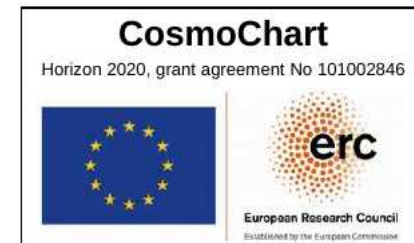


Dark matter bound states, scalar mediators and the Higgs

Kallia Petraki



Scalars 2023
Warsaw, 14/09/2023

Frontiers in particle dark matter searches

(very simplistic summary)

Past decades

Most research focused on

$$m_{\text{DM}} \sim 100 \text{ GeV} \sim m_{\text{W,Z}}$$

(e.g. prototypical
WIMP scenario)

Current frontiers

Heavy dark matter

$$m_{\text{DM}} \gtrsim \text{TeV}$$

Not constrained by colliders.

→ Experimentally probed by
existing / upcoming **telescopes**
e.g. HESS, IceCube, CTA, Antares

Light dark matter

$$m_{\text{DM}} \lesssim \text{few GeV}$$

Not constrained by older direct
detection experiments

→ Development of new generation
of **direct detection** experiments

Heavy ($m_{\text{DM}} \gtrsim \text{TeV}$) dark matter

How does the phenomenology of dark matter look like?
(in popular scenarios, e.g. thermal-relic DM)



New type of dynamics emerges:
Long-range interactions

$$\lambda_B \sim \frac{1}{\mu v_{\text{rel}}}, \quad \frac{1}{\mu \alpha} \lesssim \frac{1}{m_{\text{mediator}}} \sim \text{interaction range}$$

μ : reduced mass ($m_{\text{DM}}/2$)

Heavy ($m > \text{TeV}$)

Does this occur in models we care about?

- WIMPs with $m > \text{few TeV}$
 - WIMPs with $m < \text{TeV}$ co-annihilating with coloured/charged particles
 - Self-interacting DM
- not so heavy DM!*

type of interactions merges:
Long-range interactions

$$\lambda_B \sim \frac{1}{\mu v_{\text{rel}}}, \quad \frac{1}{\mu \alpha} \lesssim \frac{1}{m_{\text{mediator}}} \sim \text{interaction range}$$

μ : reduced mass ($m_{\text{DM}}/2$)

Heavy ($m > \text{TeV}$)

Does this occur in models we care about?

- WIMPs with $m > \text{few TeV}$
 - WIMPs with $m < \text{TeV}$ co-annihilating with coloured/charged particles
 - Self-interacting DM
- not so heavy DM!*

type of interactions merges:
Long-range interactions

$$\lambda_B \sim \frac{1}{\mu v_{\text{rel}}}, \quad \frac{1}{\mu \alpha} \lesssim \frac{1}{m_{\text{mediator}}} \sim \text{interaction range}$$

μ : reduced mass ($m_{\text{DM}}/2$)

What changes
when the interactions are long-ranged?

Sommerfeld

Distortion of scattering-state wavefunctions

⇒ **affects all cross-sections**

e.g. annihilation, elastic scattering

- Production in early universe, e.g. freeze-out
⇒ changes correlation of parameters (mass – couplings)
- Indirect detection signals
- Elastic scattering

Unstable bound states (positronium-like)

⇒ **extra annihilation channel**

• Production in early universe,
e.g. freeze-out **von Harling, Petraki 1407.7874**

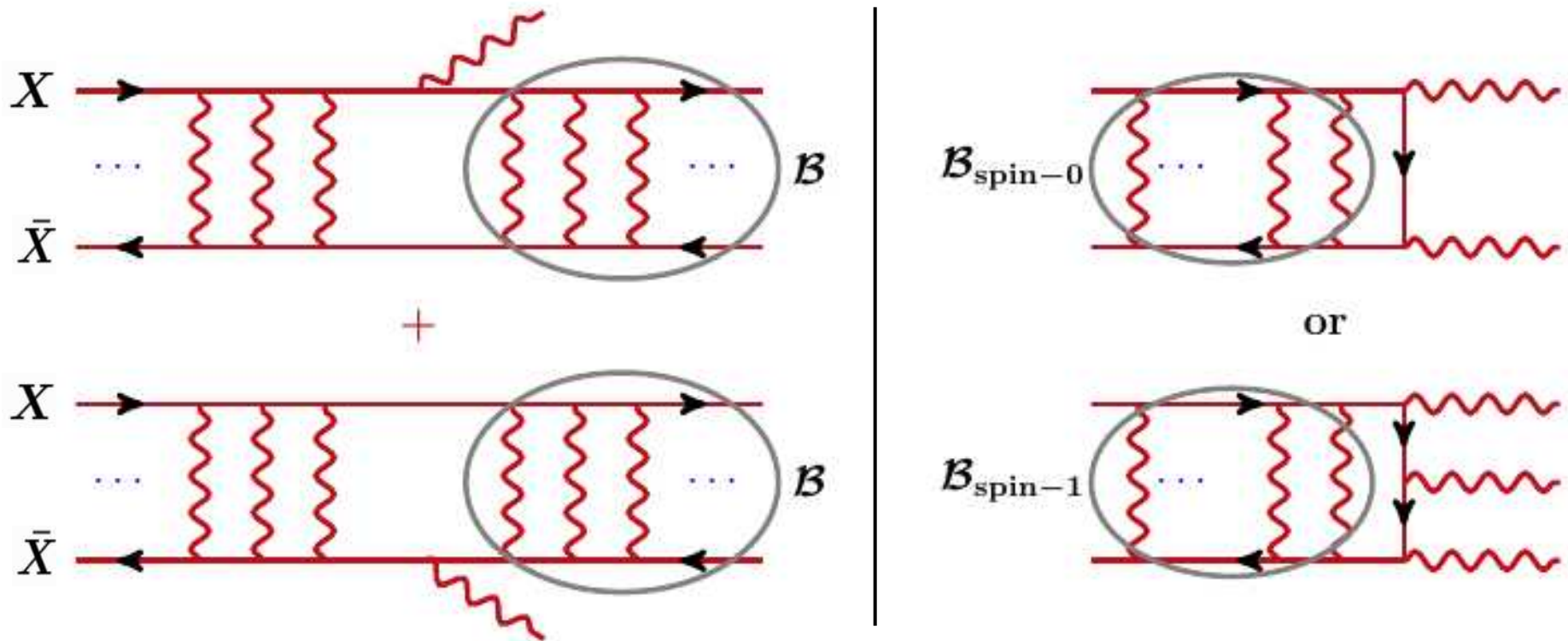
- Indirect detection
- Novel low-energy indirect detection signals
- Colliders

**Bound
states**

Stable bound states

- Elastic scattering (usually screening)
- Novel low-energy indirect detection signals
- Inelastic scattering in direct detection experiments (?)

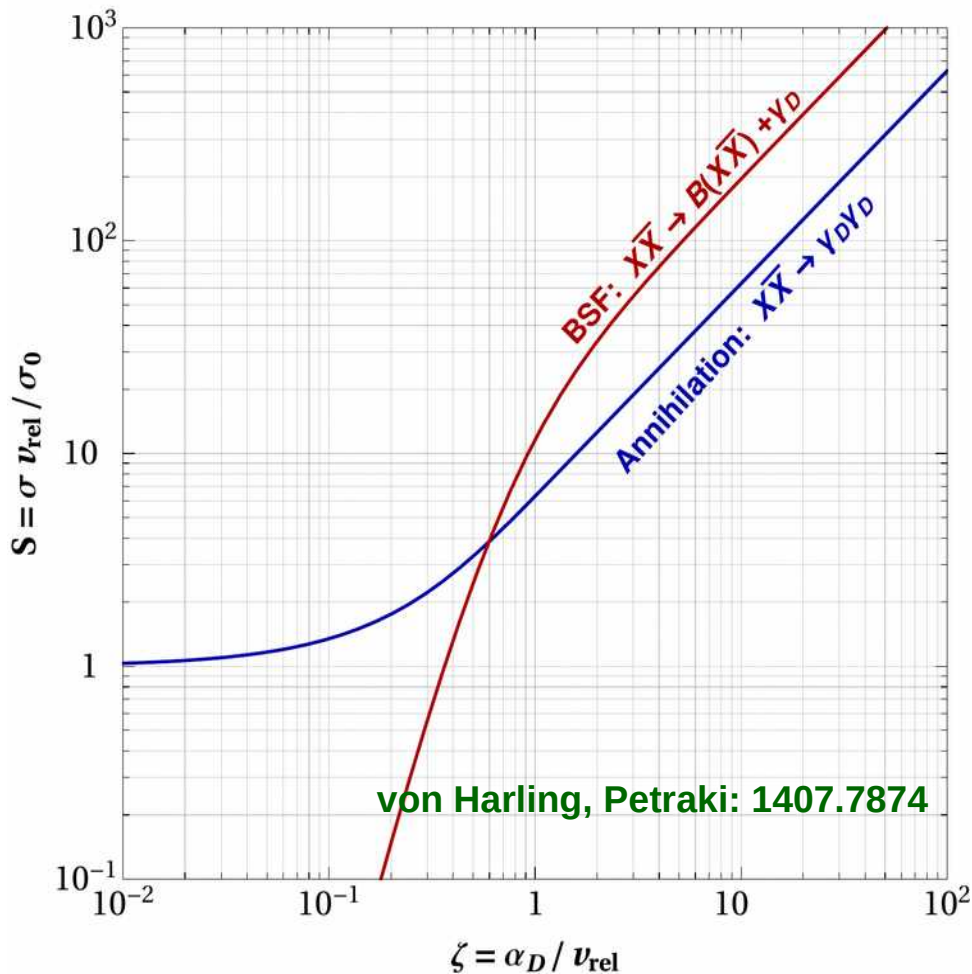
Radiative bound-state formation and decay



Opens a new annihilation channel !

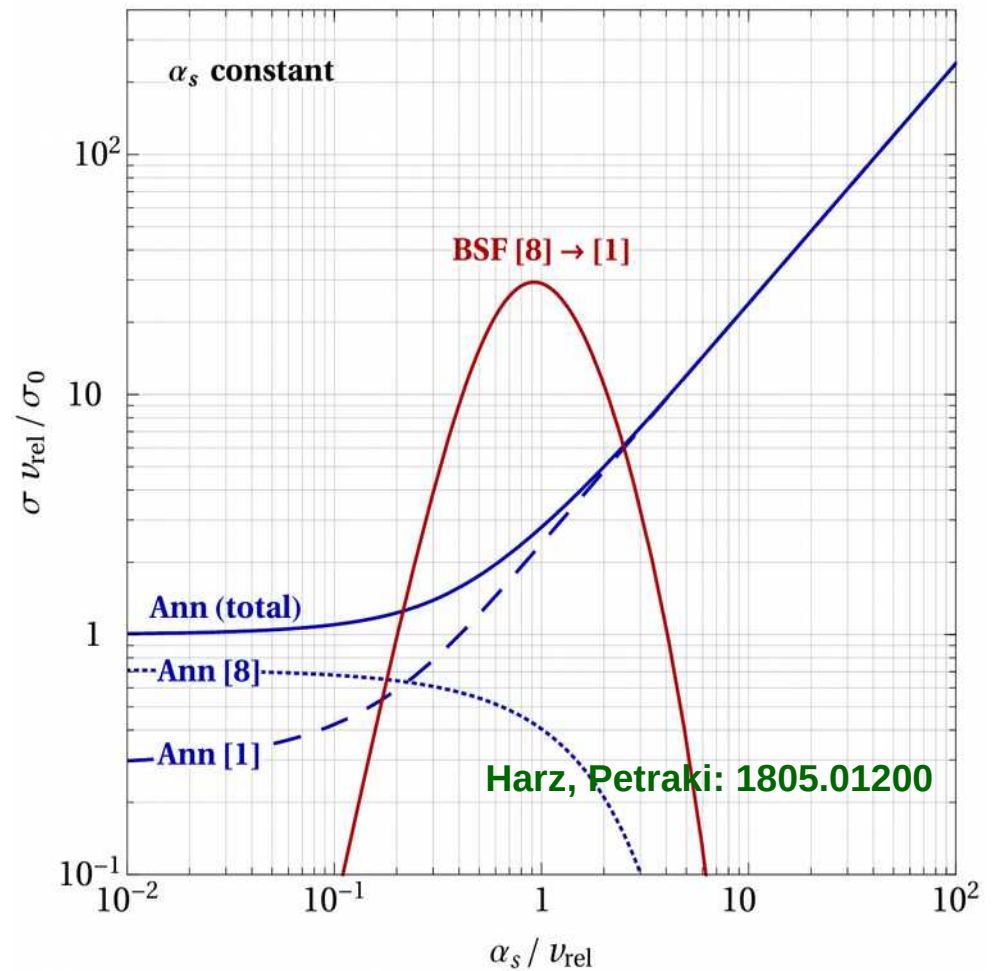
Annihilation vs Bound State Formation

U(1) model



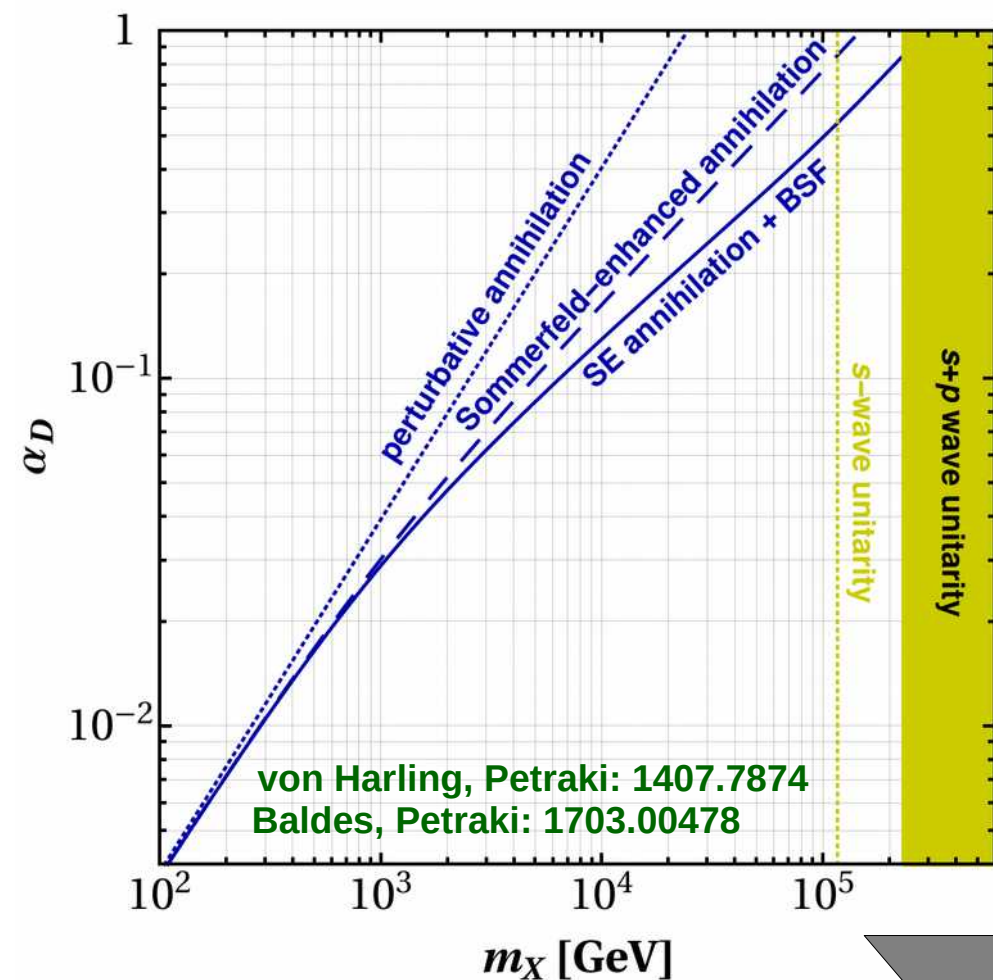
SU(3) model

e.g. neutralino-squark coann in SUSY



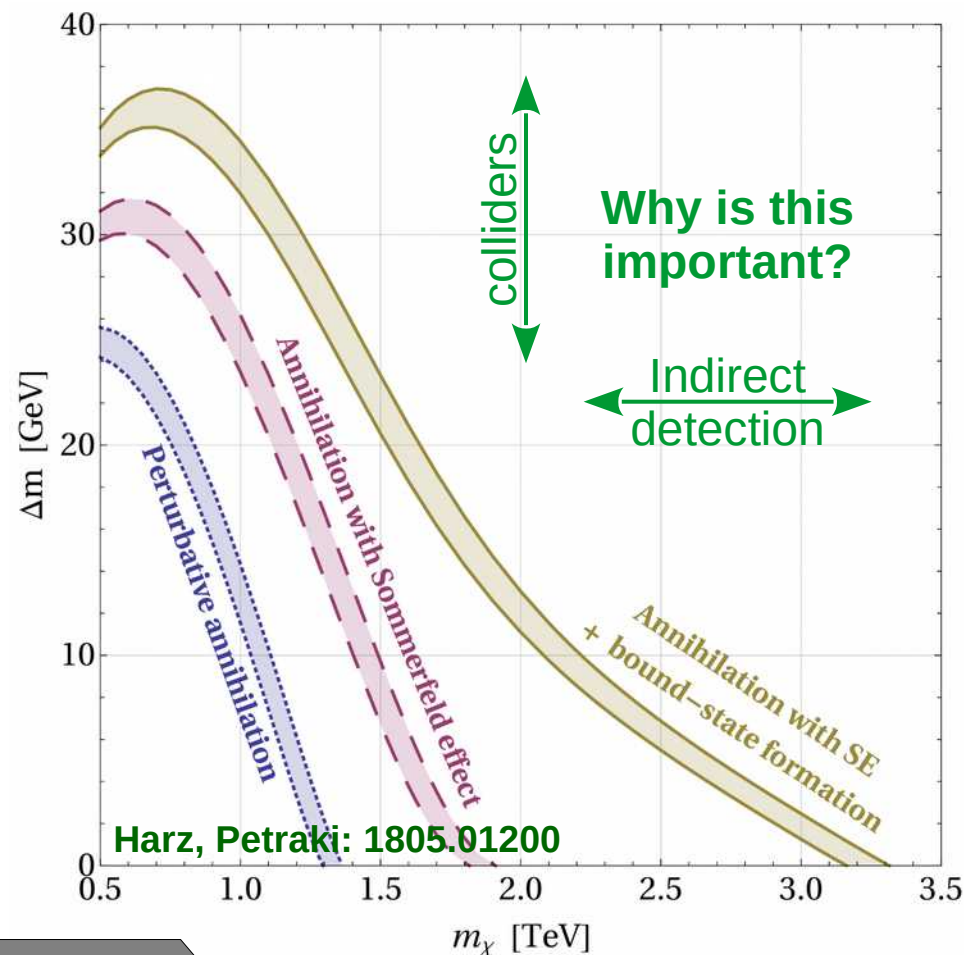
Annihilation vs Bound State Formation

U(1) model



SU(3) model

neutralino-squark coann in SUSY



Not the
final picture!

Scalars 2023

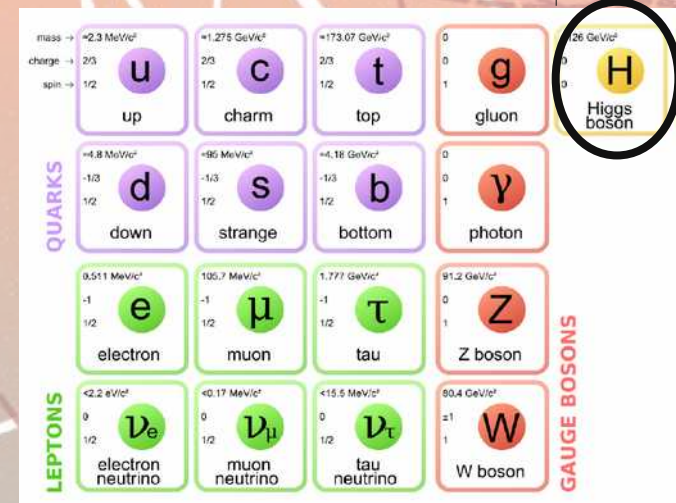
An opportunity to discuss various aspects of scalar particles.

13-16 September 2023
Warsaw (Ochota Campus)

scalars2023.fuw.edu.pl

What about light scalar mediators? And why do we care?

Hidden sector models,
e.g. self-interacting DM



A chart of the Standard Model of particle physics. It is organized into rows and columns. The first three rows are labeled 'QUARKS' on the left. The first two rows are labeled 'LEPTONS' on the left. The last two rows are labeled 'GAUGE BOSONS' on the right. Each particle is represented by a colored circle with its symbol and name. The Higgs boson (H) is circled in black. The chart includes mass, charge, and spin for each particle.

mass →	~2.3 MeV/c ²	~1.275 GeV/c ²	~173.07 GeV/c ²	0	126 GeV/c ²
charge →	2/3	2/3	2/3	0	0
spin →	1/2	1/2	1/2	1	0
	u up	c charm	t top	g gluon	H Higgs boson
	d down	s strange	b bottom	γ photon	
	e electron	μ muon	τ tau	Z Z boson	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	

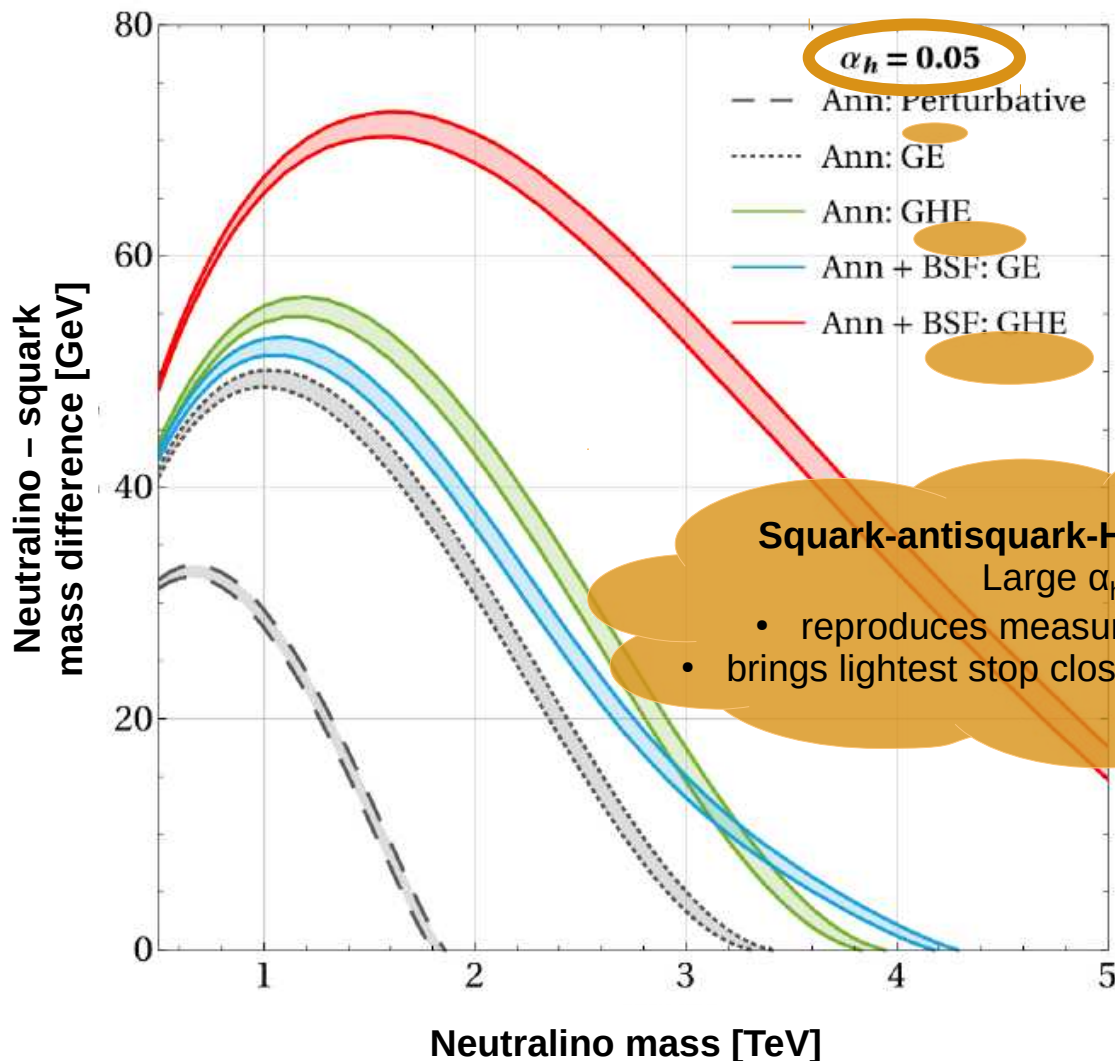
What do *light* force mediators do?

- Generate potential → Sommerfeld effect, bound states.
- Provide channel for transitions via on/off-shell emission. Spin of the emitted mediator determines selection rules.

In many realistic models, including WIMPs:
several mediators present

DM coannihilation with scalar colour triplet

The effect of the Higgs-mediated potential



$$2\text{PI} = \text{gluon} + \text{Higgs}$$

$$V(r) = -\frac{\alpha_g}{r} - \frac{\alpha_h}{r} e^{-m_h r}$$

Glue exchange Higgs exchange

Squark-antisquark-Higgs coupling

Large α_h

- reproduces measured Higgs mass
- brings lightest stop close in mass with LSP

Not the final picture!

Bound-state formation via emission of a light scalar

- Capture via emission of **neutral scalar** suppressed,
due to selection rules: **quadruple transitions**
March-Russel, West 0812.0559
Petraki, Postma, Wiechers: 1505.00109
An, Wise, Zhang: 1606.02305
Petraki, Postma, de Vries: 1611.01394
- Capture via emission of **charged scalar** [or its Goldstone mode]
extremely rapid: **monopole transitions !**
Ko, Matsui, Tang: 1910.04311
Oncala, Petraki: 1911.02605
Oncala, Petraki: 2101.08666
Oncala, Petraki: 2101.08667

Sudden change in effective Hamiltonian precipitates transitions.
Akin to atomic transitions precipitated by β decay of nucleus.

**Bound-state formation
via emission of a **charged** scalar**

Scalar DM $X, X^\dagger$ coupled to doubly charged light scalar mediator Φ

$$\mathcal{L} \supset -igX^\dagger V^\mu (\partial_\mu X) - i2g\Phi^\dagger V^\mu (\partial_\mu \Phi) - \frac{ym_X}{2} XX\Phi^\dagger + h.c.$$

$$m_X \gg m_\Phi$$

$$\begin{array}{c} X \rightarrow \quad \quad \rightarrow X \\ X^\dagger \leftarrow \quad \leftarrow X^\dagger \end{array} \quad \mathcal{A}_{XX^\dagger}^{2PI} = \begin{array}{c} X \rightarrow \quad \quad \rightarrow X \\ X^\dagger \leftarrow \quad \leftarrow X^\dagger \end{array} \quad V + \begin{array}{c} X \rightarrow \quad \quad \rightarrow X^\dagger \\ X^\dagger \leftarrow \quad \leftarrow X \end{array} \quad \Phi$$

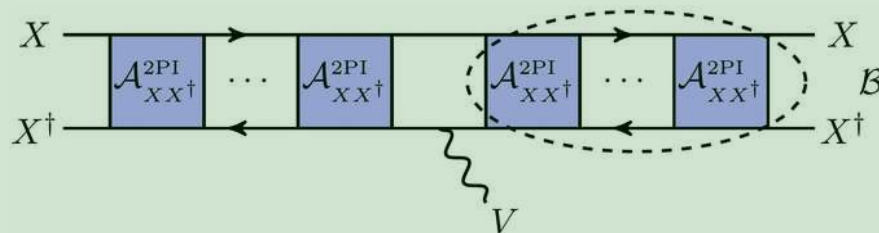
$$\begin{array}{c} X \rightarrow \quad \quad \rightarrow X \\ X \rightarrow \quad \quad \rightarrow X \end{array} \quad \mathcal{A}_{XX}^{2PI} = \begin{array}{c} X \rightarrow \quad \quad \rightarrow X \\ X \rightarrow \quad \quad \rightarrow X \end{array} \quad V$$

$$U_{XX^\dagger}(r) = -\frac{\alpha_V}{r} - (-1)^\ell \frac{\alpha_\Phi}{r} e^{-m_\Phi r}$$

$$U_{XX}(r) = +\frac{\alpha_V}{r}$$

Potential differs even in the global symmetry limit $\alpha_V \rightarrow 0$

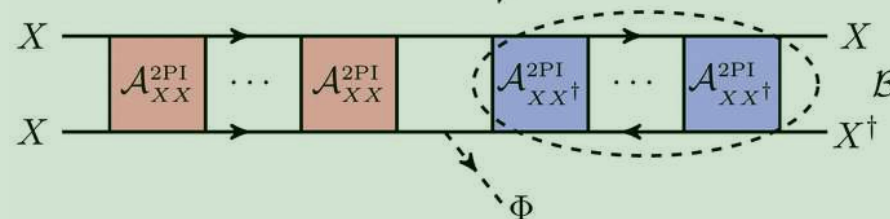
BSF_V



dipole transition

$$\mathcal{M} \sim 2g \int d^3p \psi_{n\ell m}^*(r) \nabla \phi_k(r) \sim 2g \int d^3p \psi_{n\ell m}^*(r) r \phi_k(r)$$

BSF_Φ

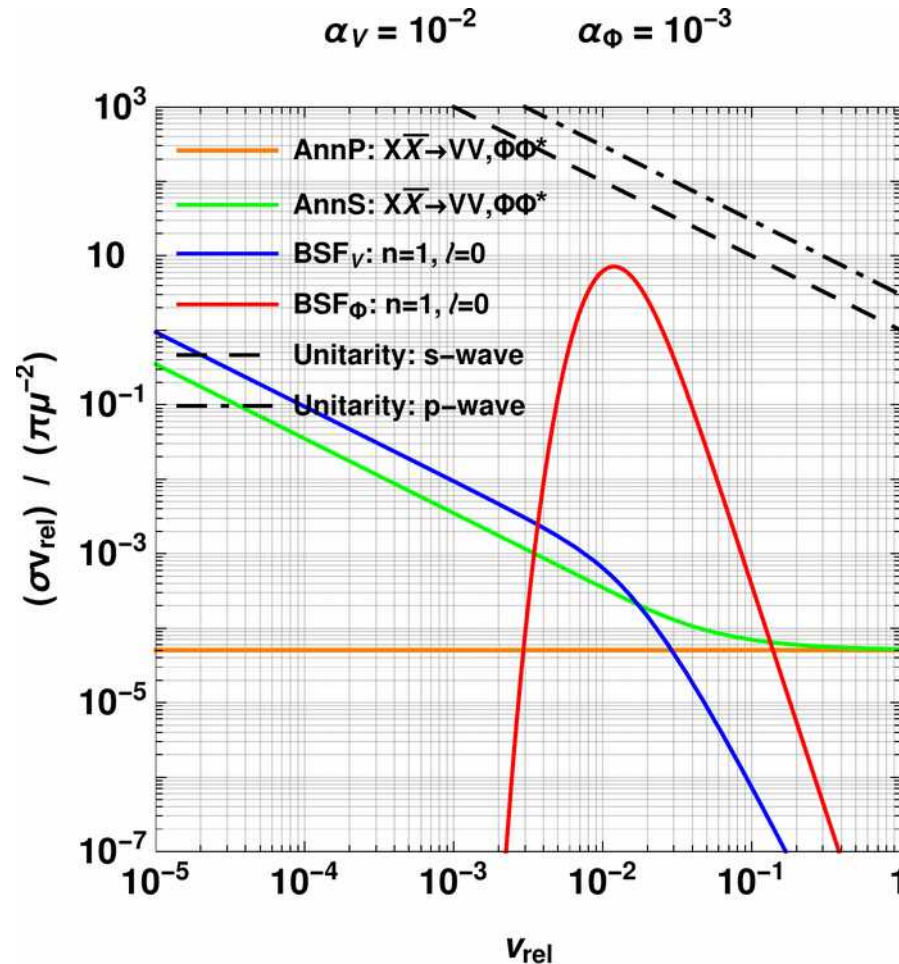


Change in effective Hamiltonian
⇒ monopole transition

$$\mathcal{M} \sim 2y \int d^3p \psi_{n\ell m}^*(r) \phi_k(r)$$

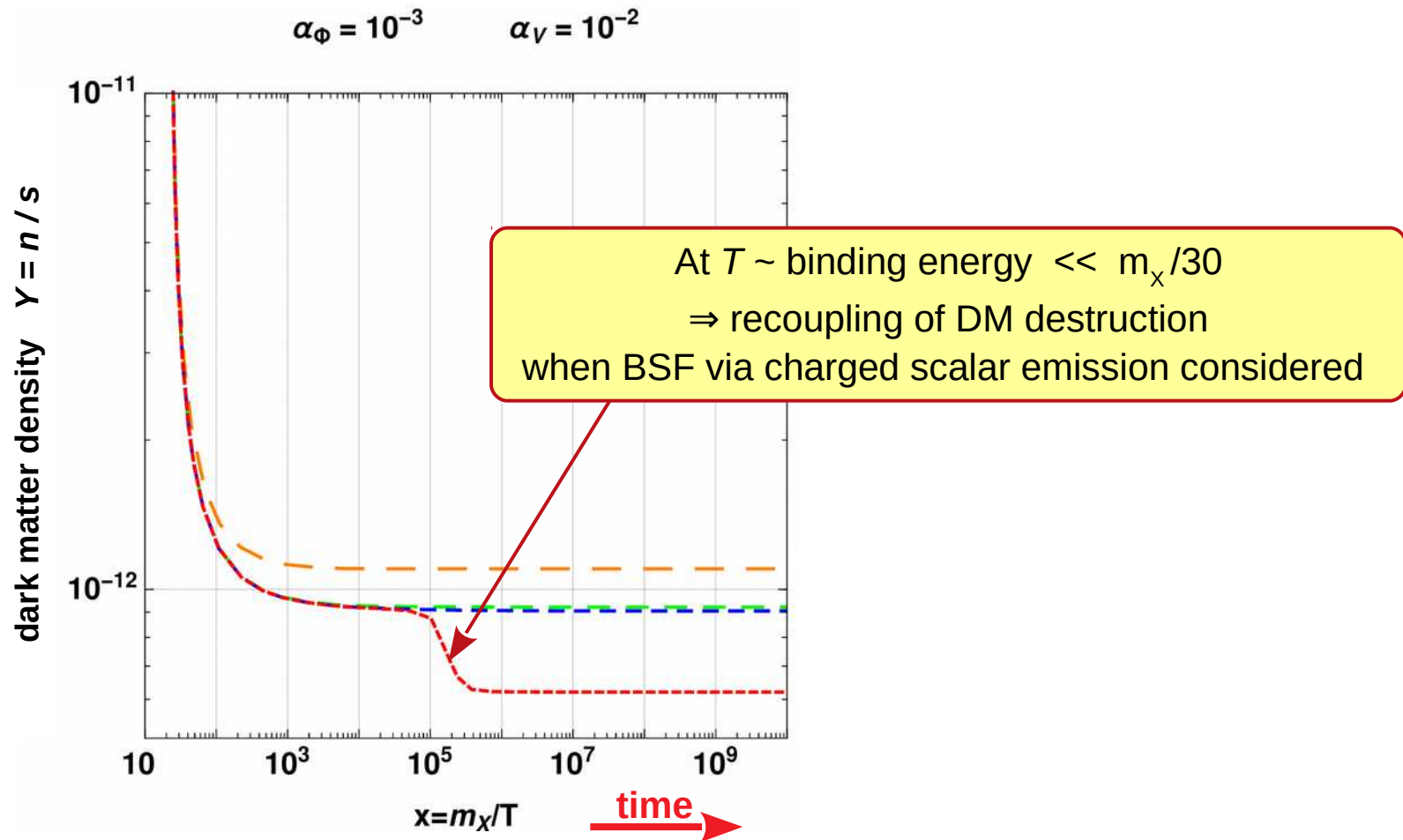
extremely fast!

Scalar DM $X, X^\dagger$ coupled to doubly charged light scalar mediator ϕ

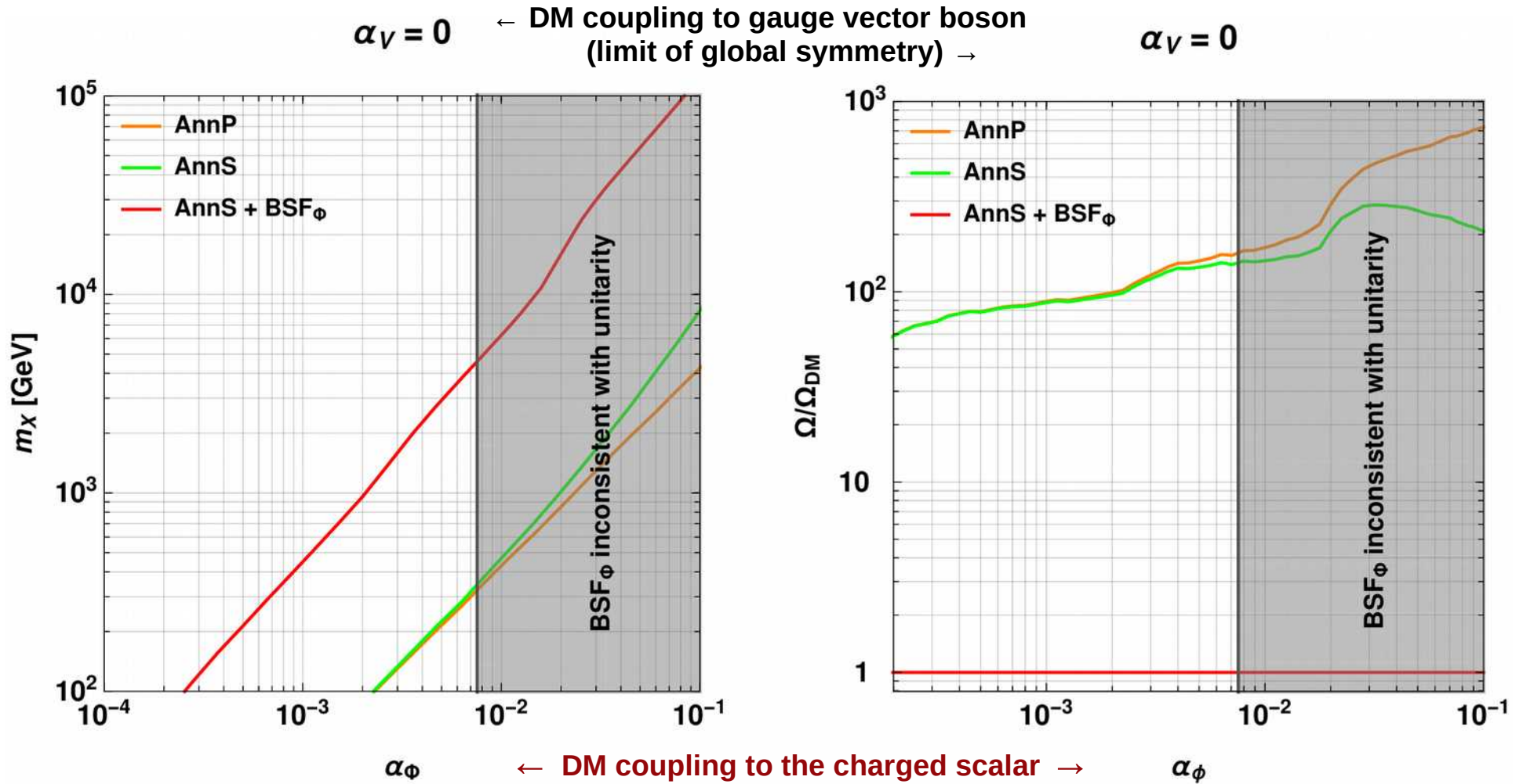


**BSF_ϕ very large,
even for small values
of α_ϕ, α_V !**

Scalar DM $X, X^\dagger$ coupled to doubly charged light scalar mediator ϕ

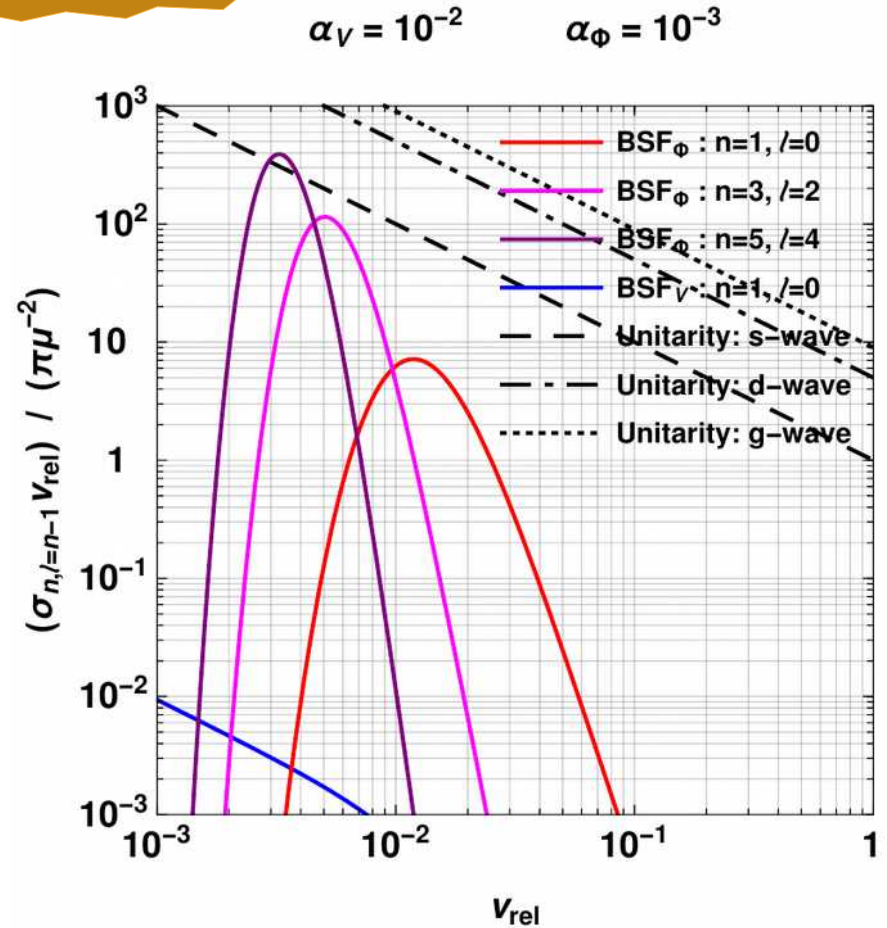
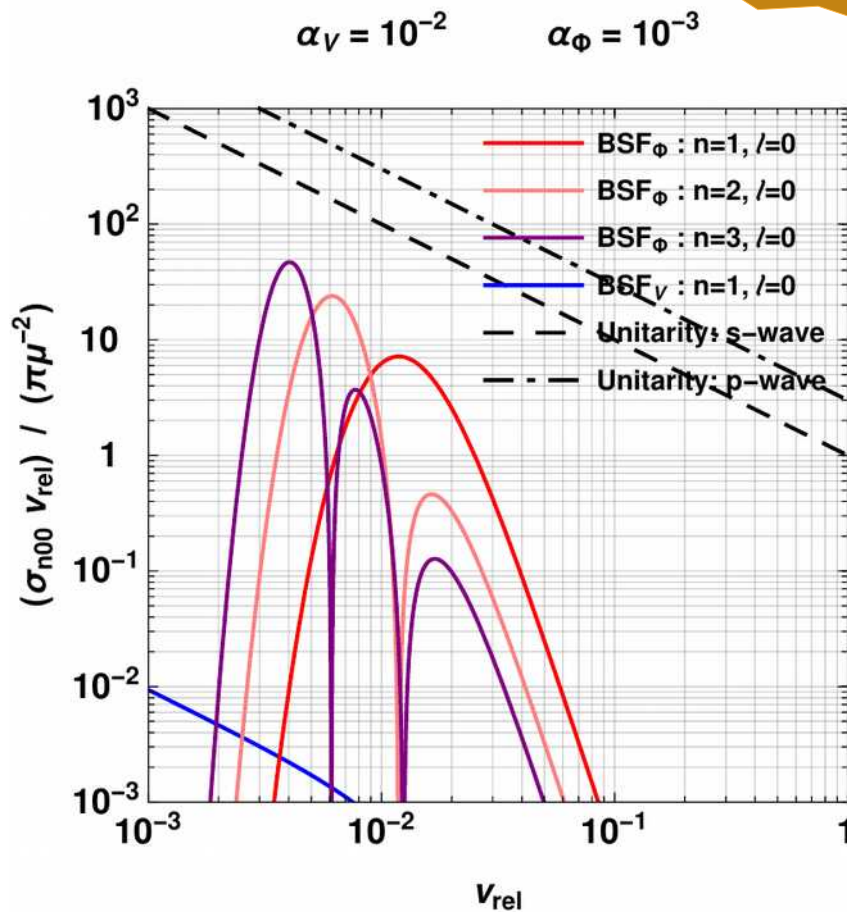


Scalar DM $X, X^\dagger$ coupled to doubly charged light scalar mediator ϕ



Scalar DM $X, X^\dagger$ coupled to doubly charged light scalar mediator ϕ

Capture into excited states important!



**Bound-state formation
via **Higgs doublet** emission**

Renormalisable WIMP models with coupling to the Higgs

In some of the archetypal WIMP DM models,
DM is the lightest linear combination of the neutral component of
SU(2) multiplets that couple to the Higgs

$$\delta\mathcal{L} \supset -y\bar{X}_n H X_{n+1} + \text{h.c.}$$

Trilinear
Higgs doublet-DM-DM,
independent of Higgs VEV

Includes many supersymmetric scenarios,
e.g. Wino-Higgsino, coloured co-annihilation

If $m > 5$ TeV, DM freeze-out begins before electroweak phase transition.

Doublet-Singlet coupled to the Higgs

$$\delta\mathcal{L} = \frac{1}{2}\bar{S}(i\not{\partial} - m_S)S + \bar{D}(i\not{\partial} - m_D)D - (y_L\bar{D}_L H S + y_R\bar{D}_R H S) + \text{h.c.}$$

for simplicity, set: $y_L = y_R \equiv y$, $m_S = m_D \equiv m$

field	$SU_L(2)$	$U_Y(1)$	$\mathbb{Z}_2$
S	1	0	-1
D	2	1/2	-1
H	2	1/2	+1

**Calculate all cross-sections and freeze-out
in the symmetric electroweak phase.**

Doublet-Singlet coupled to the Higgs

Non-relativistic potentials

$$\begin{array}{c} p_1 \\ \rightarrow \\ D_i^{s_1} \\ \leftarrow \\ \bar{D}_j^{s_2} \\ p_2 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} p_1' \\ \rightarrow \\ S^{s_1'} \\ \leftarrow \\ S^{s_2'} \\ p_2' \end{array} = \frac{1}{2} \left(\begin{array}{c} p_1 \\ \rightarrow \\ D_i^{s_1} \\ \leftarrow \\ \bar{D}_j^{s_2} \\ p_2 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} p_1' \\ \rightarrow \\ S^{s_1'} \\ \leftarrow \\ S^{s_2'} \\ p_2' \end{array} + \begin{array}{c} p_1 \\ \rightarrow \\ D_i^{s_1} \\ \leftarrow \\ \bar{D}_j^{s_2} \\ p_2 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} p_1' \\ \rightarrow \\ S^{s_2'} \\ \leftarrow \\ S^{s_1'} \\ p_1' \end{array} \right)$$

$$\begin{array}{c} p_1 \\ \rightarrow \\ D_i^{s_1} \\ \leftarrow \\ \bar{D}_j^{s_2} \\ p_2 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} p_1' \\ \rightarrow \\ D_{i'}^{s_1'} \\ \leftarrow \\ \bar{D}_{j'}^{s_2'} \\ p_2' \end{array} = \begin{array}{c} p_1 \\ \rightarrow \\ D_i^{s_1} \\ \leftarrow \\ \bar{D}_j^{s_2} \\ p_2 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} p_1' \\ \rightarrow \\ D_{i'}^{s_1'} \\ \leftarrow \\ \bar{D}_{j'}^{s_2'} \\ p_2' \end{array}$$

Need to project on

- gauge,
- spin,
- orbital angular momentum eigenstates

Neglect (T -dependent) H mass
in the potential
but consider it in
phase-space suppression
for BSF via H emission
($\rightarrow$ analytical cross-sections)

$$\begin{array}{c} p_1 \\ \rightarrow \\ D_i^{s_1} \\ \leftarrow \\ D_j^{s_2} \\ p_2 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} p_1' \\ \rightarrow \\ D_{i'}^{s_1'} \\ \leftarrow \\ D_{j'}^{s_2'} \\ p_2' \end{array} = \frac{1}{2} \left(\begin{array}{c} p_1 \\ \rightarrow \\ D_i^{s_1} \\ \leftarrow \\ D_j^{s_2} \\ p_2 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} p_1' \\ \rightarrow \\ D_{i'}^{s_1'} \\ \leftarrow \\ D_{j'}^{s_2'} \\ p_2' \end{array} + \begin{array}{c} p_1 \\ \rightarrow \\ D_i^{s_1} \\ \leftarrow \\ D_j^{s_2} \\ p_2 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} p_1' \\ \rightarrow \\ D_{j'}^{s_2'} \\ \leftarrow \\ D_{i'}^{s_1'} \\ p_1' \end{array} \right)$$

$$\begin{array}{c} p_1 \\ \rightarrow \\ D_i^{s_1} \\ \leftarrow \\ S^{s_2} \\ p_2 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} p_1' \\ \rightarrow \\ D_{i'}^{s_1'} \\ \leftarrow \\ S^{s_2'} \\ p_2' \end{array} = \begin{array}{c} p_1 \\ \rightarrow \\ D_i^{s_1} \\ \leftarrow \\ S^{s_2} \\ p_2 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} p_1' \\ \rightarrow \\ S^{s_2'} \\ \leftarrow \\ D_{i'}^{s_1'} \\ p_1' \end{array}$$

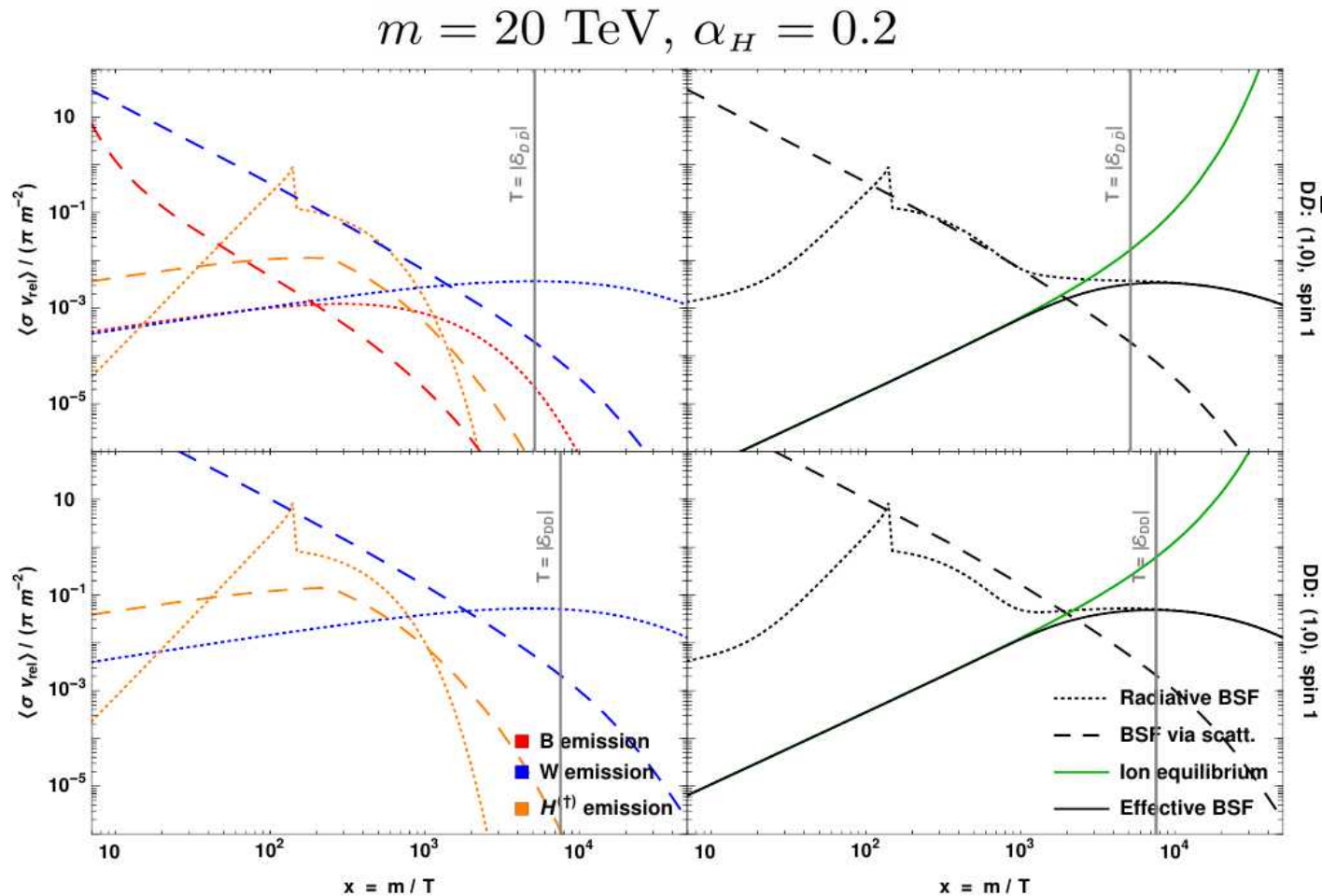
Doublet-Singlet coupled to the Higgs

Bound-state species ($n=0$, $\ell=0$)

Bound state ($\mathcal{B}$)	$U_Y(1)$	$SU_L(2)$	Spin	dof (g_B)	Bohr momentum (κ_B)	Decay rate (Γ_B)
$SS/DD\bar{D}$	0	$\mathbb{1}$	0	1	$\frac{m\alpha_A}{2}$	$\frac{m\alpha_A^3(\alpha_1^2 + 3\alpha_2^2)}{16} \left(\frac{\alpha_A}{\alpha_A + \alpha_R} \right)$
$D\bar{D}$	0	$\mathbb{1}$	1	3	$\frac{m(\alpha_1 + 3\alpha_2)}{8}$	$\frac{m(\alpha_1 + 3\alpha_2)^3[(\alpha_1 + 2\alpha_H)^2 + 40\alpha_1^2]}{2^{11} \cdot 3}$
DD	1	$\mathbb{1}$	1	3	$\frac{m(-\alpha_1 + 3\alpha_2)}{8}$	$\frac{m(-\alpha_1 + 3\alpha_2)^3\alpha_H^2}{2^7 \cdot 3}$
DS	1/2	2	0	2	$\frac{m\alpha_H}{2}$	Transition to $\mathcal{B}(SS/DD\bar{D})$:

Doublet-Singlet coupled to the Higgs

Bound-state formation cross-sections

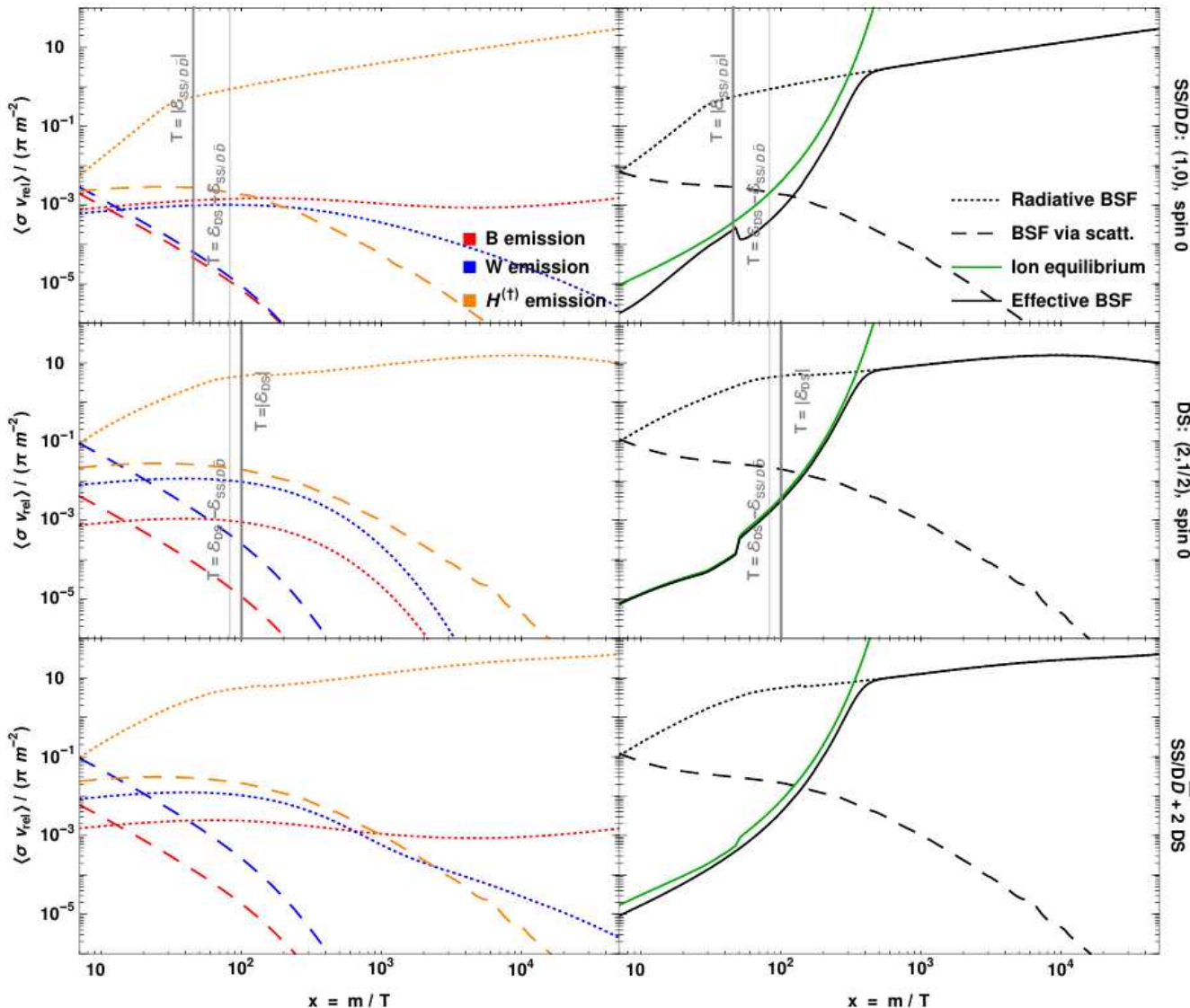


- BSF via scattering on plasma has no effect
- Departure from ionisation equilibrium at $T > \text{binding energy}$

Doublet-Singlet coupled to the Higgs

Bound-state formation cross-sections

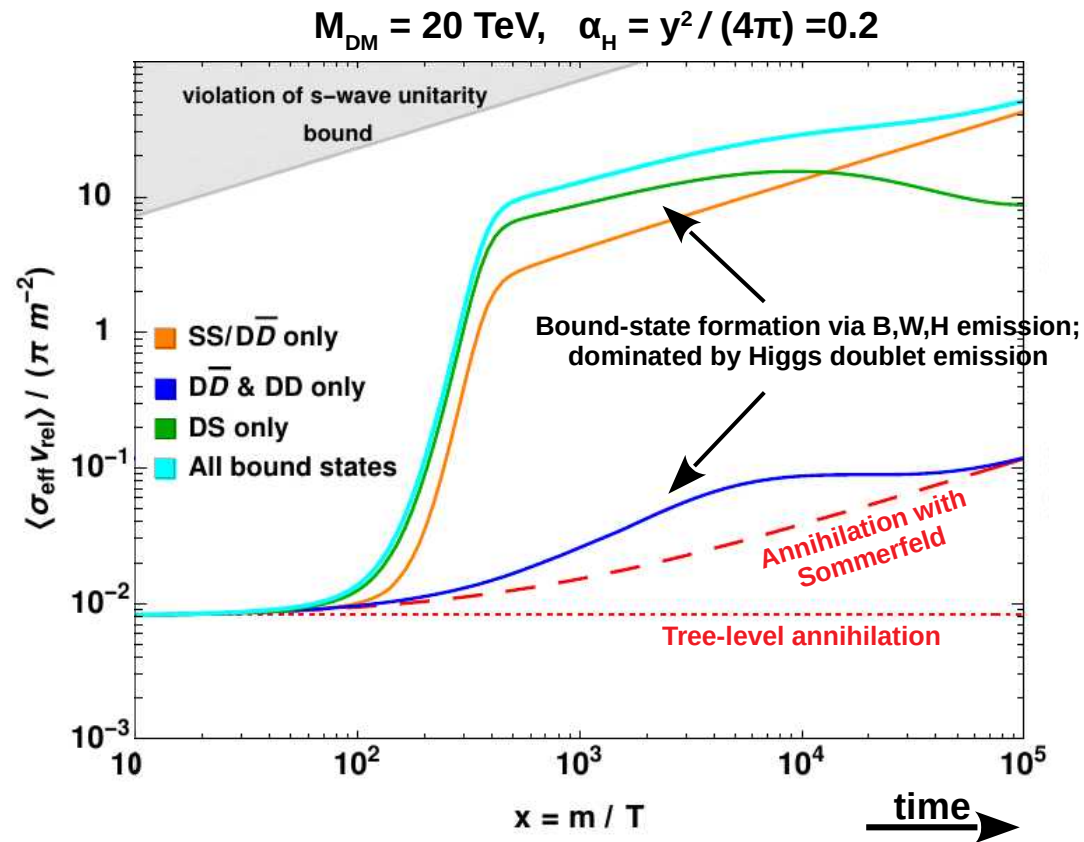
$$m = 20 \text{ TeV}, \alpha_H = 0.2$$



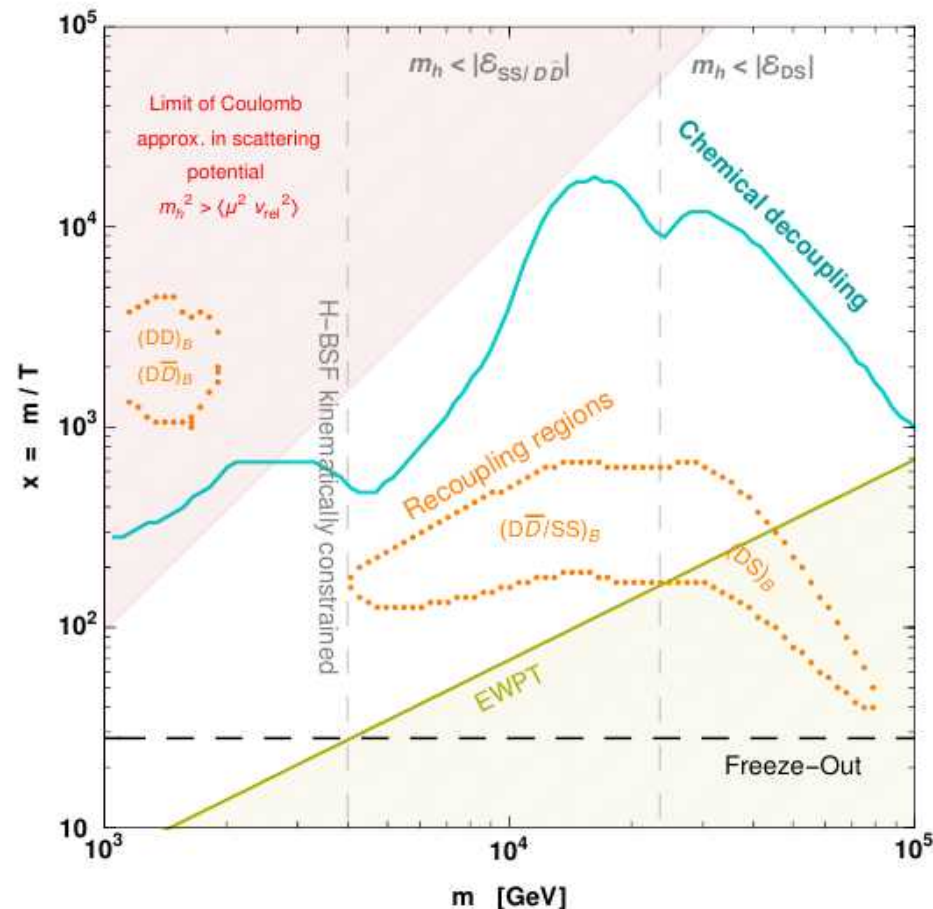
- **BSF via scattering on plasma has no effect**
(scattering kind of important for bound-to-bound transitions)
- **Departure from ionisation equilibrium at $T \ll$ binding energy**
(because of largeness of monopole BSF cross-sections)

Doublet-Singlet coupled to the Higgs

Effective cross-sections

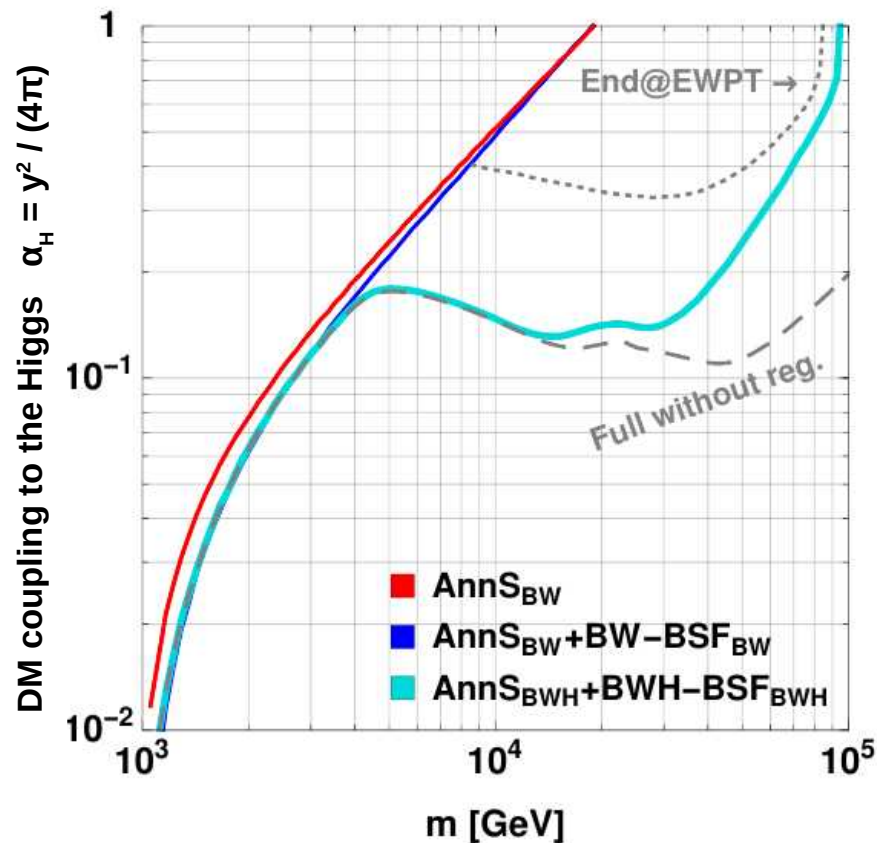


Doublet-Singlet coupled to the Higgs Timeline



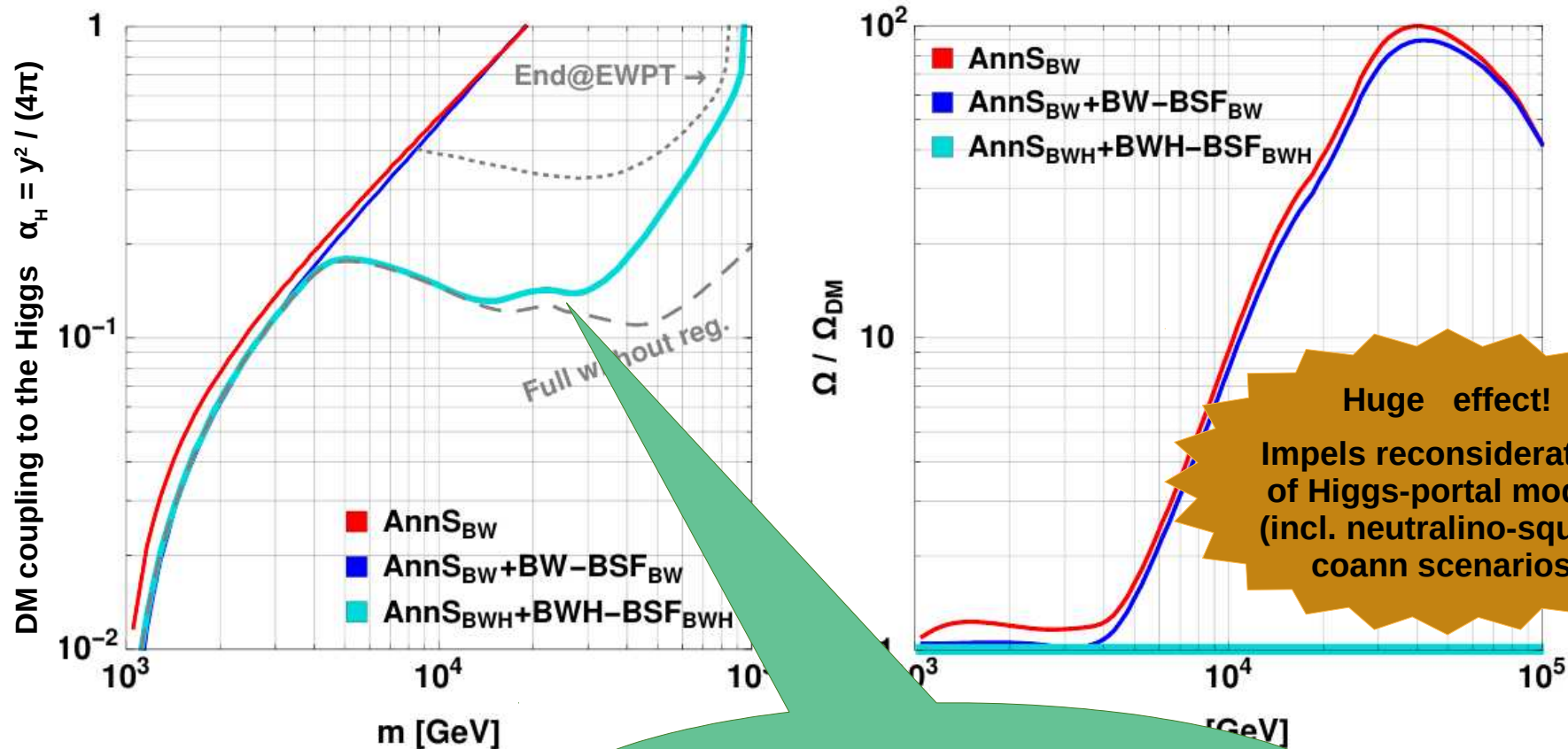
Doublet-Singlet coupled to the Higgs

Relic density



Doublet-Singlet coupled to the Higgs

Relic density



Effect only at large α_H , due to phase-space suppression in Higgs emission.

For higher multiplets, important effect expected also at lower α_H .

Conclusions

- **Bound states impel complete reconsideration of thermal decoupling at / above the TeV scale: *emergence of a new type of inelasticity*** von Harling, KP: 1407.7874

Model-independent unitarity arguments Baldes, KP: 1703.00478

- **Role of the Higgs doublet or other scalars can be extremely important.**

Ko, Matsui, Tang: 1910.04311; Oncalá, KP: 1911.02605, 2101.08666, 2101.08667

- **Experimental implications:**

- **DM heavier than anticipated:** multi-TeV probes very important
⇒ build the 100 TeV collider :)

- **Indirect detection:**

Enhanced rates due to BSF. Cirelli, Panci, KP, Sala, Taoso: 1612.07295

Novel signals: low-energy radiation emitted in BSF Baldes, Calore, KP, Poireau, Rodd: 2007.13787

Indirect detection of asymmetric DM Baldes et al.: 1703.0478, 1712.07489

- **Colliders:** improved detection prospects due increased mass gap in coannihilation scenarios

- **Effects *not* limited to freeze-out scenario:**

freeze-in, asymmetric DM, self-interacting DM, stable bound states