





#### Sneutrino dark matter

G. Bélanger

LAPTH Annecy-le-Vieux

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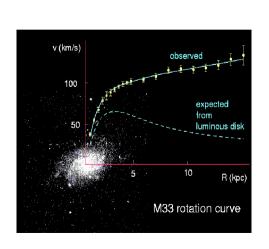
#### Introduction

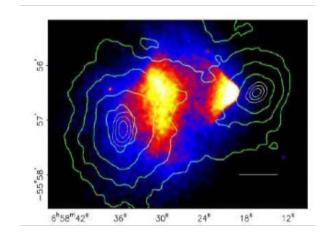
Strong evidence for dark matter from astrophysical and cosmological observations

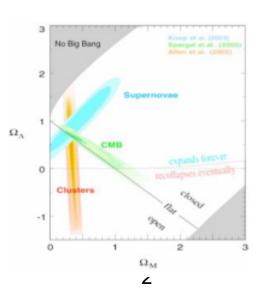
Motivation for new particles beyond standard model

Implication of precise determination of amount of CDM on DM particle properties

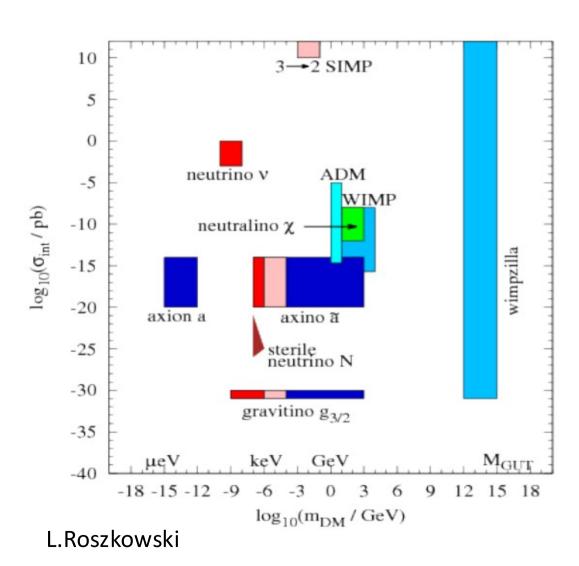
 $\Omega_{\rm cdm}$  h<sup>2</sup>=0.1196+/-0.0031







#### A wide variety of DM candidates



WIMPs

**FIMPs** 

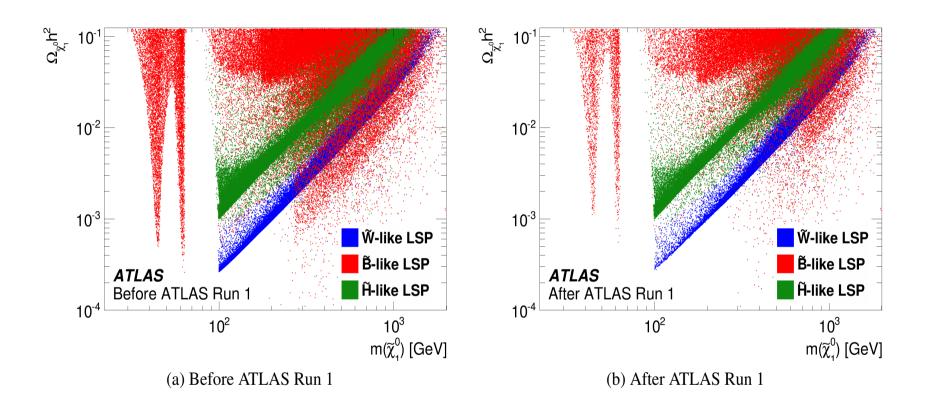
**SIMPs** 

Asymmetric

- Supersymmetry one of best motivated extension of SM
- No sign at LHC → does that mean that most popular WIMP model (neutralino) is ruled out?
- Strong constraints from LHC + direct detection especially if below TeV scale
- Properties of neutralino DM: strong dependence on its nature: partner of gauge boson (B,W) or Higgs
  - SU(2) number: efficient annihilation into WW-> relic density prefers TeV scale (higgsino) or 2TeV (wino)
  - U(1) only: bino need light sfermions LHC disfavoured
  - Mixed: satisfies relic density for any scale mixed binohiggsino strongly constrained from direct detection (binowino allowed)

#### What's left after LHC

ATLAS 1508.06608



Still large area of parameter space to be explored by LHC and (in)direct searches

What about other supersymmetry candidates?

# Sneutrino DM

- Another neutral particle in SUSY: the sneutrino
- Partner of LH neutrino NOT a good DM candidate
  - Very large contribution to direct detection through Z exchange (Falk,Olive, Srednicki, PLB354 (1995) 99)+ efficient annihilation
- Neutrino have masses RH neutrino + supersymmetric partner well-motivated if LSP then can be dark matter
- Thermalized?
  - Non-negligible L-R mixing Arkani-Hamed et al PRD61 (2001), Borzumati, Namura PRD64 (2002) 053002
  - New interactions Gauge: MSSM+U(1) (GB et al JCAP 1112:014) or scalar eg NMSSM (Cerdeno, Seto, JCAP0908:032)
  - Both cases are viable with respect to LHC constraints and feature new signatures leptons (same-sign, monoleptons) (Arina, Cabrera, 1311.6549, Arina et al, 1503.02960, GB et al, 1505.06243)

### Sneutrino DM

- Or not thermalized
  - abundance from decay of other particles 'next to lightest dark' particle which has long lifetime,
  - NLSP freeze-out as usual then decays to feebly interacting sneutrino

### MSSM+RH neutrino

- The framework : MSSM + three generations ( $v_R$  + sneutrinoR).
- Assume pure Dirac neutrino masses
- Superpotential  $W = y_{\nu} \hat{H}_{u} \cdot \hat{L} \hat{\nu}_{R}^{c} y_{e} \hat{H}_{d} \cdot \hat{L} \hat{\ell}_{R}^{c} + \mu_{H} \hat{H}_{d} \cdot \hat{H}_{u}$
- Couplings of sneutrino proportional to neutrino mass
- Lower bound on neutrino mass from fits to solar, atmospheric, accelerator neutrino data

$$|\Delta m^2| = 2.43 \pm 0.06 \times 10^{-3} \text{eV}^2 \rightarrow m_{\nu}^H > 0.049 \text{eV}$$

• For hierarchical neutrino masses

$$(y_{\nu}^{H}\sin\beta)_{\min} \simeq 2.8 \times 10^{-13}$$

• Upper limit on Yukawa couplings from cosmological bound – Planck temperature and polarisation data, lensing, supernovae, BAO

$$\sum_{i=1}^3 m_i < 0.23 \; \mathrm{eV} \; \mathrm{at} \; 95\% \; \mathrm{CL};$$
  $(y_{\nu}^H \sin \beta)_{\mathrm{max}} \simeq 4.4 \times 10^{-13}$  (for quasi-degenerate neutrinos)

#### MSSM+RH neutrino

• Sneutrino mass same order as other sfermions – can be LSP

$$-\mathcal{L}_{soft}\supset M_{ ilde{
u}_R}^2 | ilde{
u}_R|^2 + (y_
u A_
u H_u \, ilde{L} \, ilde{
u}_R^c + \, h.c.)$$

• Sneutrino mixing is very small – can be neglected

$$an 2 ilde{\Theta} = rac{2y_
u v \sineta |\coteta \mu - A_
u|}{m_{ ilde
u_L}^2 - m_{ ilde
u_R}^2}$$

- Assume mass of RH sneutrino is free parameter (even in sneu-CMSSM)
- Note that natural for sneutrinoR to be lightest particle as its mass does not evolve much with energy contrary to other sfermions.

- Sneutrino not thermalized in early universe its interactions are too weak
- One possibility for DM is production through decays of sparticles
- Consider the case where stau is the NLSP (here assume CMSSM relations, for general MSSM Heisig et al 1310.2825) neutralino NLSP no distinctive LHC signature
- Lifetime of stau (2 or 3-body decay) depends on mixing in sneutrino/stau sectors =- from a few seconds to  $10^{11}$ s.

$$\Gamma_{ ilde{ au}_1 
ightarrow ilde{
u}_R W} = rac{g^2 ilde{\Theta}^2}{32 \pi} |U_{L1}^{( ilde{ au}_1)}|^2 rac{m_{ ilde{ au}_1}^3}{m_W^2} \left[ 1 - rac{2 (m_{ ilde{
u}_R}^2 + m_W^2)}{m_{ ilde{ au}_1}^2} + rac{(m_{ ilde{
u}_R}^2 - m_W^2)^2}{m_{ ilde{ au}_1}^4} 
ight]^{3/2}$$

- Decay of NLSP (MSSM-LSP) after freeze-out
- Relic density obtained from that of the NLSP can be charged

$$\Omega^{ ext{fo}}_{ ilde{
u}_R} = rac{m_{ ilde{
u}_R}}{m_{ ext{MSSM-LSP}}} \, \Omega_{ ext{MSSM-LSP}}$$

## Model parameters and constraints

- CMSSM + RH neutrino
- Scan range

$$m_0 < 2500 \,\text{GeV}$$
;  $m_{1/2} < 2500 \,\text{GeV}$ ;  $|A_0| < 3000 \,\text{GeV}$ 

and at elevtroweak scale

$$0 < m_{\tilde{\nu}_R} < m_{\tilde{\tau}_1} \; ; \quad 5 < \tan \beta < 40$$

- M gluino>1.8TeV
- Collider constraints Higgs mass and couplings;
- Flavour constraints b-sγ, Bs-μμ, B-τυ;
- Susy searches (mostly not valid because stau is collider stable and charged);
- Charged stable stau m>340 GeV (from CMS Run 1 search)
- Constraints from BBN: lifetime of stau can be long enough for decay around or after BBN→ impact on abundance of light elements

# Big Bang Nucleosynthesis

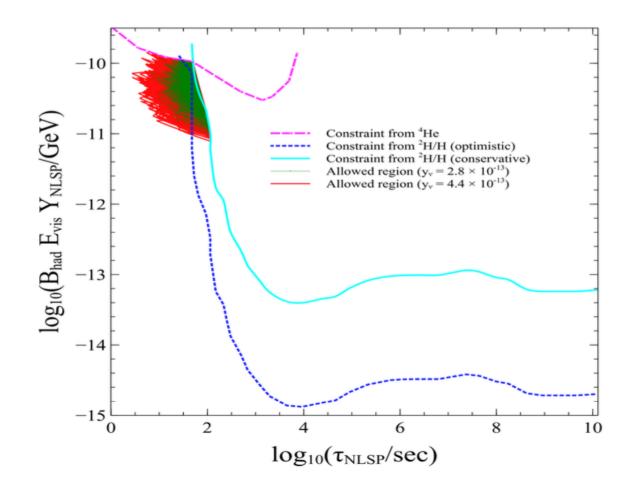
- BBN (T~MeV-10keV, t~0.1-10<sup>4</sup>s)allow to predict abundances of light elements  $D, He^3, He^4, Li$ .
- Depends on photon to baryon ratio
- In early Universe, energy density dominated by radiation
- At lower T : weak interactions fall out of equilibrium
- Freeze-out when interaction rate  $\Gamma_{\text{weak}} < H$ , species decouple
- When T approaches freeze-out (around 0.8MeV)

$$n/p \approx exp^{-\Delta m/T} \approx 1/6$$

- Nucleosynthesis begins with formation of Deuterium
- Number of photons>> number of nucleons the reverse process occurs much faster, deuterium production is delayed, starts only at T~0.1MeV  $p+n \rightarrow D+\gamma$
- ... and the chain continues with production of heavier elements
- Relationship between expansion rate of Universe (relate to total matter density) and density of p and n (baryonic matter density) determine abundance of light elements  $Y \approx \frac{2n/p}{1+n/p} \approx 0.25$
- Main product of BBN <sup>4</sup>He
- Other elements produced in lesser amounts D, <sup>3</sup>He, <sup>7</sup>Li

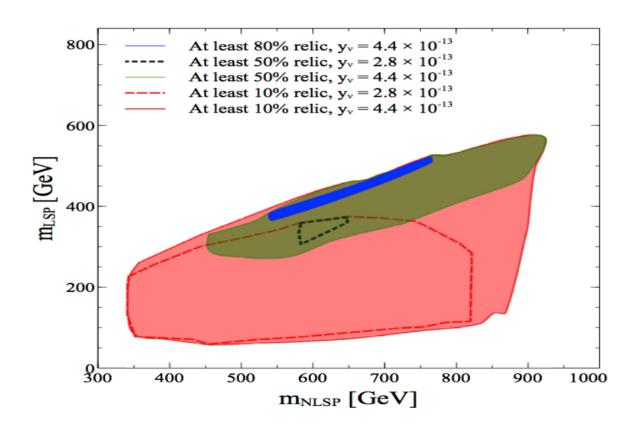
- If particle with lifetime > 0.1s decays can cause non-thermal nuclear reaction during or after BBN spoiling predictions in particular if new particle has hadronic decay modes
  - Kawasaki, Kohri, Moroi, PRD71, 083502 (2005)
- Alteration of n/p ratio for example .  $\pi^- + p \to \pi^0 + n$ 
  - -> overproduction He<sup>4</sup>
- Hadrodissociation of He<sup>4</sup> causes overproduction of D
  - $n+He^4 -> He^3+D, 2D+n, D+p+n$

- Key elements:
  - Bhad: hadronic BR of stau (nuR+W)
  - Evis : net energy carried away by hadrons
  - Ystau: yield



## Allowed region

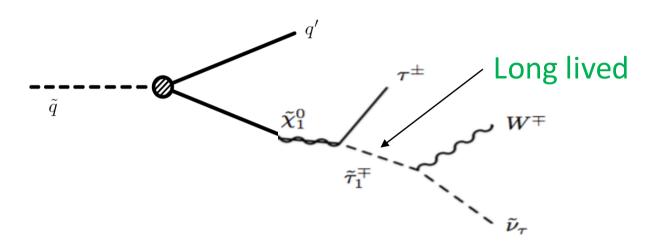
- After all constraints room for sneutrinoR DM (even in CMSSM)
- Can constitute dominant dark matter component



### LHC signatures

- Characteristic signature : stable charged particle NOT MET
- Staus live from sec to min: decay outside detector
- Searches
  - Cascades: coloured sparticles decay into jets + SUSY→ N
    jets + stau
  - Pair production of two stable staus
  - Passive search for stable particles
- Stable stau behaves like « slow » muons  $\beta=p/E<1$ 
  - Use ionisation properties and time of flight measurement to distinguish from muon
  - kinematic distribution

## Charged tracks from cascades



- Dominant contribution from squark pairs (heavy gluinos)
- Signal computed with Spheno+ Madgraph5aMC@NLO + Pythia+Delphes3+prospino k-factors
- Background: tt,µµ+jets, WW,WZ strongly suppressed with cuts
- Use approach suggested in Gupta et al PRD75075007 (2007)

# Charged tracks from cascades (2)

- $p_T^{\mu_{1,2}} > 200 \text{ GeV}, |y(\mu_{1,2})| < 2.4,$
- $p_T^{j_{1,2}} > 200 \text{ GeV}, |\eta(j_{1,2})| < 5.0,$
- $\sum |p_T^{vis.}| > 1000 \text{ GeV},$
- $\Delta R(\mu_1, \mu_2) > 0.2$ ,
- $\Delta R(j, j) > 0.4$ ,
- $\Delta R(\mu, j) > 0.4$ ,
- $M_{\mu_1,\mu_2} > 1000 \text{ GeV}$ ,

# Charged tracks from cascades (2)

Benchmark point	$\mathcal{L} \text{ for } 5\sigma \text{ [fb}^{-1}\text{]}$	$N_S$	$N_B$	$N_S/N_B$
$357~{ m GeV}$	9.1	25	0.35	72
$400~{ m GeV}$	2.5	25	0.09	265
$442~{ m GeV}$	68.5	27	2.7	10
$600~{ m GeV}$	1100	48	43	1.1

- Fairly easy to discover if mass stau < 400 GeV
- Luminosity 1ab<sup>-1</sup> can probe mass ~580GeV
- Dependence on mass of squarks

## Pair production

- No model dependence only mass of stau
- Smaller cross section (EW only)
- Background: muon pairs
- Best cuts close to current ATLAS analysis -JHEP1501 (2015) 068
- Lower reach than previous channel

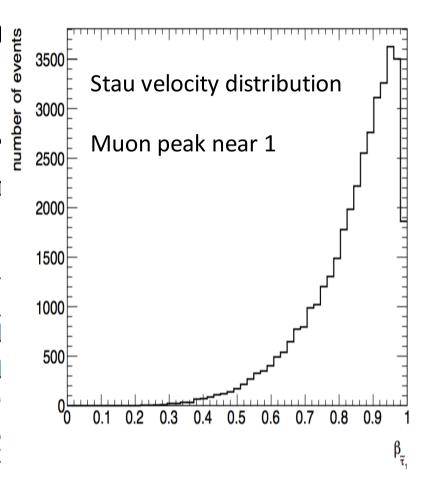
$$\mathcal{L} = 3000 \; \mathrm{fb^{-1}}$$

Cut	Benchmark	$N_S$	$N_B$	$N_S/N_B$	$\mathcal{S}$
$\Delta R(\mu\mu) > 0.4$	$357~{ m GeV}$	1543		0.44	21.8
$\beta < 0.95$	$400~{ m GeV}$	1014	3481	0.29	15.1
$p_T^{\mu_{1,2}} > 70 \text{GeV}$	$442~{ m GeV}$	715		0.21	11.0
$ y(\mu_{1,2})  < 2.5$	$600~{ m GeV}$	211		0.06	3.5

# Pair production

- No model dependence only mass of stau
- Smaller cross section (EW only)
- Background: muon pairs
- Best cuts close to current ATL
- Lower reach than previous chan

Cut	Benchmark	
$\Delta R(\mu\mu) > 0.4$	$357~{ m GeV}$	1
$\beta < 0.95$	$400~{ m GeV}$	1
$p_T^{\mu_{1,2}} > 70 \text{GeV}$	$442~{ m GeV}$	
$ y(\mu_{1,2})  < 2.5$	$600~{ m GeV}$	



#### MoEDAL detector

- Passive detector
- Array of nuclear track detector stacks
- Surrounds intersection region point 8
- Sensitive to highly ionising particles
- Does not require trigger, one detected event is enough
- Major condition: ionizing particle has velocity  $\beta$ <0.2

Benchmark point	Cascade	Pair	
$357~{ m GeV}$	45	2.5	
$400~{ m GeV}$	296	1.5	
$442~{ m GeV}$	24	1.1	
$600~{ m GeV}$	6	0.5	

CERN-LHO MISEDAL LHOB

B. Acharya et al, 1405.7662

Banerjee, et al, 1603.08834

Number of  $\tilde{\tau}_1$ 's with  $\beta \leq 0.2$  with  $\mathcal{L} = 3000 \, \mathrm{fb}^{-1}$ 

#### CONCLUSION

Sneutrino viable very weakly interacting DM candidate in supersymmetry

BBN constraints are important

LHC has unique potential to probe a whole class of DM models that predict heavy stable charged particles