

SIMPLIFIED DM MODELS WITH Z'

A STORY WITH A DARK HIGGS

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1605.09382, 1610.03063 and 1705.01105



Overview

1. Introduction
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5. Summary

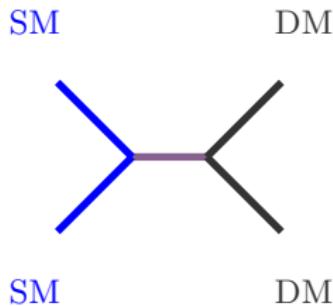
INTRODUCTION

Simplified Models and Unitarity

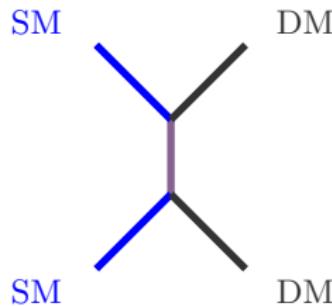
Simplified models suggested for the LHC Run 2 search 1507.00966

- a handful of new particles and interactions
- can be seen as a limit of various UV complete models

s-channel



t-channel



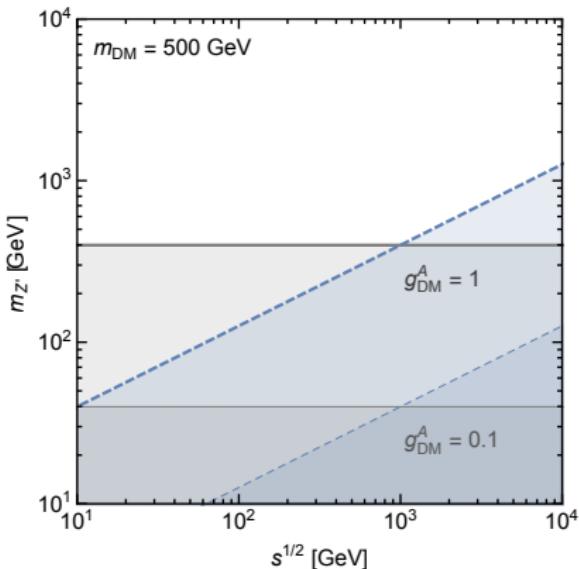
Each simplified model features **one** type of mediator.

Unitarity issues if SM or dark gauge symmetry broken 1503.07874

A Simple Example

e.g.: a spin-1 mediator [Kahlhoefer et. al. 1510.02110, 1606.07609](#)

$$\Delta\mathcal{L} \sim Z'^{\mu} \bar{f} \left[g_f^V \gamma_{\mu} + g_f^A \gamma_{\mu} \gamma^5 \right] f + Z'^{\mu} \bar{\chi} \left[g_{\chi}^V \gamma_{\mu} + g_{\chi}^A \gamma_{\mu} \gamma^5 \right] \chi$$



- Axial vector couplings break unitarity
 - DM self-scattering
 - DM annihilation
- Z' mass breaks $U(1)'$
- Extra scalar needed for AV to restore unitarity and also generate mass

Recap: Z' and a dark Higgs

Axial vector mediator implies the existence of a dark Higgs field.

- A simplified model with only a DM candidate and an axial vector mediator is oversimplified.
- This leads to a dual-mediator model: both spin-1 and spin-0 mediators.
- The presence of both a spin-1 and spin-0 mediator leads to interesting new phenomenology, not captured by a single mediator scenario.

A MODEL WITH Z' AND A DARK HIGGS

Maj. DM with Z' mediator in a Hidden Sector Scenario

- $U(1)_\chi$ with Majorana DM χ , Z' & complex scalar S
- let $Q'(S) = 1$, Yukawa interaction sets $Q'(\chi) = -\frac{1}{2}$

$$-\frac{1}{4}g_\chi \bar{\chi} \gamma_5 \gamma_\mu \chi Z'^\mu - \frac{1}{2}y_\chi \bar{\chi}_L^C \chi_L S$$

- Hidden sector scenario

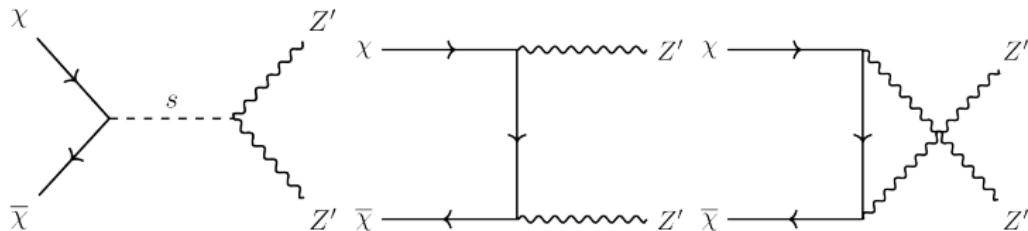
$$-\lambda_{HS}(S^\dagger S)(H^\dagger H) - \frac{1}{2}\sin\epsilon Z'_{\mu\nu}B^{\mu\nu}$$

small mixing allows decay of S and Z'

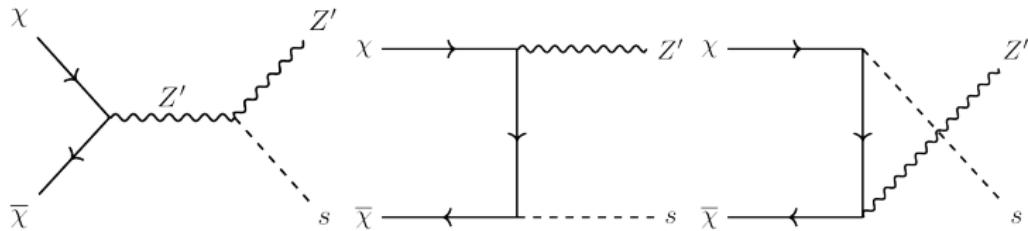
- Dark-visible sector couplings can be small, to satisfy collider and direct detection constraints
- Large indirect detection signals possible

What is New for the Dual-Mediator Model?

- A new contribution to $\chi\chi \rightarrow Z'Z'$
prevents unphysical high energy behavior



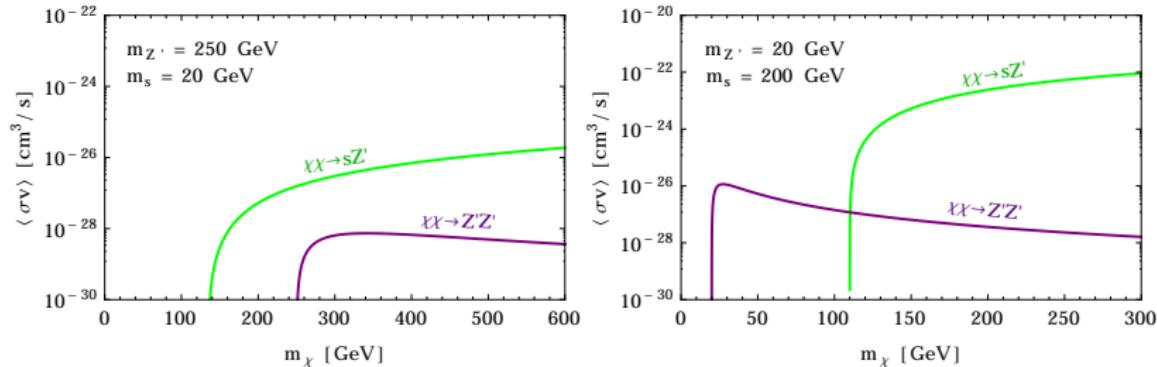
- A brand new channel of interests: **dark Higgstrahlung**



Bell, Cai and Leane, 1605.09382

Comparison of the Cross Sections

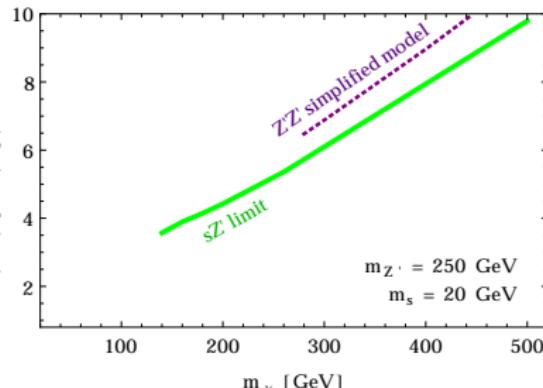
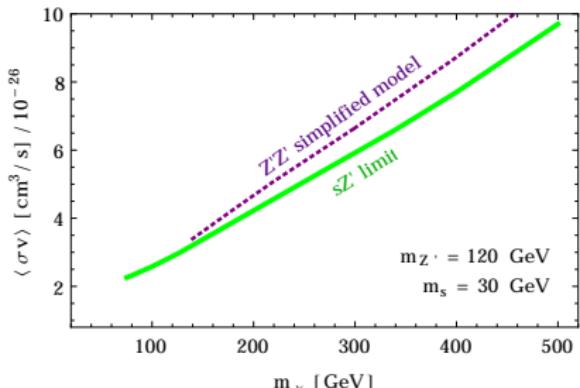
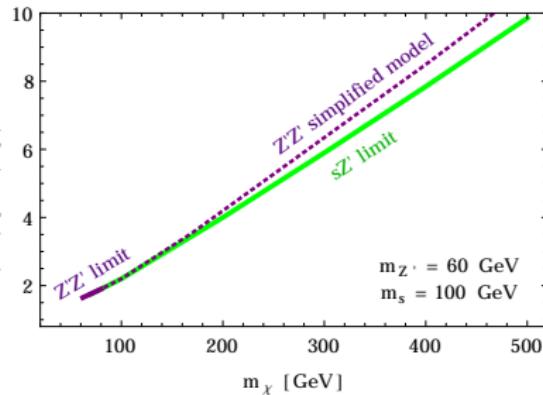
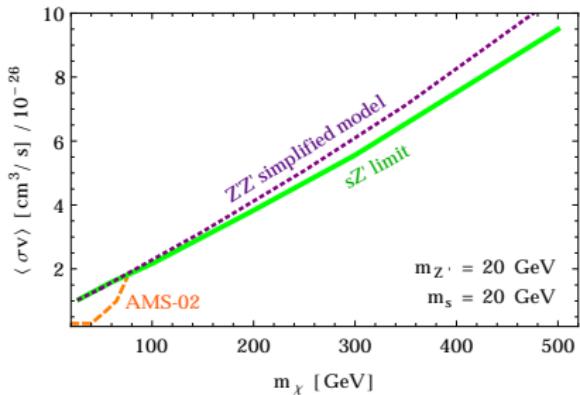
Parameters: m_s , $m_{Z'}$, m_χ , g_χ . With $m_s < 2m'_{Z'}$: $Z', s \Rightarrow \text{SM}$



- Both cross sections scale as g_χ^4 , $g_\chi = 0.1$
- Large cross section for $\chi\chi \rightarrow sZ'$
enhanced by $Z'_L \sim \frac{m_\chi^2}{m_{Z'}^4}$
- With $m_s + m_{Z'} < 2m_\chi < 2m'_{Z'}$, sZ' is the only channel
- Inaccurate limits if only $Z'Z'$ considered

Results for ID

Fermi-LAT and AMS-02



DARK SECTOR MASS GENERATION

Fermionic Dark Matter

Majorana DM

- only axial-vector couplings to a Z' allowed
- dark Higgs mechanism gives mass to both Z' and DM.

DIRAC DM

Both vector and axial-vector couplings to Z' allowed

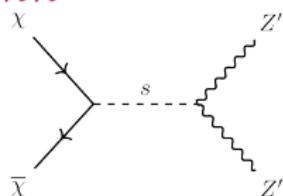
- If Z' has pure vector couplings
 - Z' mass: either dark Higgs or Stueckelberg mechanism
 - DM mass: bare mass or dark Higgs mechanism
 - **mass mechanisms not necessarily connected**
- If Z' has non-zero axial coupling
 - Dark Higgs gives mass to both Z' and DM (like Majorana)

Dark Sector Mass Generation: Dirac DM

	m_χ	$m_{Z'}$	$\chi - Z'$ coup.	Ann.	Z' pol.
I	Bare mass	Stueckelberg	Vector	$Z'Z'$	Z'_T
II	Dark Higgs	Dark Higgs	Axial $\neq 0$	$Z'Z' \& sZ'$	$Z'_T \& Z'_L$
III	Dark Higgs	Stueckelberg	Vector	$Z'Z' \& sZ'$	Z'_T
IV	Bare mass	Dark Higgs	Vector	$Z'Z' \& sZ'$	Z'_T

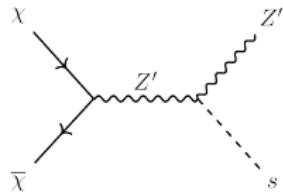
Bell, Cai & Leane 1610.03063

$$\chi\bar{\chi} \rightarrow Z'Z'$$

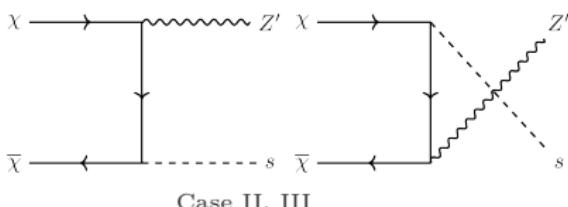
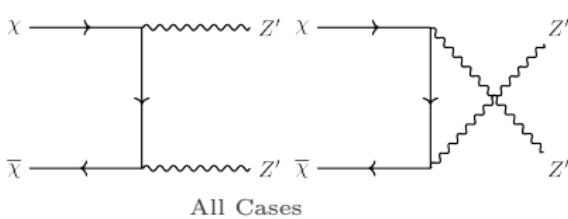


Case II only

$$\chi\bar{\chi} \rightarrow sZ'$$

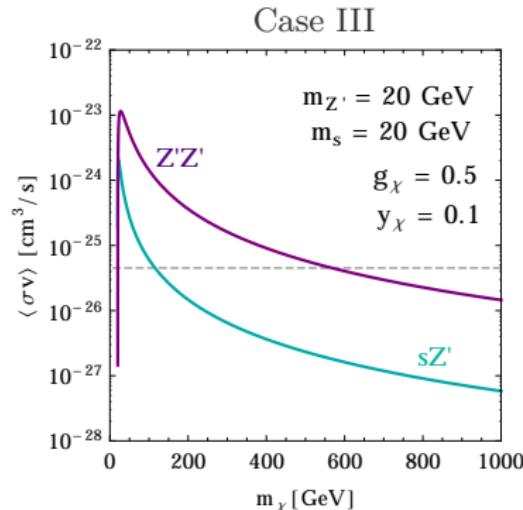
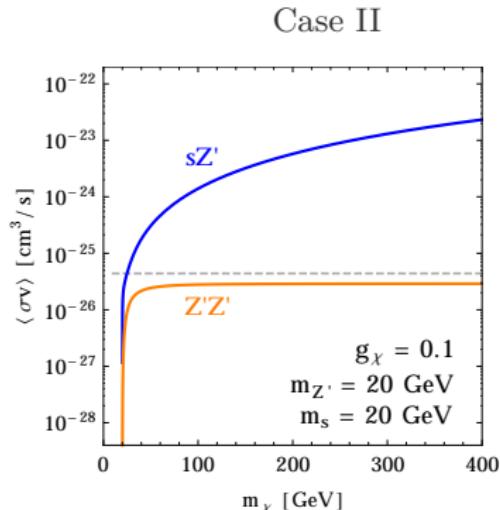


Case II, IV



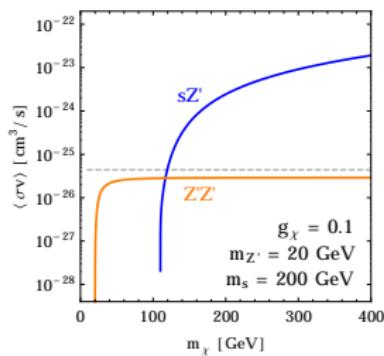
Case II & III: a quick comparison

- Case II: Z' and DM mass from dark Higgs
 - sZ' dominates over $Z'Z'$ when kinematically allowed
 - Cross sections enhanced by longitudinal Z'
- Case III: DM mass from dark Higgs, Z' from Stu. mechanism
 - y_χ and g_χ not related: relative strength freely arranged

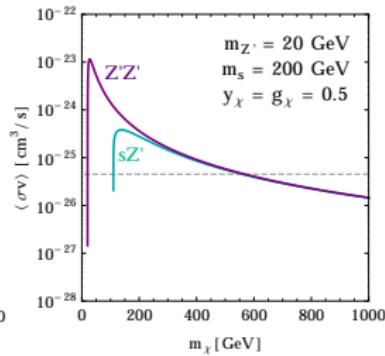


Polarization of Z'

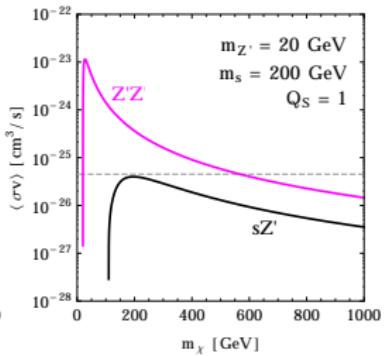
Enhancement from longitudinal Z'
only for axial couplings



Case II



Case III



Case IV

DARK ISR ENHANCES ANNIHILATION

DM Annihilation: $\chi\bar{\chi} \rightarrow f\bar{f}$

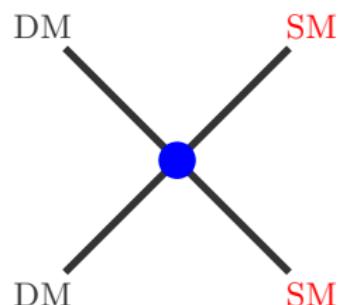
$$(\bar{\chi}\Gamma_\chi\chi)(\bar{f}\Gamma_f f)$$

Annihilation rates often suppressed

- p -wave suppressed: $\sim v^2$
 $A \otimes V, S \otimes S, S \otimes P$
- helicity suppressed: $\sim m_f^2/m_\chi^2$
 $A \otimes A$

Possible treatment

- Final state radiation
- Internal Bremsstrahlung
- Initial state radiation



Dark ISR: A simple equivalence

only true for $(\bar{\chi}\Gamma_\chi\chi)(\bar{f}\Gamma_f f)$

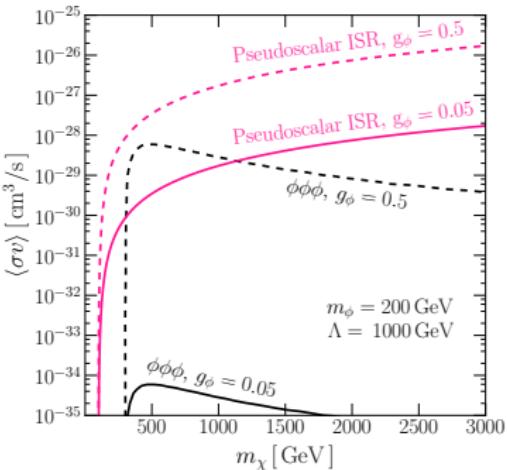
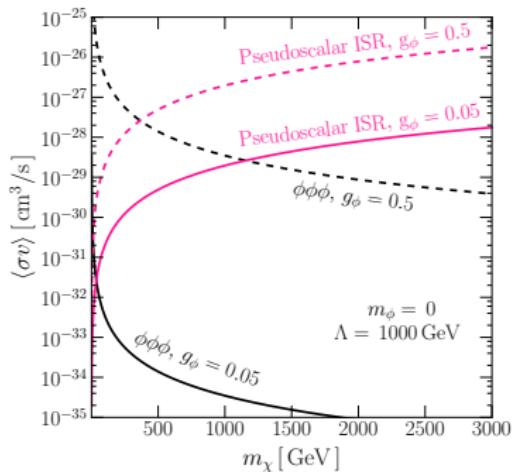
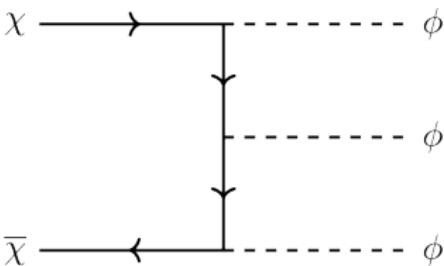
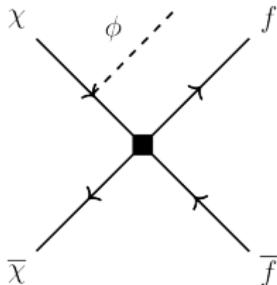


$\Gamma_{M_1} \otimes \Gamma_{M_2}$	$S \otimes S$	$S \otimes P$	$P \otimes P$	$V \otimes V$	$V \otimes A$	$A \otimes A$	$S \otimes V$	$S \otimes A$	$P \otimes V$	$P \otimes A$
$\bar{\chi}\chi \rightarrow M_1 M_2$	v^2	1	v^2	1	1	1	1	v^2	1	v^2

- p -wave suppressed: $\sim v^2$
 $A \otimes V$: V or A , $S \otimes S$: V or P , $S \otimes P$: V or P
- helicity suppressed: $\sim m_f^2/m_\chi^2$
 $A \otimes A$: V or A

Example: $S \otimes S$ with pseudoscalar

$$\Delta L \sim i g_\phi \bar{\chi} \gamma_5 \chi \phi$$



SUMMARY

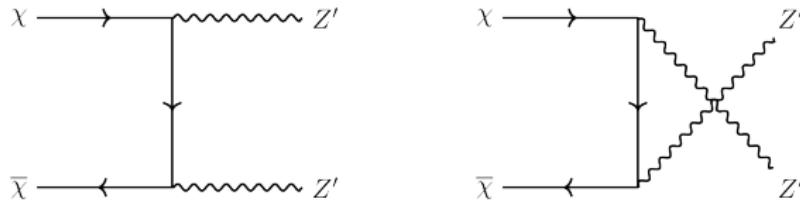
Summary

- Single mediator Simplified Models may not be self consistent
 - Two mediators can be required by gauge invariance
 - Phenomenology not captured by single-mediator model
- Axial vector Z' requires dark Higgs to unitarize Z'_L
 - New, dominant, s-wave annihilation channel $\chi\chi \rightarrow sZ'$
- Dark sector mass generation mechanisms should not be ignored
 - Choice of mass generation mechanism dictates the allowed coupling structure and annihilation processes
- Dark ISR is pretty useful to lift suppressed annihilation cross sections

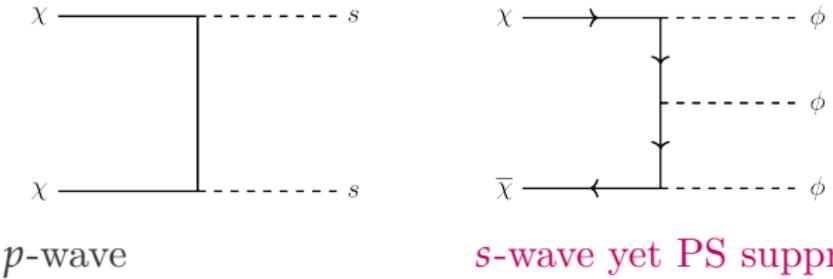
BACKUP

Dark Matter Annihilation

Simplified model with vector mediator:**s-wave**



Simplified model with scalar mediator



Mediation parameters: ϵ, λ_{HS}

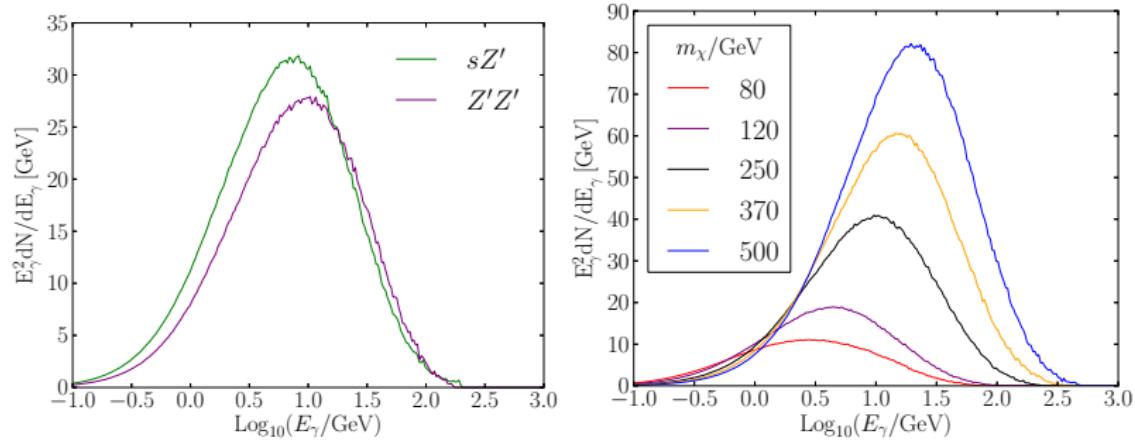
- ✓ small value consistent with non-observation of DD
- ✓ BBN requires $\tau < 1$ s, fairly easy to satisfy

COSMIC MICROWAVE BACKGROUND

- ✓ constraints on annihilation cross section weaker than AMS-02 and dSphs

Gamma Spectra

Spectra per annihilation generated with PYTHIA



$$m_{Z'} = 60 \text{ GeV}$$

$$m_s = 100 \text{ GeV}$$

$$m_\chi = 200 \text{ GeV}$$

$$\chi\chi \rightarrow sZ'$$

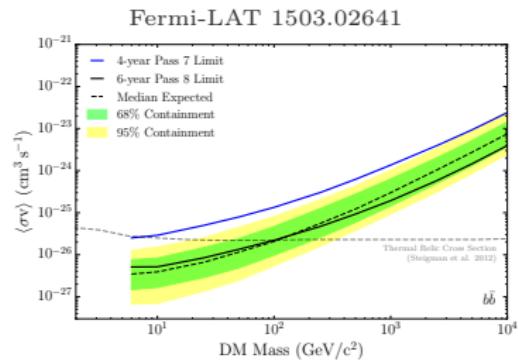
$$m_{Z'} = 120 \text{ GeV}$$

$$m_s = 30 \text{ GeV}$$

Limits from Indirect Detection

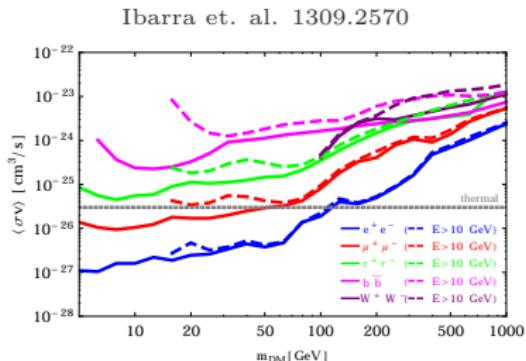
○ From Fermi-LAT dwarf spheroidal results

- No results for s and Z'
- s and Z' with different masses
- Likelihood functions of 15 dSphs used to derive limits
- J -factor as nuisance parameter
- Limits reproduced within 10%



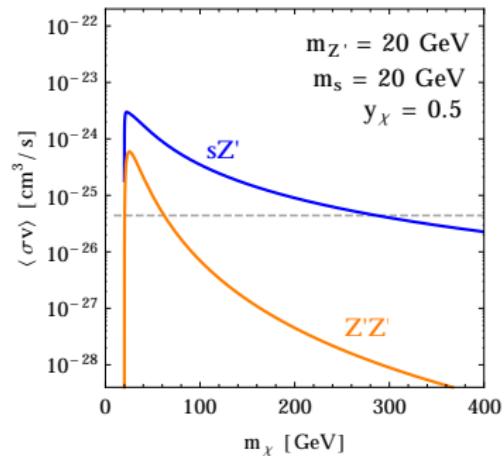
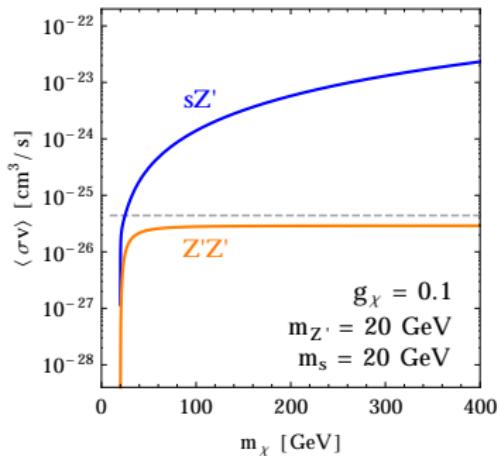
○ From AMS-02

- Consider e^\pm final states as they are the strongest
- Cascade decay via s negligible
- Only competitive for low m_χ
- Crude estimate with cross section of Z' rescaled with Br_{e^\pm}



Case II: Z' and DM mass from dark Higgs

- Couplings related: $\frac{y_\chi}{g_\chi} = \frac{\sqrt{2}m_\chi}{m_{Z'}}$
- $Q'_S = Q'_{\chi L} - Q'_{\chi R} \equiv 2Q'_A$ while Q'_V unconstrained
- sZ' dominates over $Z'Z'$ when kinematically allowed
- Cross sections enhanced by longitudinal Z'
(for $Z'Z'$ this only occur when Q'_V , Q'_A both nonzero)

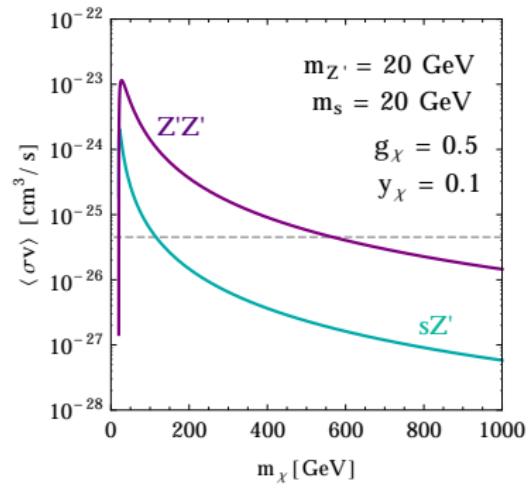
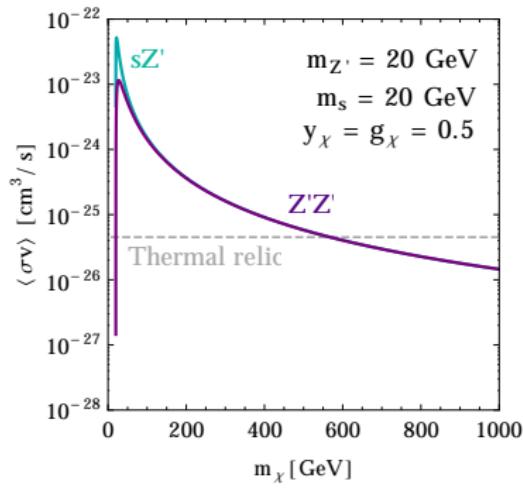


$$Q'_V = Q'_A = Q'_S/2 = 1$$

Case III

DM mass from dark Higgs and Z' mass from Stueckelberg mechanism

- y_χ and g_χ are not related: relative strength of the two ann. cross sections can be freely arranged
- Only transverse polarized Z'



Case IV: Bare DM Mass & Z' Mass from Dark Higgs

- $U(1)_\chi$ charge of Z' and S unrelated
relative strength of the two ann. cross sections can be freely arranged
- Only transverse polarized Z'

