

Tests of CP violation in multi-lepton signals at the LHC

arXiv:2212.09433

Phys. Rev. D, 103(7):075031

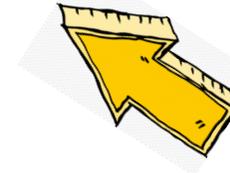
Scalars 2023

An opportunity to discuss various aspects of scalar particles.

13-16 September 2023
Warsaw (Ochota Campus)
scalars2023.fuw.edu.pl

WARSAW
POLAND

Y. Afik, S. Bar-Shalom, K. Pal,
A. Soni, J. Wudka



(presenter)

UC RIVERSIDE
UNIVERSITY OF CALIFORNIA

CP violation & new physics

SM sources:

- CKM matrix.
 - PMNS matrix.
- } small effects

... and the $\tilde{F} \cdot F$ term – apparently zero.

A good window into possible new physics (NP)



We **NEED** stronger CPV: there's no us without it (BAoU)

Many forms of NP readily supply the need.



- I assume CPT invariance holds.
- Unitarity of the S matrix:

$$\langle f | \mathcal{T} | i \rangle - \langle i | \mathcal{T} | f \rangle^* = i \underbrace{\sum_n \langle n | \mathcal{T} | f \rangle^* \langle n | \mathcal{T} | i \rangle}_{\text{Final state interactions (FSI)}} \quad (\mathcal{S} = \mathbb{1} + i\mathcal{T})$$

Final state interactions (FSI)

\mathcal{T} is Hermitian up-to FSI corrections.

useful for characterizing some CPV effects – no fundamental significance

Define 'naïve' time reversal:

$$T_N: t \rightarrow -t$$

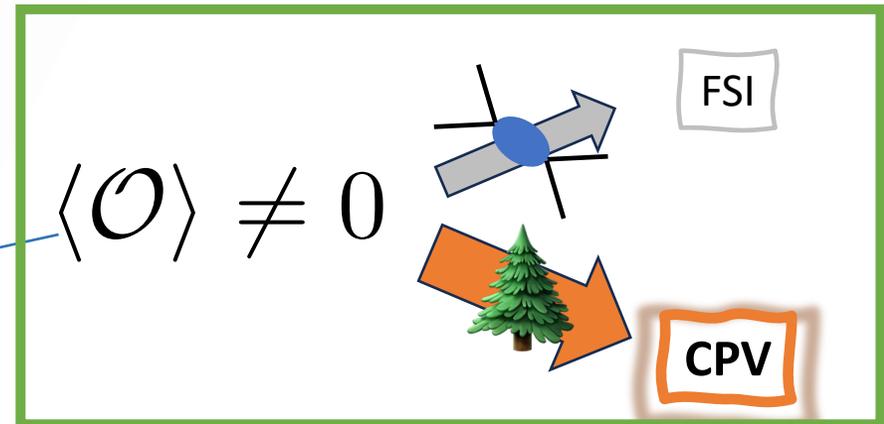
If

$$\mathcal{O} = \mathcal{O}(\mathbf{p}_f, \mathbf{s}_f) \xrightarrow{T_N} -\mathcal{O}$$

Final state momenta & spin

$$\langle \mathcal{O} \rangle = \frac{\int_R d\sigma \mathcal{O}}{\int_R d\sigma}$$

T_N -invariant phase-space region



A promising LHC process



Of the many final states,

$$pp \rightarrow n \ell + X$$

High p_T

are useful for measuring SM properties and NP searches.

Also useful in **CP_v** searches. Will look at

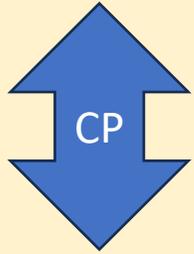
$$pp \rightarrow \ell'^- \ell^+ \ell^- + X_3$$

$$pp \rightarrow \ell'^+ \ell^- \ell^+ + \bar{X}_3$$

$$pp \rightarrow \ell'^+ \ell'^- \ell^+ \ell^- + X_4$$

$(\ell, \ell' = e, \mu, \tau)$

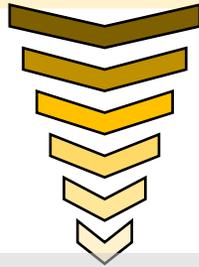
δ_i : CP even (FSI)
 ϕ_i : CP odd



$$\mathcal{M}_{ab \rightarrow \ell'^- \ell^+ \ell^-} = M_1 e^{i(\phi_1 + \delta_1)} + M_2 e^{i(\phi_2 + \delta_2)}$$

$$\bar{\mathcal{M}}_{\bar{a}\bar{b} \rightarrow \ell'^+ \ell^- \ell^+} = \bar{M}_1 e^{i(-\phi_1 + \delta_1)} + \bar{M}_2 e^{i(-\phi_2 + \delta_2)}$$

Parton-level amplitudes



$$\Delta\delta = \delta_1 - \delta_2$$

$$\Delta\phi = \phi_2 - \phi_1$$

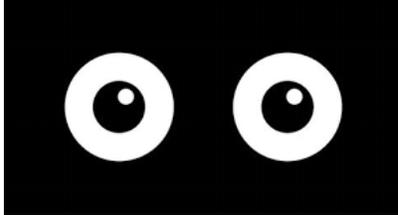
T_N even

$$T_N \text{ odd: } \begin{aligned} V' &= V \mathcal{O}_{\text{CP}} \\ \bar{V}' &= V \overline{\mathcal{O}_{\text{CP}}} \quad (V: T_N\text{-even}) \end{aligned}$$

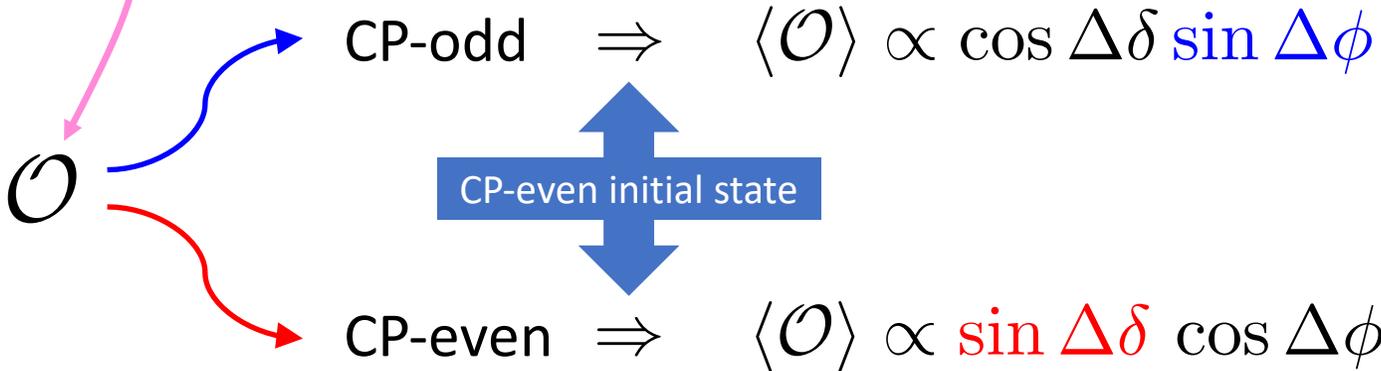
Parton-level cross sections

$$d\hat{\sigma}_{ab \rightarrow \ell'^- \ell^+ \ell^-} = W + U \cos(\Delta\delta + \Delta\phi) + V' \sin(\Delta\delta + \Delta\phi)$$

$$d\hat{\sigma}_{\bar{a}\bar{b} \rightarrow \ell'^+ \ell^- \ell^+} = W + U \cos(\Delta\delta - \Delta\phi) + \bar{V}' \sin(\Delta\delta - \Delta\phi)$$



T_N odd observables



Sensitive to **tree-level CP_v**
O(10%) CP asymmetries

Requires CP-even phases (FSI)
→ higher-order effect
→ O(0.1%) CP asymmetries

Chose

$$\mathcal{O}_{\text{CP}} = \vec{p}_{\ell'-} \cdot (\vec{p}_{\ell+} \times \vec{p}_{\ell-}) \quad \overline{\mathcal{O}_{\text{CP}}} = \vec{p}_{\ell'+} \cdot (\vec{p}_{\ell-} \times \vec{p}_{\ell+})$$

that are T_N -odd, and

$$\mathcal{O}_{\text{CP}} \longleftrightarrow_{\text{CP}} -\overline{\mathcal{O}_{\text{CP}}}$$

Construct the T_N -odd asymmetries

$$A_T = \frac{N(\mathcal{O}_{\text{CP}} > 0) - N(\mathcal{O}_{\text{CP}} < 0)}{N(\mathcal{O}_{\text{CP}} > 0) + N(\mathcal{O}_{\text{CP}} < 0)} \quad \bar{A}_T = \frac{N(-\overline{\mathcal{O}_{\text{CP}}} > 0) - N(-\overline{\mathcal{O}_{\text{CP}}} < 0)}{N(-\overline{\mathcal{O}_{\text{CP}}} > 0) + N(-\overline{\mathcal{O}_{\text{CP}}} < 0)}$$

And an observable sensitive to CP ν :

$$A_{CP} = \frac{1}{2} (A_T - \bar{A}_T)$$

May be affected by CP-even contributions.

More explicitly:

$$A_T = \mathcal{I}_{ab} \sin(\Delta\delta + \Delta\phi)$$

$$\bar{A}_T = \mathcal{I}_{\bar{a}\bar{b}} \sin(\Delta\delta - \Delta\phi)$$

$$\mathcal{I}_{ab} \propto \frac{\int_R d\Phi \cdot \overbrace{f_a f_b V}^{\text{PDF's}} |\mathcal{O}_{\text{CP}}|}{\int_R d\Phi \cdot f_a f_b \cdot (W + U \cdot \cos(\Delta\delta + \Delta\phi))}$$

T_N even phase-space region

$$A_{CP} = \frac{\mathcal{I}_{ab} + \mathcal{I}_{\bar{a}\bar{b}}}{2} \cos \Delta\delta \sin \Delta\phi + \frac{\mathcal{I}_{ab} - \mathcal{I}_{\bar{a}\bar{b}}}{2} \sin \Delta\delta \cos \Delta\phi$$

'Bona fide' CPv contribution

Vanishes if $\Delta\phi=0$

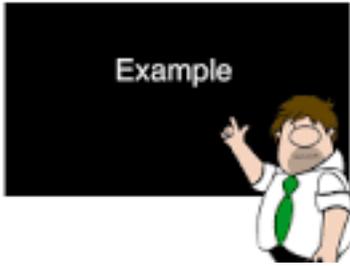
Insensitive to small CP-even phases $\Delta\delta$

Correction when initial state is not CP-even (e.g. LHC)

Survives even if $\Delta\phi=0$

Suppressed for small CP-even phases $\Delta\delta$ (no resonances)

Example



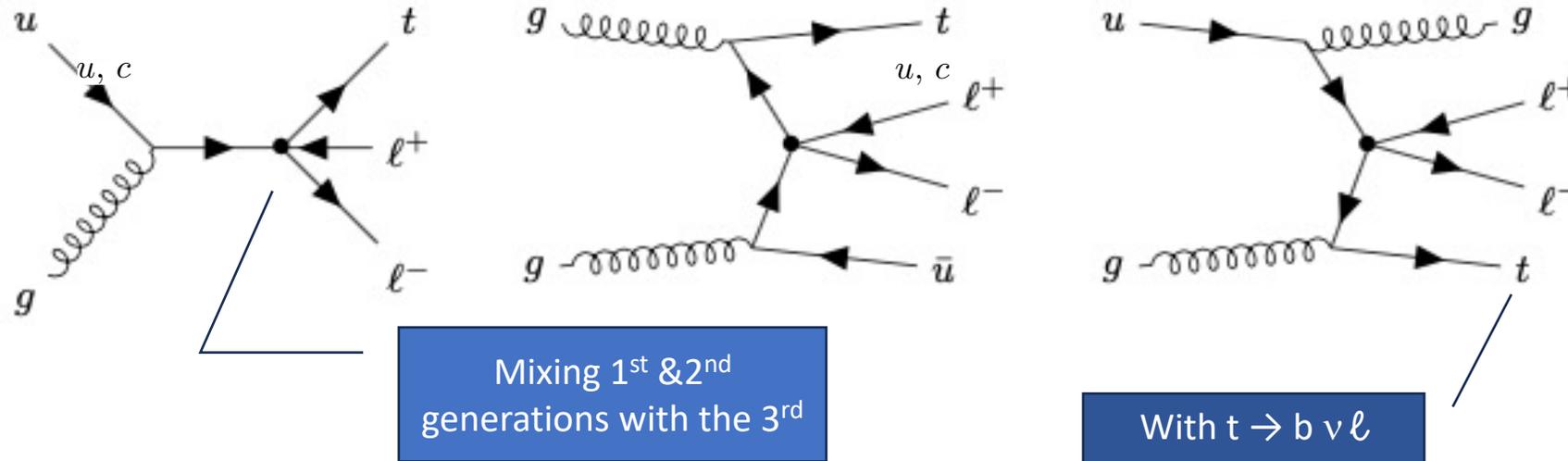
New physics:

- Decoupling
- Weakly coupled
- Not directly observed

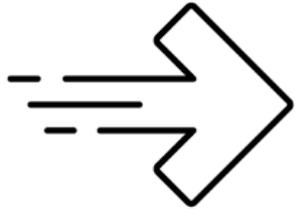
$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_{n=5}^{\infty} \frac{1}{\Lambda^{n-4}} \sum_i f_i O_i^{(n)} \subset \frac{f_S}{\Lambda^2} (\bar{l}^j e) \epsilon_{jk} (\bar{q}^k u) + \frac{f_T}{\Lambda^2} (\bar{l}^j \sigma_{\mu\nu} e) \epsilon_{jk} (\bar{q}^k \sigma^{\mu\nu} u) + \text{H.c.}$$

Can be generated by
 $S_1 \sim (3, 1, -1/3)$
 $R_2 \sim (3, 2, 7/6)$
leptoquarks, in which case
 $|f_S| = 4 |f_T|$

Reaction of interest (representative diagrams):



There is no NP-SM interference \rightarrow CP ν is purely NP-generated

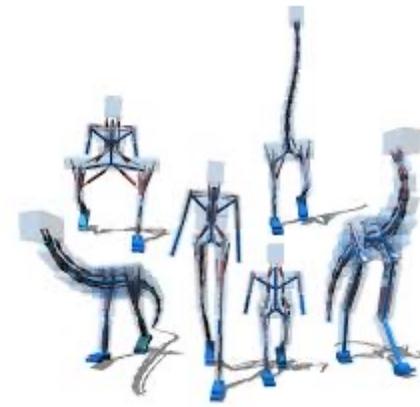


$$A_{CP} \sim \frac{\epsilon(p_u, p_{\ell'+}, p_{\ell+}, p_{\ell-}) \cdot \frac{1}{\Lambda^4} \mathbf{Im}(f_S f_T^*)}{(\text{SM}) + \frac{1}{\Lambda^4} [\text{terms} \propto |f_S|^2, |f_T|^2, \mathbf{Re}(f_S \cdot f_T^*)]}$$

Levi-Civita product

Mainly (inclusive): $pp \rightarrow WZ + X$
Subdominant (important when b tagging): $pp \rightarrow t\bar{t}Z, t\bar{t}W, t\bar{t}\bar{t}$

Simulation



Looked at (as an example)

$$pp \rightarrow e^{\pm} \mu^{+} \mu^{-} + X$$

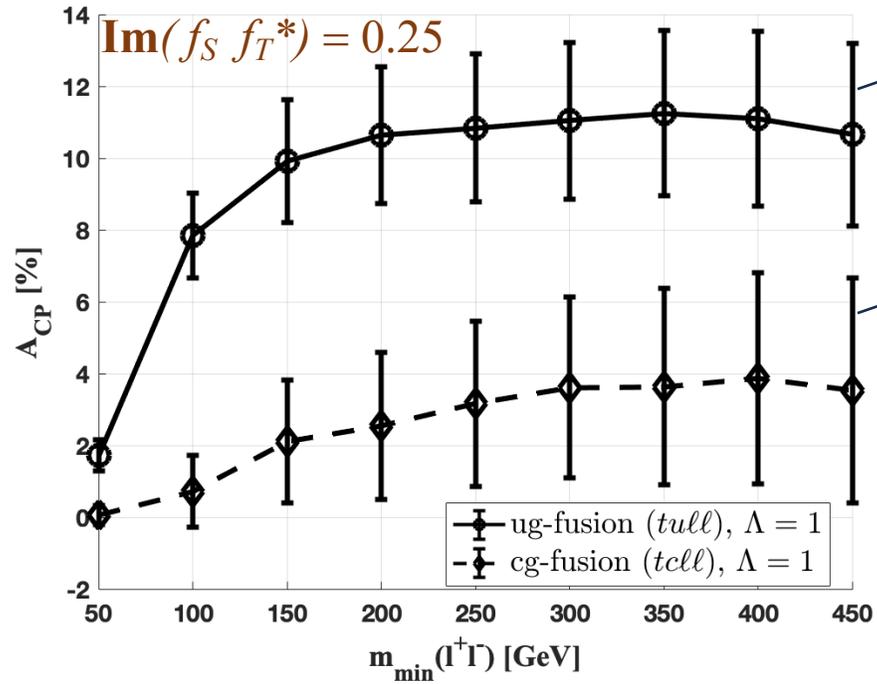
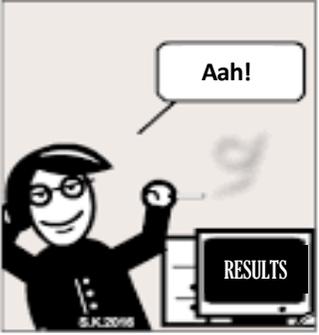
Selection:

$$m(\mu^{+} \mu^{-}) > m_{\min}(\mu^{+} \mu^{-})$$

... plus minimal acceptance and threshold restrictions (no detector modeling)

(No b-tagging, though it can suppress the main SM background.)

Results



Stat. uncertainty only ($\mathcal{L} = 1000 \text{ fb}^{-1}$)

Stat. uncertainty only ($\mathcal{L} = 3000 \text{ fb}^{-1}$)

	<i>ug</i> -fusion : $\Lambda = 1(2) \text{ TeV}$	<i>cg</i> -fusion : $\Lambda = 1(2) \text{ TeV}$
A_{CP}	11.1% (7.9)%	3.9% (0.7)%
A_T	16.4% (13.5)%	3.1% (0.5)%
\bar{A}_T	-5.8% (-2.3)%	-4.7% (-1.0)%

- $m_{\min}(\mu^+ \mu^-) = 400 \text{ GeV}$
- $\mathcal{L} = 1000 \text{ fb}^{-1}$
- $\Lambda = 1 \text{ TeV}$



~100 SM trilepton events
~10⁴ NP trilepton events

Summary



- Trilepton final states provide an $O(10\%)$ CP_v asymmetry
- A challenging probe of CP_v @ LHC
- Other T_N – odd observables can improve the asymmetry, e.g.

$$\mathcal{O}_{CP}^i = p_a^i \cdot (\vec{p}_b \times \vec{p}_c)^i$$

... but not dramatically.

