Controlled Flavour Changing Neutral Currents in the Higgs Sector

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In memoriam Maria

Scalars 2017, Warsaw, Poland
November 2017

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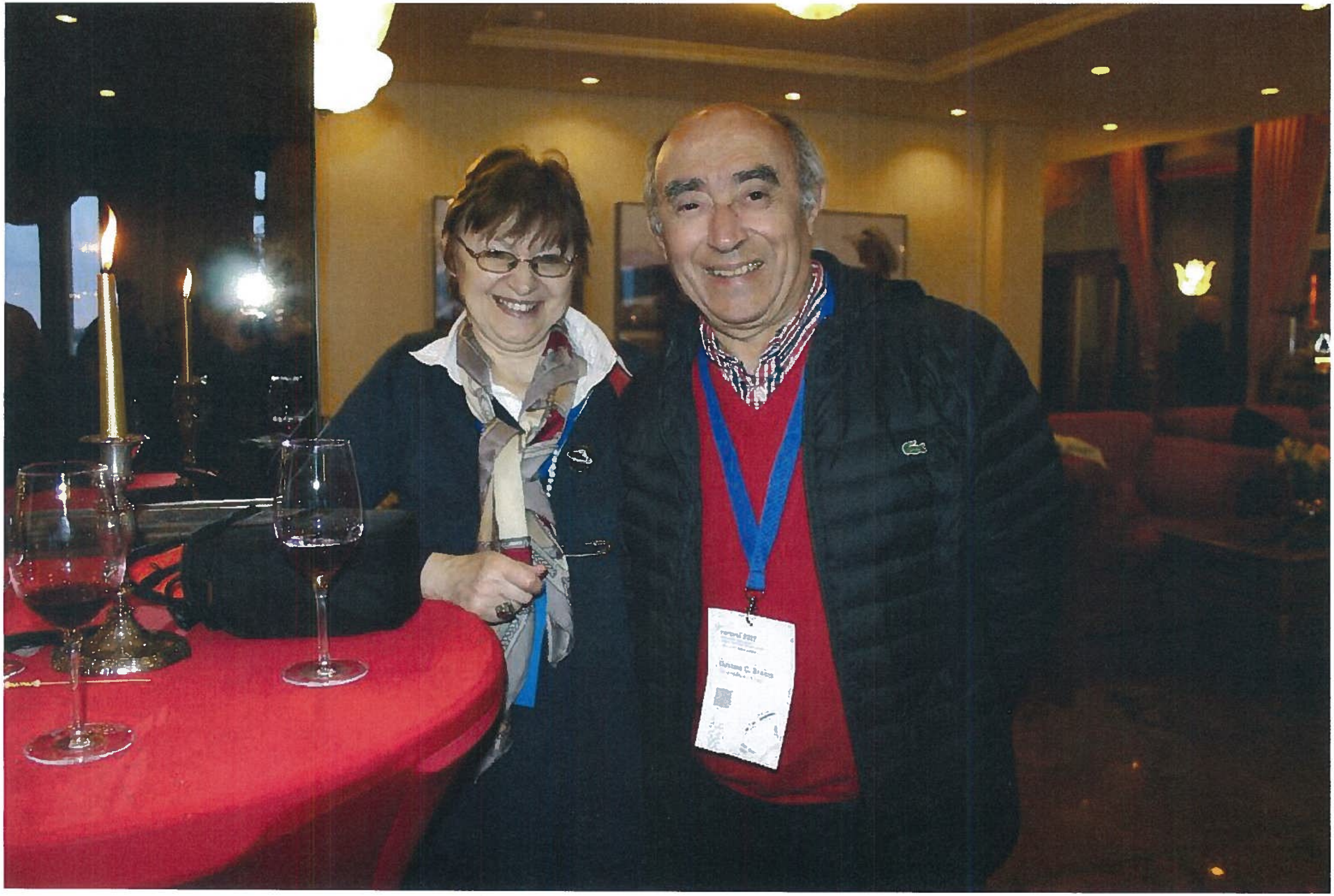
arXiv:1703.03796, arXiv:1508.05101, arXiv:1210.8163, arXiv:1102.0520, arXiv:0911.1753



Portoroz 2017

Warsaw, Discrete 2016

Work Partially supported by:

FCT Fundação para a Ciência e a Tecnologia
MINISTÉRIO DA EDUCAÇÃO E CIÊNCIA

europa.eu

European Union

QUADRO DE REFERÊNCIA ESTRATÉGICO NACIONAL
PORTUGAL 2007.2013

COMPETE
PROGRAMA OPERACIONAL FACTORES DE COMPETITIVIDADE

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Two Higgs Doublet Models

Several motivations

- New sources of CP violation

  SM cannot account for BAU

- Possibility of having spontaneous CP violation

  EW symmetry breaking and CP violation same footing

  T. D. Lee 1973, Kobayashi and Maskawa 1973

- Strong CP Problem, Peccei-Quinn

- Supersymmetry

LHC important role

In general two Higgs doublet models have FCNC

Neutral currents have played an important rôle in the construction and experimental tests of unified gauge theories

EPS Prize in 2009 to Gargamelle, CERN

In the Standard Model Flavour changing neutral currents (FCNC) are forbidden at tree level

- in the gauge sector, no ZFCNC

- in the Higgs sector, no HFCNC

Models with two or more Higgs doublets have potentially large HFCNC

Strict limits on FCNC processes!

In two Higgs Doublet Models (2HDM) Flavour Changing Neutral Currents (FCNC) have to be eliminated at tree level or naturally suppressed, in order to conform to experiment.

- Z_2 symmetry leading to Natural Flavour Conservation (NFC)

  Glashow and Weinberg (1977)

  Paschos (1977)

- Attempt at generalising NFC : R. Gatto

Can one have a framework where there are FCNC at tree level, but naturally suppressed?

Is it possible to have a framework where the FCNC exist, but are only functions of V^{CKM} and the ratio v_2/v_1?

The suppression of FCNC could be related to the smallness of some of the V^{CKM} elements.

- Naturally suppressed FCNC as a result of a symmetry of the Lagrangian. The suppression is due to small V^{CKM} elements

G.C.B, Grimus, Lavoura (1996) (BGL)

- Extension to the leptonic sector

F. Botella, GCB, MN Rebelo
F. Botella, GCB., MNRebelo, M.Nebot

- Thorough Phenomenological analysis

F. Botella, GCB, A. Carmona, M. Nebot, L. Pedro, M.N. Rebelo

Notation

Yukawa Interactions

$$\mathcal{L}_Y = -\overline{Q_L^0}\ \Gamma_1 \Phi_1 d_R^0 - \overline{Q_L^0}\ \Gamma_2 \Phi_2 d_R^0 - \overline{Q_L^0}\ \Delta_1 \tilde{\Phi}_1 u_R^0 - \overline{Q_L^0}\ \Delta_2 \tilde{\Phi}_2 u_R^0 + \text{h.c.}$$

$$\tilde{\Phi}_i = -i\tau_2 \Phi_i^*$$

Quark mass matrices

$$M_d = \frac{1}{\sqrt{2}}(v_1 \Gamma_1 + v_2 e^{i\theta} \Gamma_2), \quad M_u = \frac{1}{\sqrt{2}}(v_1 \Delta_1 + v_2 e^{-i\theta} \Delta_2),$$

Diagonalised by:

$$U_{dL}^\dagger M_d U_{dR} = D_d \equiv \text{diag}\ (m_d, m_s, m_b)\ ,$$
$$U_{uL}^\dagger M_u U_{uR} = D_u \equiv \text{diag}\ (m_u, m_c, m_t)\ .$$

Leptonic Sector

Charged leptons

$$-\overline{L_L^0}\;\Pi_1\Phi_1\ell_R^0 - \overline{L_L^0}\;\Pi_2\Phi_2\ell_R^0 \;+ \text{h.c.}$$

$$\left(\; - \overline{L_L^0}\;\Sigma_1\tilde{\Phi}_1\nu_R^0 - \overline{L_L^0}\;\Sigma_2\tilde{\Phi}_2\nu_R^0 + \text{h.c.}\;\right)$$

Neutrino Dirac

$$\left(\;\frac{1}{2}{\nu_R^0}^T C^{-1} M_R \nu_R^0 + \text{h.c.}\right)$$

Neutrino Majorana

Expansion around the vev's

$$\Phi_j = \begin{pmatrix} \phi_i^+ \\ \frac{e^{i\alpha_j}}{\sqrt{2}}(v_j + \rho_j + i\eta_j) \end{pmatrix}, \qquad j = 1, 2$$

We perform the following transformation:

$$\begin{pmatrix} H^0 \\ R \end{pmatrix} = U \begin{pmatrix} \rho_1 \\ \rho_2 \end{pmatrix}; \quad \begin{pmatrix} G^0 \\ I \end{pmatrix} = U \begin{pmatrix} \eta_1 \\ \eta_2 \end{pmatrix}; \quad \begin{pmatrix} G^+ \\ H^+ \end{pmatrix} = U \begin{pmatrix} \phi_1^+ \\ \phi_2^+ \end{pmatrix},$$

$$U = \frac{1}{v} \begin{pmatrix} v_1 e^{-i\alpha_1} & v_2 e^{-i\alpha_2} \\ v_2 e^{-i\alpha_1} & -v_1 e^{-i\alpha_2} \end{pmatrix}; \qquad v = \sqrt{v_1^2 + v_2^2} = (\sqrt{2} G_F)^{-\frac{1}{2}} \simeq 246 \mathrm{GeV}$$

U singles out

H^0 with couplings to quarks proportional to mass matrices

G^0 neutral pseudo-Goldstone boson

G^+ charged pseudo-Goldstone boson

Physical neutral fields are combinations of H^0 R I

Neutral and charged Higgs Interactions for the quark sector

$$\begin{aligned}\mathcal{L}_Y(\text{quark, Higgs}) = & -\overline{d_L^0}\frac{1}{v}\left[M_d H^0 + N_d^0 R + i N_d^0 I\right] d_R^0 \\ & -\overline{u_L^0}\frac{1}{v}\left[M_u H^0 + N_u^0 R + i N_u^0 I\right] u_R^0 \\ & -\frac{\sqrt{2}H^+}{v}\left(\overline{u_L^0} N_d^0\, d_R^0 - \overline{u_R^0} N_u^{0\dagger}\, d_L^0\right) + \text{h.c.}\end{aligned}$$

$$N_d^0 = \frac{1}{\sqrt{2}}(v_2\Gamma_1 - v_1 e^{i\theta}\Gamma_2), \quad N_u^0 = \frac{1}{\sqrt{2}}(v_2\Delta_1 - v_1 e^{-i\theta}\Delta_2).$$

Flavour structure of quark sector of 2HDM characterised by:

four matrices $M_d,\ M_u,\ N_d^0,\ N_u^0$.

Likewise for Leptonic sector, Dirac neutrinos:

$$M_\ell,\ M_\nu,\ N_\ell^0,\ N_\nu^0.$$

Yukawa Couplings in terms of quark mass eigenstates

for H^+, H^0, R, I

$$
\begin{aligned}
\mathcal{L}_Y(\text{quark, Higgs}) = \\
& -\frac{\sqrt{2}H^+}{v}\bar{u}\left(VN_d\gamma_R - N_u^\dagger\, V\gamma_L\right) d + \text{h.c.} - \frac{H^0}{v}\left(\bar{u}D_u u + \bar{d}D_d\, d\right) - \\
& - \frac{R}{v}\left[\bar{u}(N_u\gamma_R + N_u^\dagger\gamma_L)u + \bar{d}(N_d\gamma_R + N_d^\dagger\gamma_L)\, d\right] + \\
& + i\frac{I}{v}\left[\bar{u}(N_u\gamma_R - N_u^\dagger\gamma_L)u - \bar{d}(N_d\gamma_R - N_d^\dagger\gamma_L)\, d\right]
\end{aligned}
$$

$$\gamma_L = (1-\gamma_5)/2 \qquad\qquad \gamma_R = (1+\gamma_5)/2$$

$$V = V_{CKM}$$

FCNC controlled by:

$$N_d = \frac{1}{\sqrt{2}} U_{dL}^{\dagger} \left(v_2 \Gamma_1 - v_1 e^{i\alpha} \Gamma_2 \right) U_{dR}$$

$$N_u = \frac{1}{\sqrt{2}} U_{uL}^{\dagger} \left(v_2 \Delta_1 - v_1 e^{-i\alpha} \Delta_2 \right) U_{uR}$$

For general 2HDM, N_d, N_u are are arbitrary complex 3x3 matrices.

One can rewrite N_d as:

$$N_d = \frac{v_2}{v_1} D_d - \frac{v_2}{\sqrt{2}} \left(t + \frac{1}{t} \right) \underbrace{U_{dL}^{\dagger} e^{i\alpha} \Gamma_2 U_{dR}}_{\text{leads to FCNC}}$$

$t \equiv \tan\beta = \frac{v_2}{v_1}$

The general 2HDM has a HUGE number of New Flavour Parameters encoded by the two matrices N_d^0 and N_u^0 which are arbitrary complex matrices

$$18 + 18 = 36$$ new physics parameters

Questions :

- How to parametrize N_d, N_u in a sensible way?
- Can one use a parametrization which resembles the one used in V^{CKM}?

Answer : Yes!

Any arbitrary complex matrix can be written :

$$N_d^0 = X_L^\dagger \, D^{Nd} \, X_R$$

$$N_d^0 = K_L \, \hat{V}_L^{Nd} \, D^{Nd} \, \bar{K} \, (\hat{V}_R^{Nd})^\dagger \, K_R^\dagger$$

$$K_{L,R} = \text{diag.}\left[1, \exp(i\varphi_{L,R}), \exp(i\varphi_{2L,R})\right]$$

$$\bar{K} = \text{diag.}\left[e^{i\sigma_1} \quad e^{i\sigma_2} \quad e^{i\sigma_3}\right]$$

phases : $2(K_L) + 2K_R + 1\,\hat{V}_L^{Nd} + 1\,V_R^{Nd} + 3\bar{K} = 9$

real parameters : $3\,\hat{V}_L^{Nd} + 3\,\hat{V}_R^{Nd} + 3\,D^{Nd} = 9$

There are many **new sources of flavour changing and of CP violation**

$$I_1 \equiv \mathrm{tr}\left(M_d N_d^{0\dagger}\right) = m_d (N_d^*)_{11} + m_s (N_d^*)_{22} + m_b (N_d^*)_{33}$$

where N_d is the matrix N_d^0 in the basis where it couples to quark mass eigenstates.

$\mathrm{Im}\, I_1$ is very important since it probes the phases in $(N_d)_{jj}$ which contribute to the electric dipole moment of down quarks.

Recall the CP odd invariant of the SM:

$$I^{CP} \equiv \mathrm{tr}\left[H_u, H_d\right]^3 = 6i\,(m_t^2-m_c^2)(m_t^2-m_u^2)(m_c^2-m_u^2)\times$$
$$\times (m_b^2-m_s^2)(m_b^2-m_d^2)(m_s^2-m_d^2)\,\mathrm{Im}\, Q_{uscb}$$

In the SM V^{CKM} reflects a misalignment in flavour space between the two Hermitian matrices $H_d = M_d M_d^\dagger$ and $H_u = M_u M_u^\dagger$.

In an analogous way, in 2HDM, one has New mixing matrices which reflect misalignment of various Hermitian matrices.

Examples: $I_2^{CP} \equiv tr\left[H_u, H_{N_d}\right]^3 = 6i\, \Delta_u \Delta_{N_d} \operatorname{Im} Q_2$

where Q_2 is a rephasing invariant quartet of

$V_2 \equiv U_{uL}^\dagger \, U_{N_d^0 L}$ and $\Delta_u = (m_t^2 - m_c^2)(m_t^2 - m_u^2)(m_c^2 - m_u^2)$

Δ_{N_d} defined by an analogous expression for the eigenvalues of $N_d^0 N_d^{0\dagger}$

One may also construct WB invariants sensitive to right-handed mixings

$$I_7^{CP} \equiv \mathrm{tr}\left[H'_d, H'_{N_d^0}\right]^3 = 6i\,\Delta_d\,\Delta_{N_d}\,\mathrm{Im}\,Q_7$$

where $H'_d \equiv M_d^\dagger M_d$; $H'_{N_d^0} = N_d^{0\dagger} N_d^0$

Q_7 is a rephasing quartet of $U_{dR} U^\dagger_{N^0_R}$.

One may have also CP even WB invariants:

$$I_2 \equiv \mathrm{tr}\left[M_d N_d^0, M_d M_d^\dagger\right]^2 = -2 m_d m_s (m_s^2 - m_d^2) N^*_{d12} N^*_{d21} - 2 m_d m_b (m_b^2 - m_d^2)^2 (N_d^*)_{13} (N_d^*)_{31} - 2 m_s m_b (m_b^2 - m_s^2)(N_d^*)_{23} (N_d^*)_{32}$$

I_2 can probe the strength of FCNC

One of the reasons why in the SM one cannot generate sufficient BAU is due to the smallness of CP violation :

$$\boxed{\frac{\mathrm{tr}\,[H_u, H_d]^3}{v^{12}} \approx 10^{-20}}$$

Situation is much better in 2HDM

In the SM the invariant which controls the strength of CP violation is of order M^{12}.

Even in constrained 2HDM like BGL models, the lowest order invariant sensitive to CP violation is of much lower order, namely:

$$I_g^{CP} \equiv \operatorname{Im} \operatorname{tr}\left[M_d N_d^{0\dagger} M_d M_d^\dagger M_u M_u^\dagger M_d M_d^\dagger\right]$$

For one of the most interesting BGL models (top model)

$$I_g^{CP} = -\left(\frac{v_2}{v_1}+\frac{v_1}{v_2}\right)(m_b^2-m_s^2)(m_b^2-m_d^2)(m_s^2-m_d^2)(m_c^2-m_u^2)\times$$

$$\operatorname{Im}\left(V_{22}^* V_{32} V_{33}^* V_{23}\right)$$

$V \rightarrow V^{CKM}$!!

Example of a BGL-type model: Impose the following discrete symmetry:

$$Q^0_{Lj} \to \exp(i\tau)\, Q^0_{Lj} \; ; \quad u^0_{Rj} \to \exp(2i\tau)\, u^0_{Rj} \; ;$$

$$\phi_2 \to \exp(i\tau)\, \phi_2$$

$\tau \neq 0, \pi$

Γ_j, Δ_j have the form:

$$\Gamma_1 = \begin{bmatrix} \times & \times & \times \\ \times & \times & \times \\ 0 & 0 & 0 \end{bmatrix} ; \quad \Gamma_2 = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \times & \times & \times \end{bmatrix} ; \quad \Delta_1 = \begin{bmatrix} \times & \times & 0 \\ \times & \times & 0 \\ 0 & 0 & 0 \end{bmatrix} ; \quad \Delta_2 = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \times \end{bmatrix}$$

FCNC only in the down sector

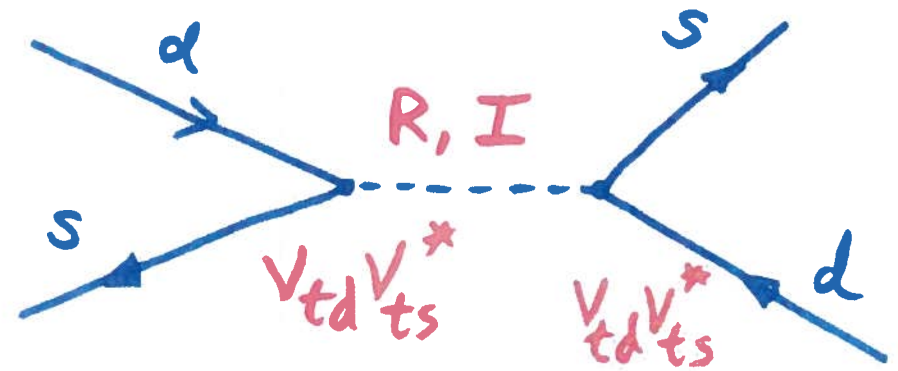
If one imposes $d^0_{Rj} \to \exp(2i\tau)\, d^0_{Rj}$ instead of $u^0_{Rj} \to \exp(2i\tau)\, u^0_{Rj}$ only FCNC in up sector

Considering only the Quark sector there are 6 different BGL type models. In the example considered, one has:

$$\left(N_d\right)_{rs} = \frac{v_2}{v_1}\left(D_d\right)_{rs} - \left(\frac{v_2}{v_1} + \frac{v_1}{v_2}\right)\left(V_{CKM}^{\dagger}\right)_{r3}\left(V_{CKM}\right)_{3s}\left(D_d\right)_{ss}$$

$$N_u = -\frac{v_1}{v_2}\,\mathrm{diag}\left(0, 0, m_t\right) + \frac{v_2}{v_1}\,\mathrm{diag.}\left(m_u, m_c, 0\right)$$

Strong and natural suppression of K^0–$\overline{K^0}$ transitions



$\rightarrow \lambda^{10}$ suppression

Neutral couplings in BGL models

$$N_u = -\frac{v_1}{v_2}\mathrm{diag}(0,0,m_t) + \frac{v_2}{v_1}\mathrm{diag}(m_u,m_c,0)$$

Explicitely

$$N_d = \frac{v_2}{v_1}\begin{pmatrix} m_d & 0 & 0 \\ 0 & m_s & 0 \\ 0 & 0 & m_b \end{pmatrix} - \left(\frac{v_1}{v_2}+\frac{v_2}{v_1}\right)\begin{pmatrix} m_d|V_{31}|^2 & m_s V_{31}^* V_{32} & m_b V_{31}^* V_{33} \\ m_d V_{32}^* V_{31} & m_s|V_{32}|^2 & m_b V_{32}^* V_{33} \\ m_d V_{33}^* V_{31} & m_s V_{33}^* V_{32} & m_b|V_{33}|^2 \end{pmatrix}$$

<u>It all comes from the symmetry</u>

BGL models have some features in Common with the Minimal Flavour Violation Framework

Buras, Gambino, Gorbahn, Jager, Silvestrini (2001)

D'Ambrosio, Giudice, Isidori, Strumia (2002)

Namely, Flavour dependence of New Physics is completely controlled by V^{CKM}, with no other flavour parameters.

Note: MFV is an "Hypothesis" not a model!!!

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An important question:

Can one introduce other discrete symmetries leading to other models with FCNC, completely controlled by V_{CKM}?

Answer: In the framework of 2HDM with Abelian symmetries and the constraint that FCNC only depend on V_{CKM}, BGL models are unique!

Ferreira and Silva 2010.

The explicit example of a BGL model was written in a weak-basis chosen by the symmetry. How to recognize a BGL model when written in a different WB?

The following relations

$$\Delta_1^\dagger \Delta_2 = 0 \ ; \quad \Delta_1 \Delta_2^\dagger = 0 \ ; \quad \Gamma_1^\dagger \Delta_2 = 0 \ ; \quad \Gamma_2^\dagger \Delta_1 = 0$$

are necessary and sufficient conditions for a set of Yukawa matrices Γ_i, Δ_i to be of the BGL type, with FCNC in the down sector.

- In a certain sense, **BGL models** are rather unique. They have FCNC either in the up or the down sectors but not in both.

- If one restricts oneself to **abelian symmetries**, BGL models are the only 2HDM with **FCNC** at tree level, but no new flavour parameters, apart from V_{CKM}.

- **Question** – Can one generalize **BGL models** and construct a 2HDM with non-vanishing but controlled **FCNC** in both the up and down sectors? These gBGL models would contain BGL models as special cases, corresponding to specific values of the parameters of gBGL

Answer: Yes!!
gBGL allowing for HFCNC both in up and down sectors

Symmetry: (not flavour blind!!)

$$Q_{L_3} \mapsto -Q_{L_3},$$
$$d_R \mapsto d_R, \qquad \Phi_1 \mapsto \Phi_1,$$
$$u_R \mapsto u_R, \qquad \Phi_2 \mapsto -\Phi_2.$$

- drastic reduction in number of free parameters

-no NFC

one may say that *the principle leading to gBGL constrains the Yukawa couplings so that each line of* Γ_j, Δ_j *couples only to one Higgs doublet.*

$$\Gamma_1 = \begin{pmatrix} \times & \times & \gamma_{13} \\ \times & \times & \gamma_{23} \\ 0 & 0 & 0 \end{pmatrix}, \quad \Gamma_2 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \gamma_{31} & \gamma_{32} & \times \end{pmatrix},$$
$$\Delta_1 = \begin{pmatrix} \times & \times & \delta_{13} \\ \times & \times & \delta_{23} \\ 0 & 0 & 0 \end{pmatrix}, \quad \Delta_2 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \delta_{31} & \delta_{32} & \times \end{pmatrix},$$

- renormalisable;

- FCNC both in up and down sectors;

- no longer of MFV type, four additional flavour parameters;

- both up and down type BGL appear as special limits;

gBGL verify:

$$\Gamma_2^\dagger \Gamma_1 = 0, \qquad \Gamma_2^\dagger \Delta_1 = 0,$$
$$\Delta_2^\dagger \Delta_1 = 0, \qquad \Delta_2^\dagger \Gamma_1 = 0.$$

Structure of Yukawa Couplings

$$\Gamma_1 = \begin{bmatrix} \times & \times & \gamma_{13} \\ \times & \times & \gamma_{23} \\ 0 & 0 & 0 \end{bmatrix} \quad ; \quad \Gamma_2 = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \gamma_{31} & \gamma_{32} & \times \end{bmatrix}$$

$$\Delta_1 = \begin{bmatrix} \times & \times & \delta_{13} \\ \times & \times & \delta_{23} \\ 0 & 0 & 0 \end{bmatrix} \quad ; \quad \Delta_2 = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \delta_{31} & \delta_{32} & \times \end{bmatrix}$$

For $\delta_{ij} = 0$ one obtains uBGL models, with

$$\Gamma_1 = \begin{bmatrix} \times & \times & \times \\ \times & \times & \times \\ 0 & 0 & 0 \end{bmatrix} \quad ; \quad \Gamma_2 = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \times & \times & \times \end{bmatrix}$$

$$\Delta_1 = \begin{bmatrix} \times & \times & 0 \\ \times & \times & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad ; \quad \Delta_2 = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \times \end{bmatrix}$$

Similarly for $\gamma_{ij} = 0$, one obtains dBGL

It can be shown that N_d, N_u can be parametrized as :

$$N_d = \left[t_\beta \mathbb{1} - (t_\beta + t_\beta^{-1}) V^\dagger \mathcal{U} P_3 \mathcal{U}^\dagger V \right] M_d$$

$$N_u = \left[t_\beta \mathbb{1} - (t_\beta + t_\beta^{-1}) \mathcal{U} P_3 \mathcal{U}^\dagger \right] M_u$$

$V \equiv V^{CKM}$. It is clear that in gBGL one has more freedom, due to the presence of the arbitrary matrix $\mathcal{U}$. Nevertheless, there is much less freedom than one might expect, since the only quantities involving $\mathcal{U}$ are :

$$\left[\mathcal{U} P_3 \mathcal{U}^\dagger \right]_{ij} = \mathcal{U}_{i3} \mathcal{U}^*_{j3}$$

Phenomenological Consequences

- New Physics contributions to:

  $K^0 - \bar{K}^0$, $B_d - \bar{B}_d$, $B_s - \bar{B}_s$ mixings

- Rare b-decays
- Rare top decays
- Flavour changing Higgs decays

  $h \to \mu\tau$

  $h \to bd, bs$

etc, etc

For detailed phenomenological analysis see talk by

Miguel Nebot

Conclusions

- The 2HDM is one of the most plausible extensions of the SM

- There are 2HDM with naturally suppressed FCNC at tree level. examples: BGL models; gBGL models

- If extra scalars, either neutral or charged are discovered , an important task will be to measure the matrices Nu and Nd

These measurements may help to uncover a possible

family symmetry

- There is no scientific reason to accept the $\left\{\begin{array}{l}\text{Dogma}\\ \text{Myth}\end{array}\right.$ that the question of flavour can only be understood at the Planck scale
- It is crucial to construct a

Higgs/Top Factory