A global view on the Higgs self coupling

Planck 2017, Warsaw

Thibaud Vantalon
DESY - IFAE

Based on:

ArXiv:1704.01953, S. Di Vita, C. Grojean, G. Panico, M.Riembau, T. Vantalon

And

Gorbahn et al ’16
arXiv:1607.03773 [hep-ph]

Degrassi et al ’16
arXiv:1607.04251 [hep-ph]

Bizon et al ’16
arXiv:1610.05771 [hep-ph]

The Higgs potential

$$V_{\mathrm{SM}} = \frac{1}{2} m_h^2 + \lambda_3^{\mathrm{SM}} h^3 + \lambda_4^{\mathrm{SM}} h^4$$

$$\lambda_3^{\mathrm{SM}} = \frac{m_h^2}{2v} \qquad\qquad \lambda_4^{\mathrm{SM}} = \frac{m_h^2}{8v^2}$$

Standard model Higgs potential depend on only 2 parameters and is precisely measured

Direct measurements of h^3 and h^4 are challenging but an important consistency check.
- Stability of EW vacuum
- Baryogenesis through first order phase transition?

h^3 challenging to measure at LHC

h^4 out of reach of LHC

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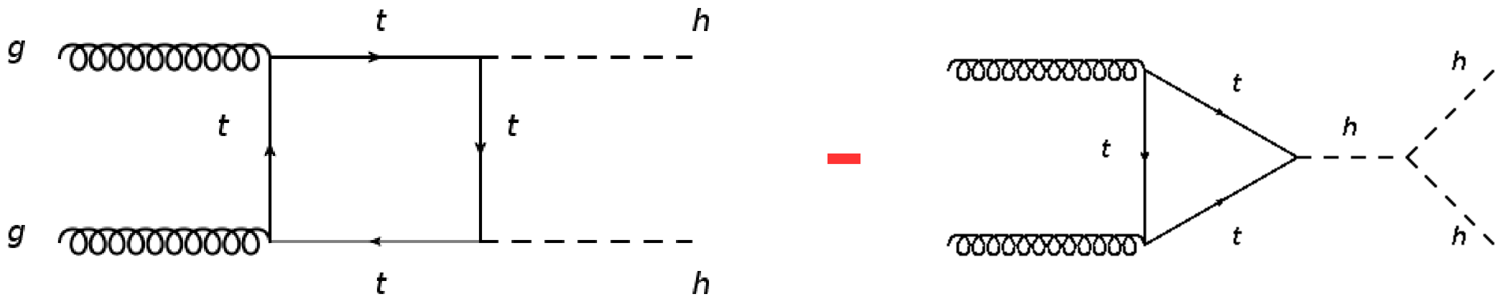
Double Higgs production



Negative interference decrease cross section by 50%

$$\frac{\sigma(pp \to hh)}{\sigma(pp \to h)} \sim 10^{-3}$$

Need a trade off between cleanness and statistic

$$\mathrm{Br}(h \to b\bar{b}) \times \mathrm{Br}(h \to \gamma\gamma) \sim 60\% \times 0.1\%$$

HL-LHC @ 3 ab^{-1} $\kappa_\lambda \in [-0.8\,,\,8.8]$ ATL-PHYS_PUB_2017-001

Idea, since the bounds are so loose and trilinear enter at NLO in single Higgs process

Can single Higgs process help?

McCullough, 1312.3322
Gorbahn, Haisch 1607.03773
Degrassi, et al. 1607.04251
Bizon, et al. 1610.05771
Degrassi, et al. 1702.01737

LHC from discovery to high precision

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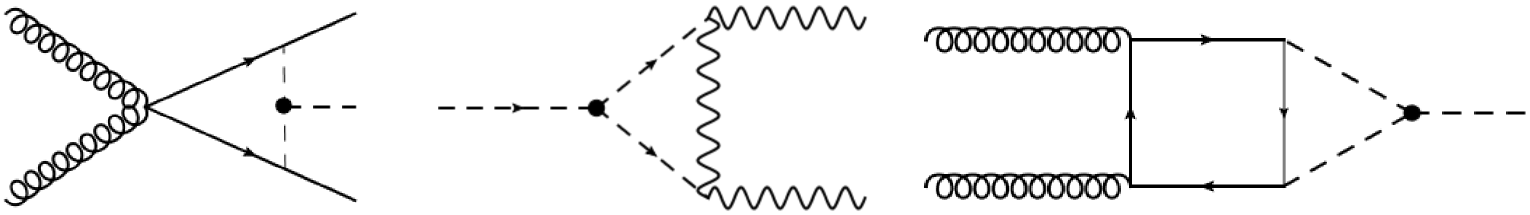
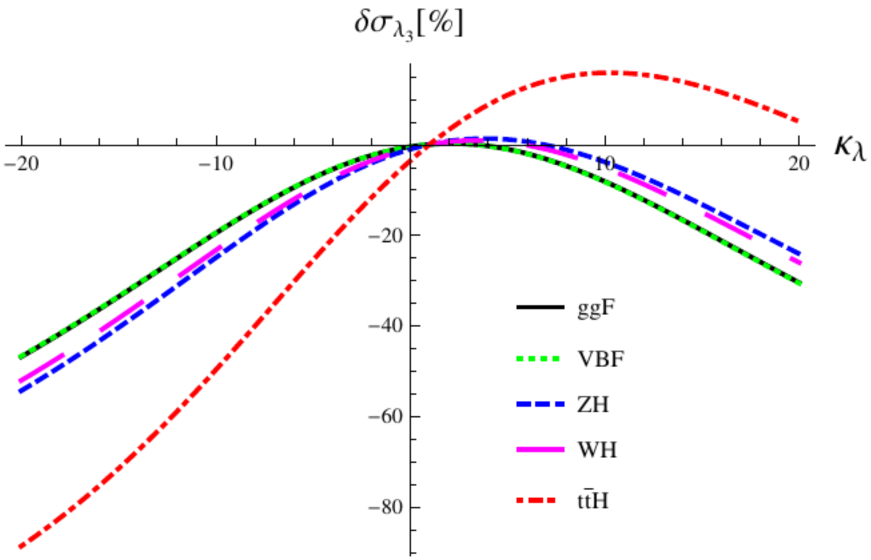
The trilinear coupling enter at loop level in single Higgs observables





Degrassi, et al. 1607.04251

Only κ_λ deviate from SM:

$$\kappa_\lambda \in [-0.7, 4.2]$$

Compared to other double Higgs expected bound in $\mathrm{HH} \to b\bar{b}\gamma\gamma$

$$\kappa_\lambda \in [0, 2.8] \cup [4.5, 6.1]$$

Azatov et al. 1502.00539

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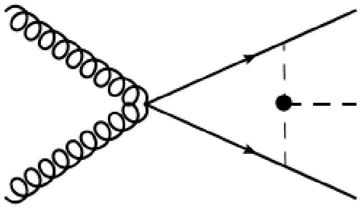

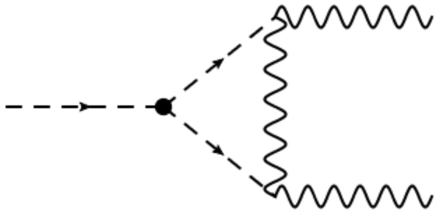

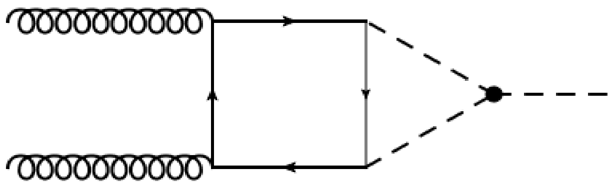

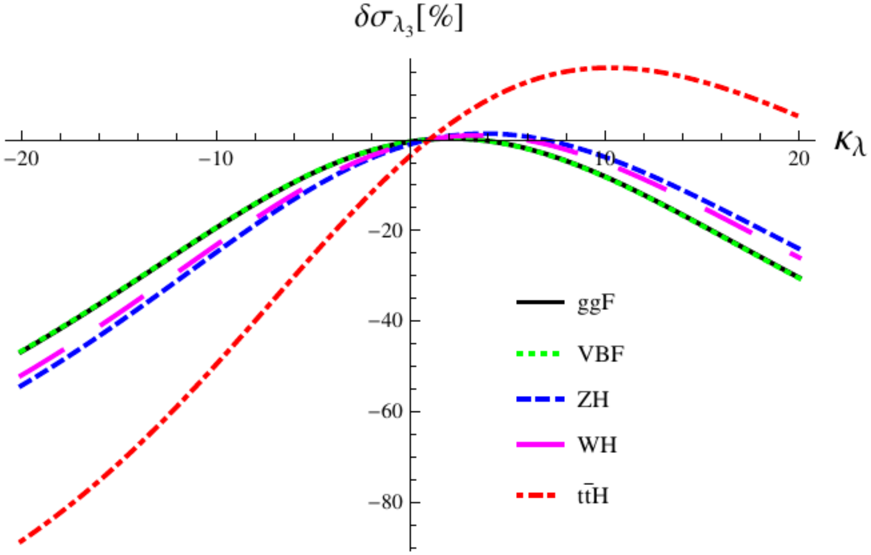

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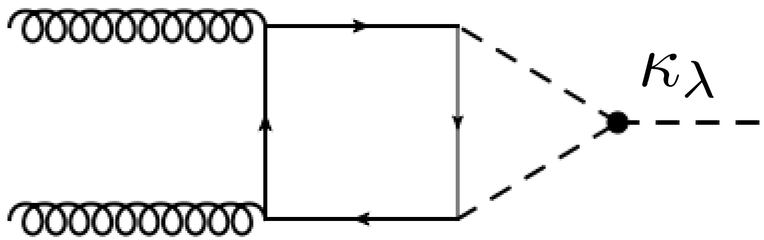
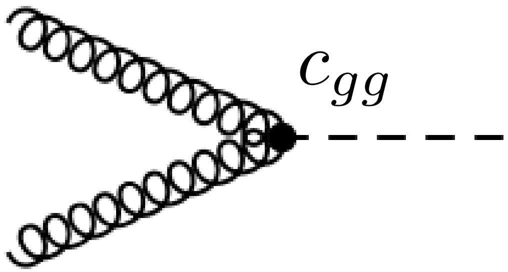
But this comparison is not fair

Other deviation

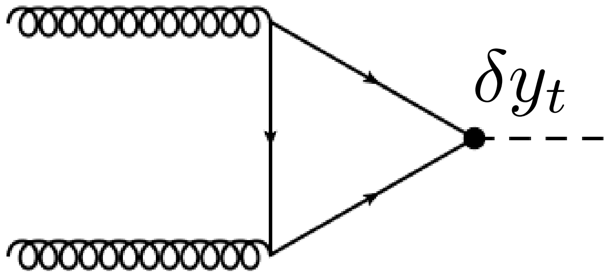
Setting on one anomalous coupling at a time is a strong assumption.



Versus



?



Is it possible to disentangle the different contributions?

The model

Parametrization of dominating BSM effects in Higgs physics using dimension 6 Lagrangian in the "Higgs basis"

Assuming flavour universality and no CP violating operator

Tested in TGC

8 (+2) Independent operators that affect Higgs physics at leading order and have not been tested in existing precision measurements

6 parameters controlling deformations of the couplings to the SM gauge bosons

$$\delta c_z \; , \;\; c_{zz} \; , \;\; c_{z\Box} \; , \;\; \hat{c}_{z\gamma} \; , \;\; \hat{c}_{\gamma\gamma} \; , \;\; \hat{c}_{gg} \; ,$$

3 related to the deformations of the fermion Yukawa's

$$\delta y_t \; , \;\; \delta y_d \; , \;\; \delta y_\tau \; ,$$

1 distortion to the Higgs trilinear self-coupling

$$\kappa_\lambda \; .$$

Inclusive observables

Global Chi squared fit of the signal strengths

We explore the sensitivity of HL-LHC at 3/ab, using the ATLAS projection.

ATL-PHYS-PUB-2014-016 **ATL-PHYS-PUB-2016-008**

ATL-PHYS-PUB-2016-018

+ Updated ggF uncertainties

Process		Combination	Theory	Experimental
	ggF	0.07	0.05	0.05
	VBF	0.22	0.16	0.15
$H \to \gamma\gamma$	$t\bar{t}H$	0.17	0.12	0.12
	WH	0.19	0.08	0.17
	ZH	0.28	0.07	0.27
	ggF	0.06	0.05	0.04
	VBF	0.17	0.10	0.14
$H \to ZZ$	$t\bar{t}H$	0.20	0.12	0.16
	WH	0.16	0.06	0.15
	ZH	0.21	0.08	0.20
$H \to WW$	ggF	0.07	0.05	0.05
	VBF	0.15	0.12	0.09
$H \to Z\gamma$	incl.	0.30	0.13	0.27
$H \to b\bar{b}$	WH	0.37	0.09	0.36
	ZH	0.14	0.05	0.13
$H \to \tau^+\tau^-$	VBF	0.19	0.12	0.15

We assume that in our EFT the dim 6 level is a good approximation.

Higher order therm can be neglected so we linearized the signal strength in the wilson coefficient

$$\mu = \frac{\sigma_i}{(\sigma_i)_{\text{SM}}} \times \frac{\text{BR}[f]}{(\text{BR}[f])_{\text{SM}}}$$

$$\approx 1 + \delta\sigma + \delta\text{BR}$$

Single Higgs observable without the trilinear

Run 1 channel, Observable = SM exactly

$$
\begin{pmatrix} \hat{c}_{gg} \\ \delta c_z \\ c_{zz} \\ c_{z\Box} \\ \hat{c}_{z\gamma} \\ \hat{c}_{\gamma\gamma} \\ \delta y_t \\ \delta y_b \\ \delta y_\tau \end{pmatrix} = \pm \left(\begin{array}{c|c} 0.07 & (0.02) \\ 0.07 & (0.01) \\ 0.64 & (0.02) \\ 0.24 & (0.01) \\ 4.94 & (0.65) \\ 0.08 & (0.02) \\ 0.09 & (0.02) \\ 0.14 & (0.03) \\ 0.17 & (0.09) \end{array} \right) \qquad \begin{bmatrix} 1 & -0.01 & -0.02 & 0.03 & 0.08 & 0.01 & \mathbf{-0.71} & 0.03 & 0.01 \\ & 1 & -0.45 & 0.36 & -0.61 & -0.33 & 0.18 & \mathbf{0.89} & 0.53 \\ & & 1 & \mathbf{-0.99} & 0.69 & 0.11 & 0.38 & -0.47 & \mathbf{-0.74} \\ & & & 1 & -0.58 & -0.23 & -0.42 & 0.42 & \mathbf{0.71} \\ & & & & 1 & -0.58 & 0.09 & -0.46 & -0.63 \\ & & & & & 1 & 0.14 & 0.04 & 0.04 \\ & & & & & & 1 & 0.25 & -0.08 \\ & & & & & & & 1 & 0.57 \\ & & & & & & & & 1 \end{bmatrix} .
$$

Global fit | Fit with only 1 wilson

$$
\left.\begin{array}{l} c_{zz} - c_{z\Box} \\ c_{zz} - \delta y_\tau \\ \hat{c}_{gg} - \delta y_t \\ \delta c_z - \delta y_b \\ \delta c_{z\Box} - \delta y_\tau \end{array}\right\} \text{Very correlated}
$$

$\longrightarrow$ **Global fit is needed!**

Falkowski:1505.00046

Using 8 TeV channel

$\Delta\mu/\mu$	300 fb^{-1}		3000 fb^{-1}	
	All unc.	No theory unc.	All unc.	No theory unc.
$H \to \gamma\gamma$ (comb.)	0.13	0.09	0.09	0.04
(0j)	0.19	0.12	0.16	0.05
(1j)	0.27	0.14	0.23	0.05
(VBF-like)	0.47	0.43	0.22	0.15
(WH-like)	0.48	0.48	0.19	0.17
(ZH-like)	0.85	0.85	0.28	0.27
(ttH-like)	0.38	0.36	0.17	0.12
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[redacted]				
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[redacted]				
$H \to b\bar{b}$ (comb.)	0.26	0.26	0.14	0.12
(WH-like)	0.57	0.56	0.37	0.36
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$H \to \tau\tau$ (VBF-like)	0.21	0.18	0.19	0.15
[redacted]				

ATL-PHYS-PUB-2014-016

Inclusive observables at 8 TeV

We have 10 quantities

$$\delta\sigma \qquad\qquad \delta\mathrm{BR}$$

$$\mathrm{ggF} \quad \mathrm{ttH} \qquad\qquad b\bar{b} \quad \tau^+\tau^-$$

$$\mathrm{VBF} \quad \mathrm{WH} \quad \mathrm{ZH} \qquad\qquad \gamma\gamma \quad \mathrm{WW} \quad \mathrm{ZZ}$$

$$\mu \approx 1 + \delta\sigma + \delta\mathrm{BR}$$

Receiving modifications from 9+1 parameters

$$\hat{c}_{gg},\ \delta c_z,\ c_{zz},\ c_{z\Box},\ \hat{c}_{z\gamma},\ \hat{c}_{\gamma\gamma},\ \delta y_t,\ \delta y_b,\ \delta y_\tau,\ \kappa_\lambda$$

So, we should be able to constrain them by looking at the <u>signal strengths</u>

This is not possible

Only 9 Independent signal strength combinations (at the linear level)

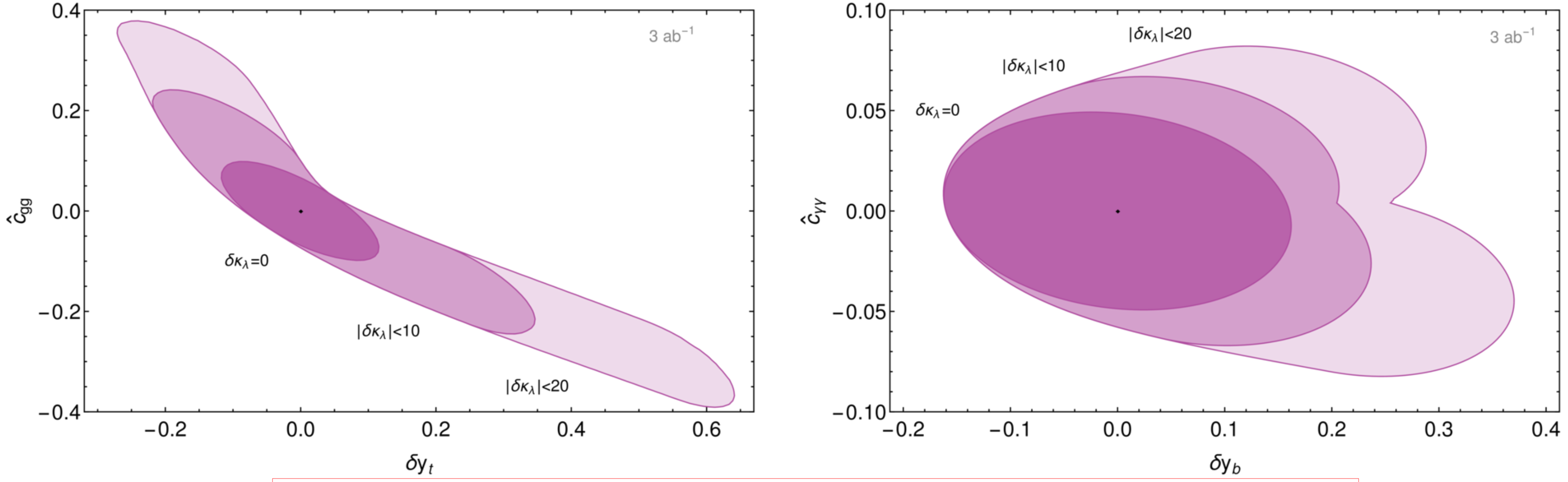
Shift in production can be compensated by opposite shift in decay

$$\delta\sigma = -\delta\mathrm{BR} \longrightarrow \boxed{\text{Unconstrained direction}}$$

Effect of the flat direction

Single Higgs without NLO effect validity

Incl. single Higgs data



Only valid for reasonable value of the trilinear coupling

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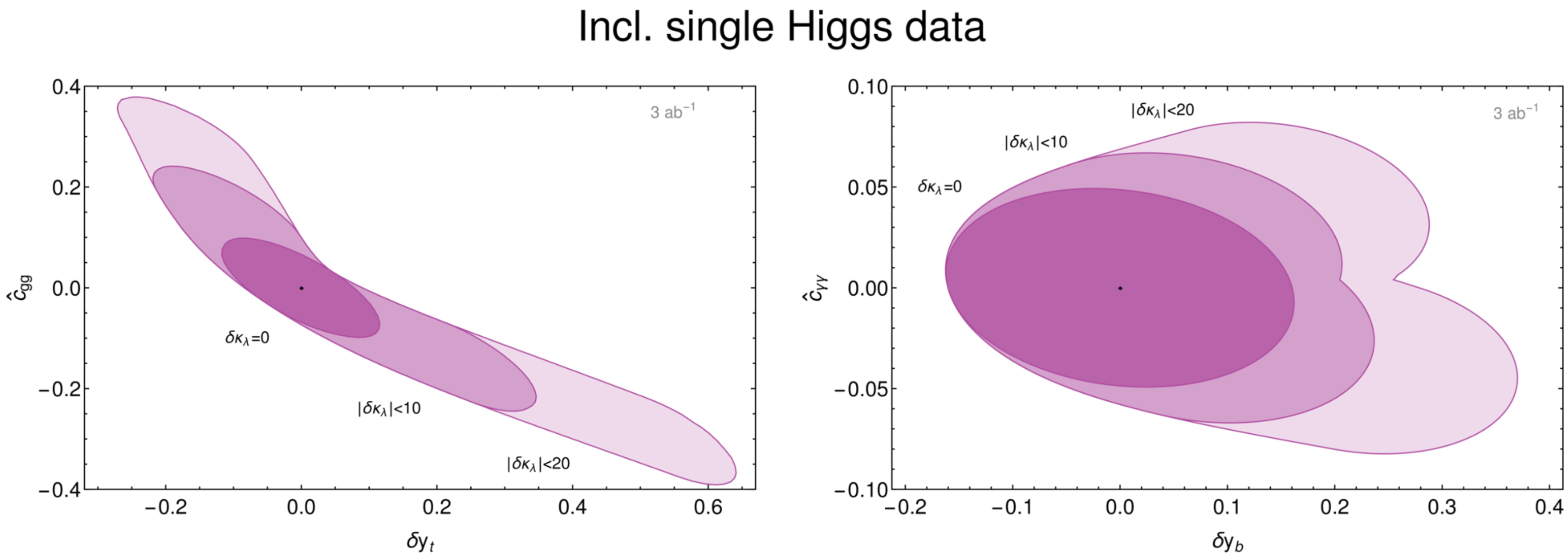

Only valid for reasonable value of the trilinear coupling

Valid in a SILH model

$$\delta c_z \sim v^2/f^2$$

$$\delta\kappa_\lambda \equiv \kappa_\lambda - 1 \sim v^2/f^2\,,$$

$$f \sim \frac{m*}{g*}$$

$$\delta c_z \sim \delta\kappa_\lambda$$

This is true for a broad class of model

A counter example

May not be valid for Higgs portal

$$\mathcal{L} \supset \theta g_* m_* H^\dagger H \varphi - \frac{m_*^4}{g_*^2} V(g_* \varphi / m_*)$$

Will generate:

$$\delta c_z \sim \theta^2 g_*^2 \frac{v^2}{m_3^2} \qquad \delta \kappa_\lambda \sim \theta^3 g_*^4 \frac{1}{\lambda_3^{\mathrm{SM}}} \frac{v^2}{m^2}$$

With a typical tuning of $\Delta \sim \frac{\theta^2 g_*^2}{\lambda_3^{SM}}$

Perturbative expansion $\varepsilon \equiv \frac{\theta g_*^2 v^2}{m_*^2} \ll 1$

$$\theta \simeq 1,\ g_* \simeq 3 \text{ and } m_* \simeq 2.5 \text{ TeV}$$

$$\varepsilon \simeq 0.1\,, \quad 1/\Delta \simeq 1.5\%$$

$$\boxed{\delta c_z \simeq 0.1\,, \quad \delta \kappa_\lambda \simeq 6}$$

Hard to have model with large deviation only in $\delta\kappa$ $\longrightarrow$ **Single Higgs fit valid for most model**

Inclusive observables

Way out:

Extra constraints

1 **- Higgs total width**
$ **- Compare different energies**
1 **- decay** $\mu\mu$
2 **- Anomalous triple gauge couplings(aTGCs)**
1 **- decay** $Z\gamma$
L **- Differential distributions**
1 **- Add double Higgs**

Not helping too much
See paper for detail

Inclusive observables

aTGCs

At dimension 6, the aTGCs can be written in terms of the Higgs basis parameters

$$\delta g_{1,z} = \frac{1}{2(g-g')}\left[c_{\gamma\gamma}e^2 g' + c_{z\gamma}\left(g^2-g'^2\right)g'^2 - c_{zz}\left(g^2+g'^2\right)g'^2 - c_{z\Box}\left(g^2+g'^2\right)g^2\right],$$

$$\delta\kappa_\gamma = -\frac{g^2}{2}\left(c_{\gamma\gamma}\frac{e^2}{g^2+g'^2} + c_{z\gamma}\frac{g^2-g'^2}{g^2+g'^2} - c_{zz}\right)$$

$H \to Z\gamma$

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$H\to WW$ (comb.)	0.13	0.08	0.11	0.05
(0j)	0.18	0.09	0.16	0.05
(1j)	0.30	0.18	0.26	0.10
(VBF-like)	0.21	0.20	0.15	0.09
$H\to Z\gamma$ (incl.)	0.46	0.44	0.30	0.27
$H\to b\bar{b}$ (comb.)	0.26	0.26	0.14	0.12
(WH-like)	0.57	0.56	0.37	0.36
(ZH-like)	0.29	0.29	0.14	0.13
$H\to\tau\tau$ (VBF-like)	0.21	0.18	0.19	0.15
$H\to\mu\mu$ (comb.)	0.39	0.38	0.16	0.12
(incl.)	0.47	0.45	0.18	0.14
(ttH-like)	0.74	0.72	0.27	0.23

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Correlation with new observables

$$
\begin{pmatrix} \hat{c}_{gg} \\ \delta c_z \\ c_{zz} \\ c_{z\Box} \\ \hat{c}_{z\gamma} \\ \hat{c}_{\gamma\gamma} \\ \delta y_t \\ \delta y_b \\ \delta y_\tau \end{pmatrix} = \pm \begin{pmatrix} 0.07 & (0.02) \\ 0.07 & (0.01) \\ 0.64 & (0.02) \\ 0.24 & (0.01) \\ 4.94 & (0.65) \\ 0.08 & (0.02) \\ 0.09 & (0.02) \\ 0.14 & (0.03) \\ 0.17 & (0.09) \end{pmatrix}
\begin{bmatrix}
1 & -0.01 & -0.02 & 0.03 & 0.08 & 0.01 & \mathbf{-0.71} & 0.03 & 0.01 \\
 & 1 & -0.45 & 0.36 & -0.61 & -0.33 & 0.18 & \mathbf{0.89} & 0.53 \\
 & & 1 & \mathbf{-0.99} & 0.69 & 0.11 & 0.38 & -0.47 & \mathbf{-0.74} \\
 & & & 1 & -0.58 & -0.23 & -0.42 & 0.42 & \mathbf{0.71} \\
 & & & & 1 & -0.58 & 0.09 & -0.46 & -0.63 \\
 & & & & & 1 & 0.14 & 0.04 & 0.04 \\
 & & & & & & 1 & 0.25 & -0.08 \\
 & & & & & & & 1 & 0.57 \\
 & & & & & & & & 1
\end{bmatrix} .
$$

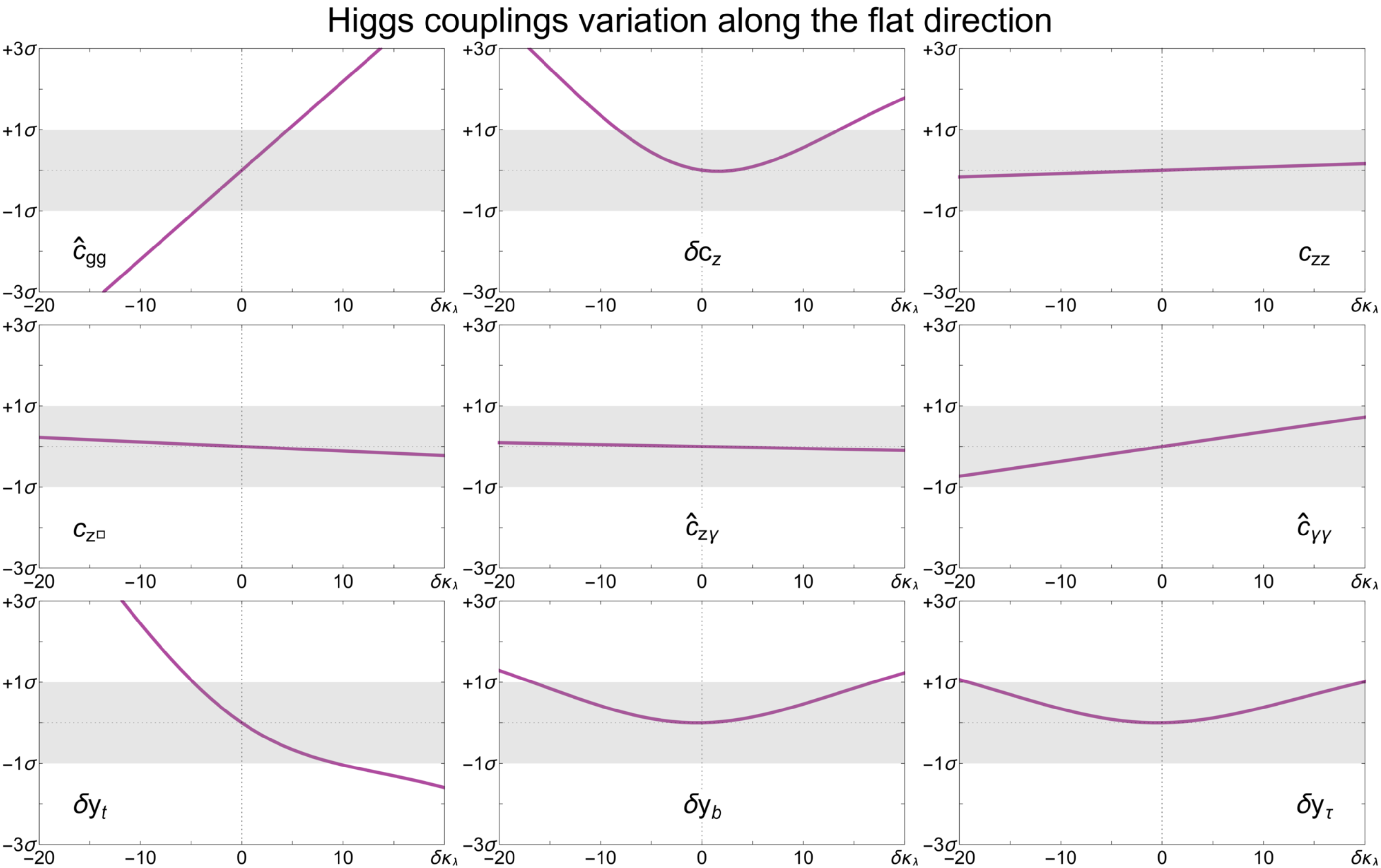
New channels help the correlations

$$
\begin{pmatrix} \hat{c}_{gg} \\ \delta c_z \\ c_{zz} \\ c_{z\Box} \\ \hat{c}_{z\gamma} \\ \hat{c}_{\gamma\gamma} \\ \delta y_t \\ \delta y_b \\ \delta y_\tau \end{pmatrix} = \pm \begin{pmatrix} 0.07 & (0.02) \\ 0.05 & (0.01) \\ 0.05 & (0.02) \\ 0.02 & (0.01) \\ 0.09 & (0.09) \\ 0.03 & (0.02) \\ 0.08 & (0.02) \\ 0.12 & (0.03) \\ 0.11 & (0.09) \end{pmatrix}
\begin{bmatrix}
1 & 0.04 & -0.01 & -0.01 & 0.04 & 0.31 & -0.76 & 0.05 & 0.02 \\
 & 1 & -0.07 & -0.26 & 0.01 & 0.01 & 0.36 & 0.88 & 0.27 \\
 & & 1 & -0.87 & 0.13 & 0.20 & 0.03 & -0.07 & -0.06 \\
 & & & 1 & -0.09 & -0.09 & -0.09 & -0.17 & 0.08 \\
 & & & & 1 & 0.05 & -0.02 & -0.02 & -0.03 \\
 & & & & & 1 & -0.32 & -0.19 & -0.12 \\
 & & & & & & 1 & 0.50 & 0.28 \\
 & & & & & & & 1 & 0.36 \\
 & & & & & & & & 1
\end{bmatrix} .
$$

Linear fit become a good approximation
(If we can constrain the trilinear!)

The flat direction

Value of all the couplings in function of $\delta\kappa_\lambda$ such that
All the $\delta\mu=0$



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Value of all the couplings in function of $\delta\kappa_\lambda$ such that

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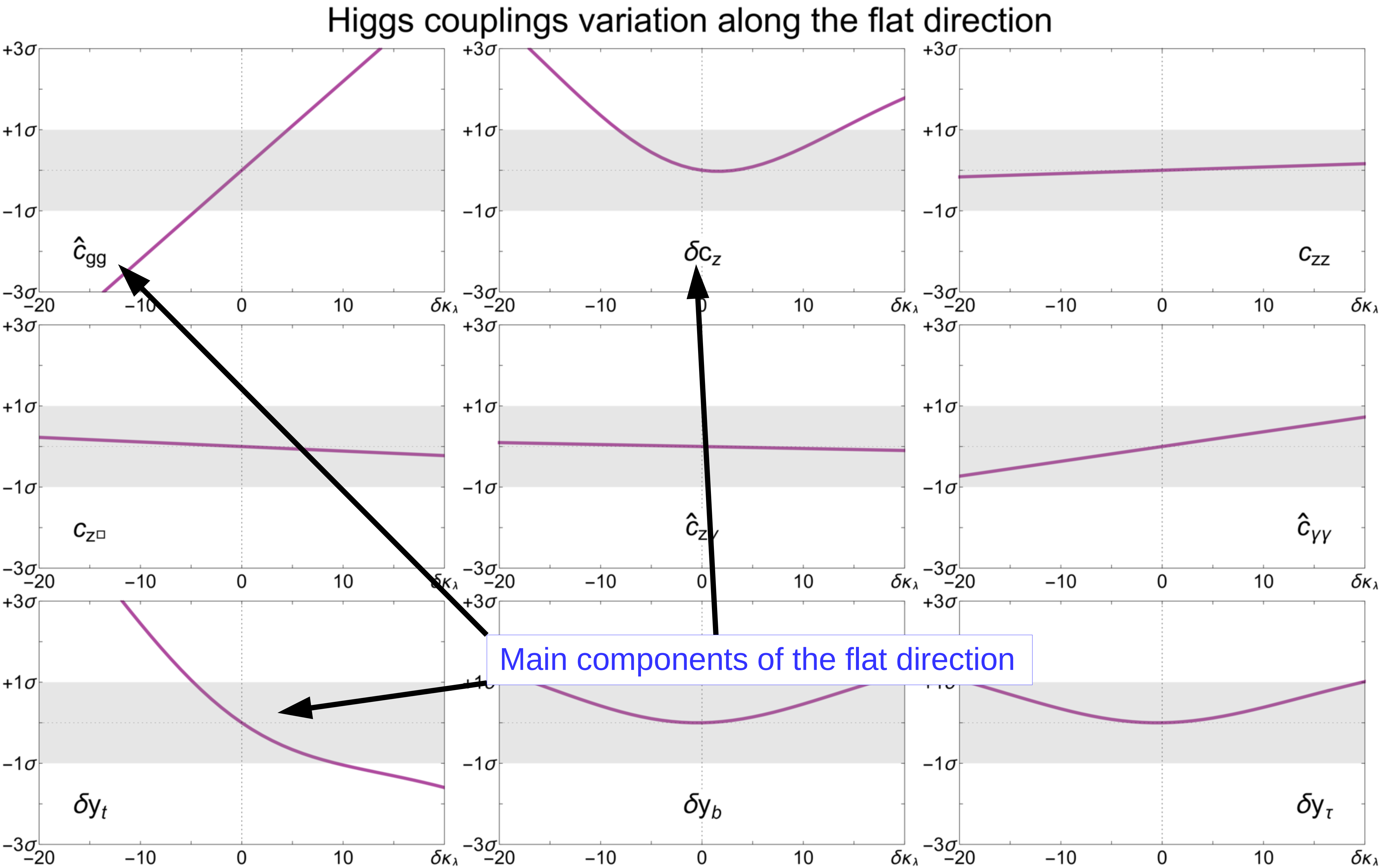
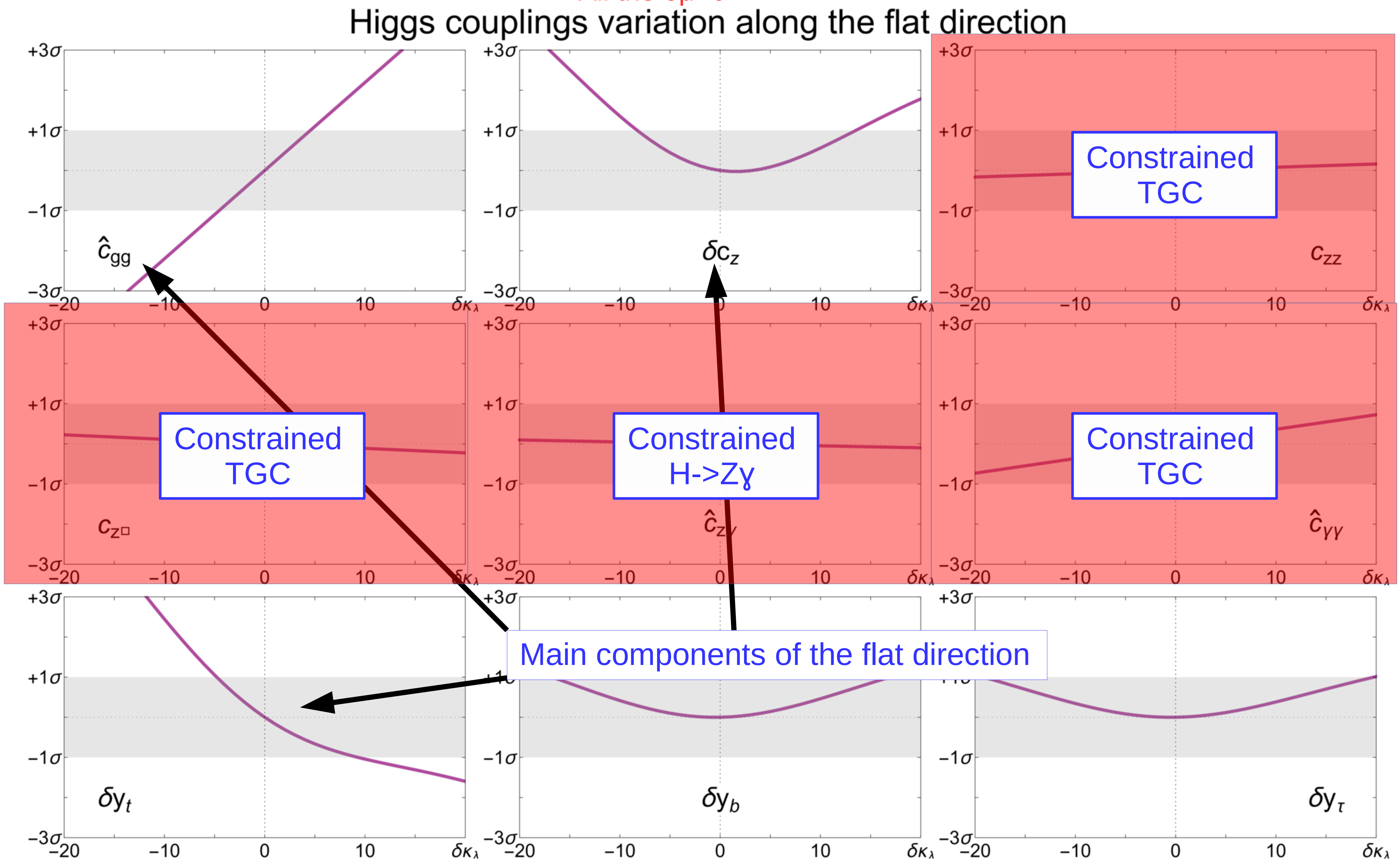


What we constrained

Value of all the couplings in function of $\delta\kappa_\lambda$ such that

All the $\delta\mu=0$



Inclusive observables

TGCs

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$$\delta g_{1,z} = \frac{1}{2(g-g')}\left[c_{\gamma\gamma}e^2 g' + c_{z\gamma}\left(g^2-g'^2\right)g'^2 - c_{zz}\left(g^2+g'^2\right)g'^2 - c_{z\Box}\left(g^2+g'^2\right)g^2\right],$$

$$\delta\kappa_\gamma = -\frac{g^2}{2}\left(c_{\gamma\gamma}\frac{e^2}{g^2+g'^2} + c_{z\gamma}\frac{g^2-g'^2}{g^2+ \text{[illegible]}}\right)$$

$H \to Z\gamma$

Do not constrain the flat direction

$\Delta\mu/\mu$	300 fb^{-1}		3000 fb^{-1}	
	All unc.	[illegible]	All unc.	No theory unc.
$H \to \gamma\gamma$ (comb.)	[illegible]	[illegible]	0.09	0.04
(0j)	[illegible]	0.12	0.16	0.05
[illegible]	[illegible]	0.14	0.23	0.05
[illegible]	[illegible]	0.43	0.22	0.15
[illegible]	0.48	0.48	0.19	0.17
[illegible]	0.85	0.85	0.28	0.27
[illegible]	0.38	0.36	0.17	0.12
[illegible] ZZ (comb.)	0.11	0.07	0.09	0.04
(VH-like)	0.35	0.34	0.13	0.12
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(ZH-like)	0.29	0.29	0.14	0.13
$H \to \tau\tau$ (VBF-like)	0.21	0.18	0.19	0.15
$H \to \mu\mu$ (comb.)	0.39	0.38	0.16	0.12
(incl.)	0.47	0.45	0.18	0.14
(ttH-like)	0.74	0.72	0.27	0.23

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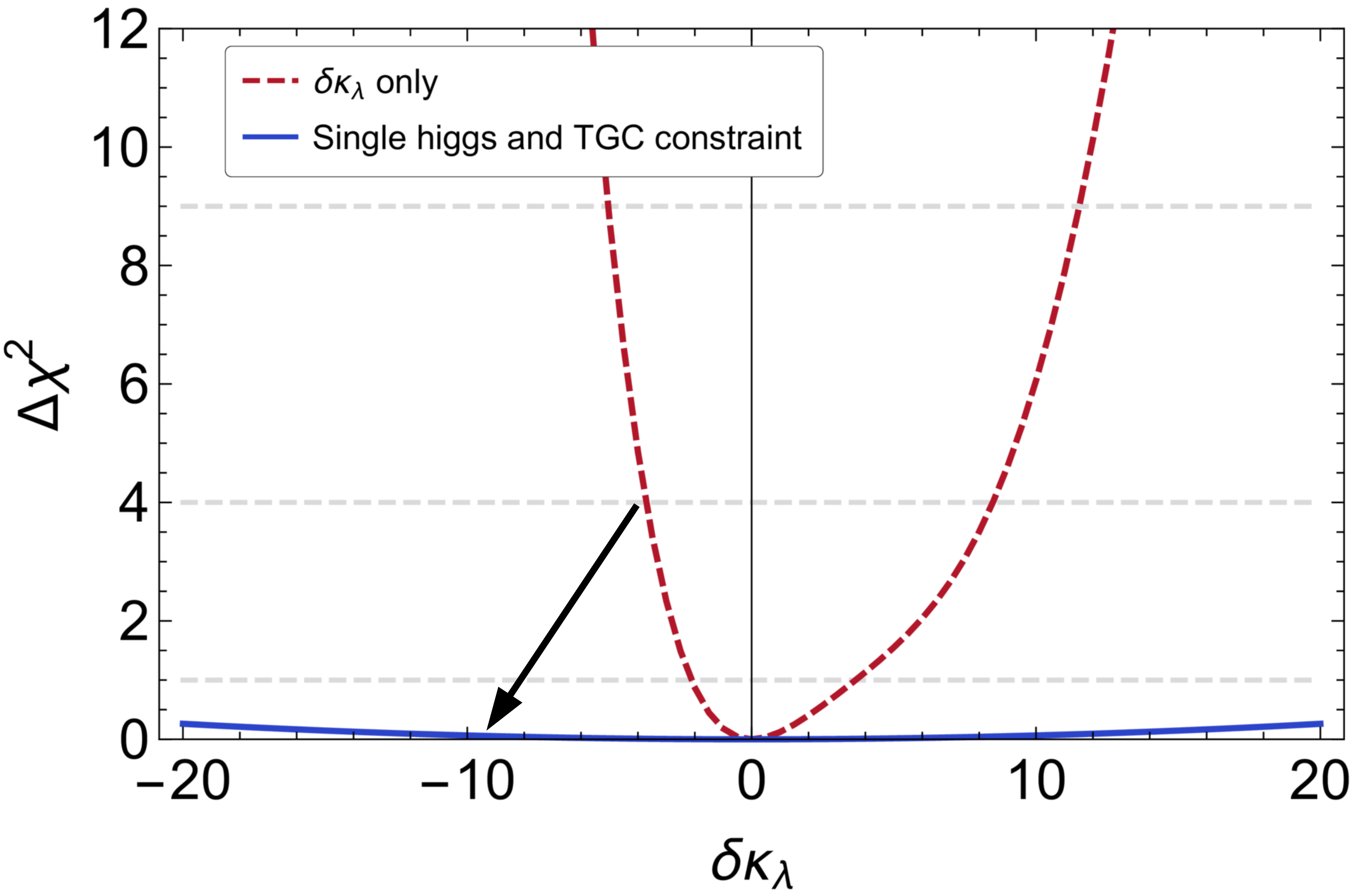
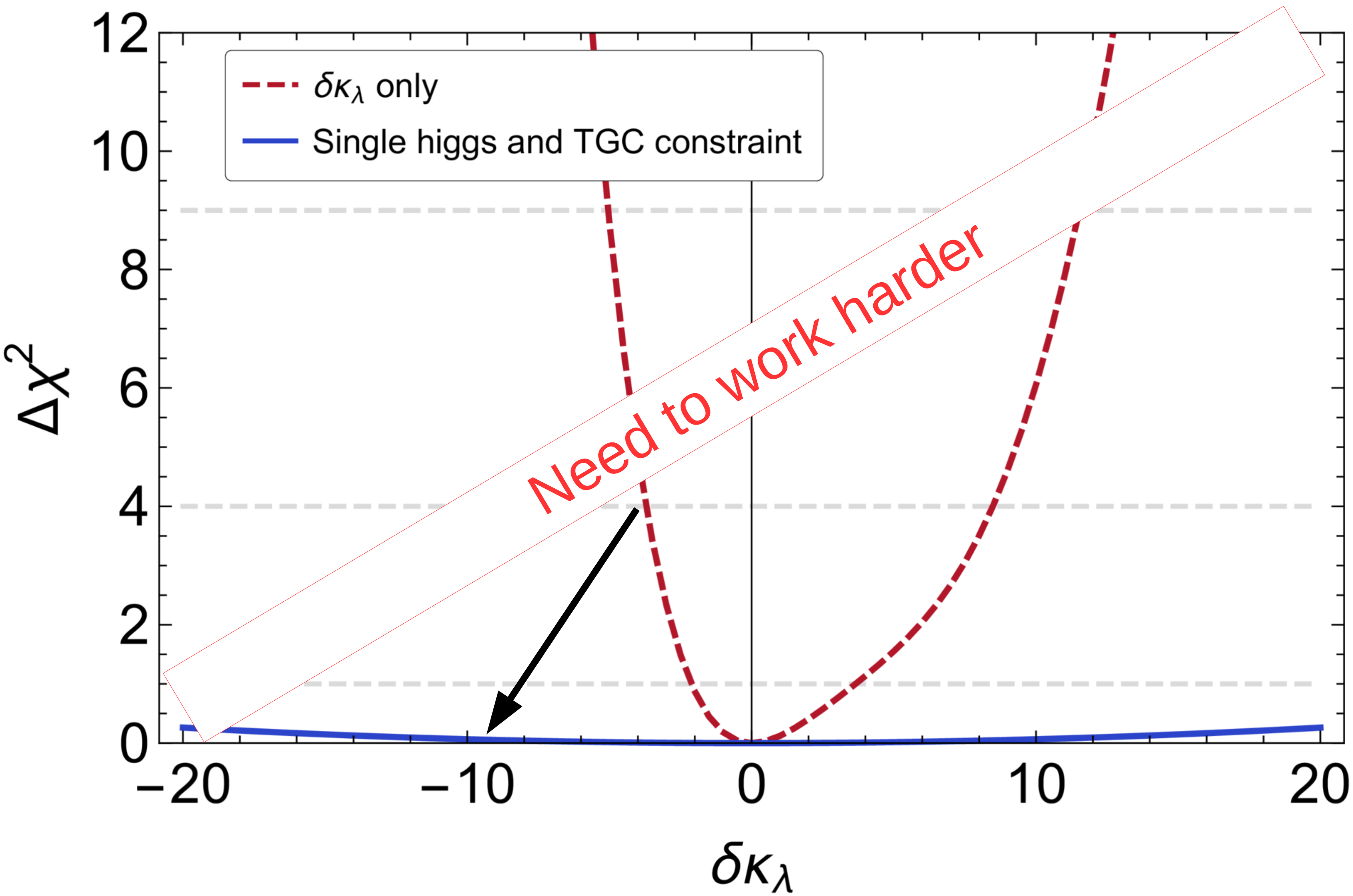
Not enough constraints



Not enough constraints



Differential Observables

<u>Rough</u> analysis looking at the prospects of differential observables

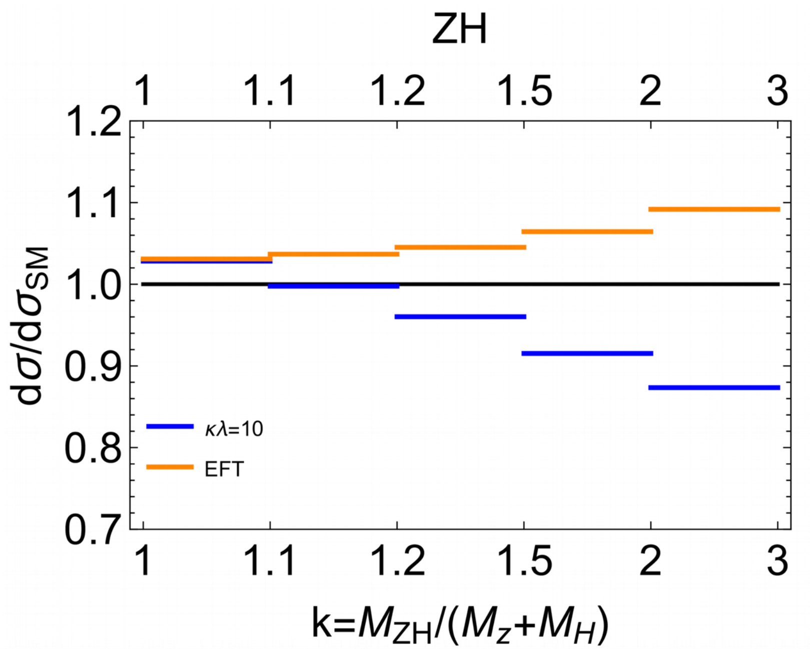
Cross section in each bin in terms of the EFT parameters computed using MadGraph.

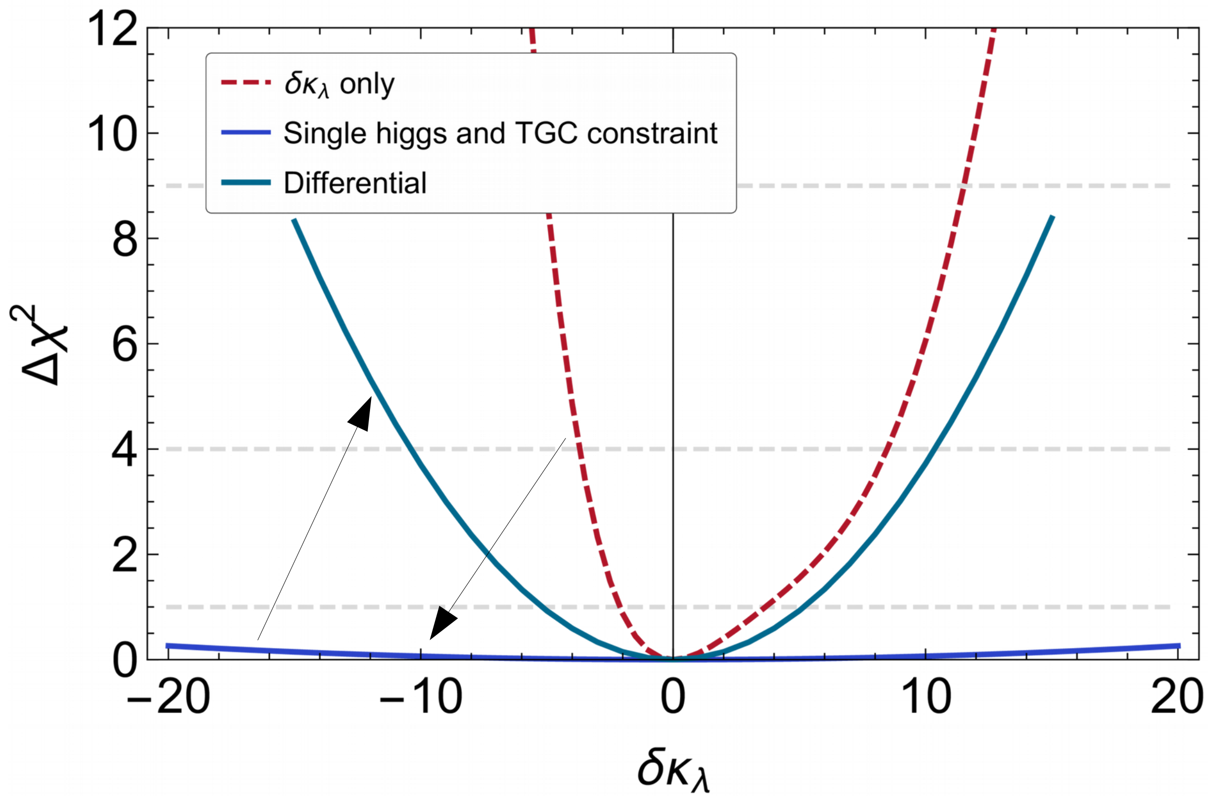
Dependence on Higgs trilinear computed in Degrassi, et al. 1607.04251

Restore some power to the method, may be seen as complement to double Higgs

Maybe other differential observable can be more powerful

$$\delta\kappa_\lambda \in [-5.4, 5.4]$$

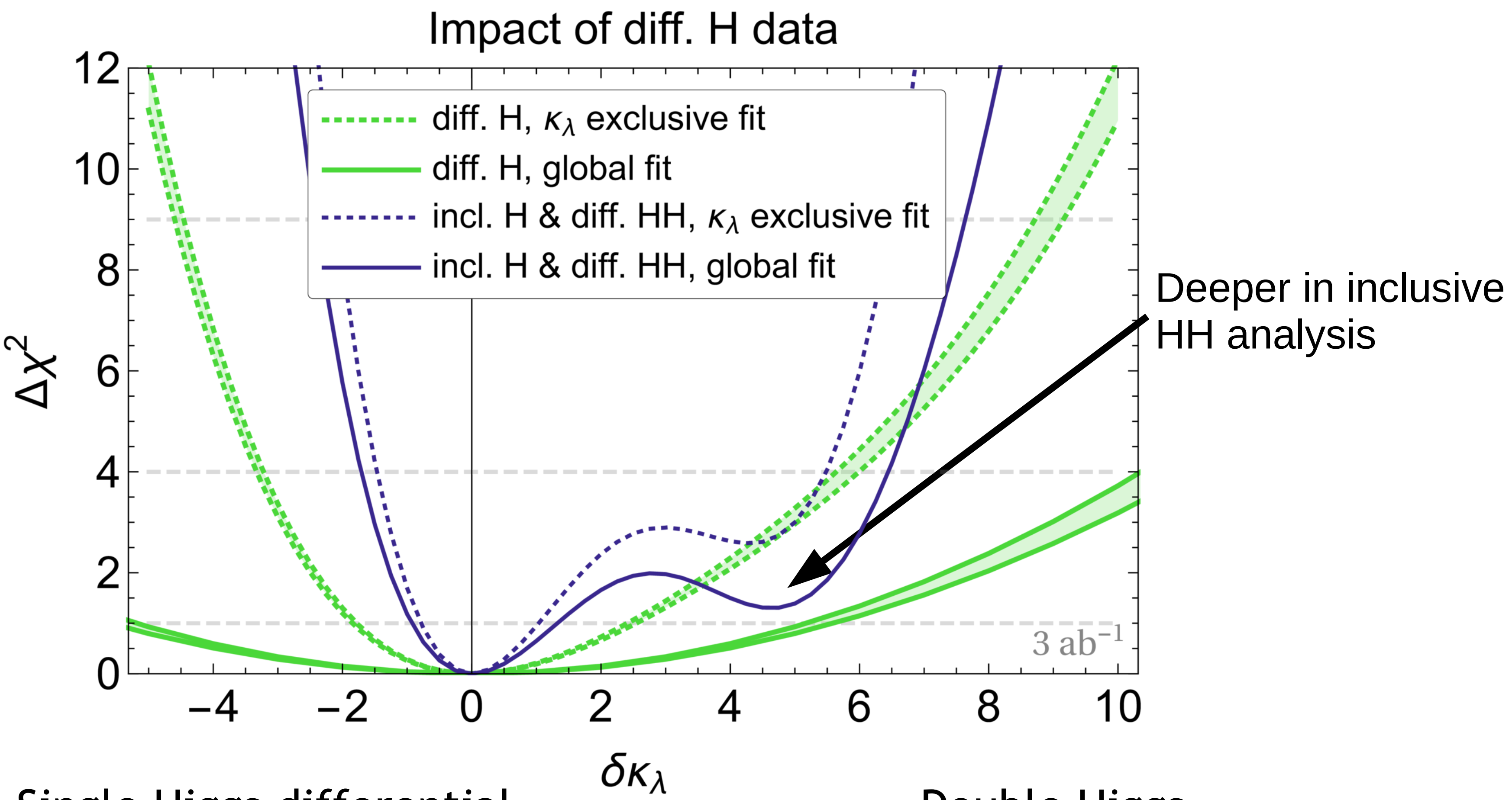




Differential Observables versus double Higgs

Double Higgs analysis more powerful

It also solves the flat direction issue in single Higgs



Single Higgs differential

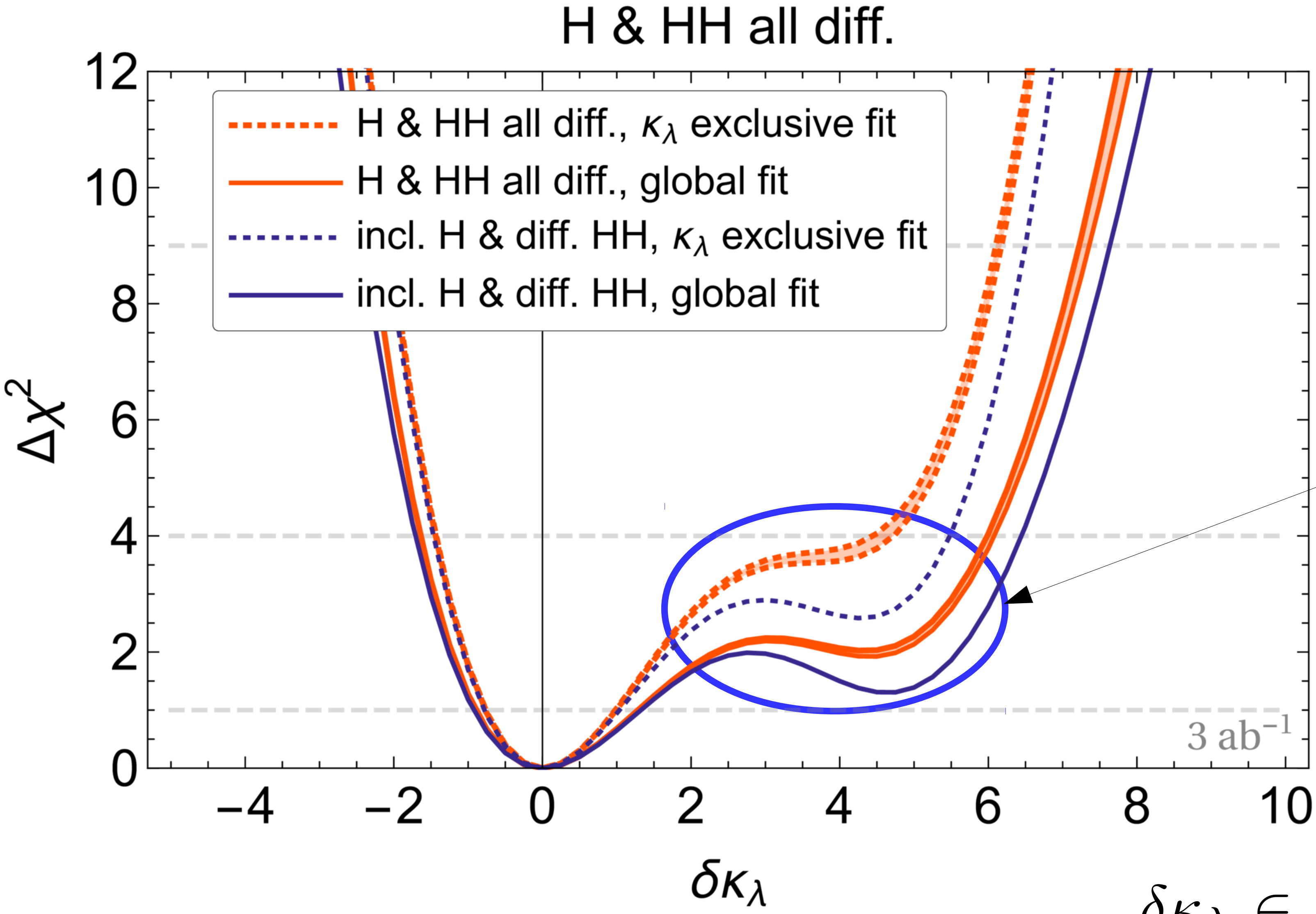
$\delta\kappa_\lambda \in [-5.4, 5.4]$

Double Higgs

$\delta\kappa_\lambda \in [-0.9, 1.2]$

Single and double Higgs together

Single Higgs help little

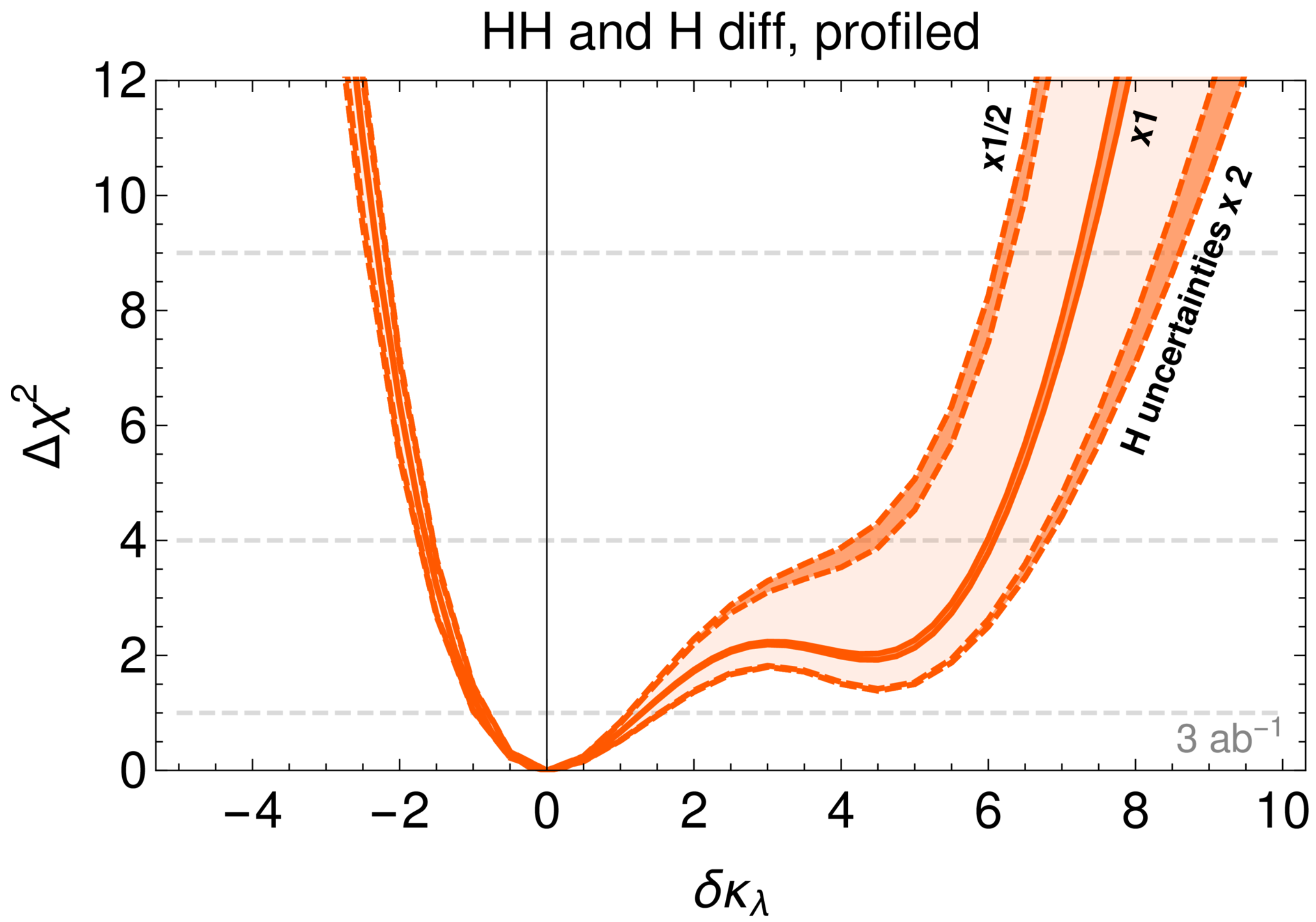


Single Higgs help lifting this minimum
(More clear for Inclusive double Higgs)

$$\delta\kappa_\lambda \in [-0.88, 1.24]$$

Robustness of the analysis

Sensibility to single Higgs uncertainties



Conclusion

- **Single Higgs observables are a complementary source of information for the Higgs trilinear.**
- **At the inclusive level the trilinear corrections to single Higgs observables introduce a flat direction in the global fit.**
- **This flat direction degrades the precision achievable on the wilson coefficients. Some control on the trilinear is needed to solve this issue.**
- **Double Higgs is still the best way to extract Higgs trilinear and to restore the control over single Higgs fit.**
- **Most promising way to remove the flat direction without using double Higgs is to use differential distribution. More work in this direction is needed.**

More results in
ArXiv:1704.01953

For preliminary results on the trilinear extraction at future lepton collider see Jiayin Gu talk tomorrow morning

Thank you

Our parametrisation:

Parametrization of dominating BSM effects in Higgs couplings couplings:

$$
\begin{aligned}
\mathcal{L}^{\rm NP} \supset \frac{h}{v}\Bigg[& \delta c_w \frac{g^2 v^2}{2} W^+_\mu W^{-\mu} + \delta c_z \frac{(g^2+g'^2)v^2}{4} Z_\mu Z^\mu \\
& + c_{ww} \frac{g^2}{2} W^+_{\mu\nu} W_{-\mu\nu} + c_{w\Box} g^2 \left(W^+_\mu \partial_\nu W_{+\mu\nu} + \text{h.c.} \right) \\
& + \hat{c}_{\gamma\gamma} \frac{e^2}{4\pi^2} A_{\mu\nu} A^{\mu\nu} + c_{z\Box} g^2 Z_\mu \partial_\nu Z^{\mu\nu} + c_{\gamma\Box} g g' Z_\mu \partial_\nu A^{\mu\nu} \\
& + c_{zz} \frac{g^2+g'^2}{4} Z_{\mu\nu} Z^{\mu\nu} + \hat{c}_{z\gamma} \frac{e\sqrt{g^2+g'^2}}{2\pi^2} Z_{\mu\nu} A^{\mu\nu} \Bigg] \\
& + \frac{g_s^2}{48\pi^2} \left(\hat{c}_{gg} \frac{h}{v} + \hat{c}^{(2)}_{gg} \frac{h^2}{2v^2} \right) G_{\mu\nu} G^{\mu\nu} \\
& - \sum_f \left[m_f \left(\delta y_f \frac{h}{v} + \delta y^{(2)}_f \frac{h^2}{2v^2} \right) \bar{f}_R f_L + \text{h.c} \right] \\
& + (\kappa_\lambda - 1) \lambda_{SM} v h^3
\end{aligned}
$$

$\delta y_\tau,\ \delta y_b,\ \delta y_t$

Only enter at loop level in single Higgs observable