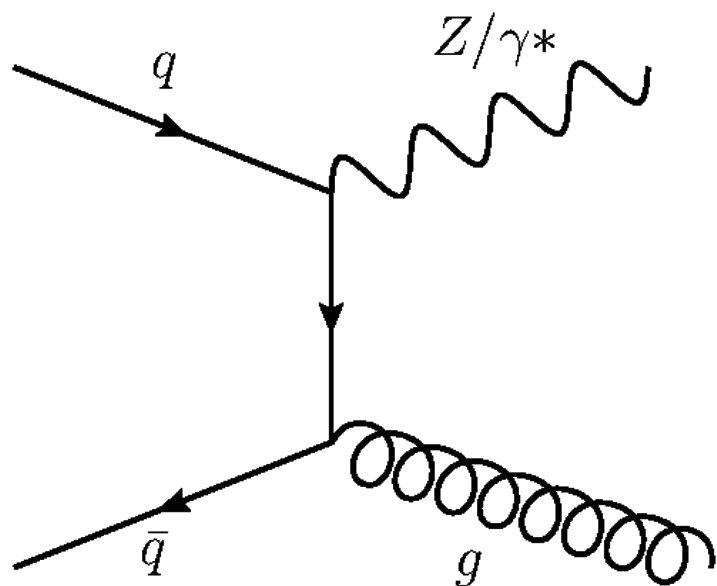
VS



Signal

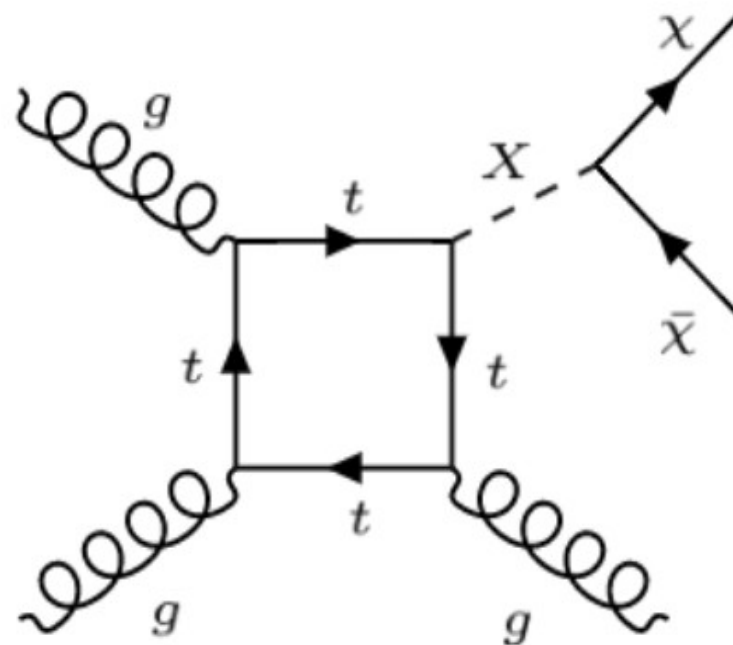
For Vector and Axial mediators not much
Only real difference is mediator mass

Where do we gain from 2nd Jet?



Background

VS

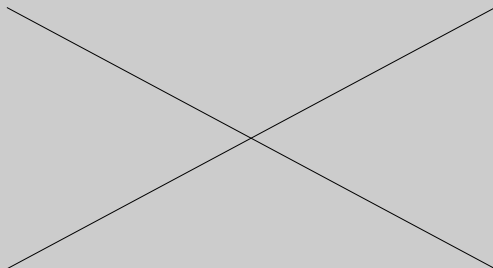


Signal

For Scalar and Pseudoscalar mediators more
 Now the production modes are different

In addition to 2nd jet can consider a quarkgluon discriminator

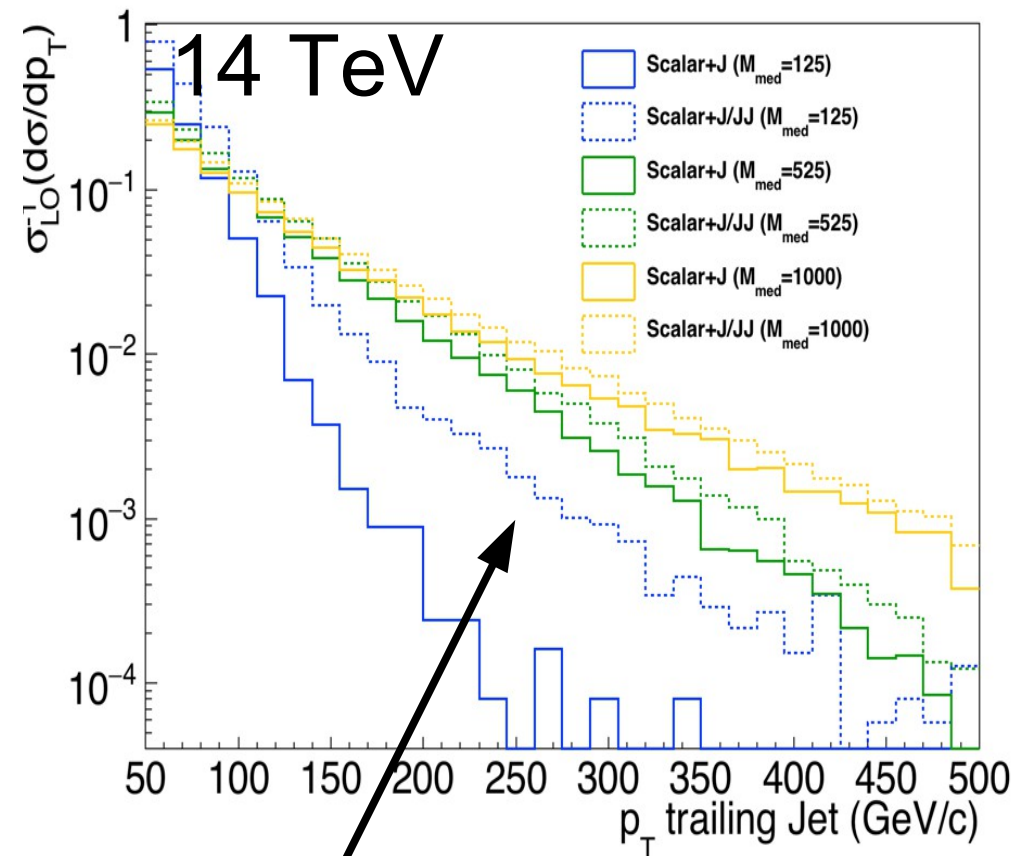
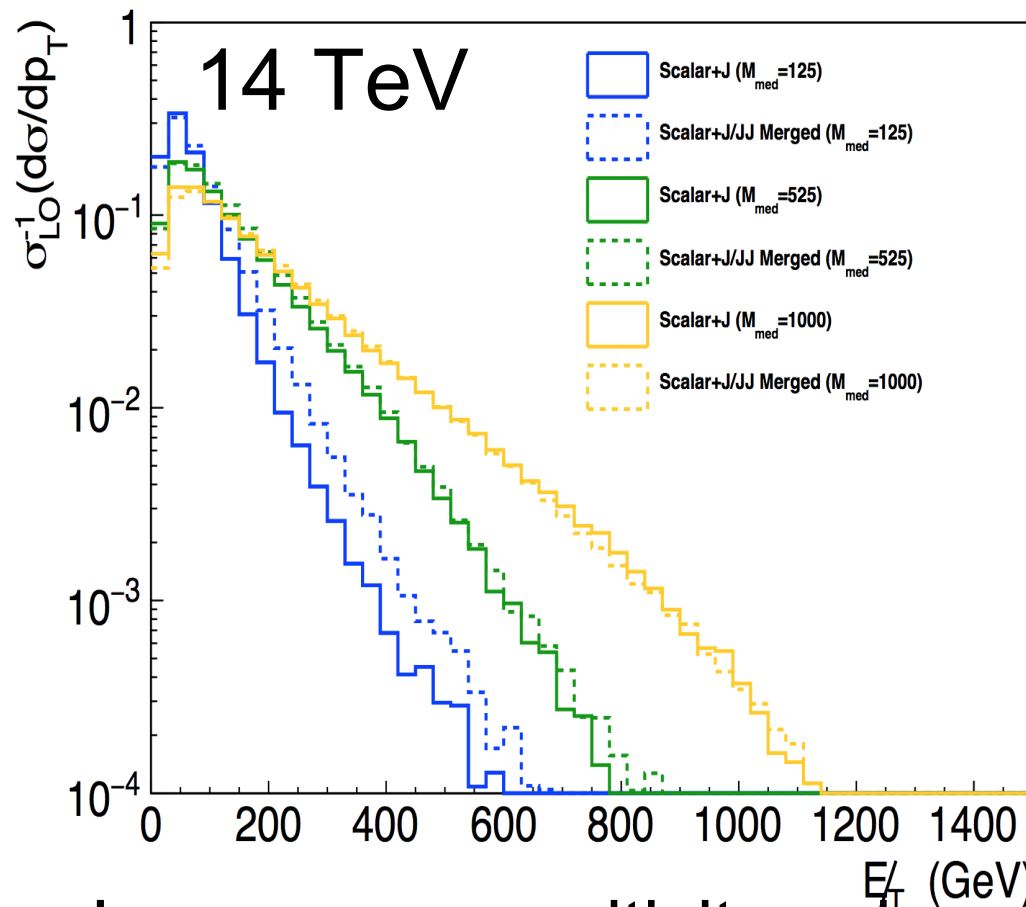
Spectrum of Signal MCs

Sample	LO/ Leading Loop	LO in 2j	NLO,1,2j
Vector/ Axial	Madgraph MCFM	Powheg	aMC@NLO
Scalar/ Pseudoscalar	Powheg MCFM aMC@NLO	VBF@NLO aMC@NLO	

aMC@NLO+MG get highest order 1/2 jets merged

Advantage of merged MC

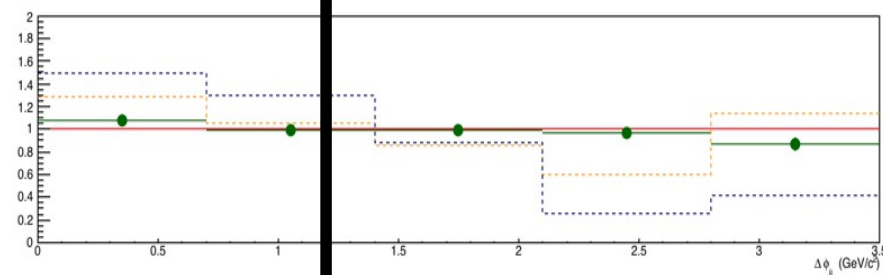
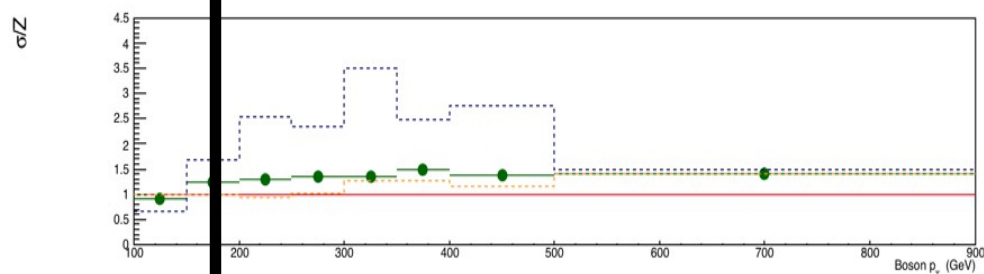
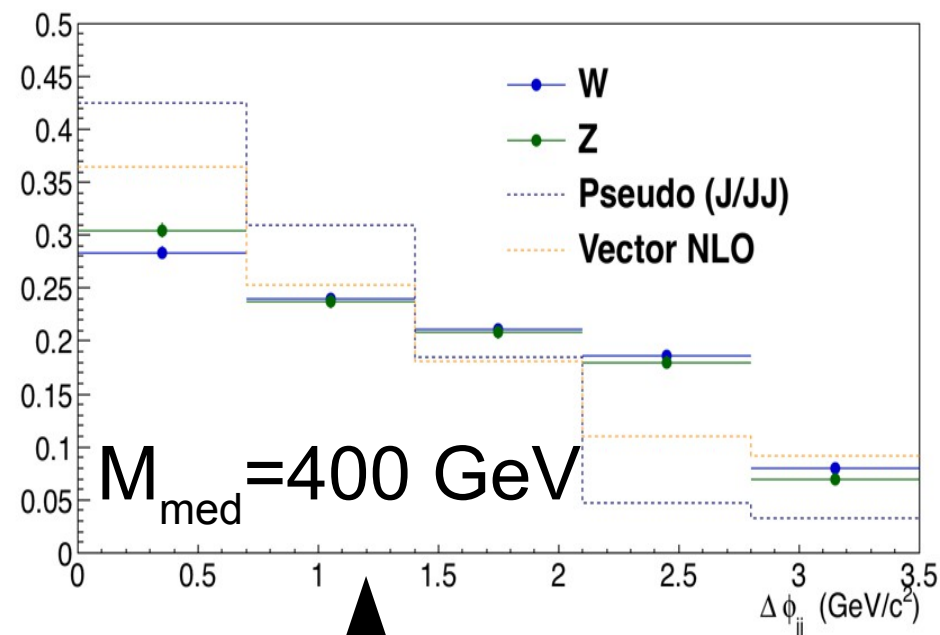
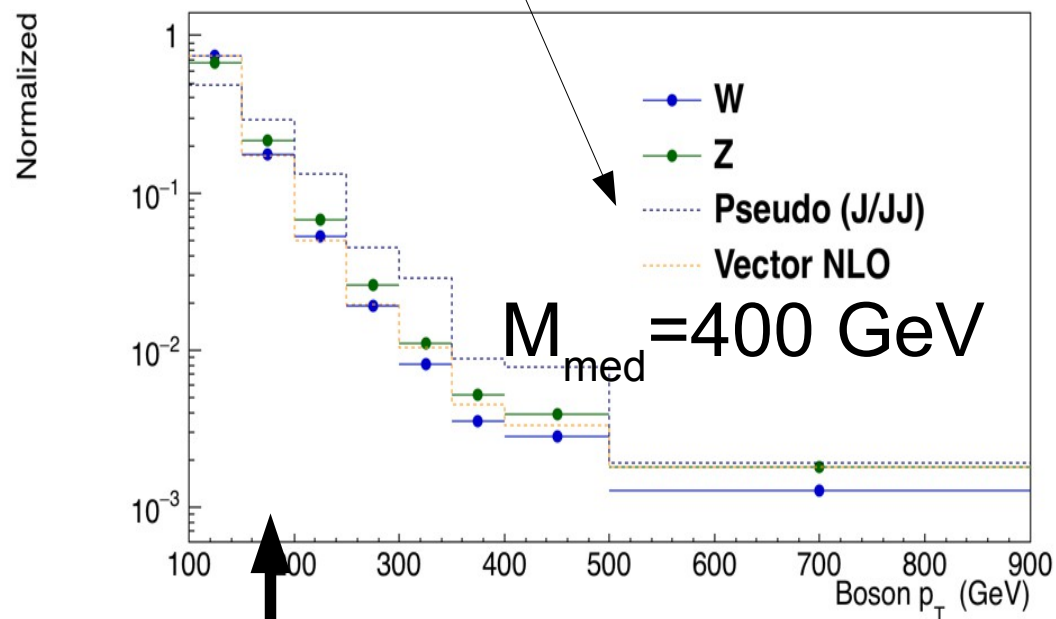
- Taking advantage of the new technology
 - Can consider exploring new regions of phase space



Low mass sensitivity enhanced in the multi jet final state

FYI aMC@NLO
merged 0,1,2jet
pseudoscalar

Basic Concept of Gains

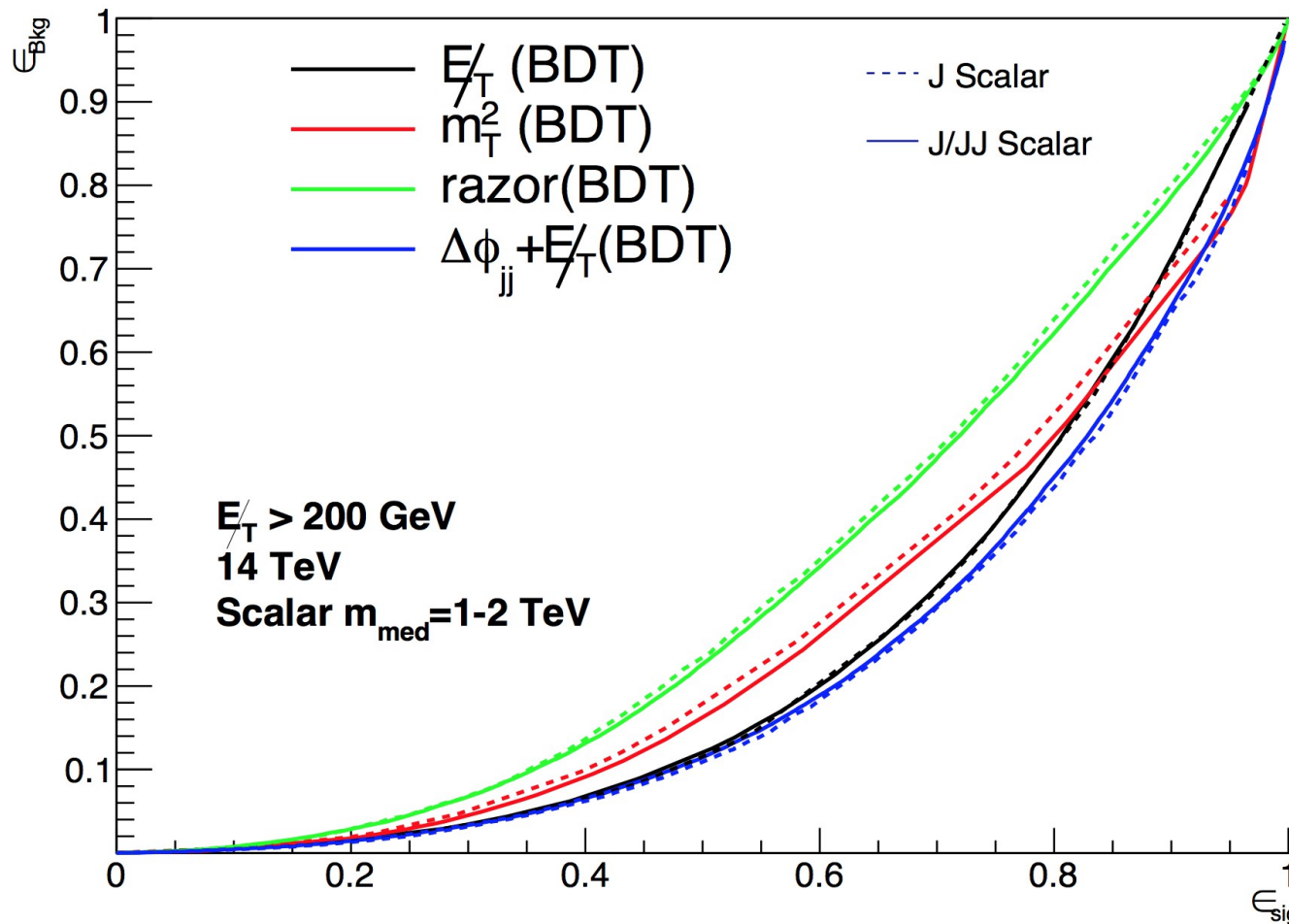


Gluon fusion induces
Higher p_T spectrum

Heavier mediator forces
jets to be closer

How do the single variables perform?

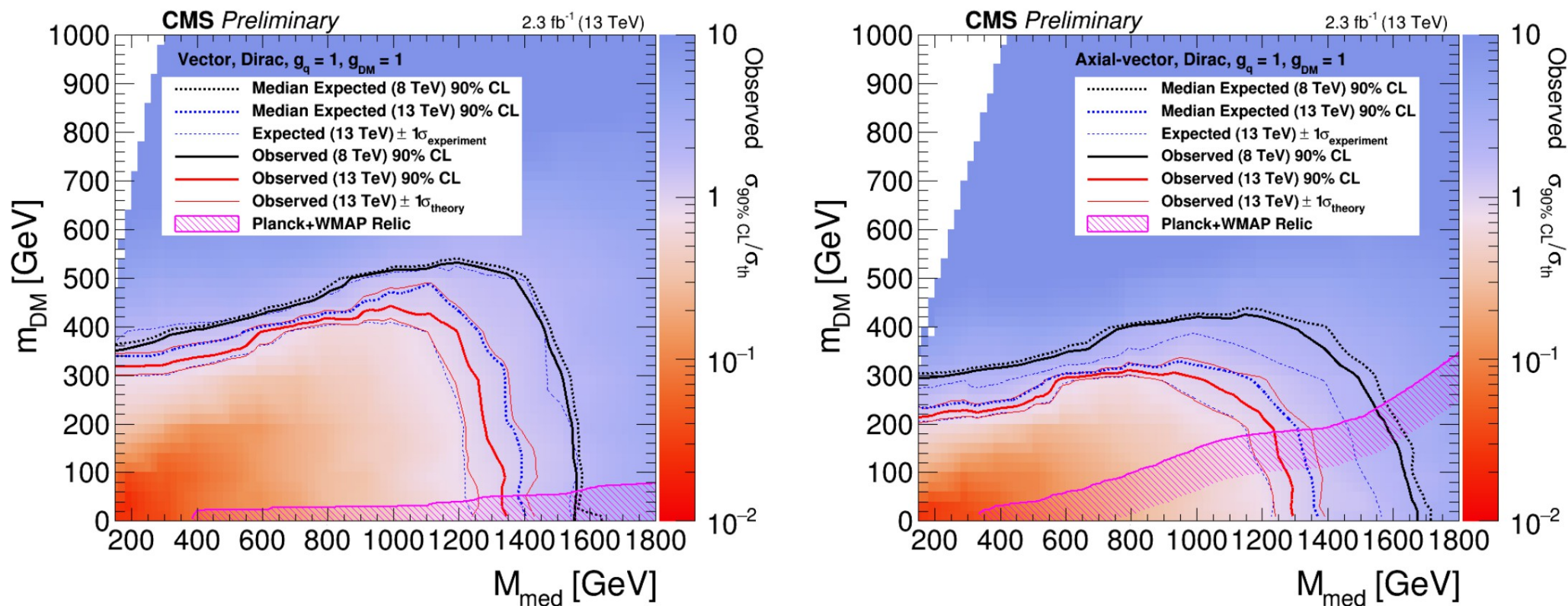
- Comparison of single variables



Gain comes from fact that light mass objects have collinear jets
 Using $\Delta\phi_{jj}$ can bring as much as 20% gain

Results

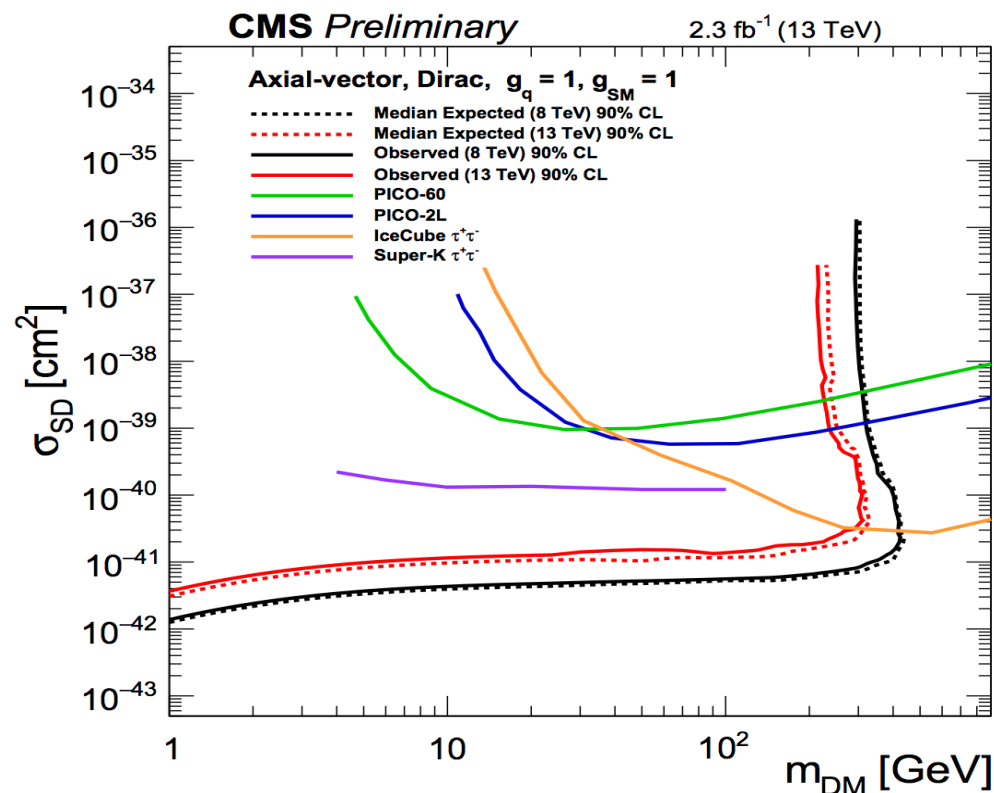
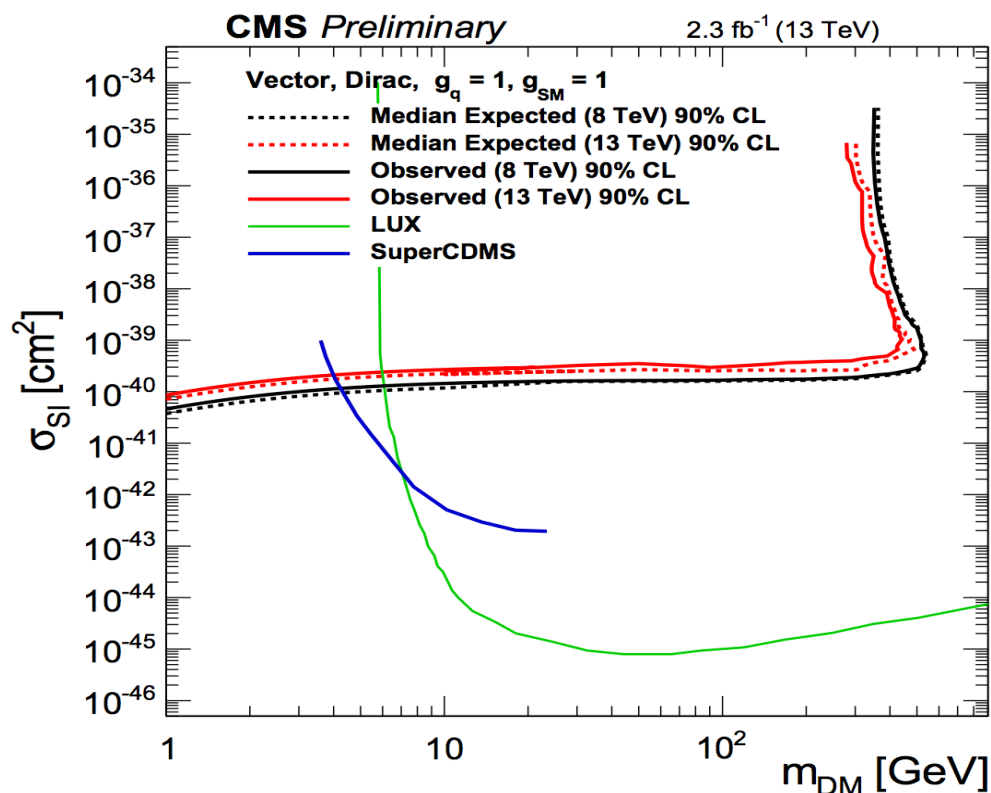
Our Current Public Results



Both 13 TeV and 8 TeV analysis treat:
mono-V and monojet on equal footing

An $1-2\sigma$ excess is present in both data sets in tail

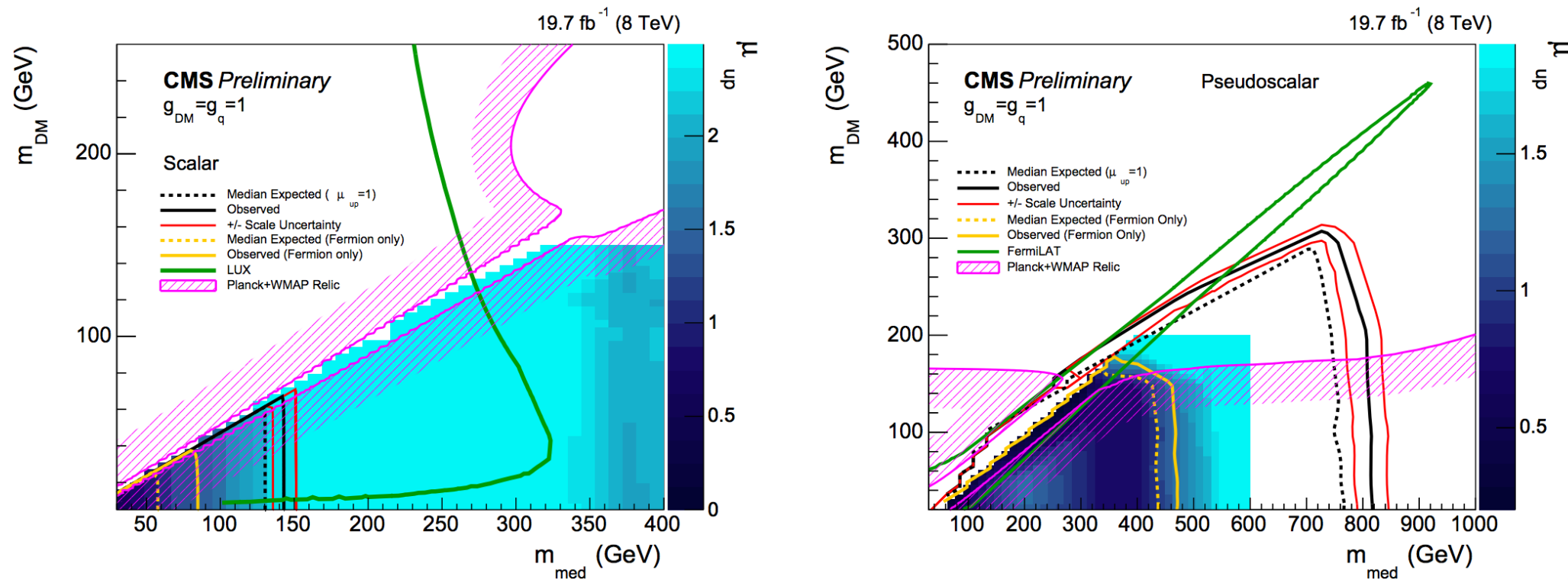
Our Current Public Results



Translation to direct detection now standardized

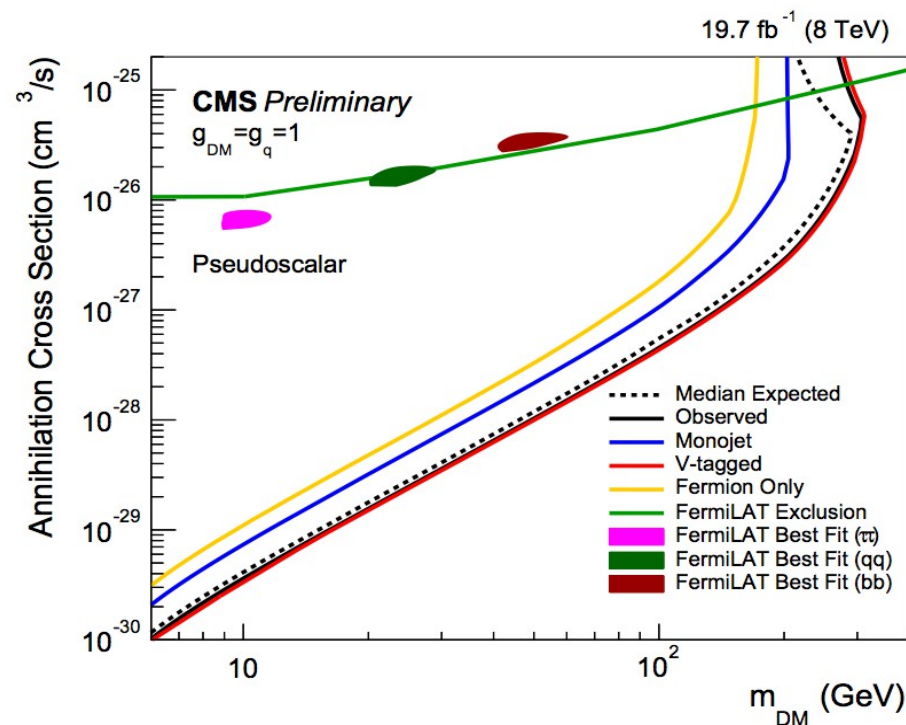
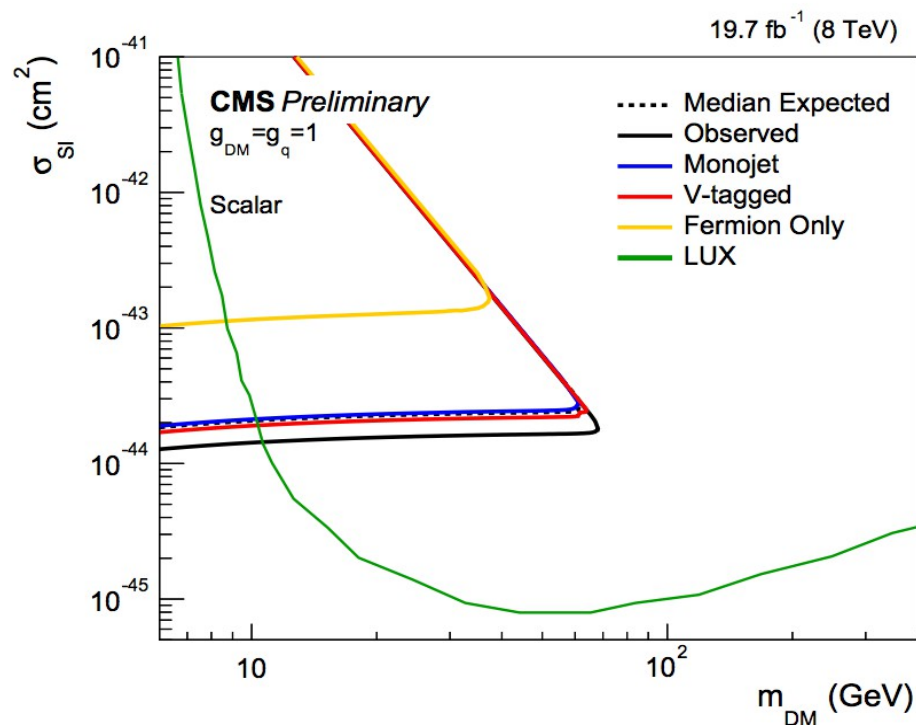
An 1-2 σ excess is present in both data sets in tail

Our Current Scalar & Psuedo results



- Currently only have 8TeV exclusion
 - Yellow line : Official simplified models
 - Black/Red (controversial) : Simplified + EWSB
 - Allows us to add Higgsstrahlung

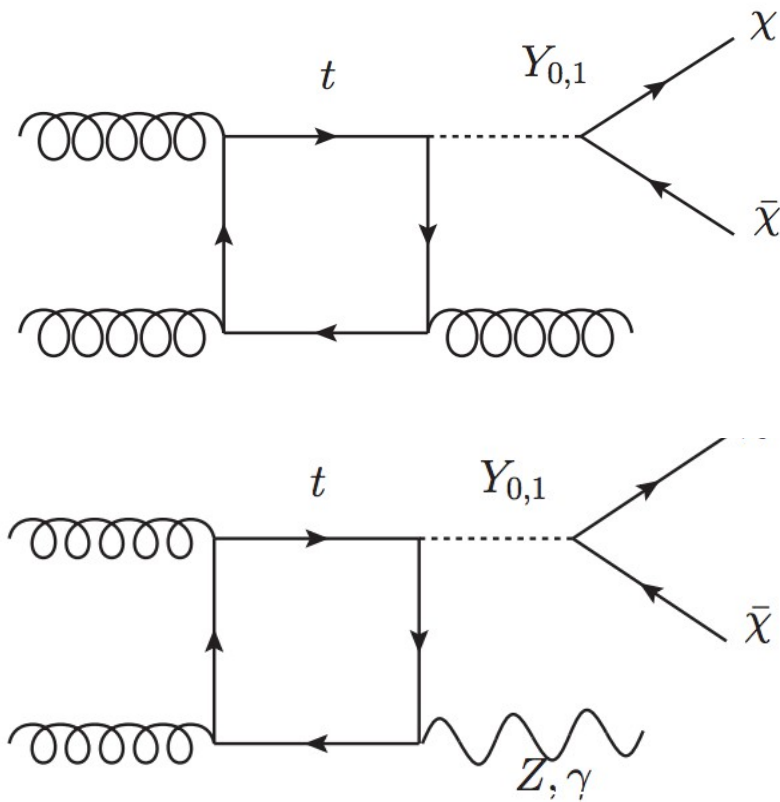
Our Current Scalar & Psuedo results



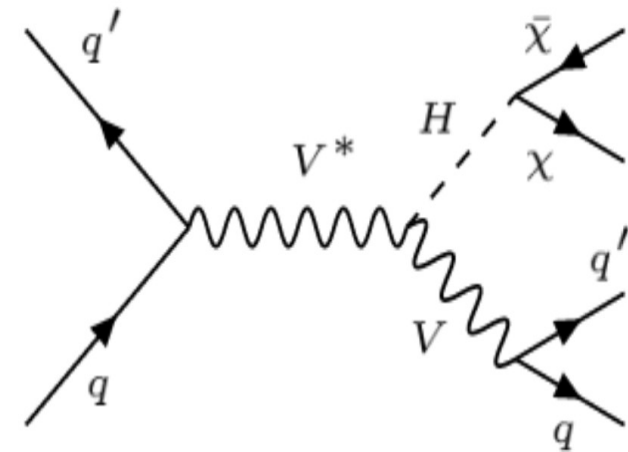
- When the dark matter is not onshell
 - Strong exclusion of pseudoscalar interpretation of LAT
 - Scalar and Direct detection are in close competition
 - Expect LHC to pass LUX this summer winter!

To break or not to break?

- EWK symmetry breaking adds lots of mono-V
 - Contribution can be very significant if pseudoscalar
- There are models that do that (e.g. 2HDM...)
 - Need physics at a higher scale (dim-7 operator)

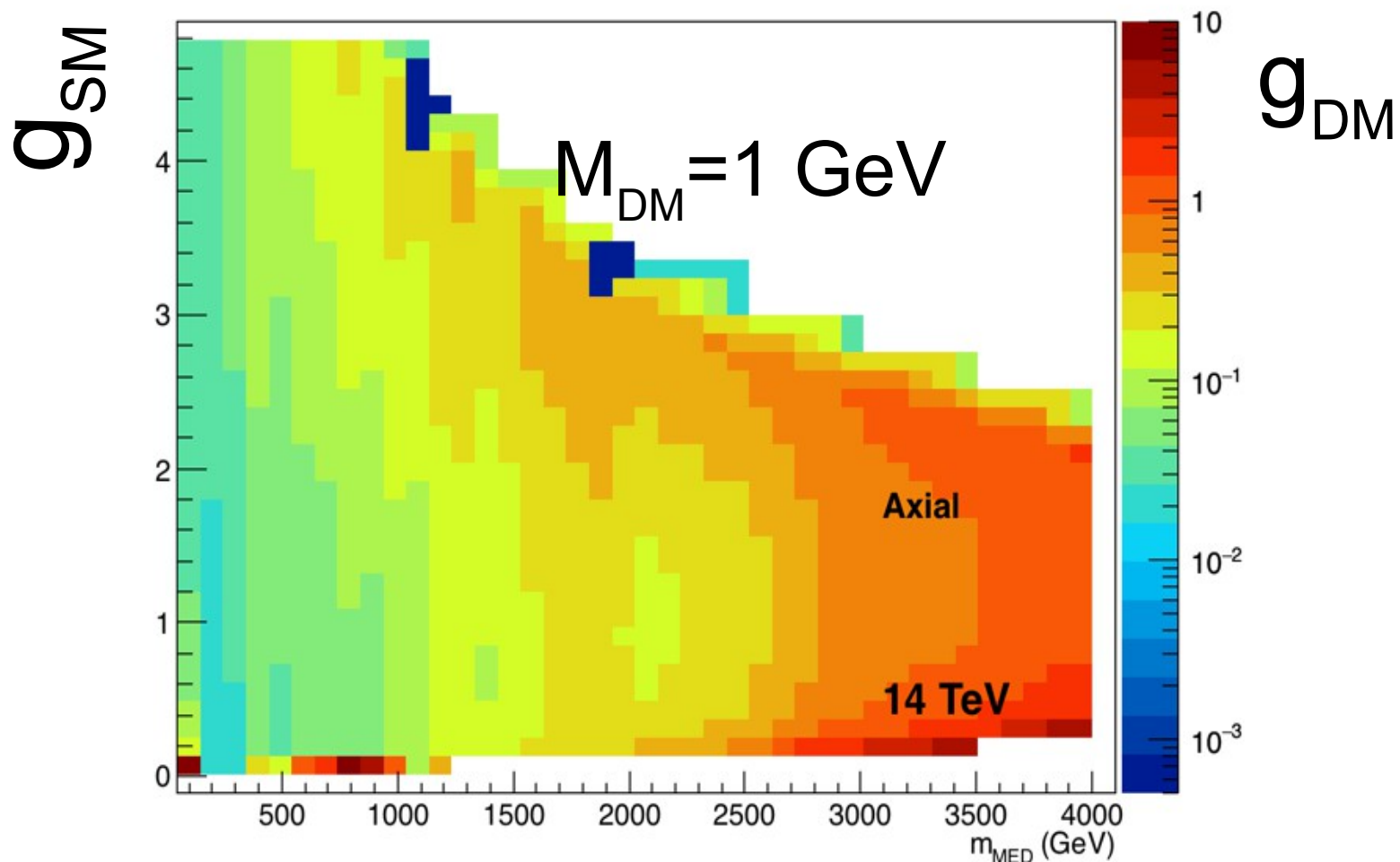


+?



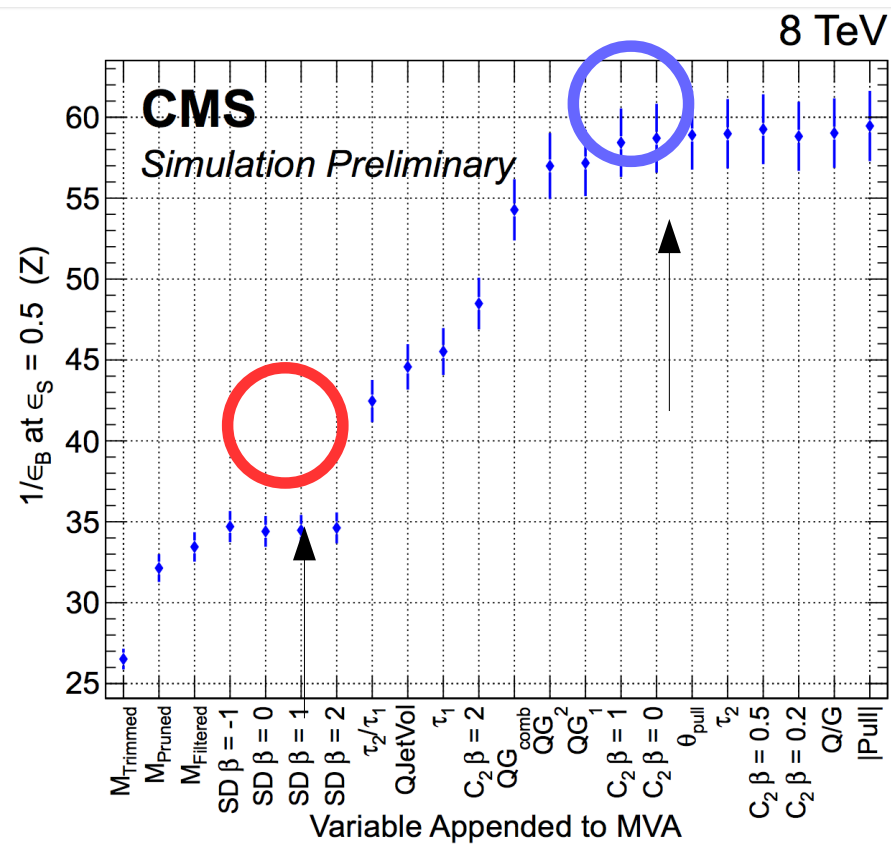
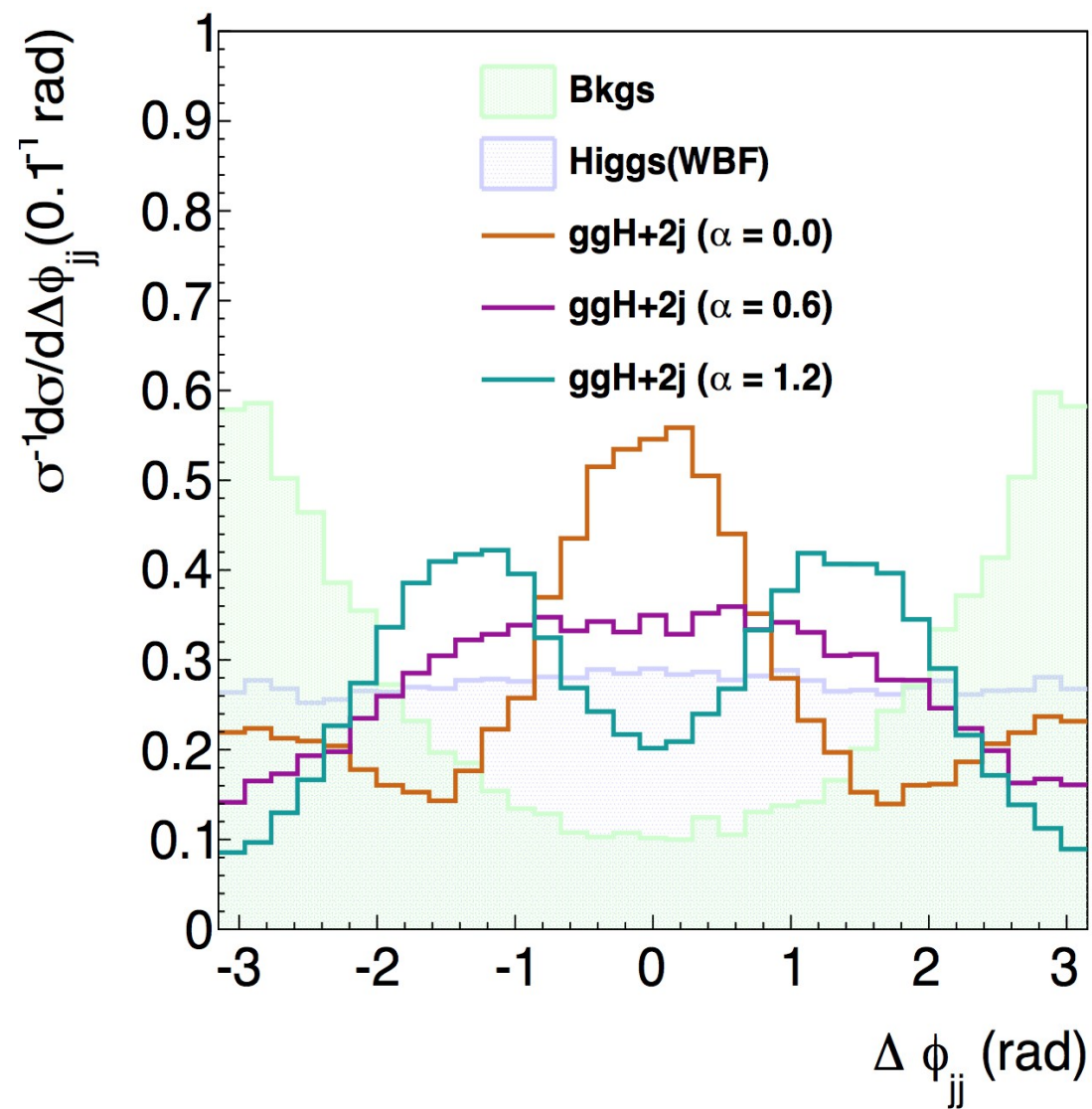
MIXING!?

Extending Our results



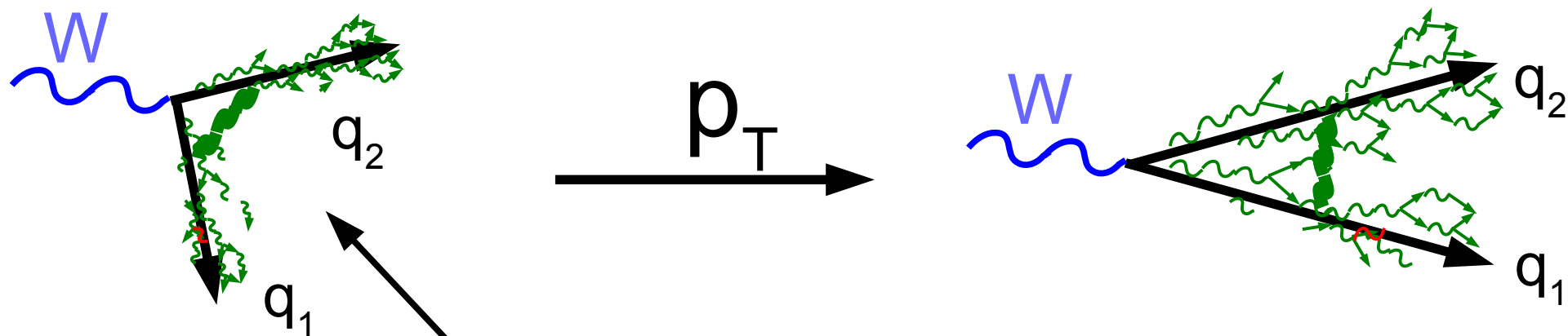
Is it time to consider varying the couplings?

$$\mathcal{L}_{\text{axial}} \supset \frac{1}{2} m_{\text{MED}}^2 Z_{\mu}'' Z''^{\mu} - g_{\text{DM}} Z_{\mu}'' \bar{\chi} \gamma^{\mu} \gamma^5 \chi - \sum_q g_{\text{SM}}^q Z_{\mu}'' \bar{q} \gamma^{\mu} \gamma^5 q.$$



Building a V-tagger

- Evolution of effects



At low p_T “resolve” two jets : resolved tag

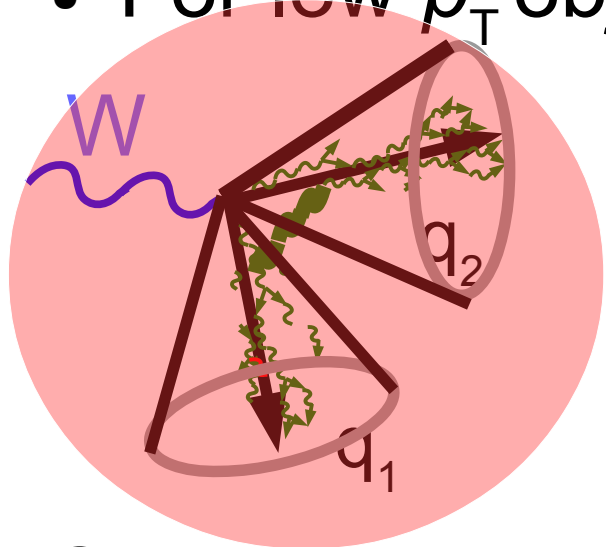
Focus on identifying two jets like a W

At high p_T we obtain one big jet

Focus on identifying one jet like a W

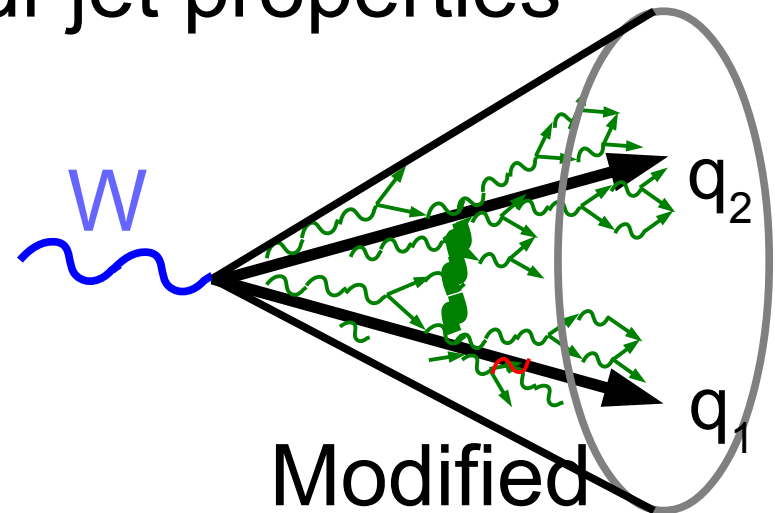
Resolved Tagger

- For low p_T objects focus on di-jet properties



Jet Quark

Gluon likelihood



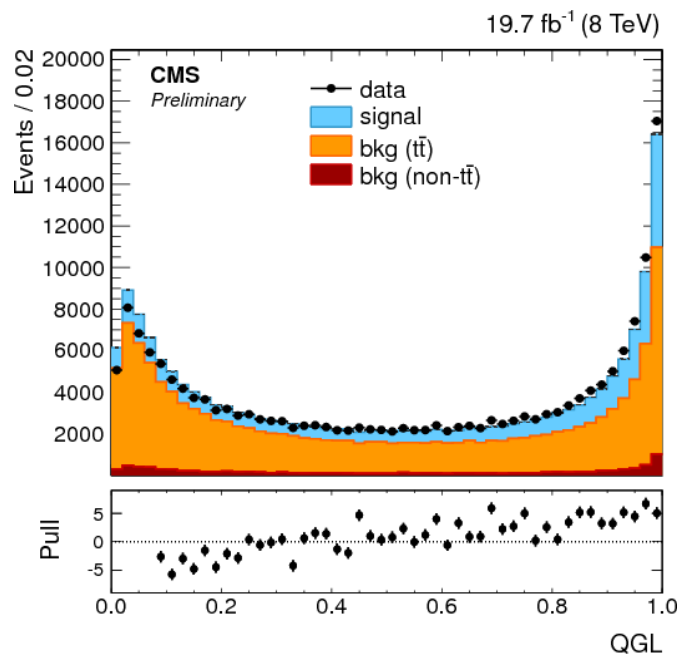
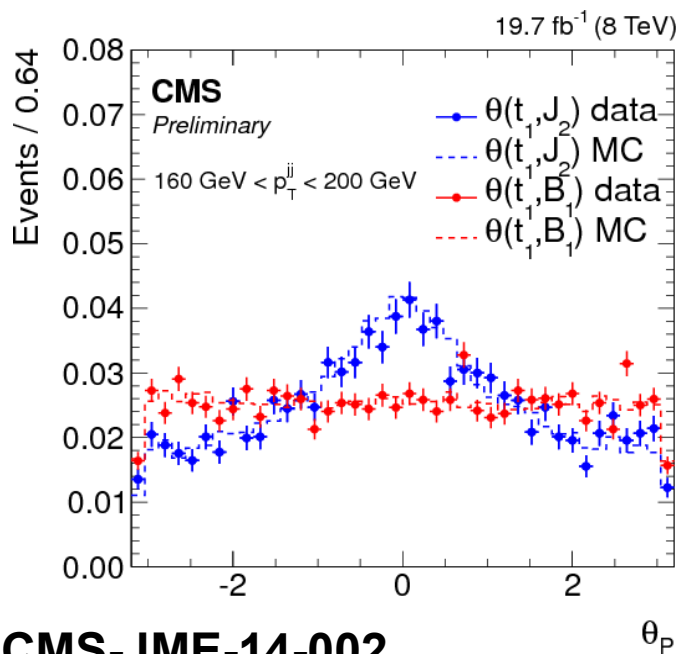
Modified

+ Mass drop

Color Flow +

Jet Quark
Gluon likelihood

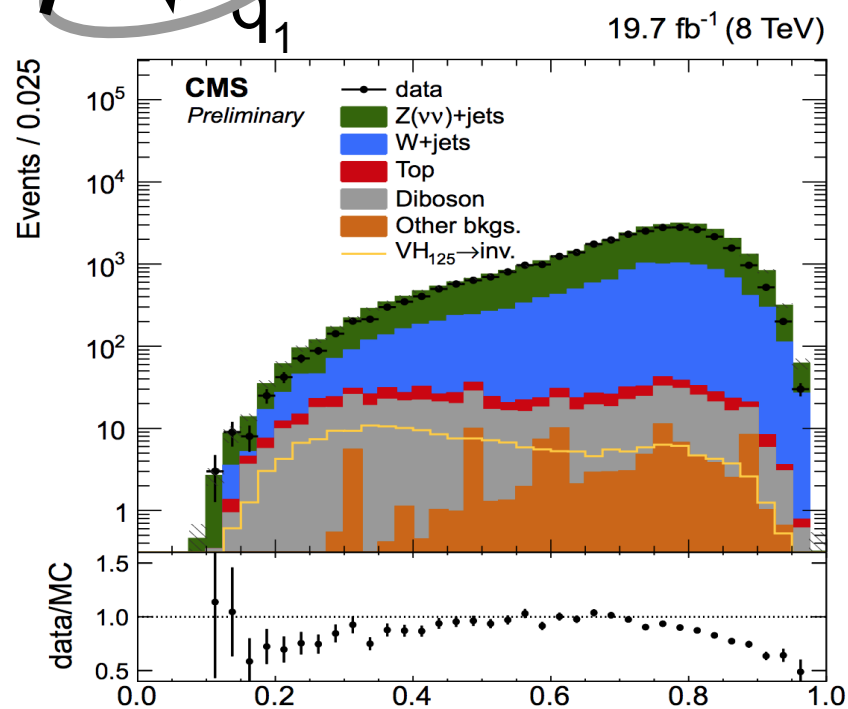
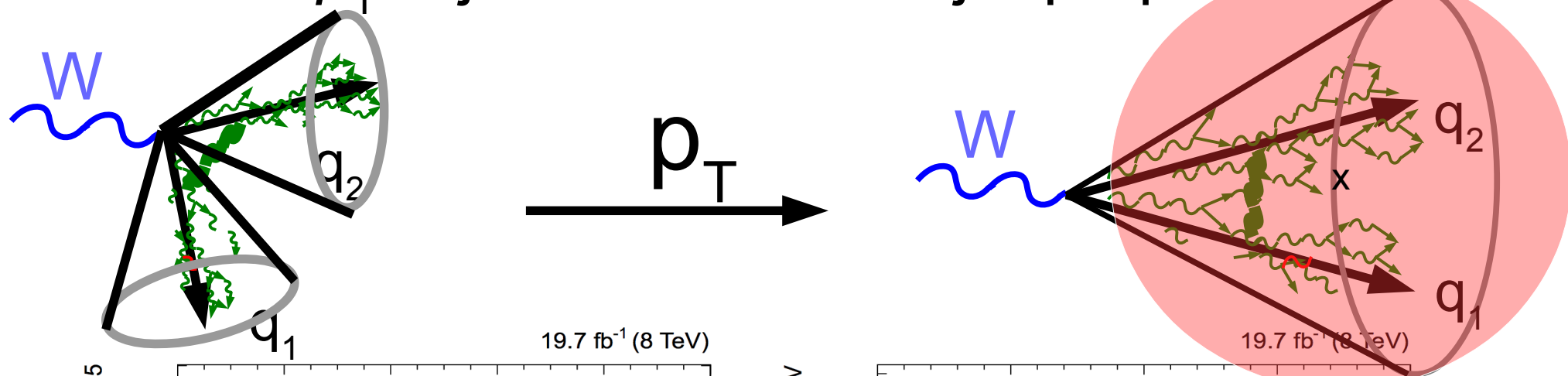
+ Mass drop



$$\zeta = \Delta R_{12} \cdot \max(m_{j1}, m_{j2}) / m_{jj}$$

Boosted Tagger

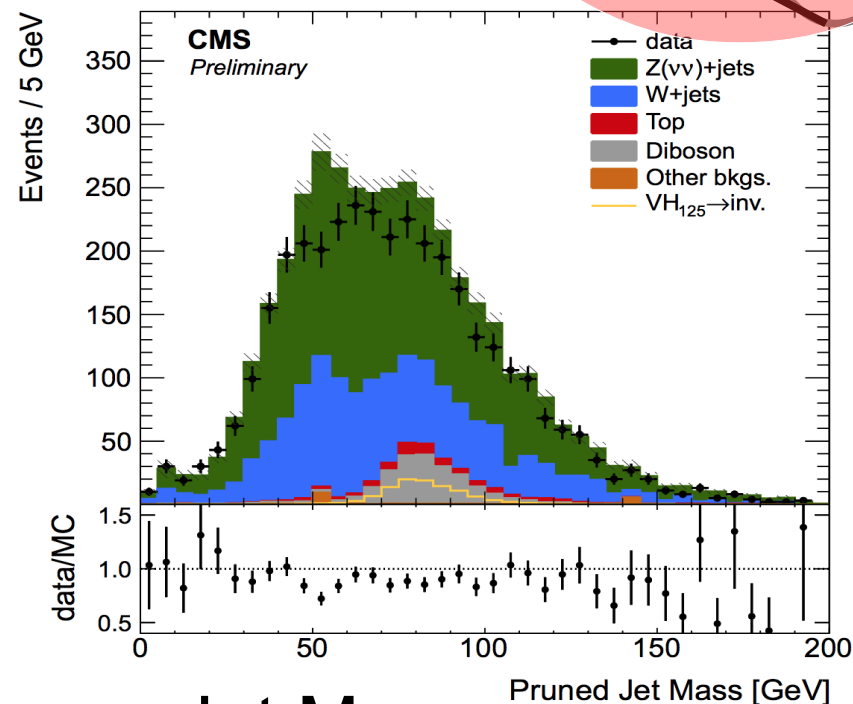
- For low p_T objects focus on di-jet properties



Likelihood for 2 prongs

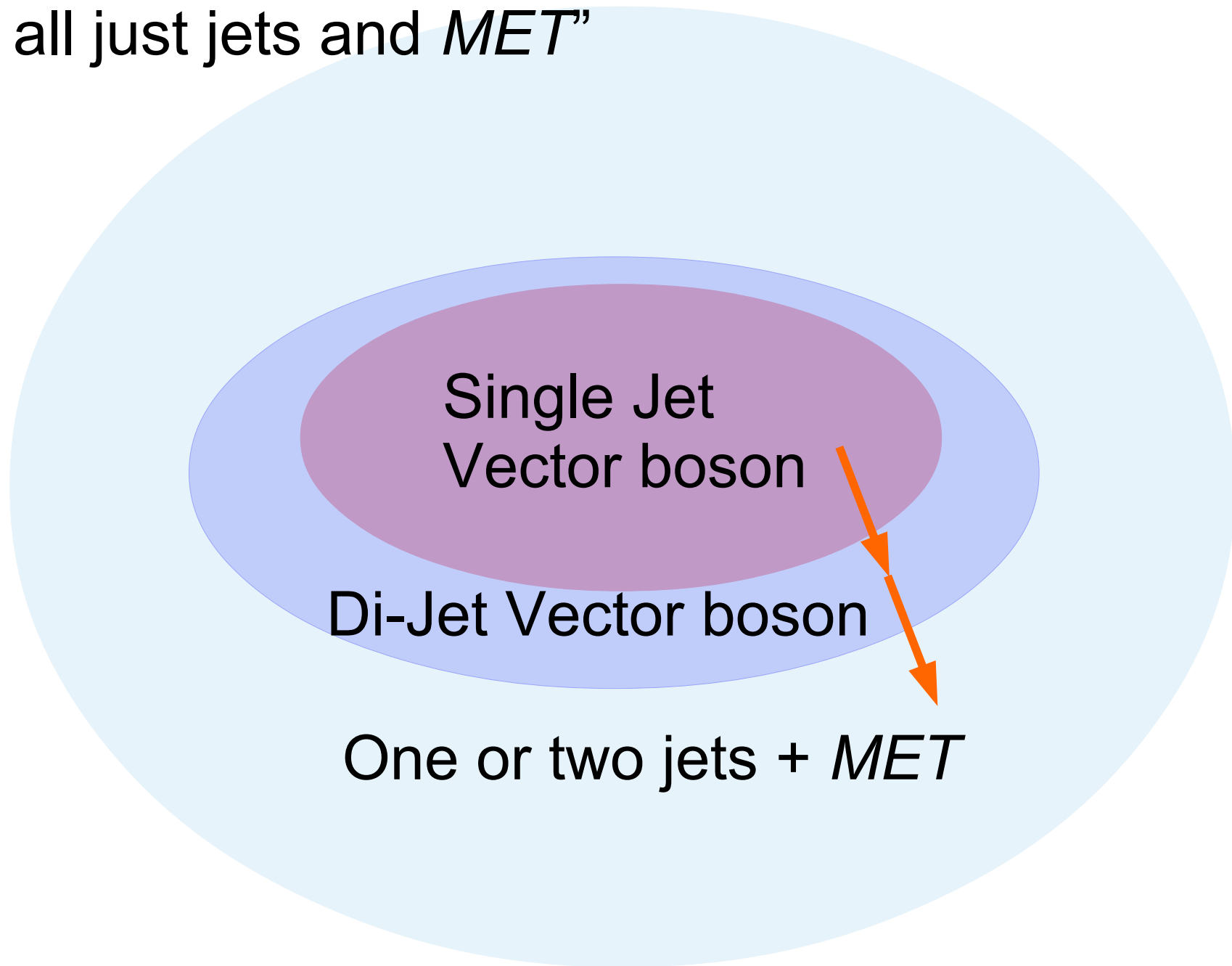
+

Jet Mass



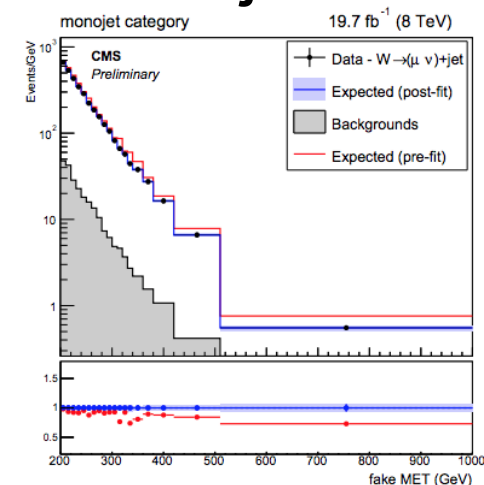
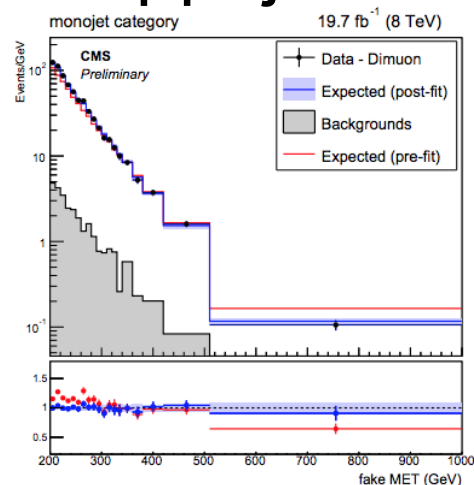
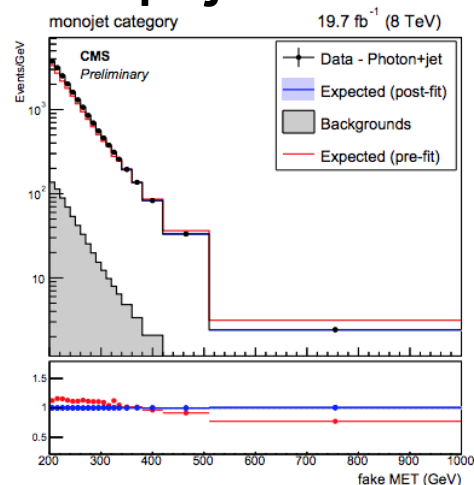
One Big Analysis

“Its all just jets and *MET*”

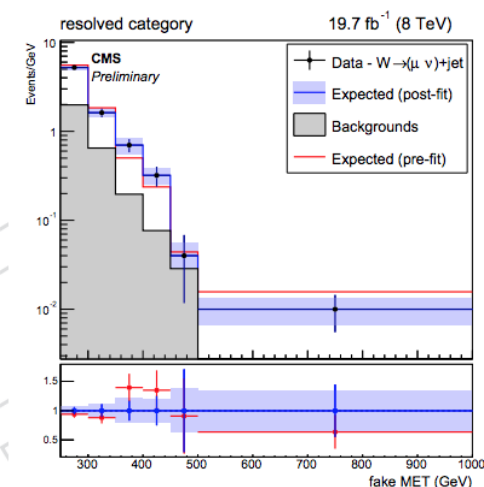
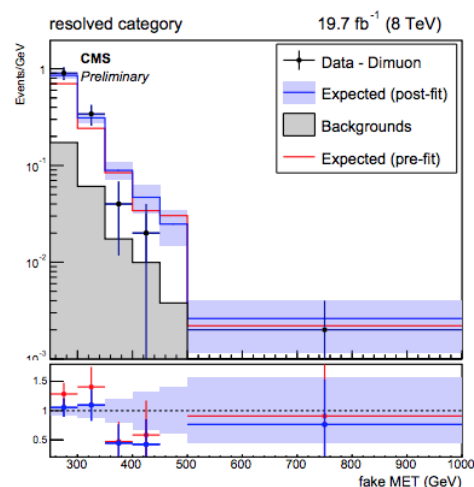
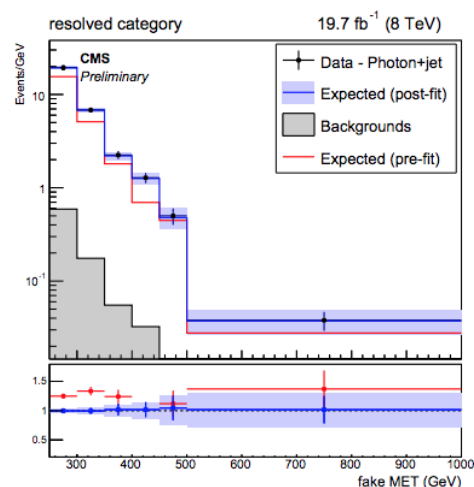


γ +jets $Z\mu\mu$ +jets $Wl\nu$ +jets

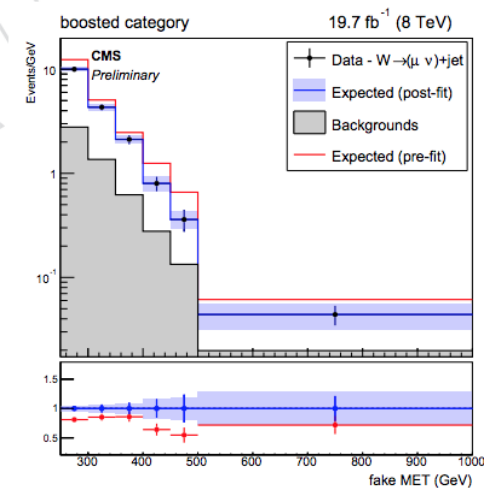
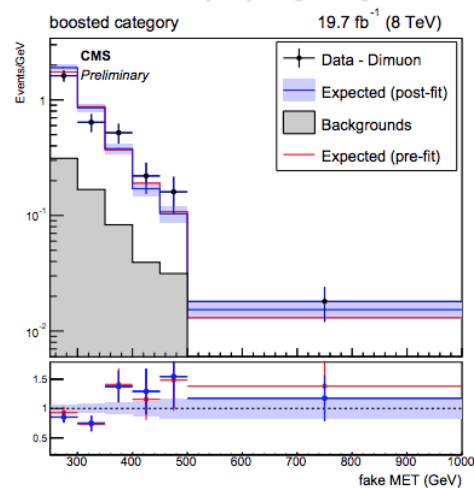
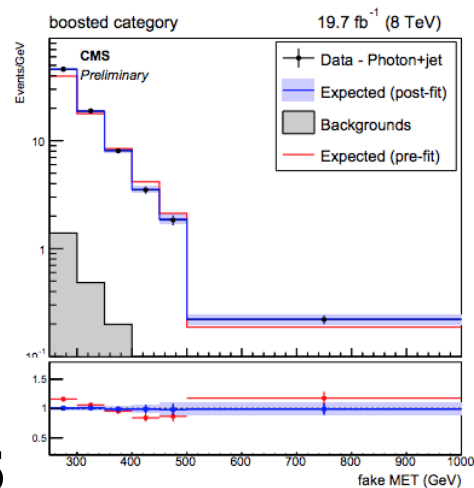
Monojet category



Di-jet category



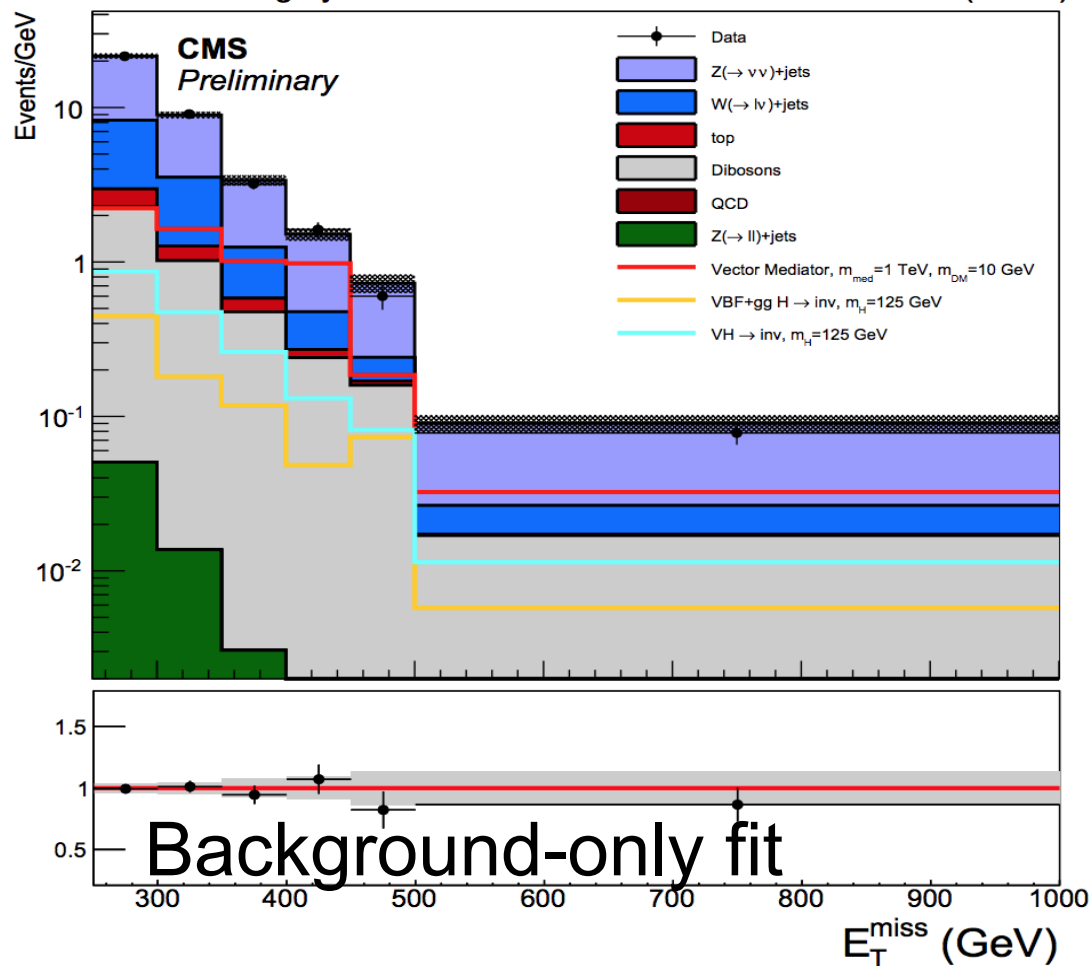
Single jet category



Single “Boosted” jet

boosted category

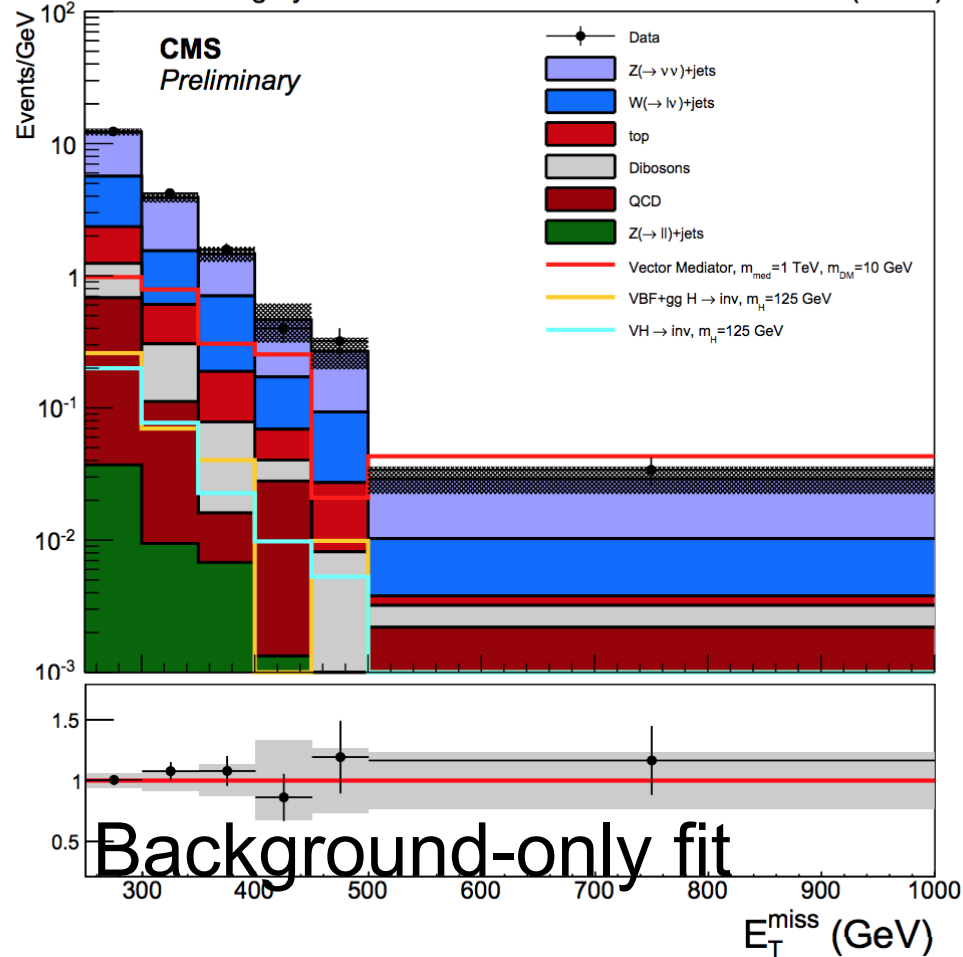
19.7 fb⁻¹ (8 TeV)



Two “Resolved” jets

resolved category

19.7 fb⁻¹ (8 TeV)

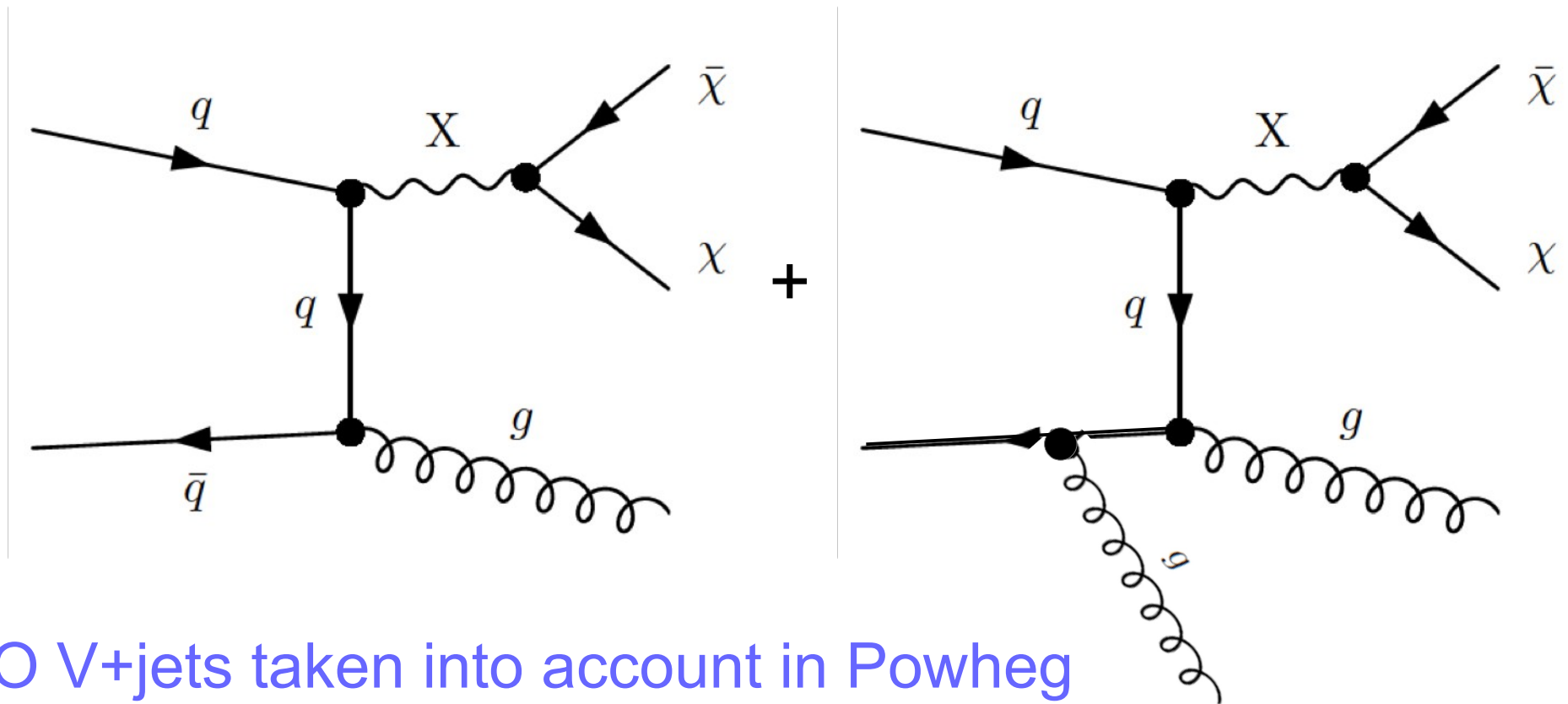


Observe a small excess in resolved MET tail (1σ)

Observe a small deficit in the mono-V MET tail (1.5σ)

2 Jets @ Vector/Axial Simplified Model

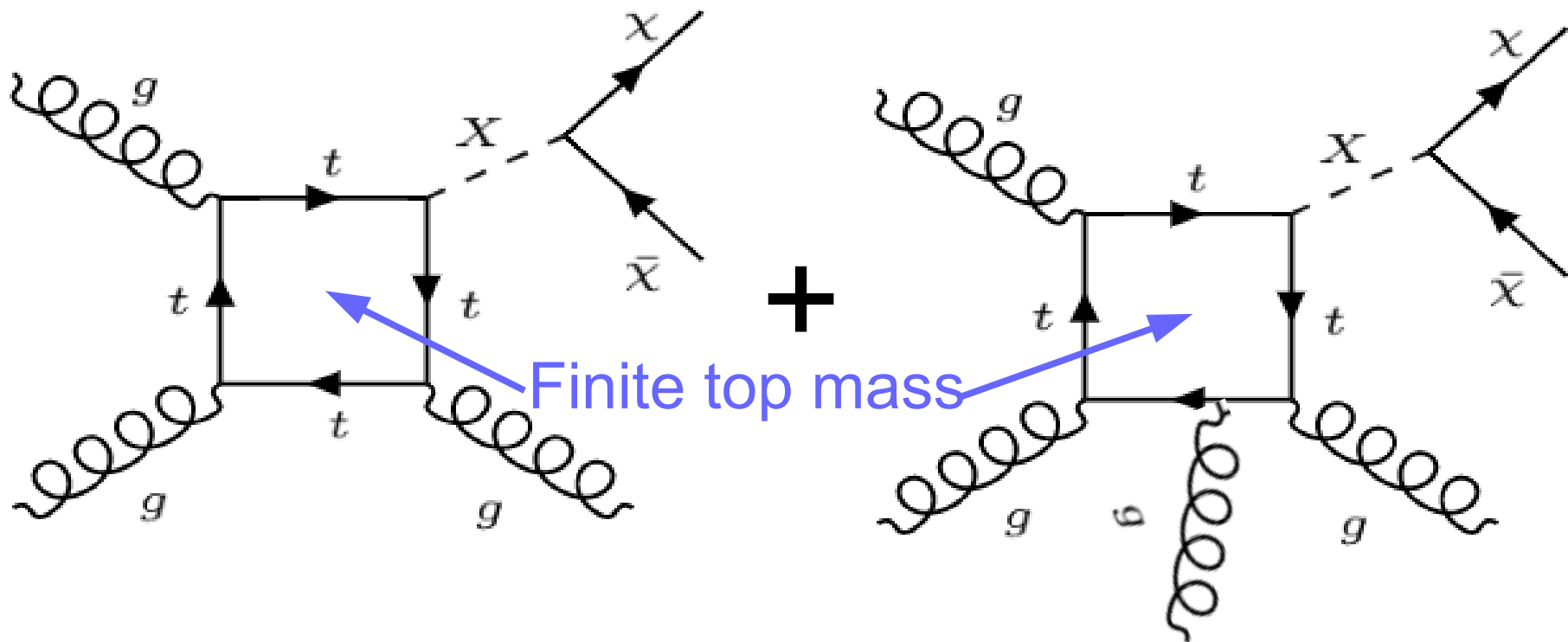
- At higher $\sqrt{s}$ multi-jet final states predominant
- In light of building on new ideas



NLO V+jets taken into account in Powheg
Now available in Madgraph as well

(Pseudo)Scalar Simplified Model

- Requires finite top mass at 2 jets order
 - Available now in now with Madgraph
 - Also can do some hacky procedure



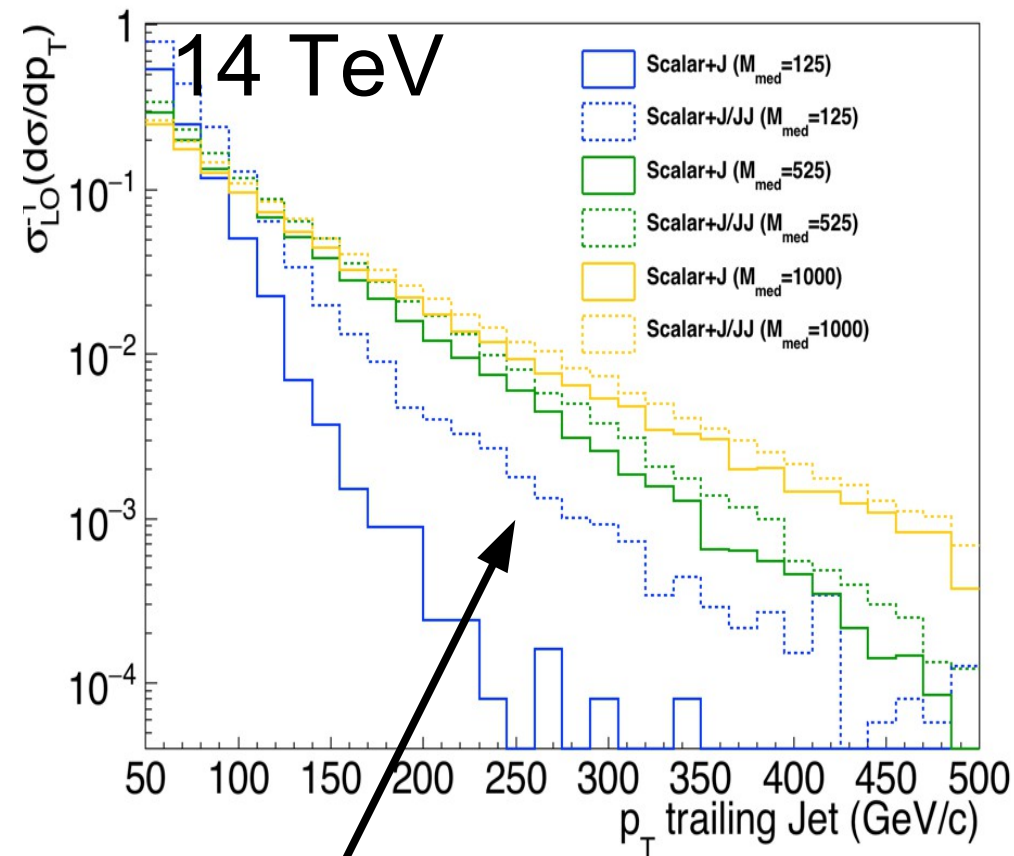
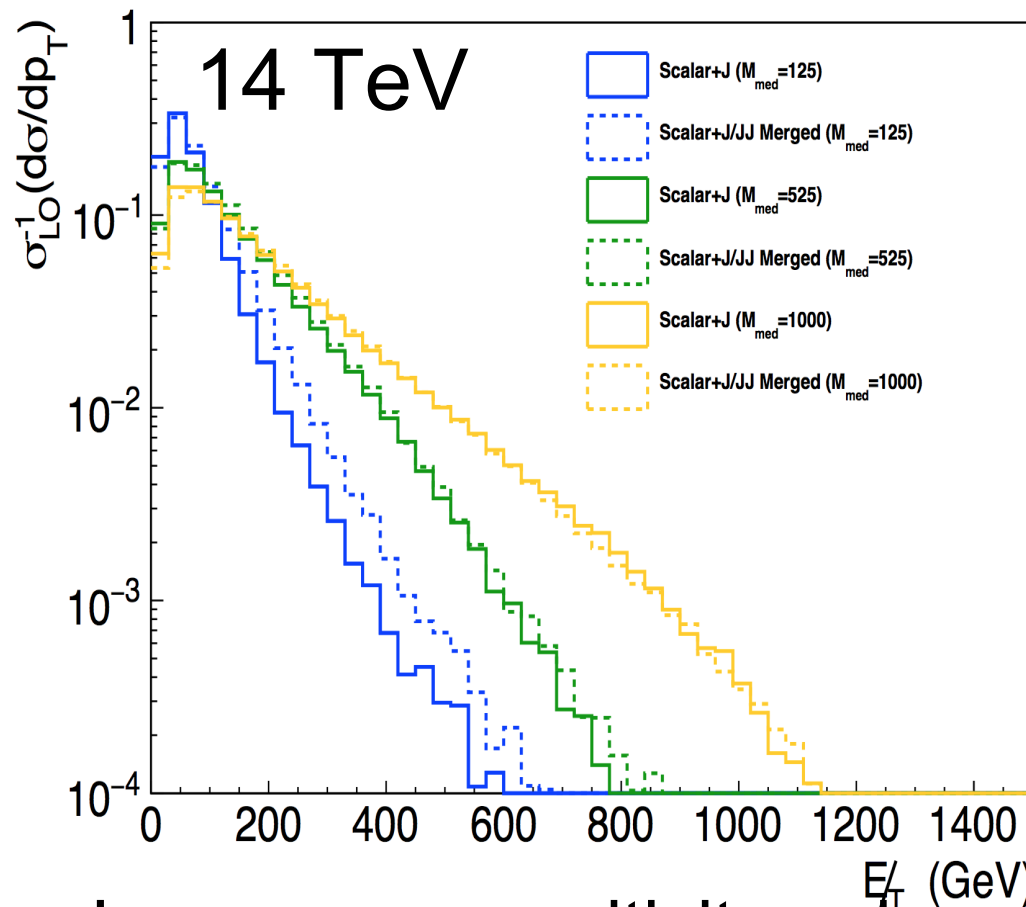
MCFM gluon fusion + 1 jet

VBF@NLO gluon fusion + 2 jet

One caveat : On-shell production is only available for 2 jet final state

Advantage of merged MC

- Taking advantage of the new technology
 - Can consider exploring new regions of phase space



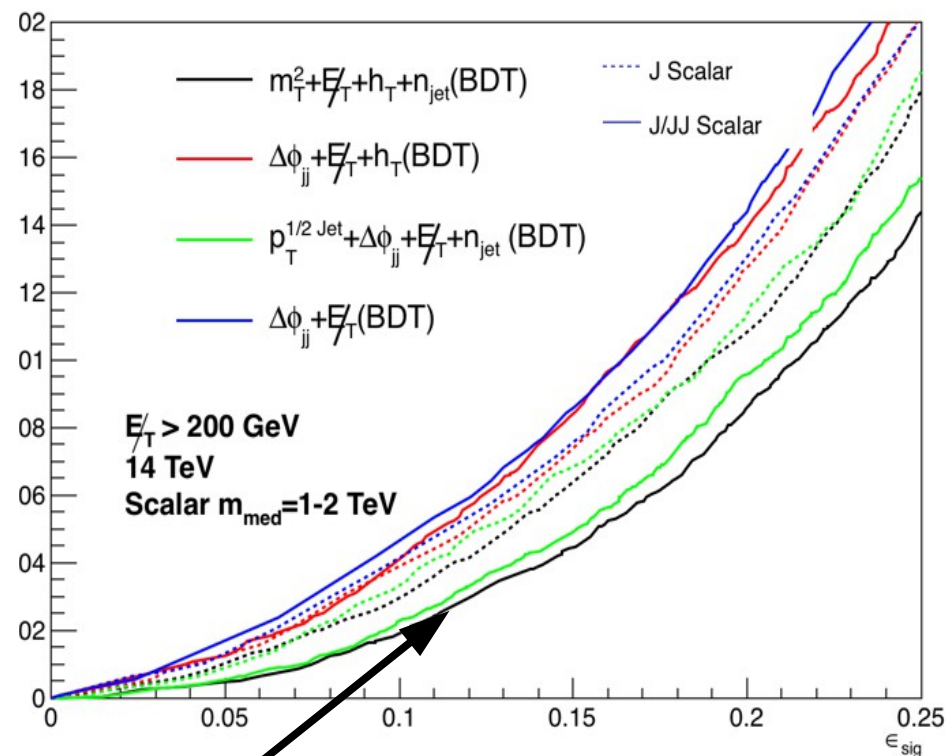
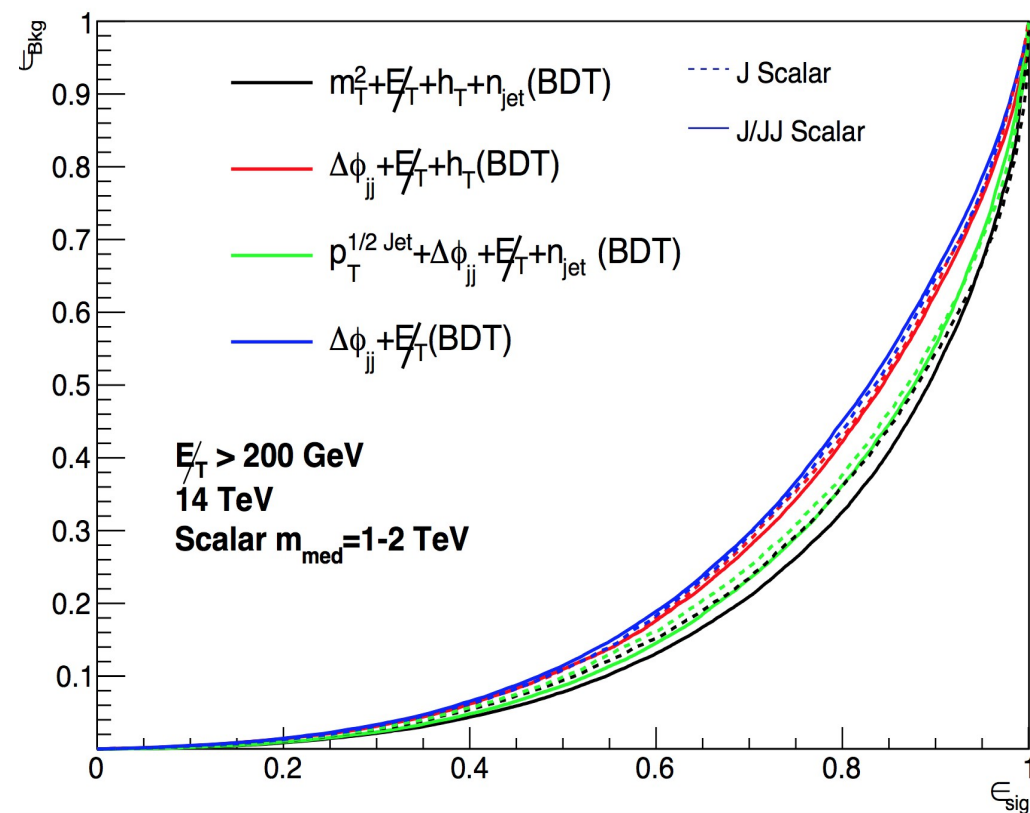
Low mass sensitivity enhanced in the multi jet final state

Question #2 : Advantage of MC

- With 2 jet MC : can now probe multijet final states
- Two questions can be answered :
 - Which variables are most sensitive with 2 jet MC?
 - Which variables are sensitive at 100 TeV?
- Considered a number of multi-jet variables :
 - M_T^2 : SUSY like variable obtained for pairwise sparticles
 - Razor variables : M_R , R : Related SUSY variables
 - MET : standard
 - $\Delta\phi_{jj}$: angle between the two jets

Whats the maximum gain?

- Making an MVA combining all information



Background drops by a factor of 2

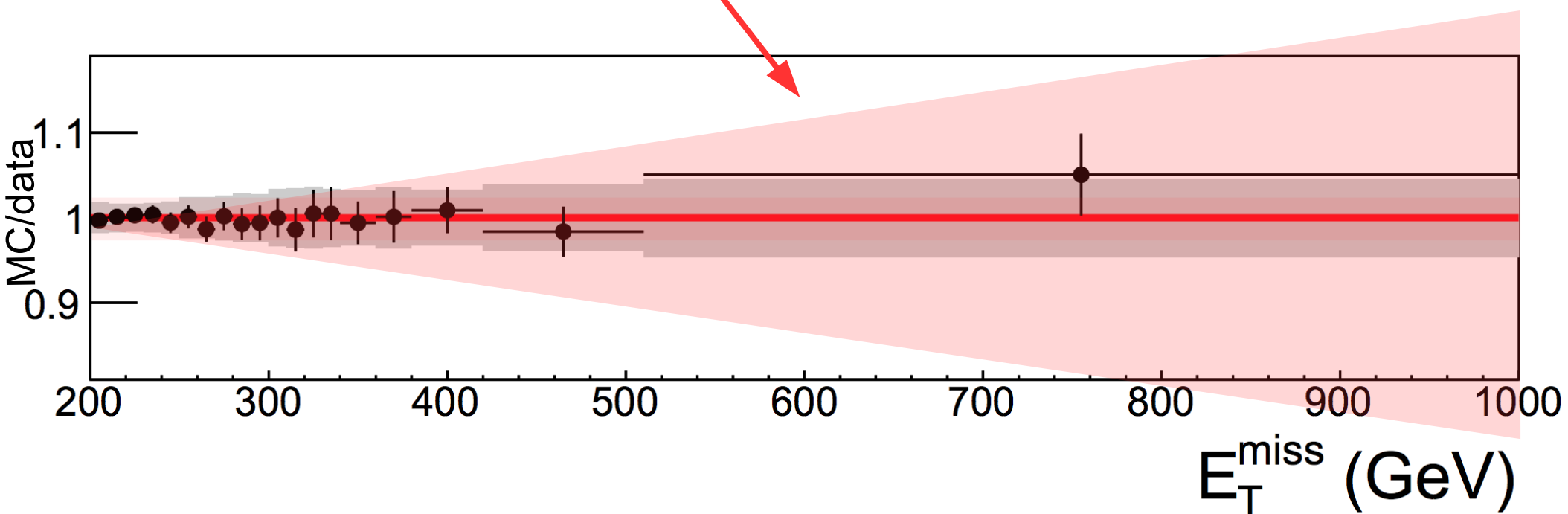
Can maximize sensitivity by an additional $\sqrt{2}$

The only other way to gain is to reduce the systematics

Strategy to fix agreement

10x less $Z \rightarrow \mu\mu$ than $Z \rightarrow \nu\nu$

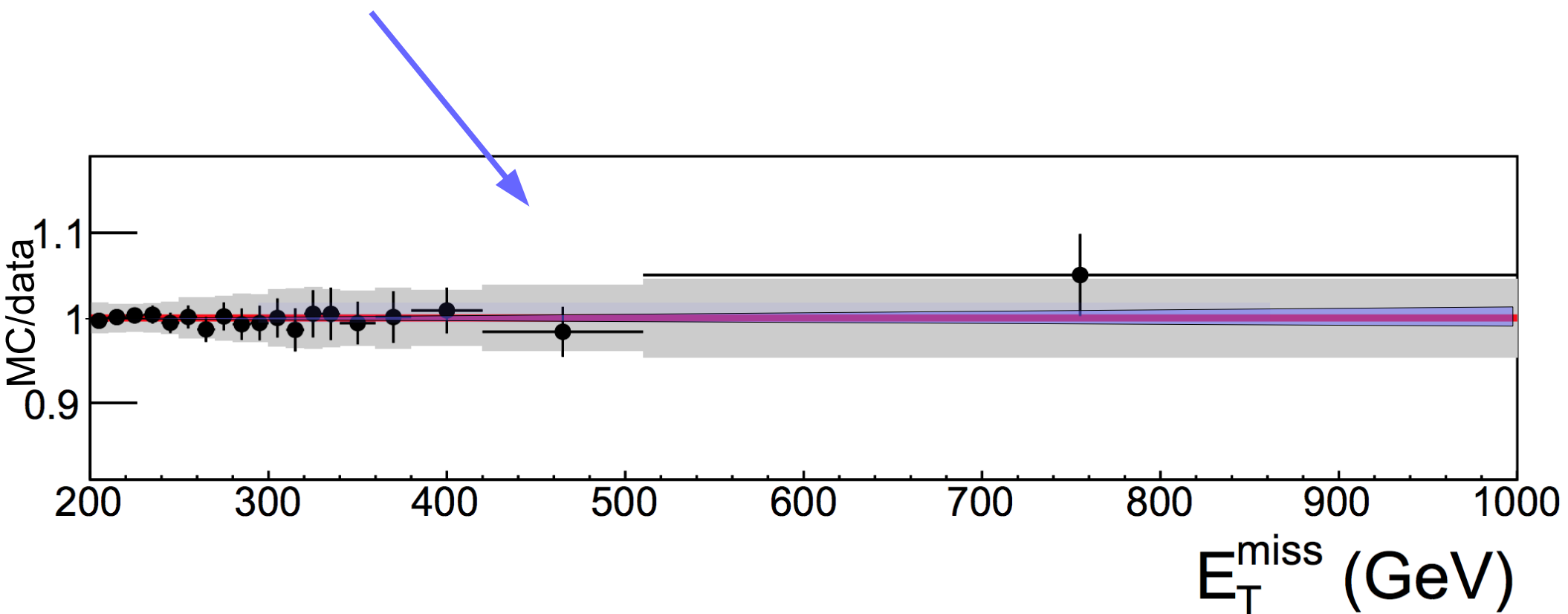
$Z \rightarrow \mu\mu + \text{Jets}$ prediction uncertainty



Statistical uncertainty too large

Strategy to fix agreement

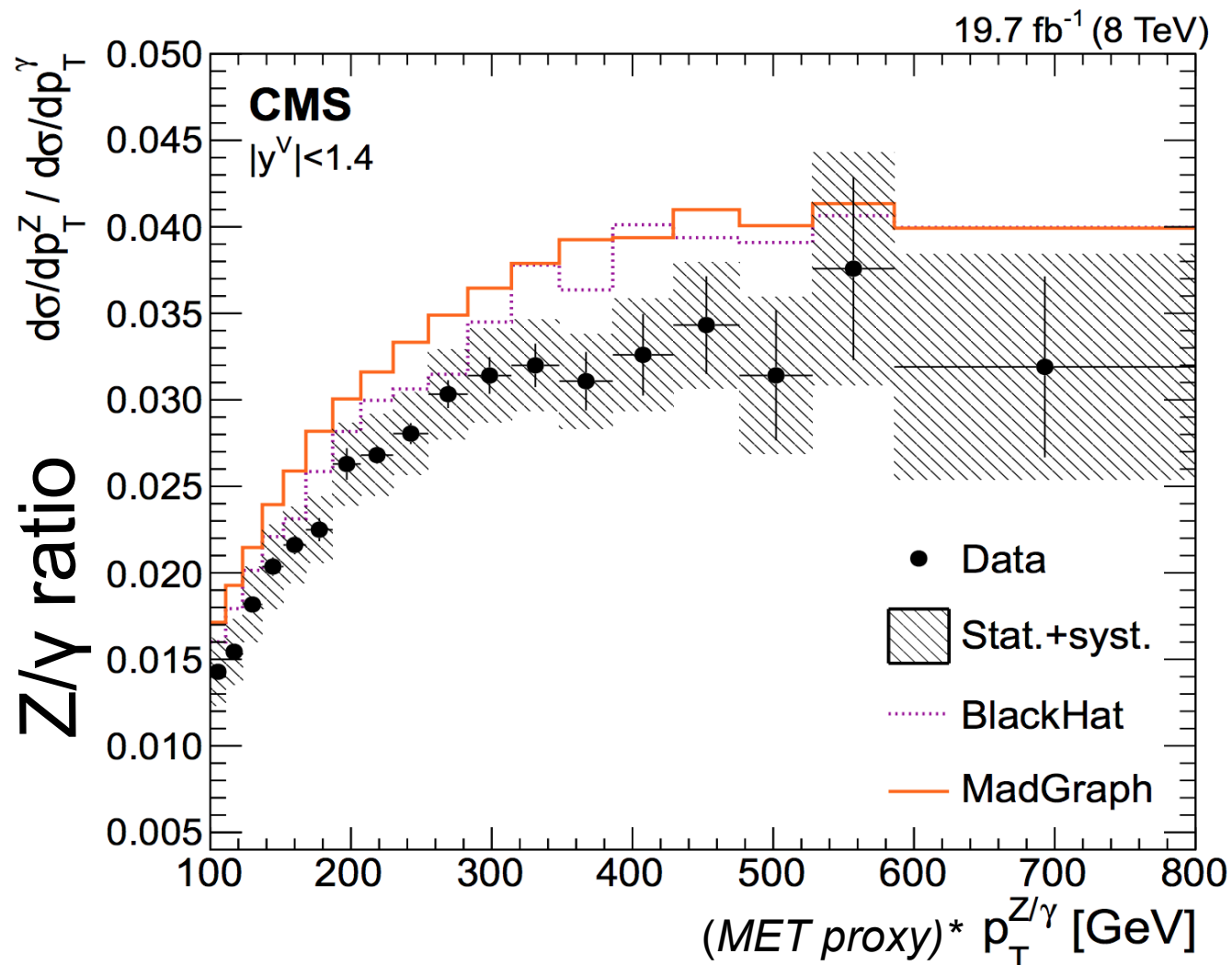
5x More γ +Jets than $Z \rightarrow \nu\nu$



Statistical uncertainty small
....However

A mystery? Understanding $Z/\gamma p_T$

Can we really use Photons to model Zs ?

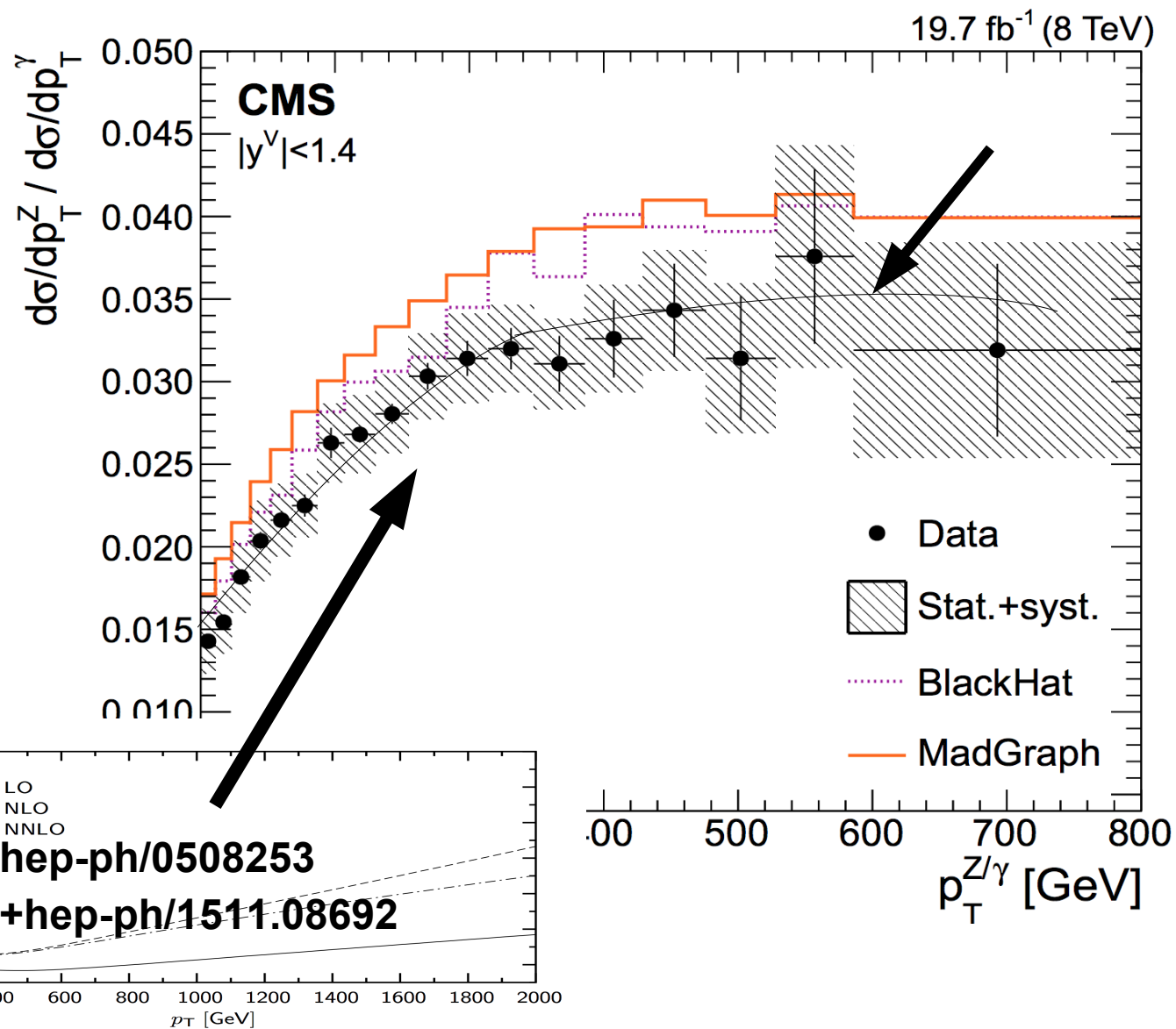


*See backup

CMS-SMP-14-005

A mystery? The $Z p_T$ spectrum

- These results are missing NLO EWK corrections!



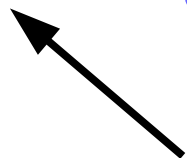
How do we fix this?

$$\frac{d\sigma^Y}{dp_T} \quad \text{and} \quad \frac{d\sigma^Z}{dp_T}$$

Before : $\sigma_{\text{tot}} = \sigma_{\text{NLO}}(0, 1\text{jet})$

After : $\sigma_{\text{tot}} = \sigma_{\text{NLO}}(0, 1, 2j)(1 + \sigma_{\text{EWK}})$ (added)

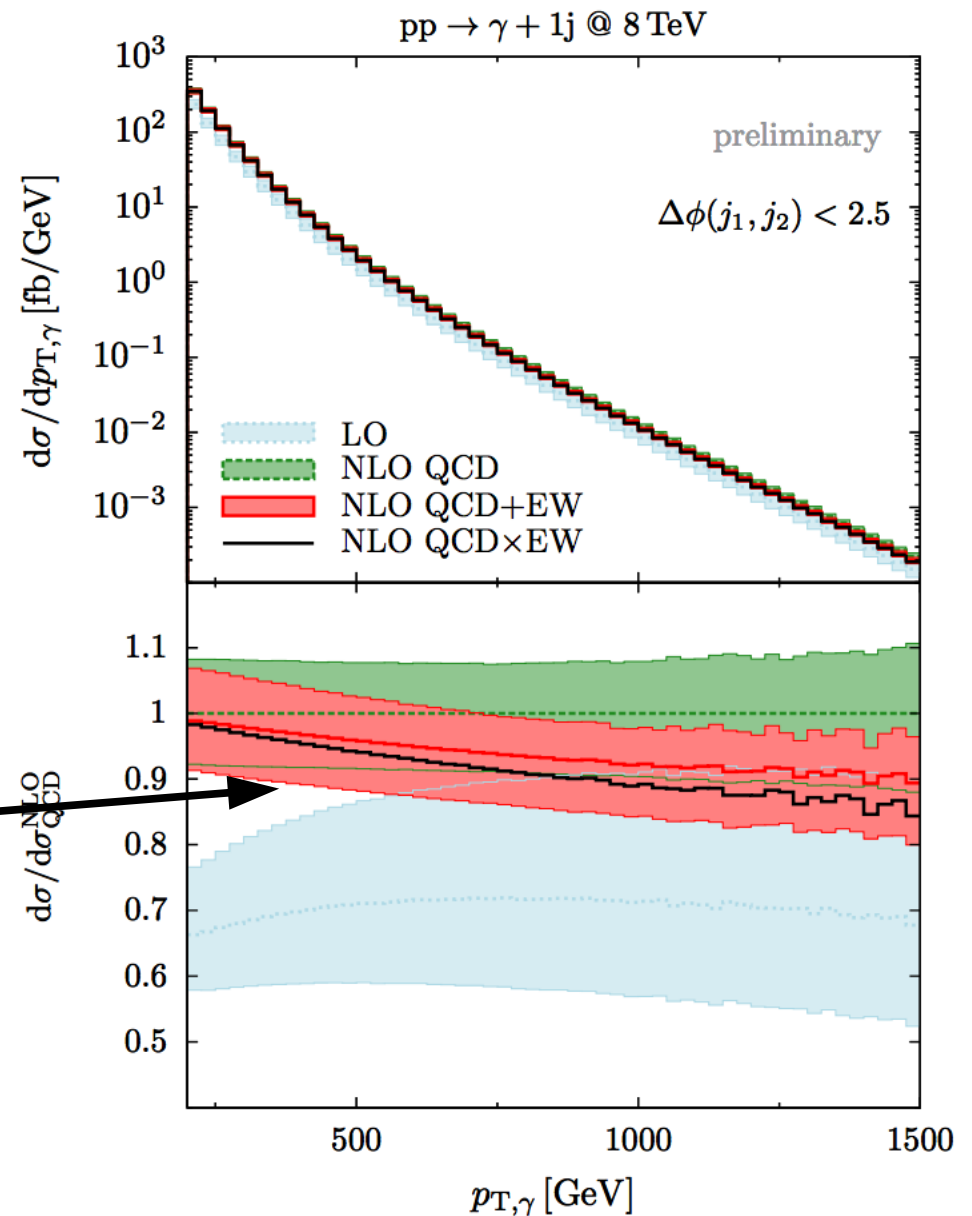
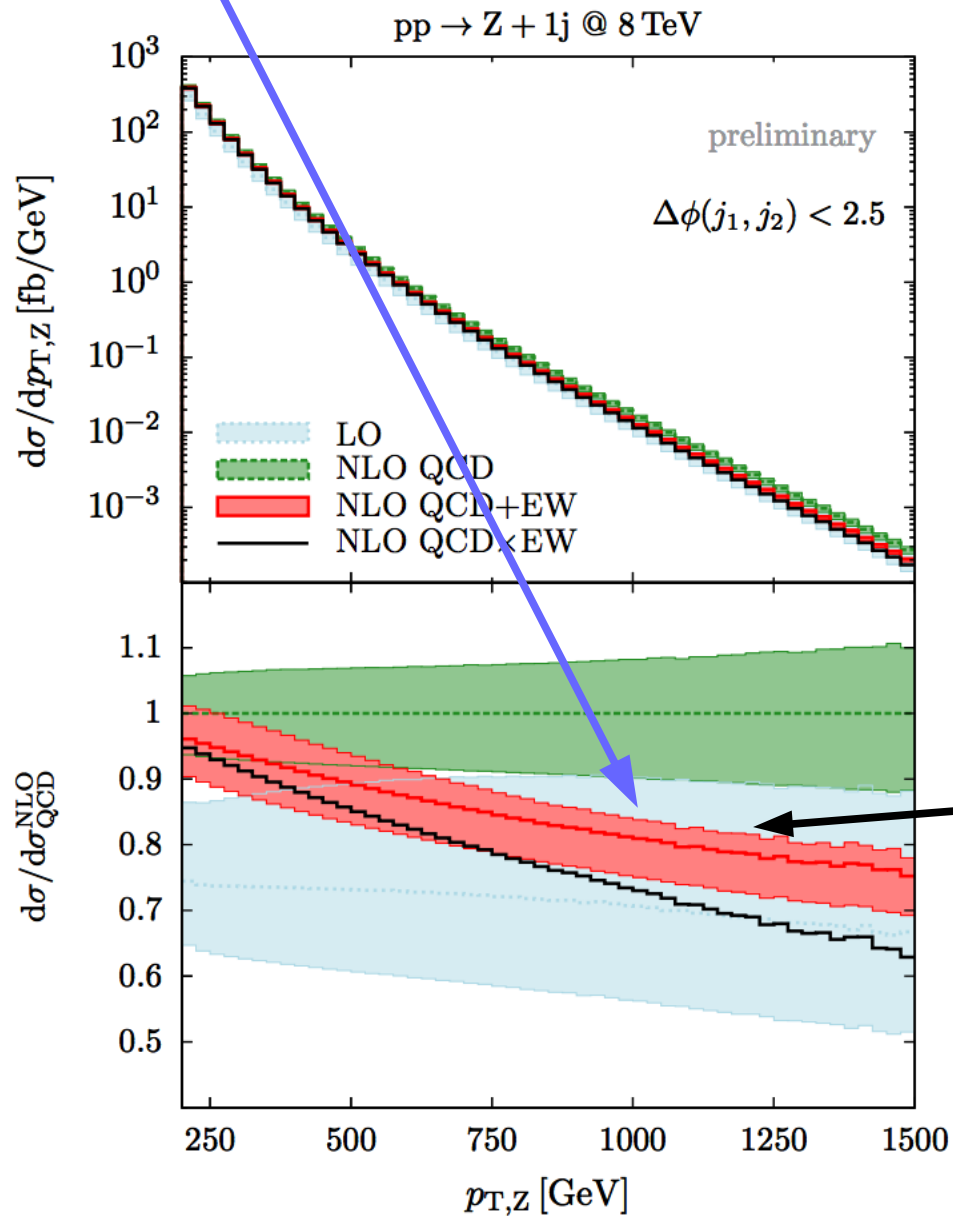
Energy leakage outside of photon which biases *MET*



This was the harder one

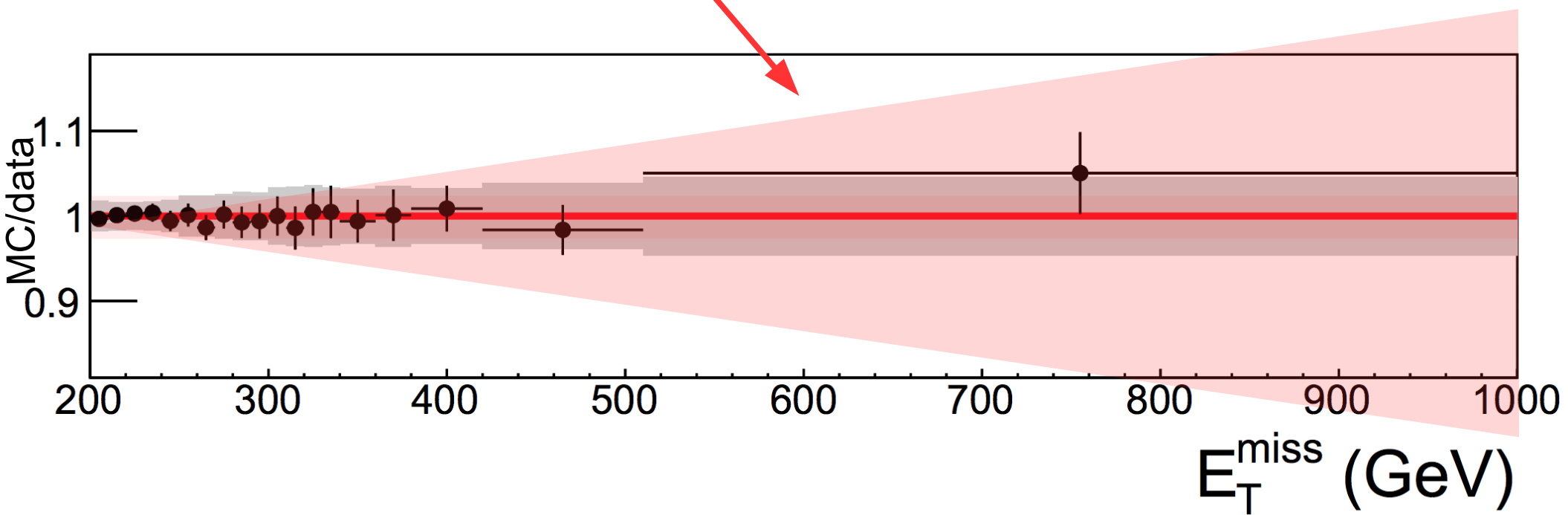
How do we fix this?

Impact of the electroweak corrections

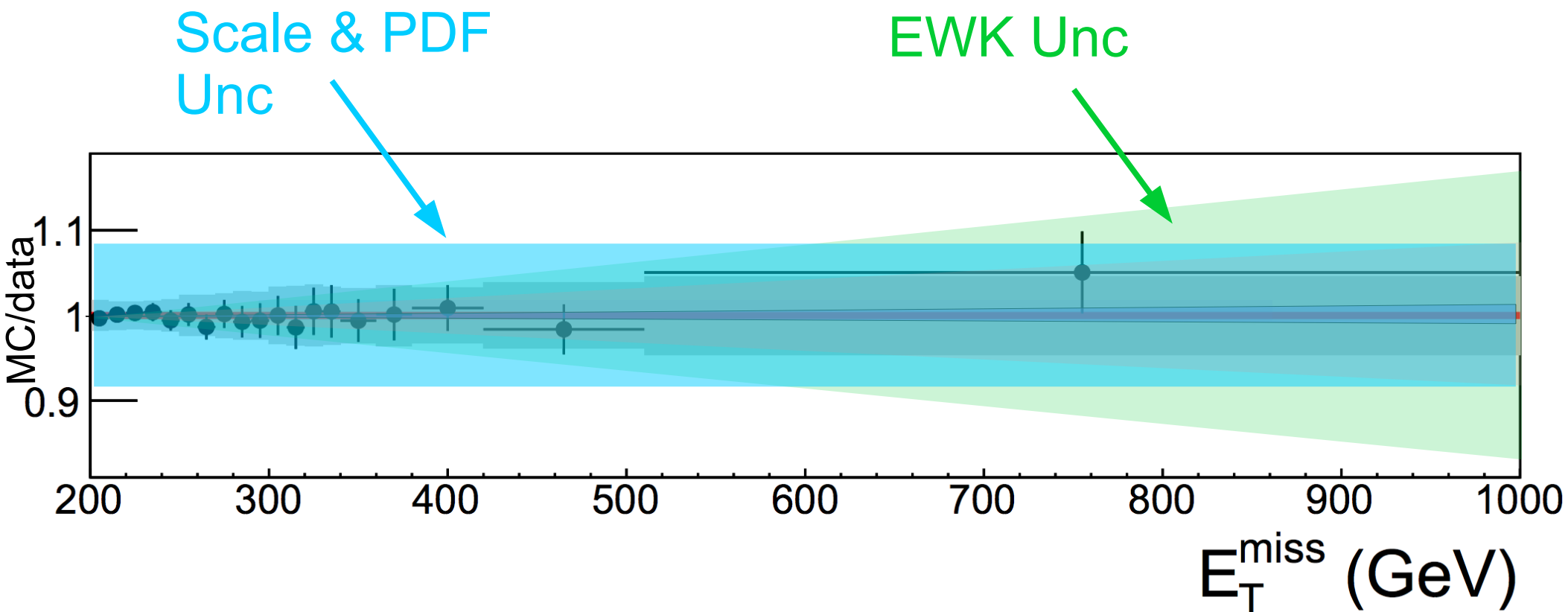


Strategy to fix agreement

$Z \rightarrow \mu\mu + \text{Jets}$ prediction
uncertainty



Strategy to fix agreement



Updated unc still too large

What is the previous unc?

Scale & PDF
Unc



$$\text{Unc.} \longrightarrow \frac{d\sigma^Y}{dp_T} / \frac{d\sigma^Z}{dp_T} = d\sigma^Y/d\sigma^Z(\mu) \quad \uparrow \quad \text{On the ratio}$$

$$\begin{pmatrix} d\sigma^Y(+\sigma) \\ d\sigma^Z(+\sigma) \end{pmatrix} = \begin{pmatrix} 1 & C \\ C & 1 \end{pmatrix} \begin{pmatrix} d\sigma^Y(\mu^{\text{up}})/d\sigma^i(\mu_0) \\ d\sigma^Z(\mu^{\text{up}})/d\sigma^i(\mu_0) \end{pmatrix}$$

Adjust C until
uncertainty is

$$d\sigma^Y/d\sigma^Z(+\sigma) < \max_i (d\sigma^i(\mu^{\text{up}})/d\sigma^i(\mu_0))$$

What is the previous unc?

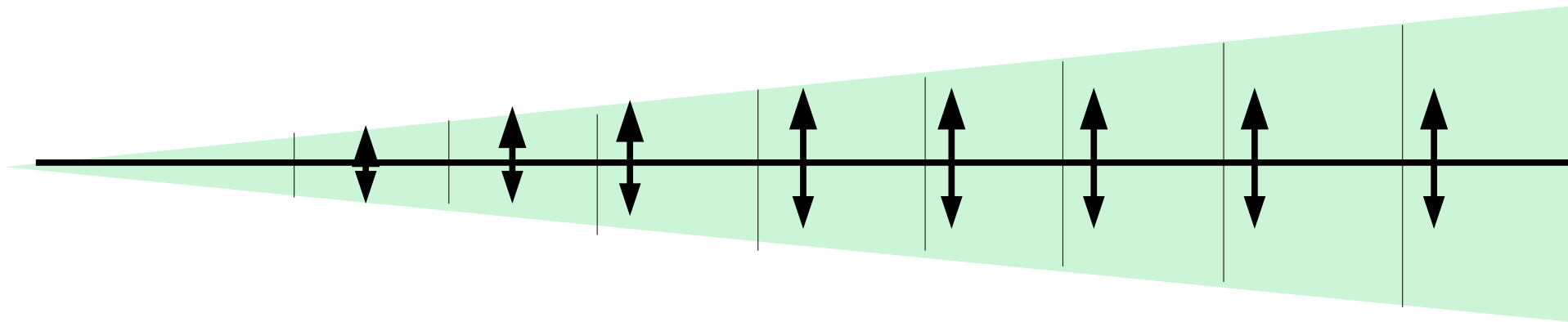
Scale & PDF
Unc

Unc. $\longrightarrow \frac{d\sigma^Y}{dp_T} / \frac{d\sigma^Z}{dp_T} = d\sigma^Y/d\sigma^Z(\mu)$ $\uparrow$
On the ratio

$$\begin{pmatrix} d\sigma^Y(+\sigma) \\ d\sigma^Z(+\sigma) \end{pmatrix} = \begin{pmatrix} 1 & C \\ C & 1 \end{pmatrix} \begin{pmatrix} d\sigma^Y(\mu^{\text{up}})/d\sigma^i(\mu_0) \\ d\sigma^Z(\mu^{\text{up}})/d\sigma^i(\mu_0) \end{pmatrix}$$

Makes No
Sense

What about the EWK uncertainty?



In light of being conservative :

Treated full correction as an uncertainty

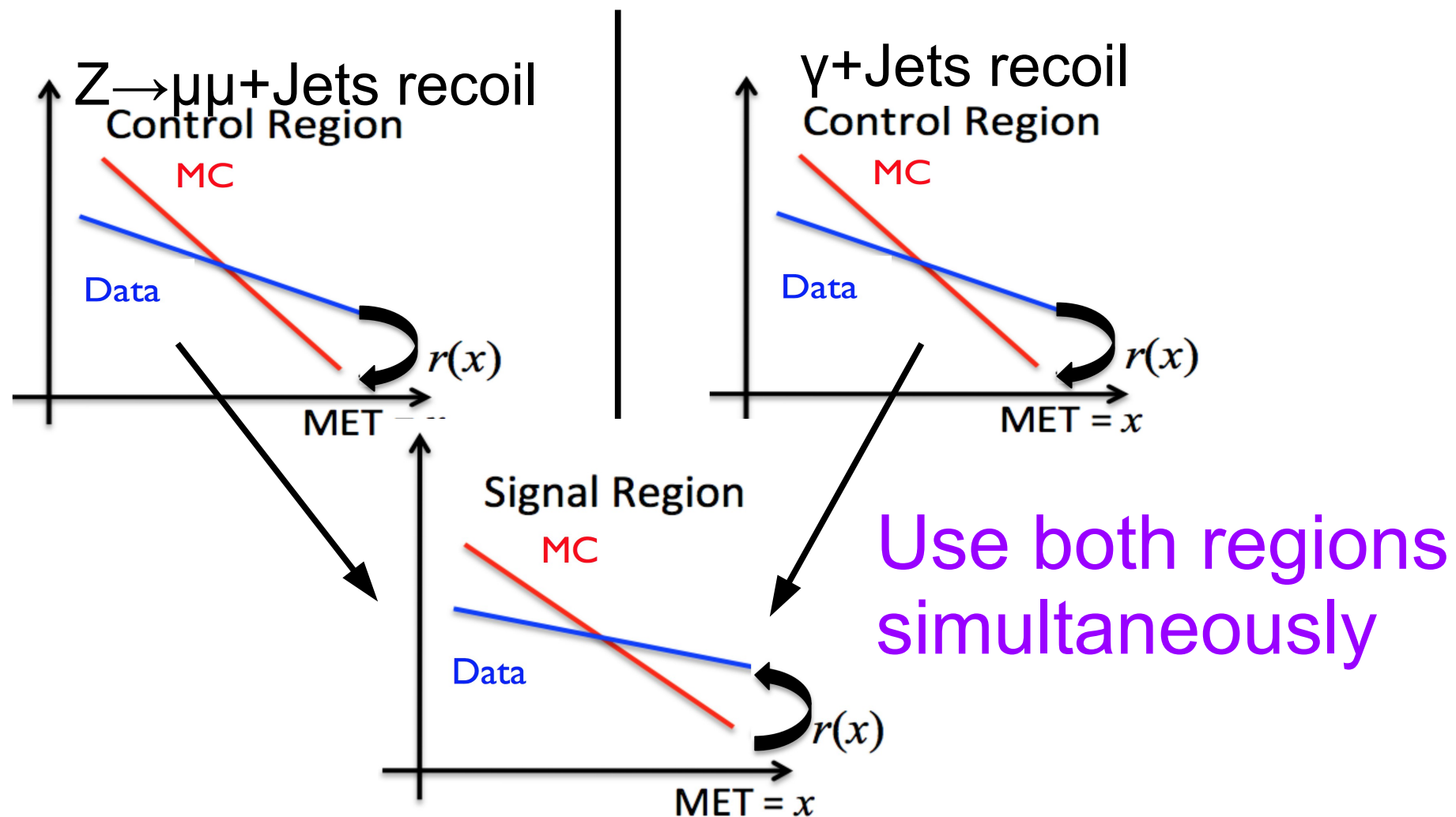
Additionally de-correlated this per bin

Avoids low MET to high MET constraints

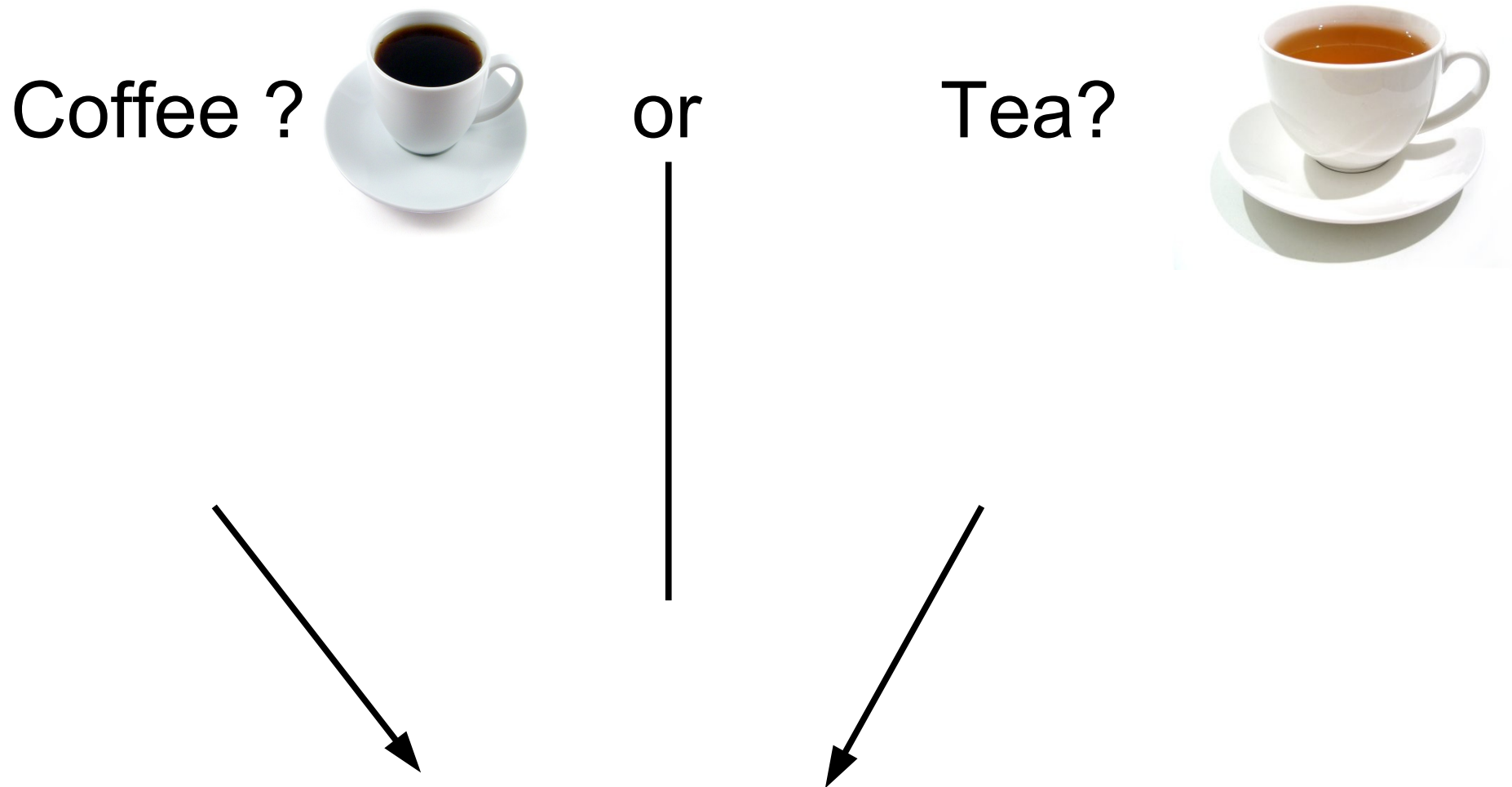
Not very logical

Other (better) schemes exist

Strategy to fix agreement



Strategy to fix agreement



Answer : Yes Please

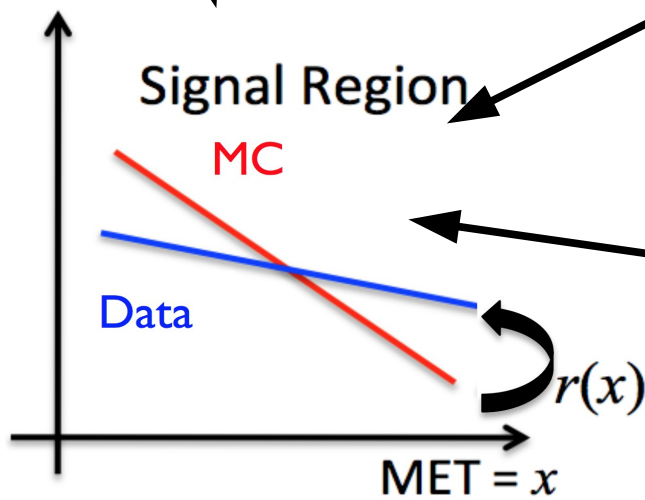
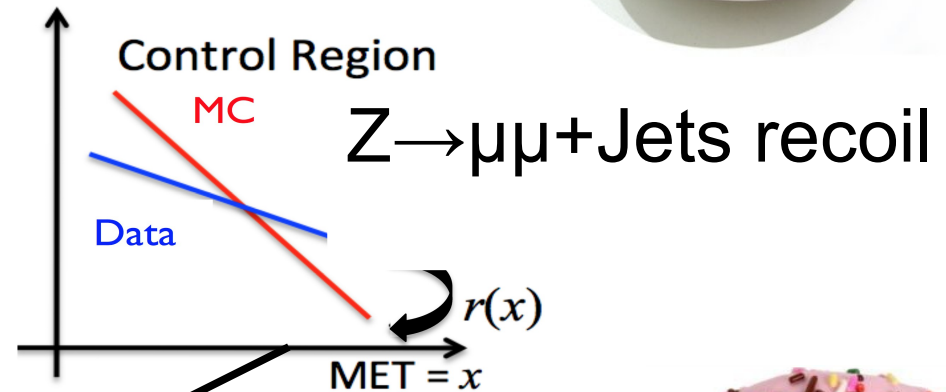
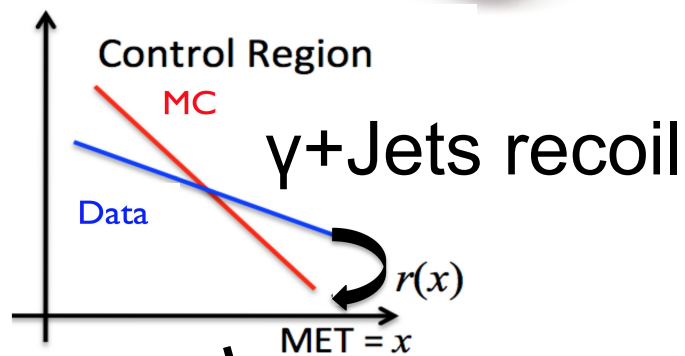
Strategy to fix agreement

Coffee ?



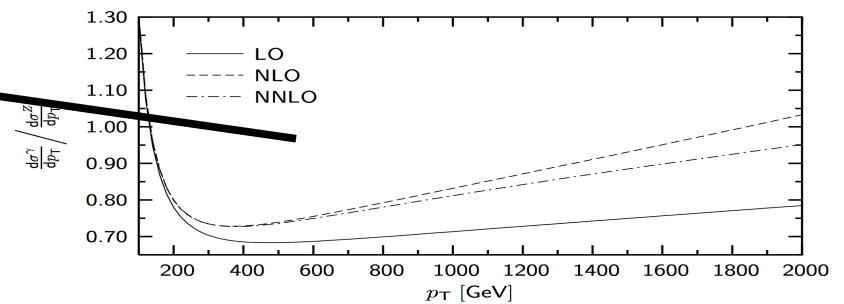
or

Tea?



and Donuts

Not enough LO



Strategy to fix agreement

Coffee ?



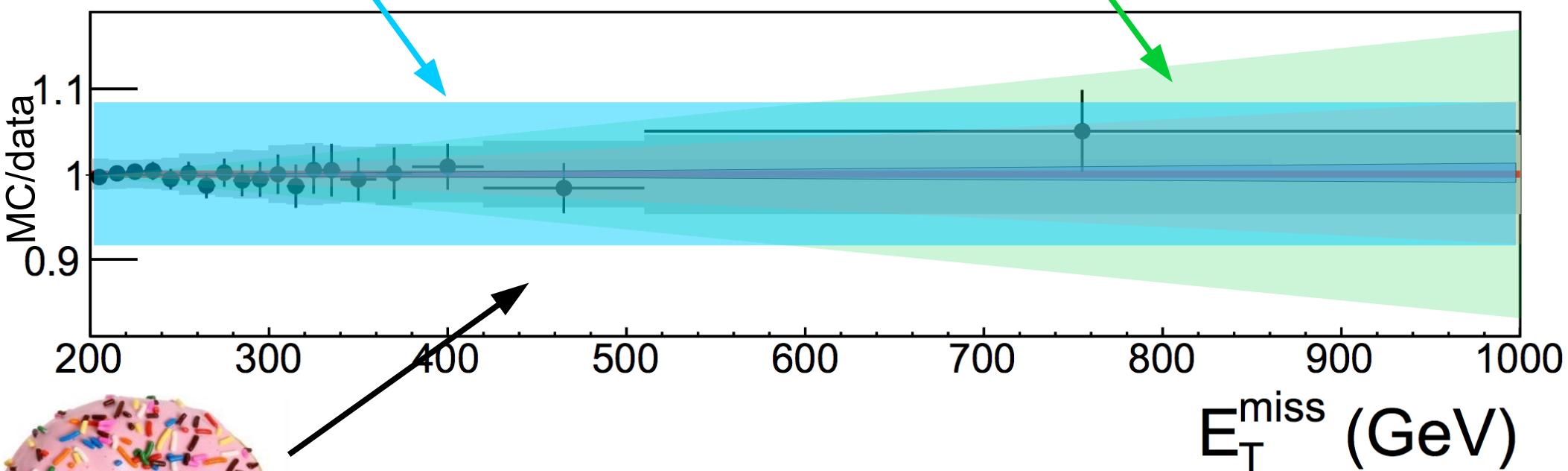
or

Tea?



Scale & PDF
Unc

EWK Unc



Donuts? **Constraint**

Strategy to fix agreement

Coffee ?

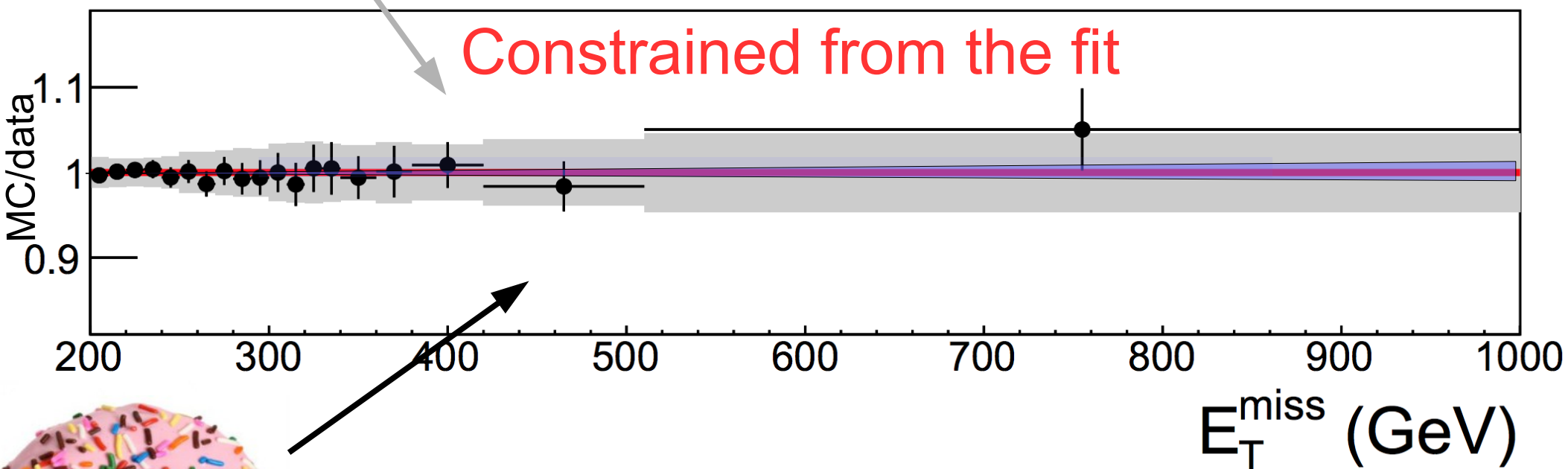


or

Tea?



Actual Uncertainty we use



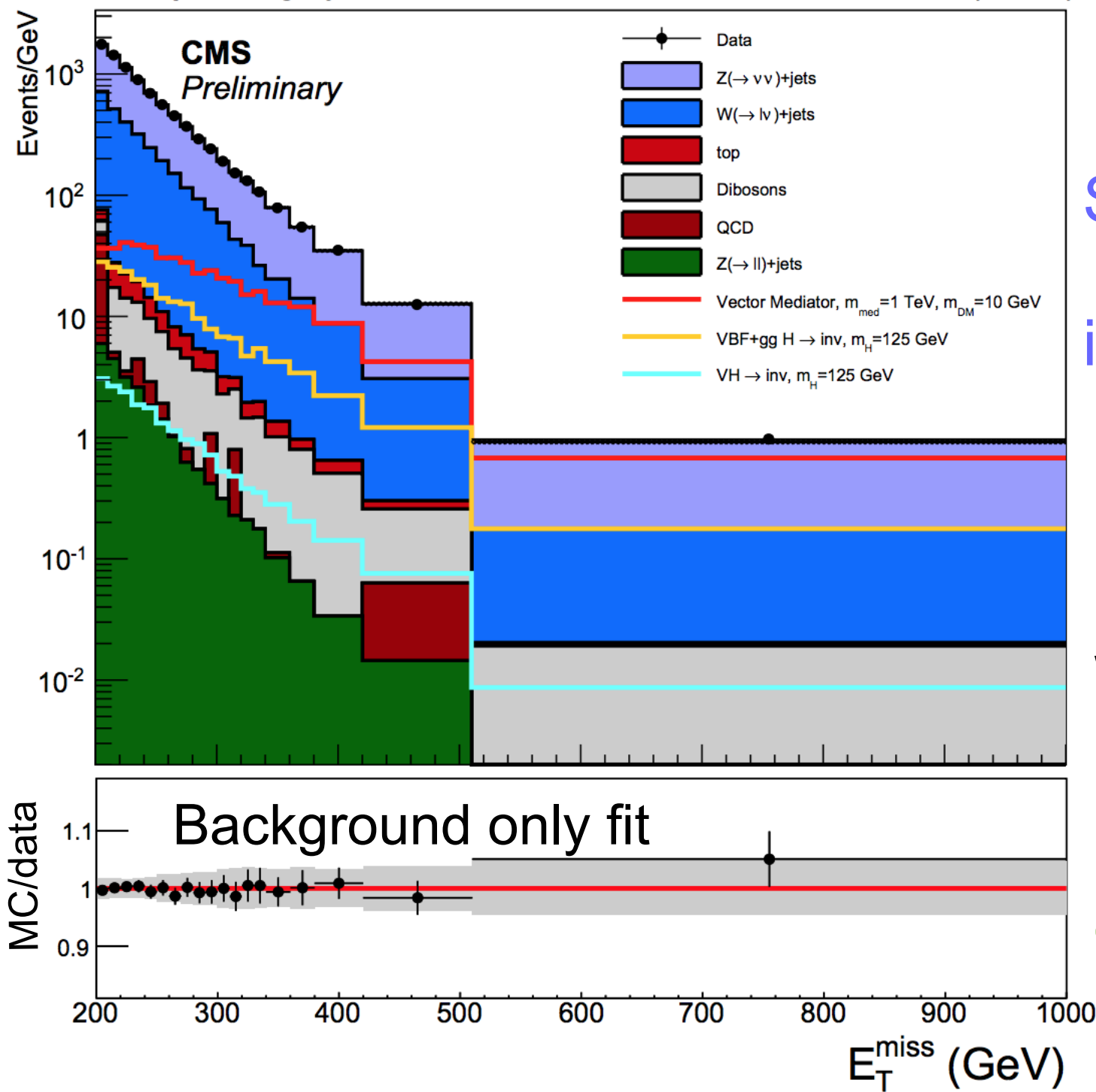
Donuts

monojet category

19.7 fb⁻¹ (8 TeV)

The Result

Small excess
(1-2 σ)
in *MET* tail

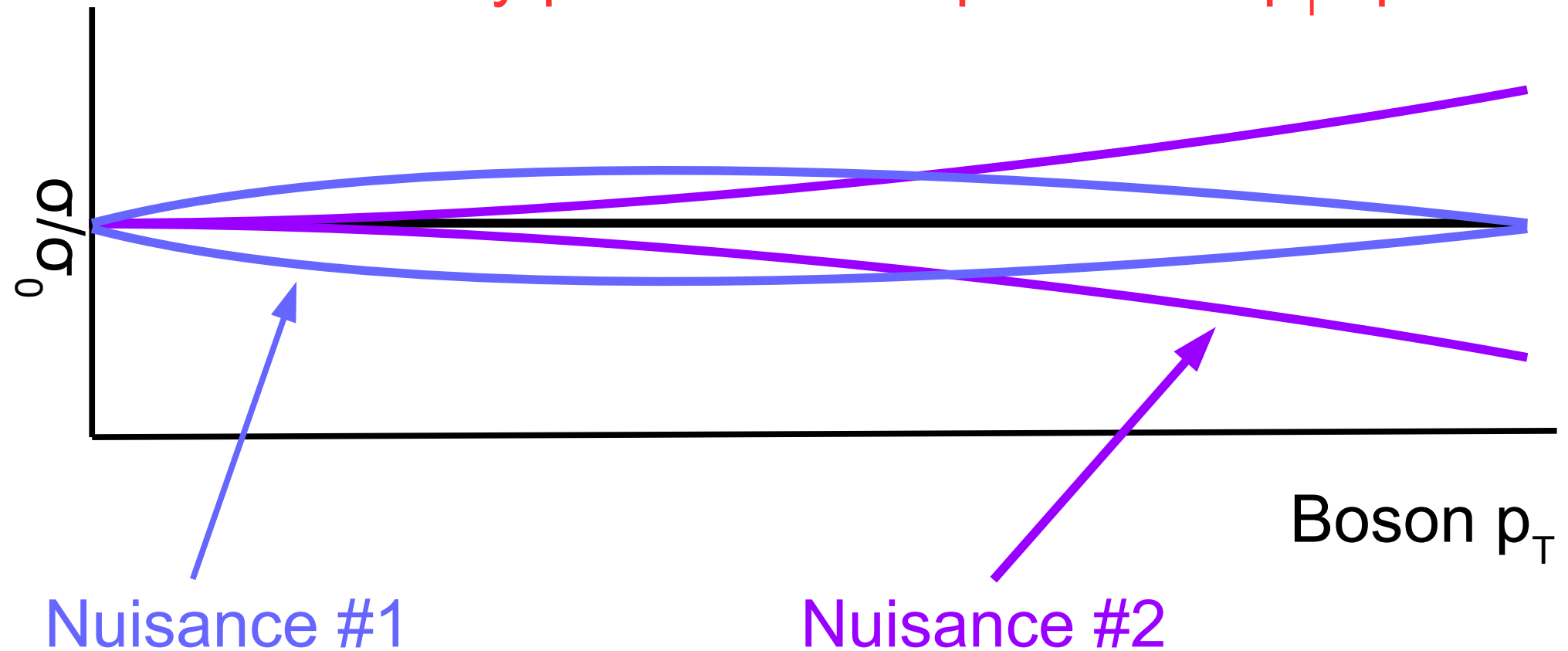


With new method
Still **systematics**
limited
EWK uncertainty
dominates

How do we do the fit?

The updated version of this search fits the
 W, Z, γ p_T simultaneously

Simultaneously profile the shapes of the p_T spectra



Can we bound our uncertainties into a class of shapes?

We are Stuck!

We are relying on

$$d\sigma^Y/d\sigma^Z(+\sigma) < \max_i (d\sigma^i(\mu^{up})/d\sigma^i(\mu_0))$$

For the uncertainty on

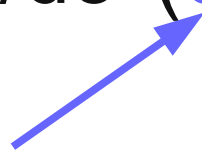
$$\frac{d\sigma^Y}{dp_T} \bigg/ \frac{d\sigma^Z}{dp_T}$$

We need a better approach

Ideally one that we can embed to the likelihood(L)

$$\text{Log}(L) = \text{Log}(L_0) + (d\sigma^Y/d\sigma^Z(\theta) - d\sigma^Y/d\sigma^Z(\mu_0))/\sigma^2$$

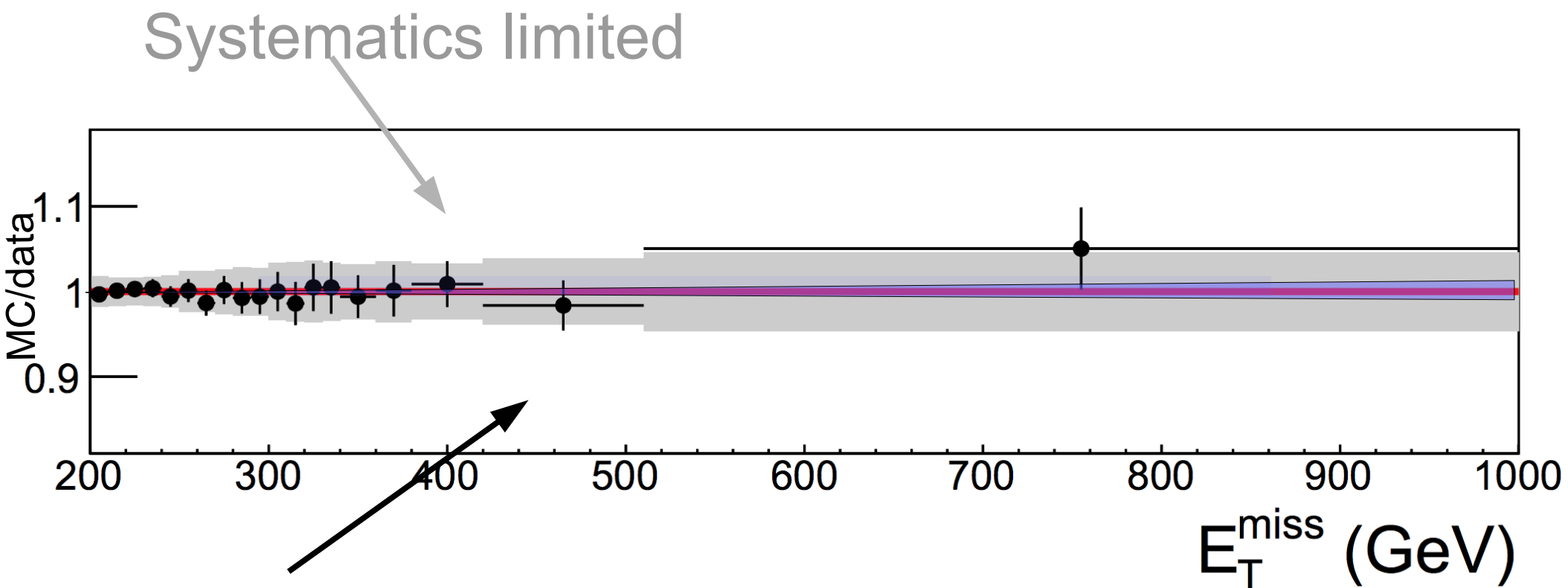
Profiled nuisance



Improved knowledge of high p_T spectrum drives search

Can we improve?

Monojet search will not improve quickly in the future



Driven by our NLO+EWK
uncertainties

Towards a complete statement on Dark Matter

Analyses presented

- Mono Jet:

Scalar Axial Higgs modified vector modified scalar mixture

Jets	Inclusive	V tag	Top	b tag	Higgs
1					
2					
3					
n					
leptons					
γ					
$X \rightarrow \gamma\gamma$					

- Mono Everything: Extending to improve scalar at 100 TeV

Scalar Axial Higgs modified vector modified scalar mixture

Jets	Inclusive	V tag	Top	b tag	Higgs
1					
2					
3					
n					
leptons					
γ					
$X \rightarrow \gamma\gamma$					

- Mono Everything:

Extending models to cover
modified simplified models

Scalar Axial Higgs modified vector modified scalar mixture

Jets	Inclusive	V tag	Top	b tag	Higgs
1					
2					
3					
n					
leptons					
γ					
$X \rightarrow \gamma\gamma$					

- Mono Everything:

Extending models to cover
modified simplified models

Scalar Axial Higgs modified vector modified scalar mixture

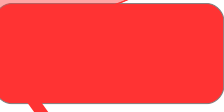
Jets	Inclusive	V tag	Top	b tag	Higgs
1					
2					
3					
n					
leptons					
γ					
$X \rightarrow \gamma\gamma$					

Presented

- Mono Everything:

Extending models to cover
modified simplified models

Scalar Axial Higgs modified vector modified scalar mixture

Jets	Inclusive	V tag	Top	b tag	Higgs
1					
2					
3					
n					
leptons					
γ					
$X \Rightarrow \gamma\gamma$					

Requires High p_T V/t-tag
Experimentally difficult
At extreme p_T

- Mono Everything:

Extending models to cover
modified simplified models

Scalar Axial Higgs modified vector modified scalar mixture

Jets	Inclusive	V tag	Top	b tag	Higgs
1					
2					
3					
n					
leptons					
γ					
$X \rightarrow YY$					

VBF final state
Very useful in higgs
portal models
Very challenging
experimentally

- Mono Everything:

Extending models to cover modified simplified models

Scalar Axial Higgs modified vector modified scalar mixture

	Jets	Inclusive	V tag	Top	b tag	Higgs
1						
2						
3						
n						
leptons						
γ						
$X \rightarrow YY$						

Compliments

$V \rightarrow jj$

$\Phi \rightarrow tt$

Both very powerful

Compliments

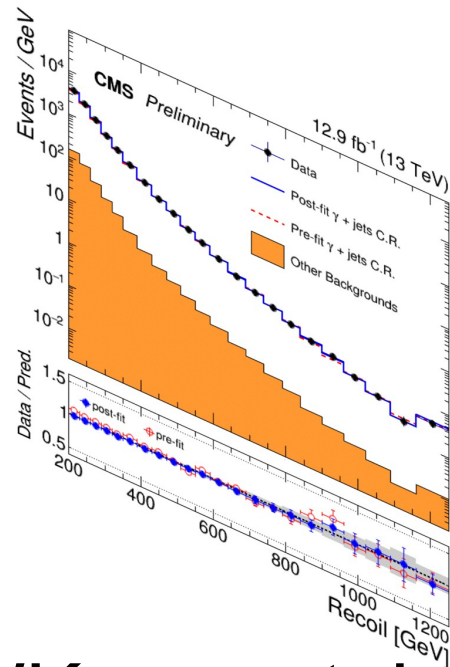
$V \rightarrow jj$

$\Phi \rightarrow tt$

Both very powerful

Common misnomers

- γ +jets control suffers from γ fragmentation model



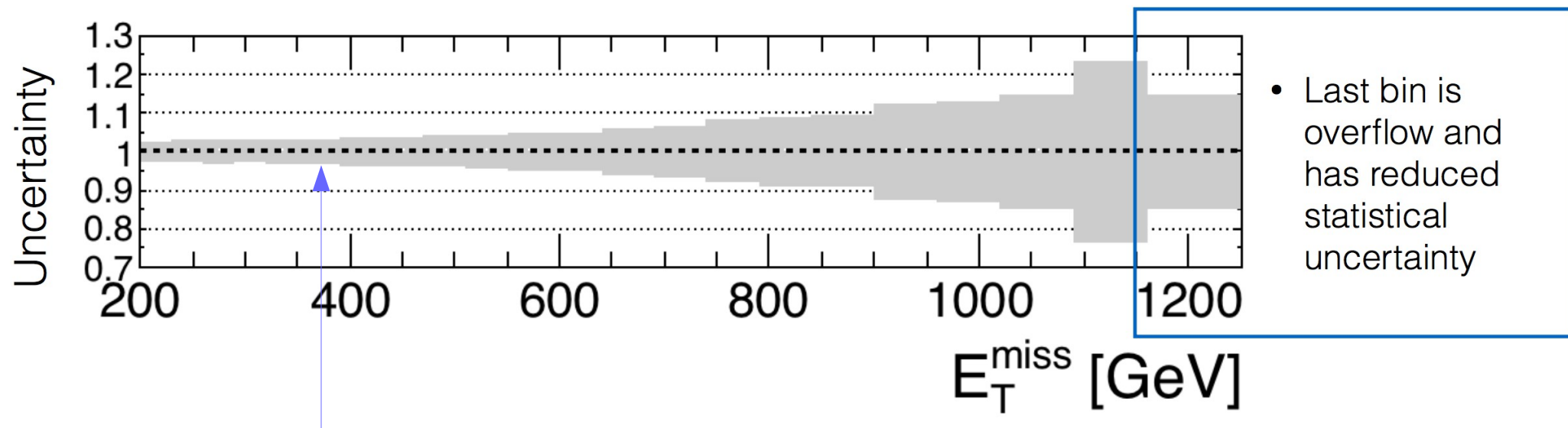
Perform 2D fit in isolation to compute fragmentation component

Photon Isolation

fit resolves issues of photon frag. component

- γ +jets EWK uncertainty too large to constrain
 - W+jets EWK corrections are much smaller
 - Point is we use both in simultaneous fit
 - Dropping one or other is a $>10\%$ reduction in sensitivity

Projecting beyond this



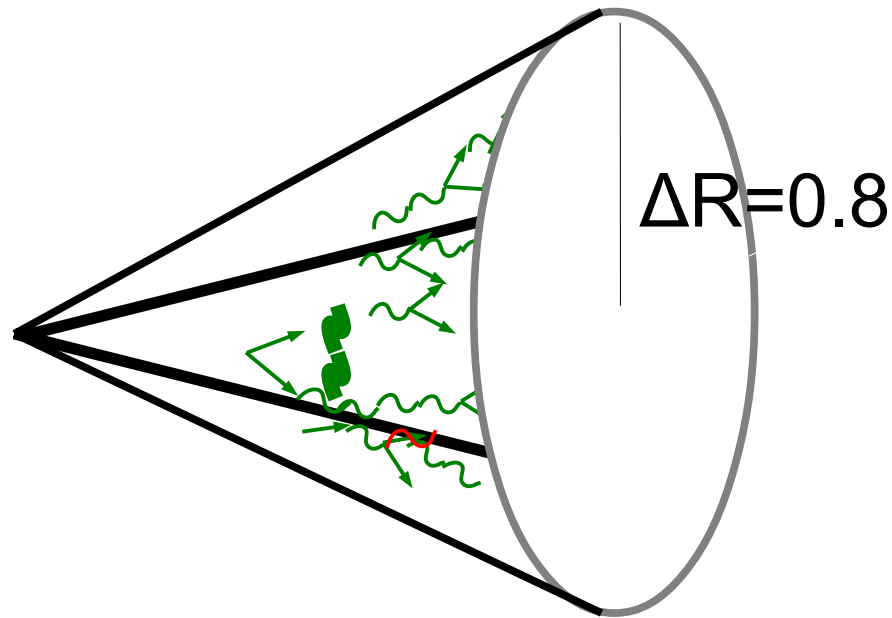
Current low *MET* region unc. comes from $Z \rightarrow \ell\ell$ control region
 Dominant uncertainty is the lepton efficiency measurement
 Will improve considerably with updated analyses

Current high *MET* region unc. from the γ & W control region
 Will improve with new scale/EWK unc. Recommendations

As lumi LHC increases correlated data approaches are critical

What differentiates them?

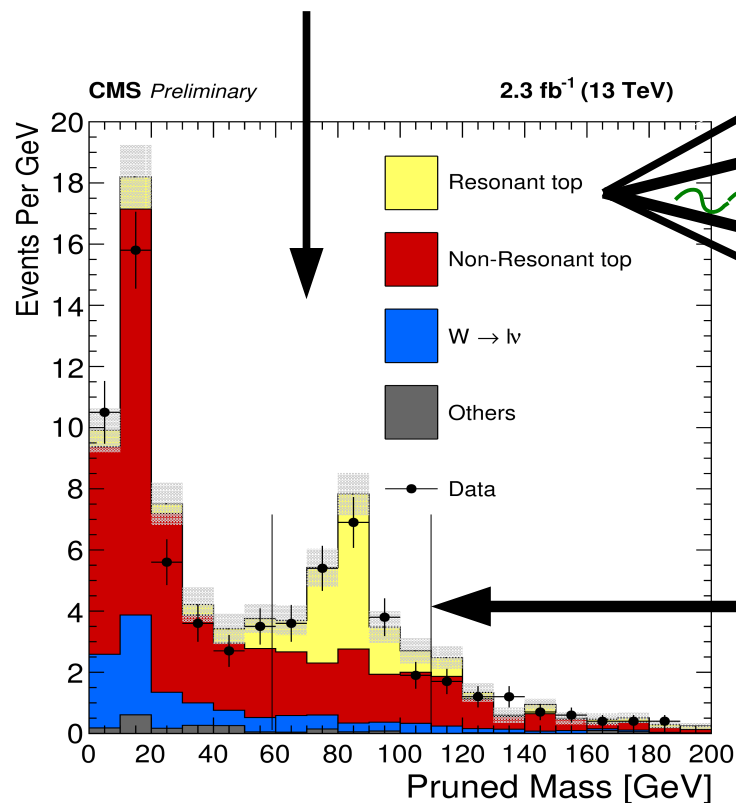
- For both the mono-top and mono-V we tag



What differentiates them?

- For both the mono-top and mono-V we tag

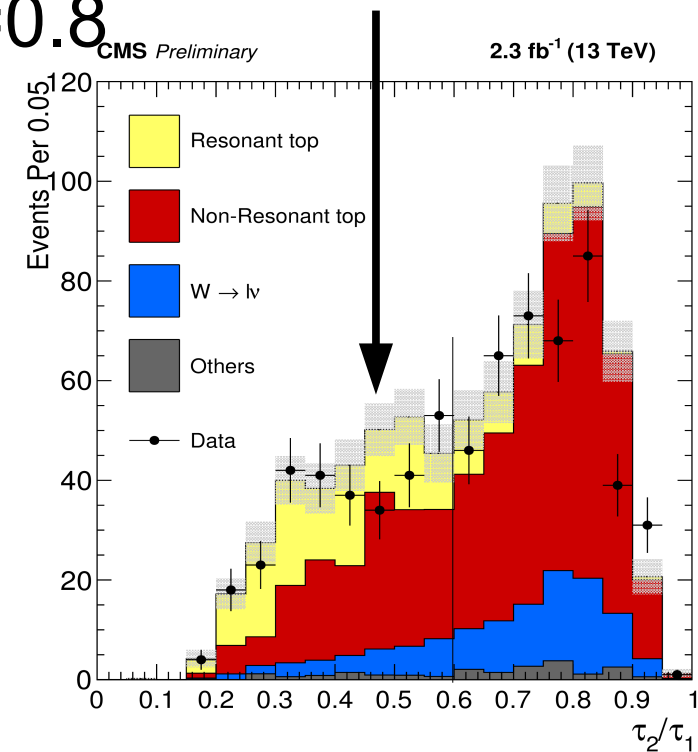
Cut on mass



Tag & Probe on this peak
To get efficiency

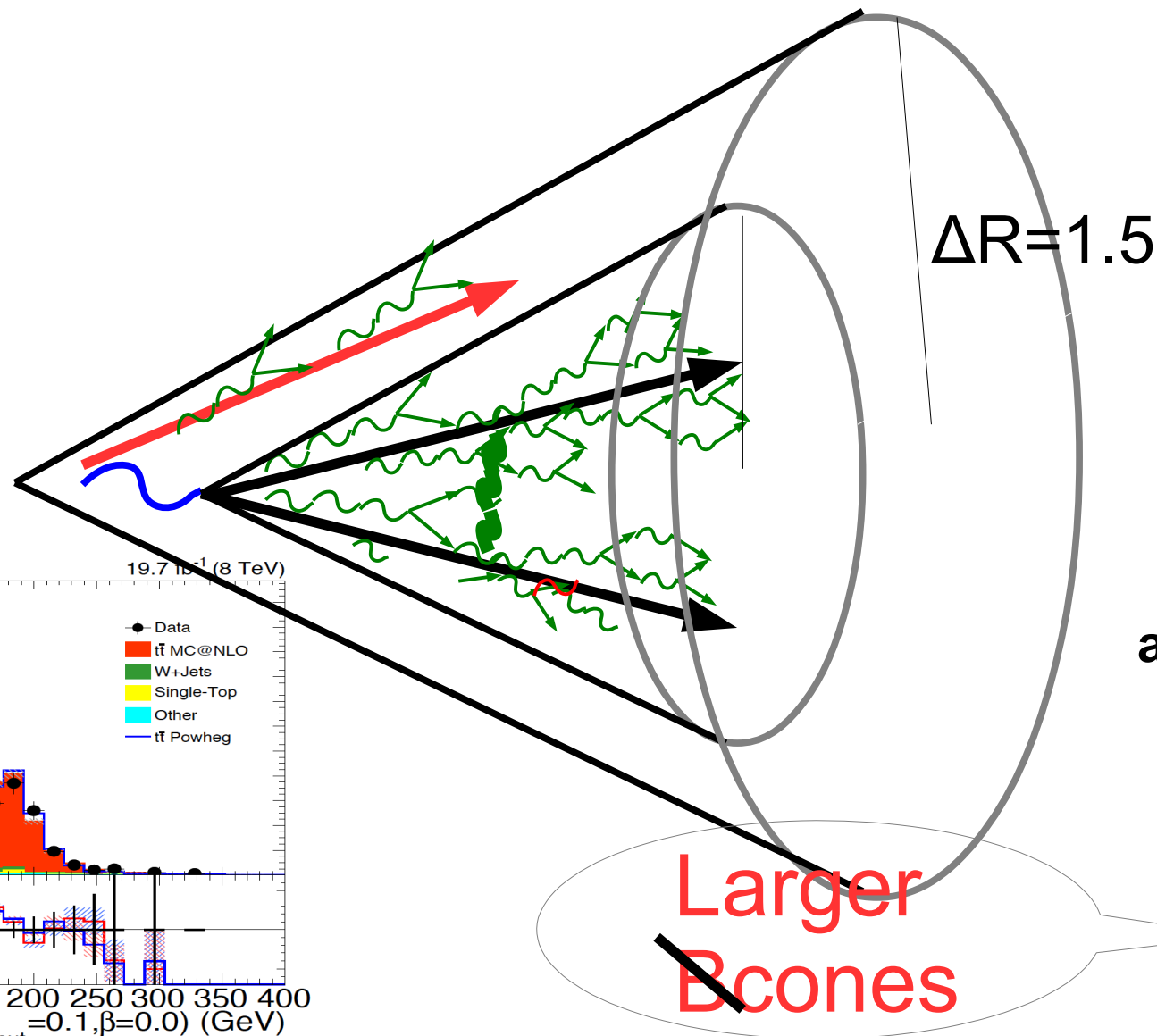
$\Delta R = 0.8$

Cut on
Likelihood of
2 prongs (τ_2/τ_1)



What differentiates them?

- For both the mono-top and mono-V we tag



arXiv:1407.6013

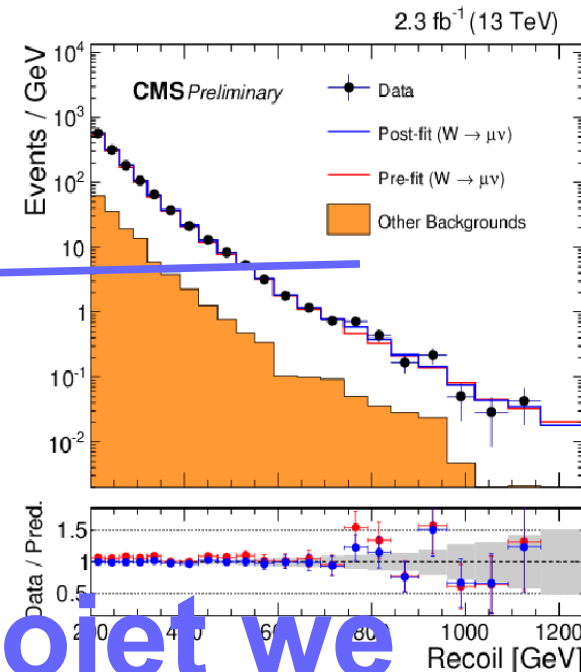
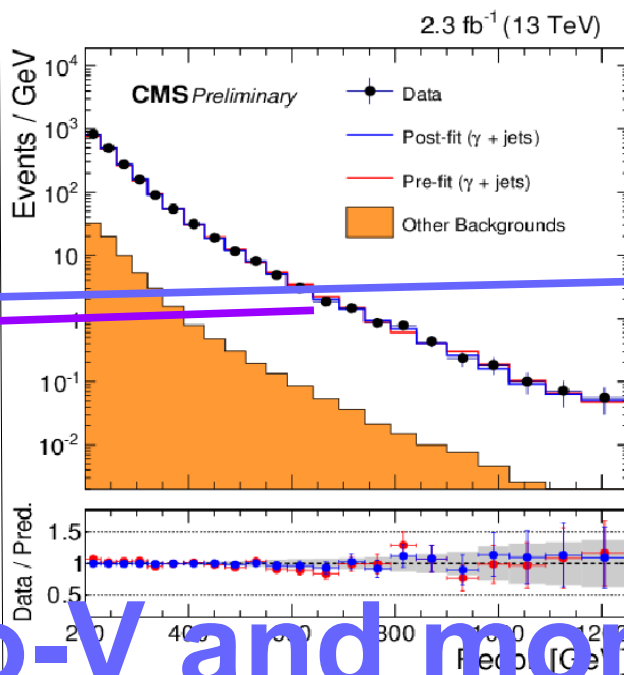
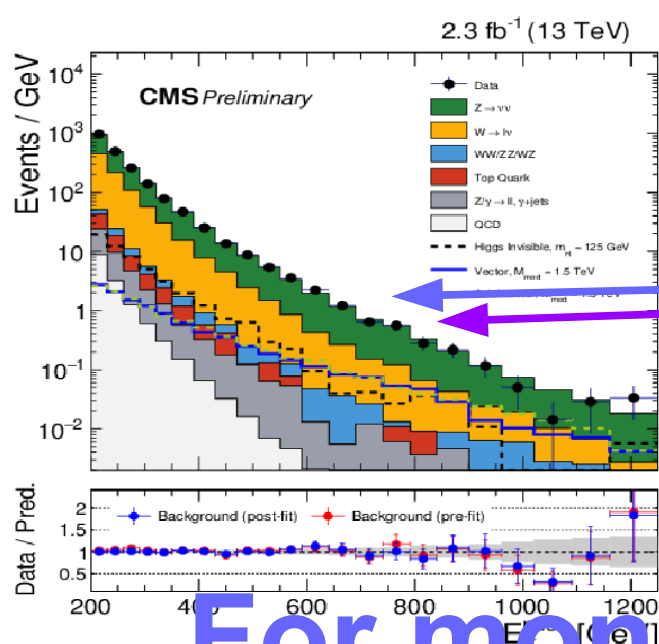


Signal

CRs: γ +jets

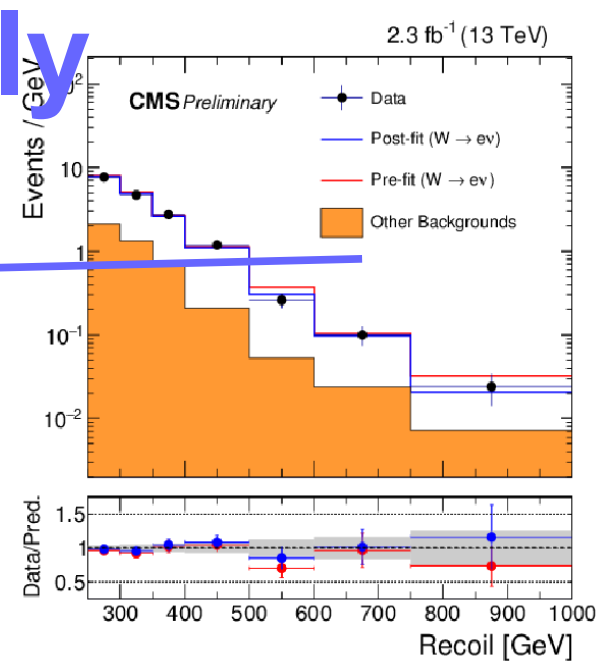
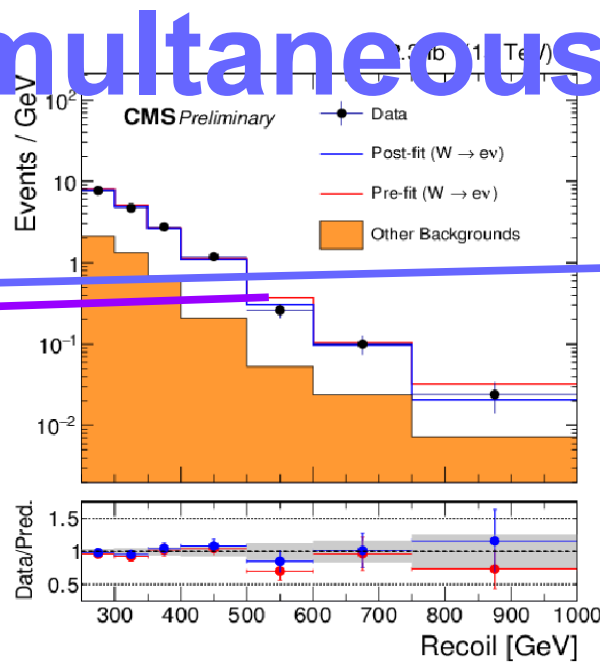
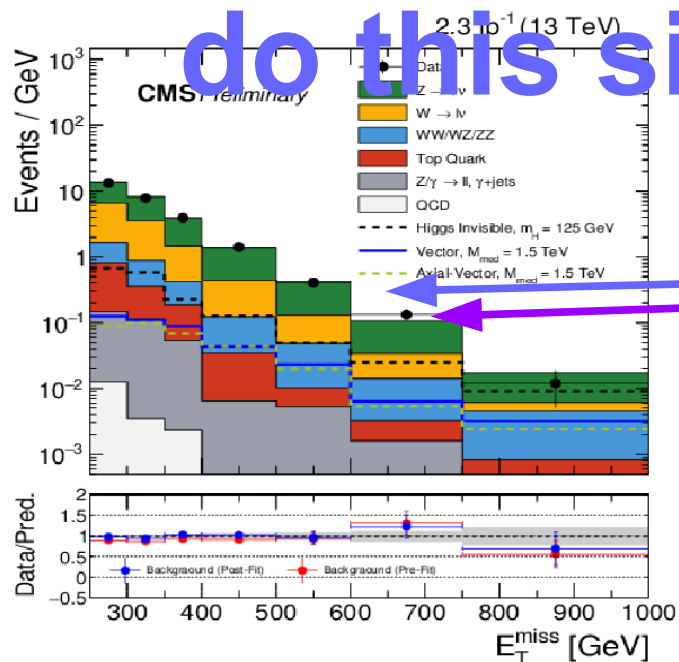
+ W +

Monojet



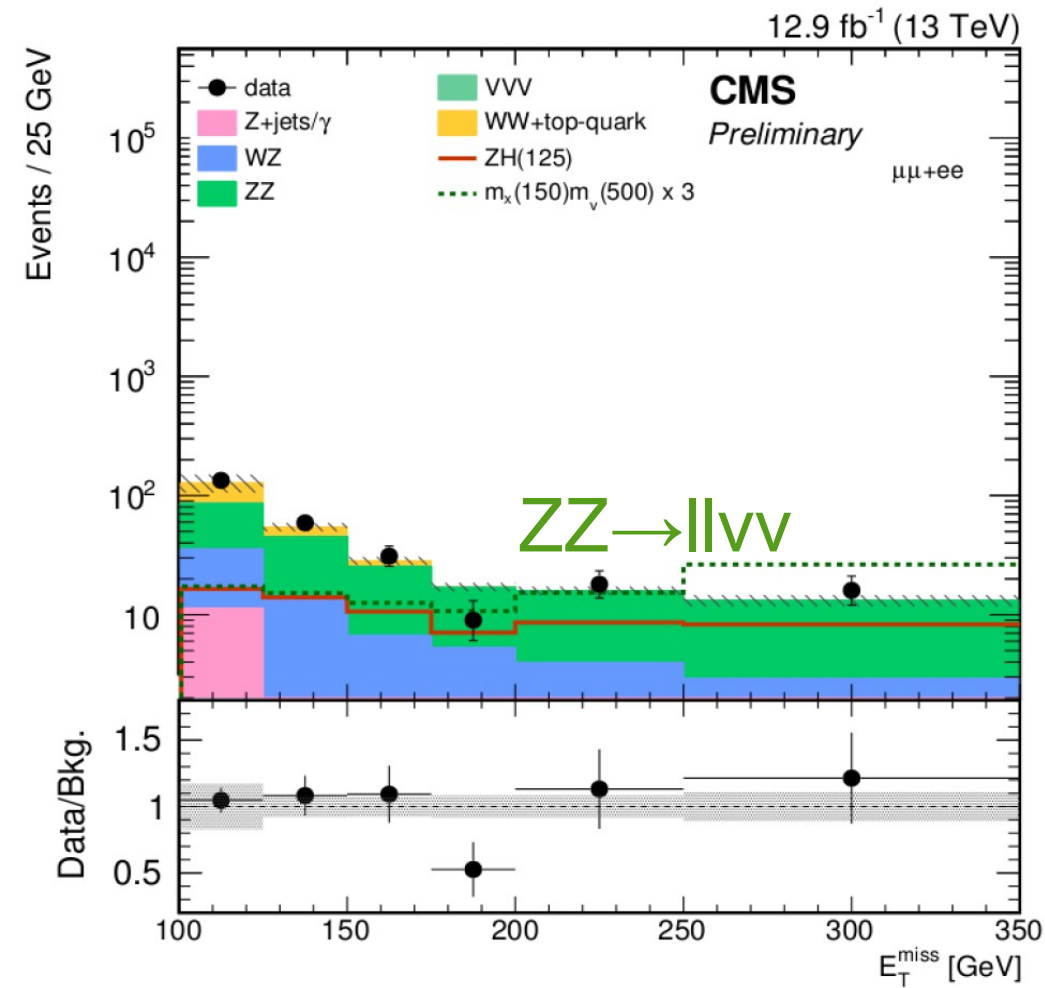
For mono-V and monojet we do this simultaneously

Mono-V

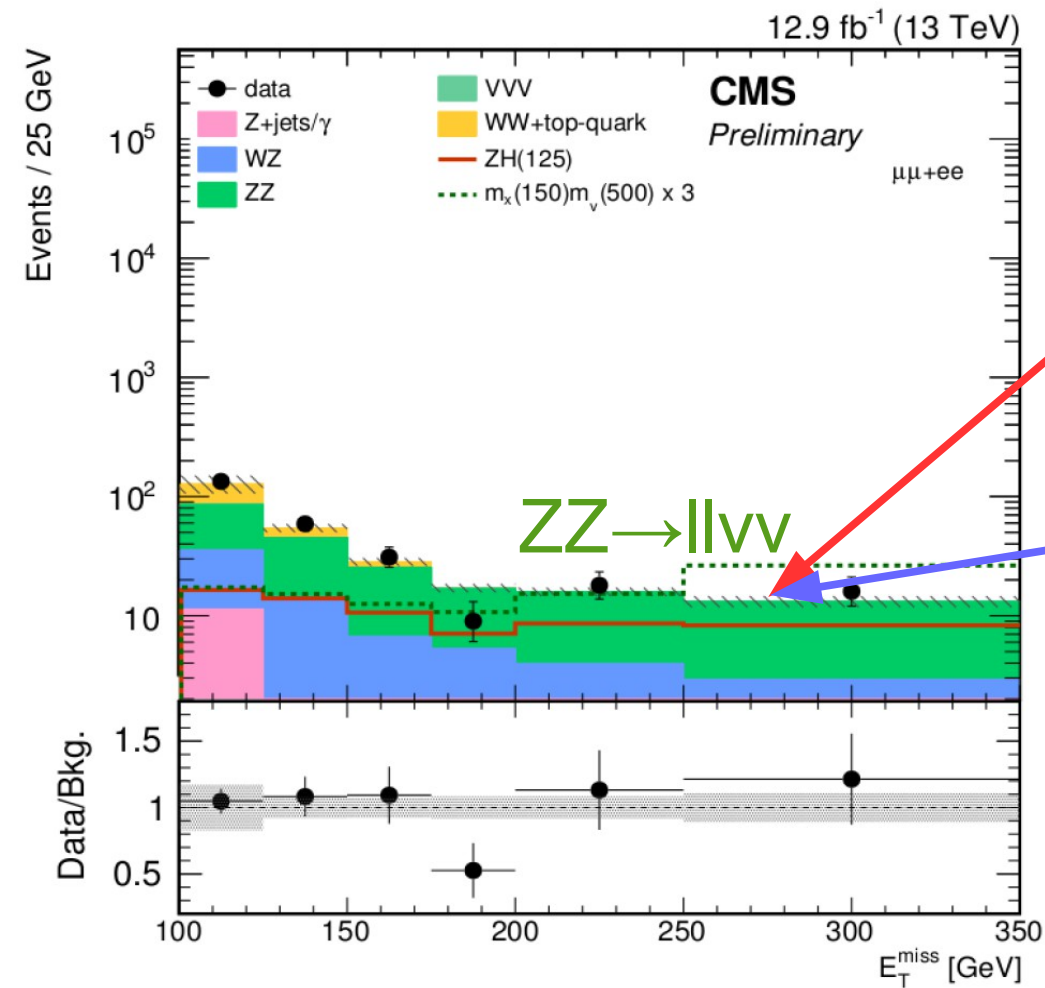


What would it be for Mono-Z?

Can modeled by



What would it be for Mono-Z?



Can modeled by

$WZ \rightarrow ll\nu$

$ZZ \rightarrow ll\nu\nu$

Simultaneously

While not done yet

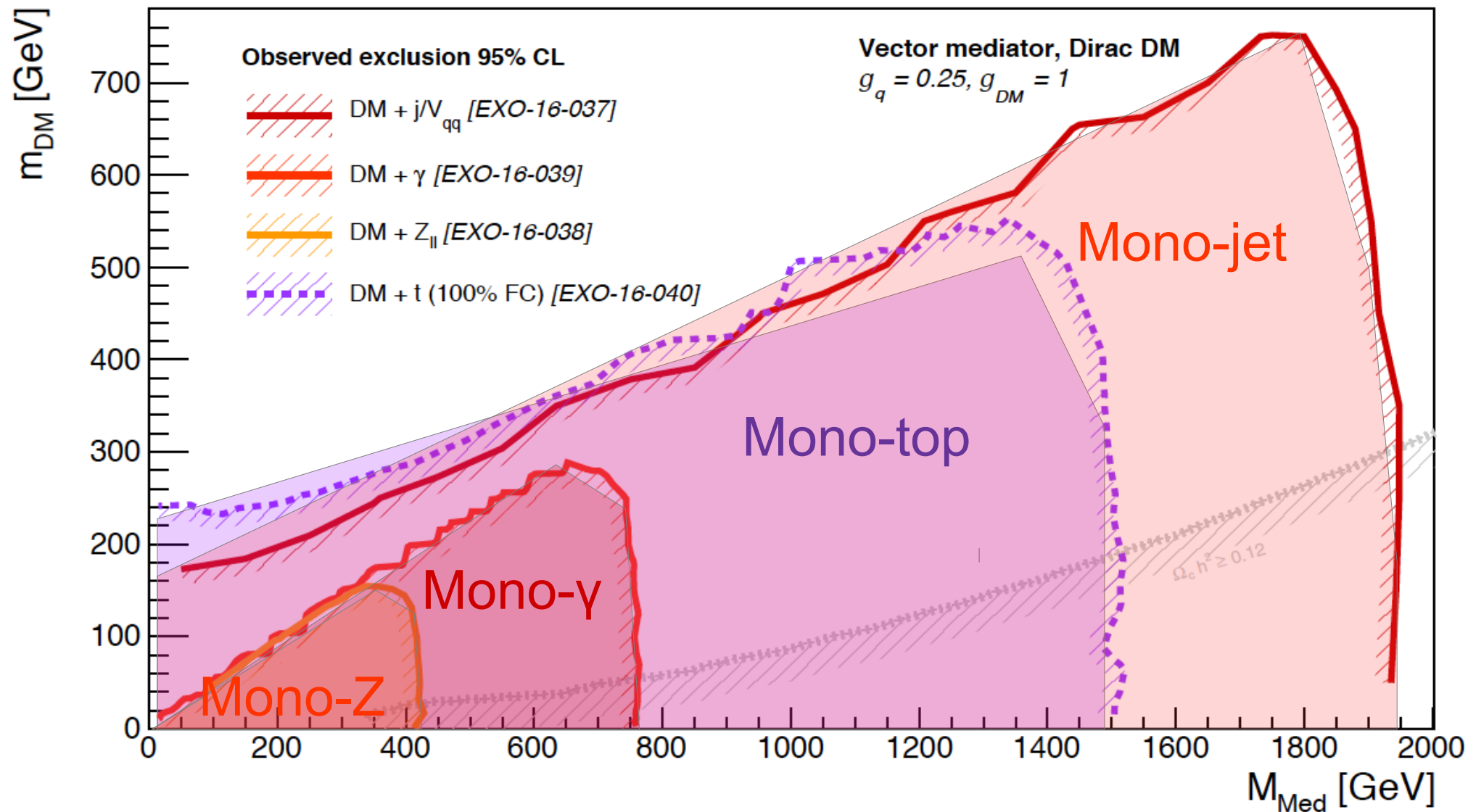
Expect to see this result coming out soon

How do the channels compare?



CMS Preliminary

Dark Matter Summary ICHEP 2016

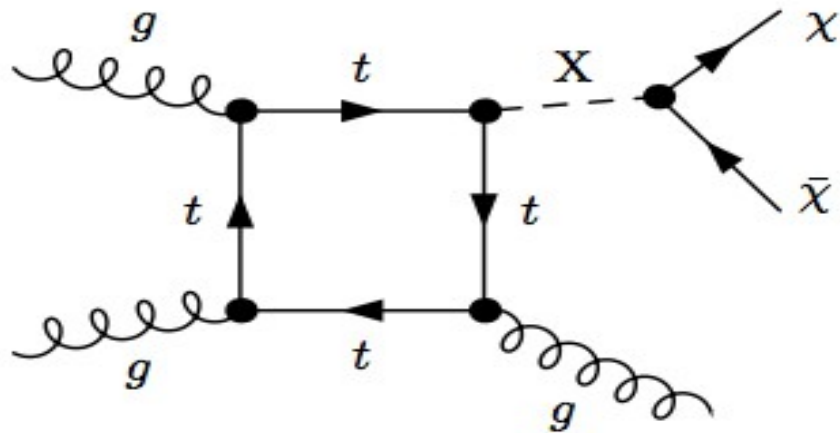


Spin 0

What can you do with Spin 0?

Basic production is gluon fusion

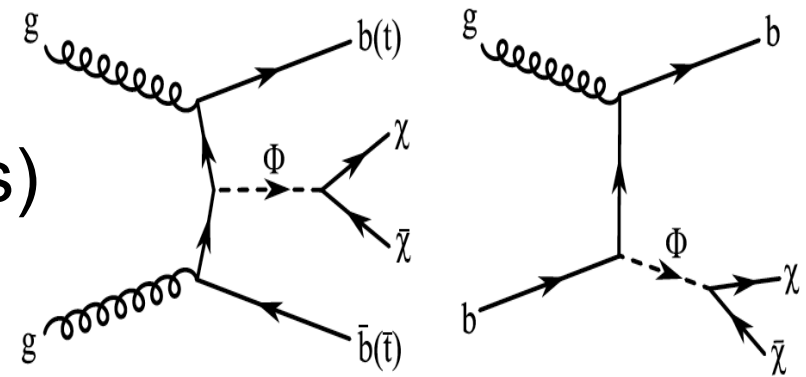
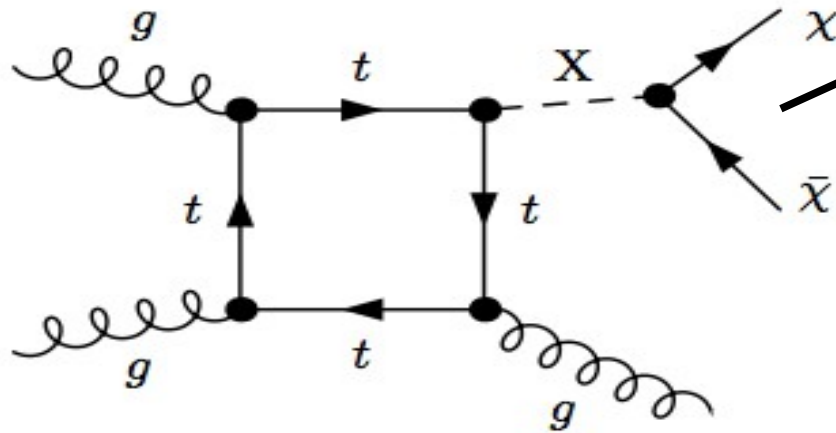
Amplitude is double for pseudoscalar mediator



What can you do with Spin 0?

Heavy flavor channels

Can be added with (same couplings)

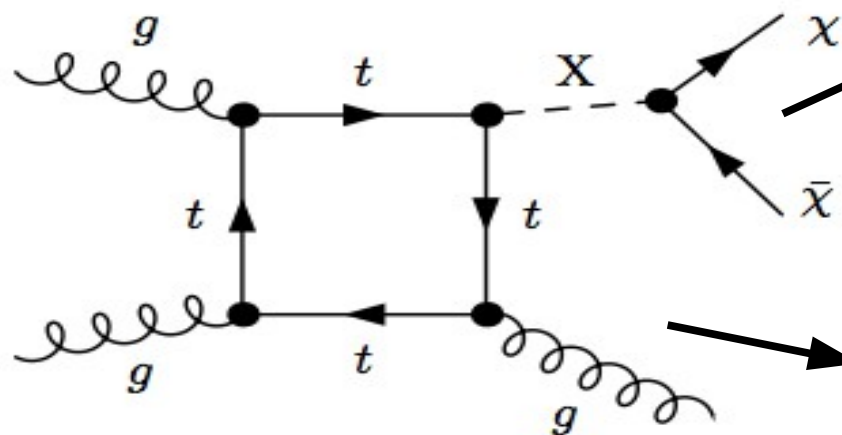


Heavy Flavor
final states

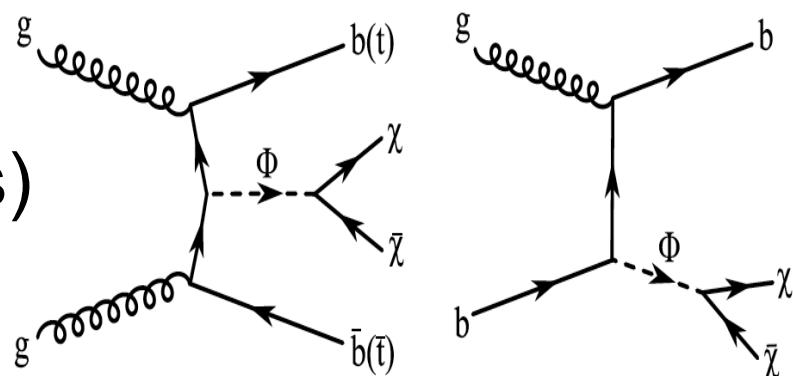
What can you do with Spin 0?

Heavy flavor channels

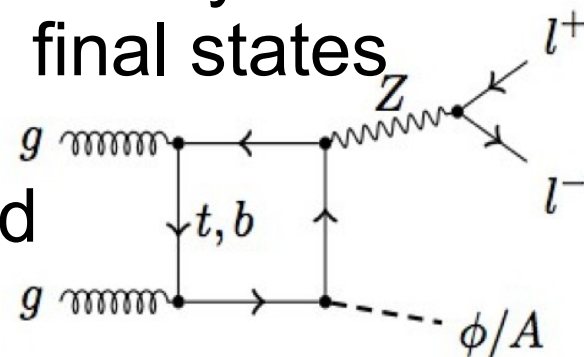
Can be added with (same couplings)



Loop induced
Z boson

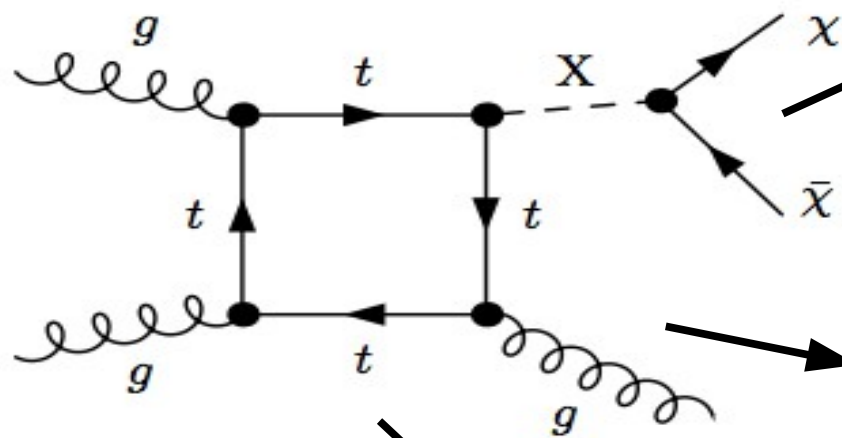


Heavy Flavor
final states

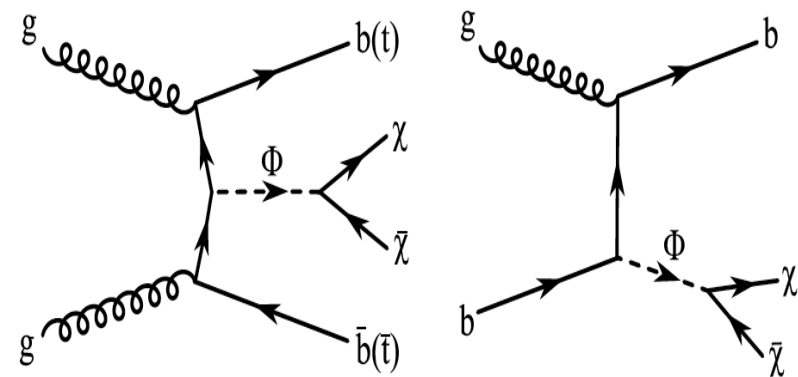


What can you do with Spin 0?

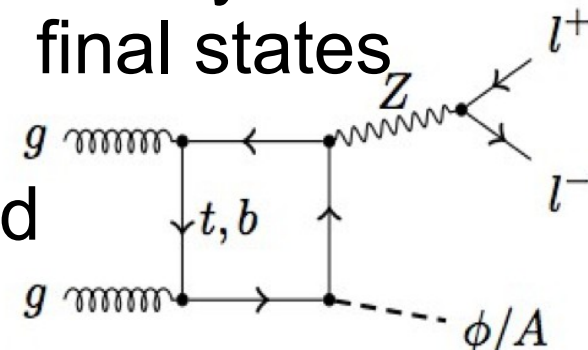
Big Assumption :
No mixing w/Higgs



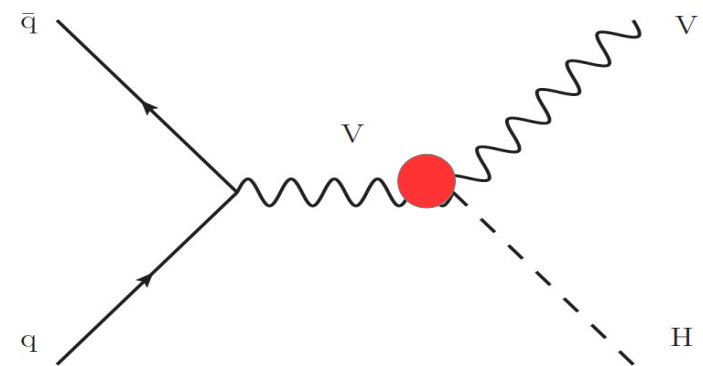
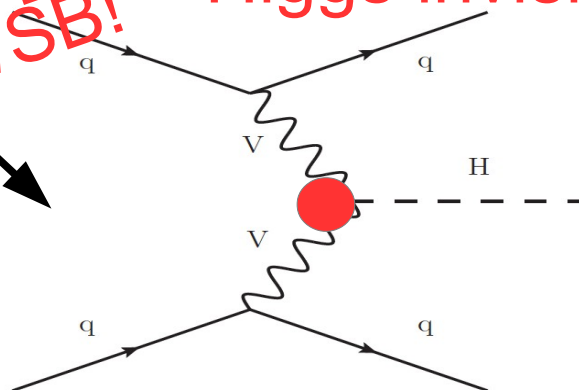
Loop induced
Z boson



Heavy Flavor
final states

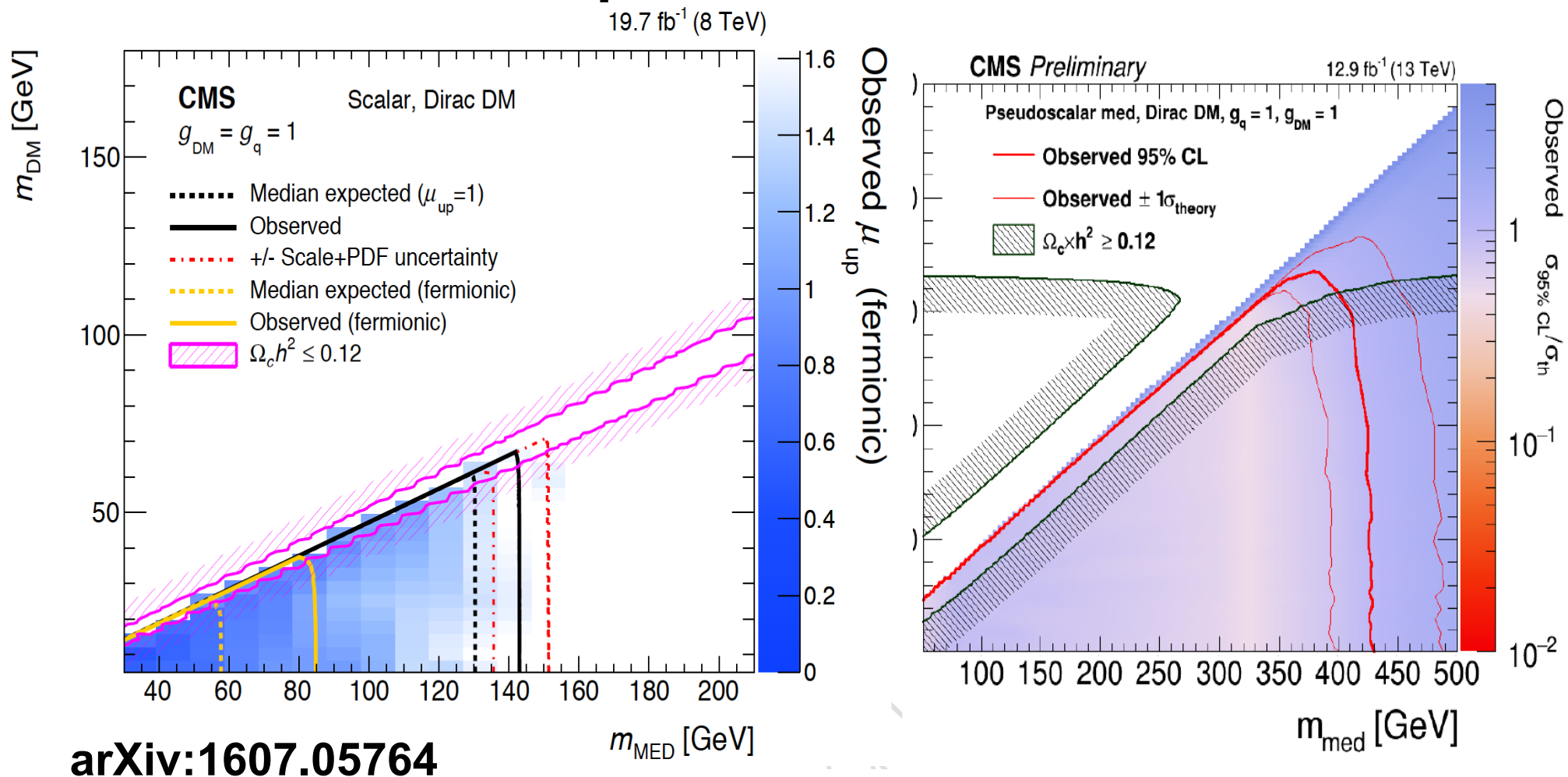


Applying EWSB! Higgs invisible or Scalar w/EWSB



Spin 0
No EWSB

Spin 0 Bounds: No EWSB

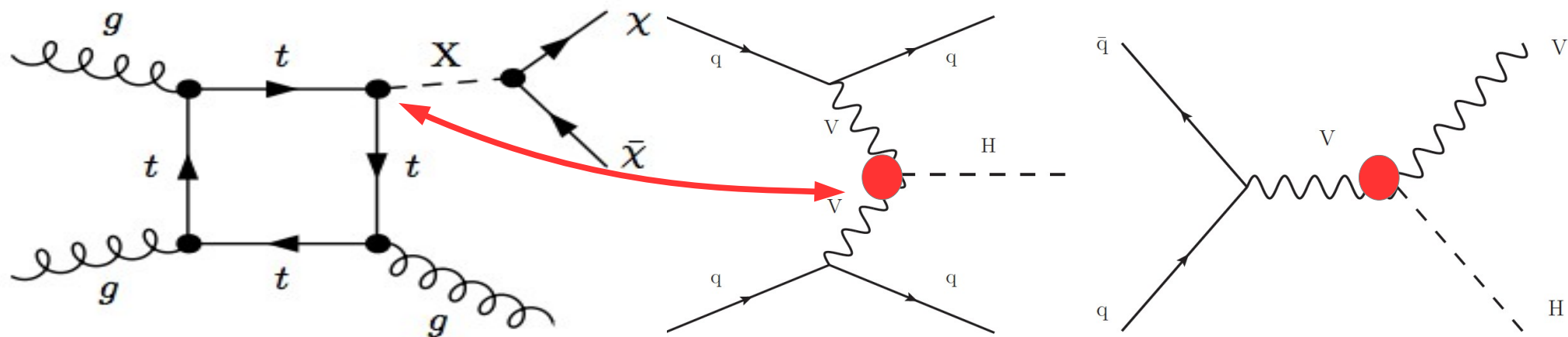


For scalar the 8TeV result slightly more sensitive
 Expect next set of results to exclude much more

Spin 0
With EWSB

Issue with adding EWSB

- Whats the right ratio of vector boson couplings
 - Do we have to deal with mixing with the Higgs?



Simplest model (not necessarily realistic) :
Higgs, however other models exist